

Beyond the Standard Model: Where do we go from here?



Marie-Helene Genest,
Howard E. Haber, and
James Olsen
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The Galileo Galilei Institute For Theoretical Physics

Centro Nazionale di Studi Avanzati dell'Istituto Nazionale di Fisica Nucleare

Arcetri, Firenze



A Theorist's perspective

1. Has the idea of naturalness run its course?



run *its* course

★ to develop and finish **naturally**:

The doctor's advice is to let the fever run its course.

I had to accept that the relationship had run its course.

— Thesaurus: synonyms and related words

Based on an image from the BackReaction Blog
of [Sabine Hossenfelder](#)

1939: Scalar fields portend an energy scale associated with new phenomena that are close at hand.

JULY 1, 1939

PHYSICAL REVIEW

VOLUME 56

On the Self-Energy and the Electromagnetic Field of the Electron

V. F. WEISSKOPF

University of Rochester, Rochester, New York

(Received April 12, 1939)

The charge distribution, the electromagnetic field and the self-energy of an electron are investigated. It is found that, as a result of Dirac's positron theory, the charge and the magnetic dipole of the electron are extended over a finite region; the contributions of the spin and of the fluctuations of the radiation field to the self-energy are analyzed, and the reasons that the self-energy is only

logarithmically infinite in positron theory are given. It is proved that the latter result holds to every approximation in an expansion of the self-energy in powers of e^2/hc . The self-energy of charged particles obeying Bose statistics is found to be quadratically divergent. Some evidence is given that the "critical length" of positron theory is as small as $h/(mc) \cdot \exp(-hc/e^2)$.

The situation is, however, entirely different for a particle with Bose statistics. Even the Coulombian part of the self-energy diverges to a first approximation as $W_{st} \sim e^2 h / (mca^2)$ and requires a much larger critical length that is $a = (hc/e^2)^{-\frac{1}{2}} \cdot h / (mc)$, to keep it of the order of magnitude of mc^2 . This may indicate that a theory of particles obeying Bose statistics must involve new features at this critical length, or at energies corresponding to this length; whereas a theory of particles obeying the exclusion principle is probably consistent down to much smaller lengths or up to much higher energies.

Weisskopf's arguments imply that there should be new physics at the scale of $m_H/g \sim 1$ TeV. But where is the new TeV-scale physics?

ATLAS SUSY Searches* - 95% CL Lower Limits

July 2018

ATLAS Preliminary

$\sqrt{s} = 7, 8, 13$ TeV

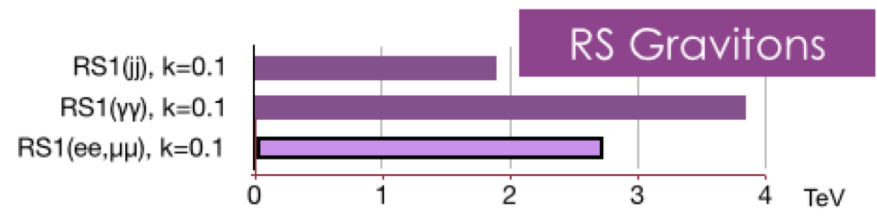
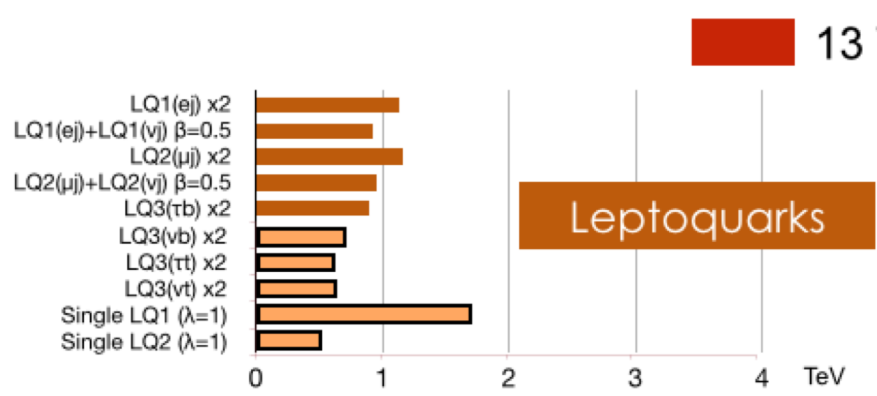
Model	e, μ, τ, γ	Jets	E_T^{miss}	$\int \mathcal{L} dt [\text{fb}^{-1}]$	Mass limit	$\sqrt{s} = 7, 8$ TeV	$\sqrt{s} = 13$ TeV	Reference	
Inclusive Searches	$\tilde{q}\tilde{q}, \tilde{q} \rightarrow q\tilde{\chi}_1^0$	0 mono-jet	2-6 jets 1-3 jets	Yes Yes	36.1 36.1	\tilde{q} [2x, 8x Degen.] \tilde{q} [1x, 8x Degen.]	0.9 0.43 0.71 1.55	$m(\tilde{\chi}_1^0) < 100$ GeV $m(\tilde{q}) - m(\tilde{\chi}_1^0) = 5$ GeV	1712.02332 1711.03301
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}\tilde{\chi}_1^0$	0	2-6 jets	Yes	36.1	\tilde{g} \tilde{g}	2.0 Forbidden 0.95-1.6	$m(\tilde{\chi}_1^0) < 200$ GeV $m(\tilde{\chi}_1^0) = 900$ GeV	1712.02332 1712.02332
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}(\ell\ell)\tilde{\chi}_1^0$	3 e, μ $ee, \mu\mu$	4 jets 2 jets	- Yes	36.1 36.1	\tilde{g} \tilde{g}	1.85 1.2	$m(\tilde{\chi}_1^0) < 800$ GeV $m(\tilde{g}) - m(\tilde{\chi}_1^0) = 50$ GeV	1706.03731 1805.11381
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow qqWZ\tilde{\chi}_1^0$	0	7-11 jets	Yes	36.1	\tilde{g}	1.8	$m(\tilde{\chi}_1^0) < 400$ GeV	1708.02794
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow qqWZ\tilde{\chi}_1^0$	3 e, μ	4 jets	-	36.1	\tilde{g}	0.98	$m(\tilde{g}) - m(\tilde{\chi}_1^0) = 200$ GeV	1706.03731
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow t\tilde{t}\tilde{\chi}_1^0$	0-1 e, μ 3 e, μ	3 b 4 jets	Yes -	36.1 36.1	\tilde{g} \tilde{g}	2.0 1.25	$m(\tilde{\chi}_1^0) < 200$ GeV $m(\tilde{g}) - m(\tilde{\chi}_1^0) = 300$ GeV	1711.01901 1706.03731
3 rd gen. squarks direct production	$\tilde{b}_1\tilde{b}_1, \tilde{b}_1 \rightarrow b\tilde{\chi}_1^0/\tilde{\chi}_1^\pm$		Multiple Multiple Multiple		36.1 36.1 36.1	\tilde{b}_1 Forbidden \tilde{b}_1 Forbidden \tilde{b}_1 Forbidden	0.9 0.58-0.82 0.7	$m(\tilde{\chi}_1^0) = 300$ GeV, BR($b\tilde{\chi}_1^0$) = 1 $m(\tilde{\chi}_1^0) = 300$ GeV, BR($b\tilde{\chi}_1^0$) = BR($\tilde{\chi}_1^\pm$) = 0.5 $m(\tilde{\chi}_1^0) = 200$ GeV, $m(\tilde{\chi}_1^\pm) = 300$ GeV, BR($b\tilde{\chi}_1^\pm$) = 1	1708.09266, 1711.03301 1708.09266 1706.03731
	$\tilde{b}_1\tilde{b}_1, \tilde{t}_1\tilde{t}_1, M_2 = 2 \times M_1$		Multiple Multiple		36.1 36.1	\tilde{t}_1 \tilde{t}_1 Forbidden	0.7 0.9	$m(\tilde{\chi}_1^0) = 60$ GeV $m(\tilde{\chi}_1^0) = 200$ GeV	1709.04183, 1711.11520, 1708.03247 1709.04183, 1711.11520, 1708.03247
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow Wb\tilde{\chi}_1^0$ or $\tilde{\chi}_1^0$	0-2 e, μ	0-2 jets/1-2 b	Yes	36.1	\tilde{t}_1	1.0	$m(\tilde{\chi}_1^0) = 1$ GeV	1506.08616, 1709.04183, 1711.11520
	$\tilde{t}_1\tilde{t}_1, \tilde{H}$ LSP				36.1 36.1	\tilde{t}_1 \tilde{t}_1 Forbidden	0.4-0.9 0.6-0.8	$m(\tilde{\chi}_1^0) = 150$ GeV, $m(\tilde{\chi}_1^\pm) - m(\tilde{\chi}_1^0) = 5$ GeV, $\tilde{t}_1 \approx \tilde{t}_L$ $m(\tilde{\chi}_1^0) = 300$ GeV, $m(\tilde{\chi}_1^\pm) - m(\tilde{\chi}_1^0) = 5$ GeV, $\tilde{t}_1 \approx \tilde{t}_L$	1709.04183, 1711.11520 1709.04183, 1711.11520
	$\tilde{t}_1\tilde{t}_1, \text{Well-Tempered LSP}$				36.1	\tilde{t}_1	0.48-0.84	$m(\tilde{\chi}_1^0) = 150$ GeV, $m(\tilde{\chi}_1^\pm) - m(\tilde{\chi}_1^0) = 5$ GeV, $\tilde{t}_1 \approx \tilde{t}_L$	1709.04183, 1711.11520
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow c\tilde{\chi}_1^0/\tilde{c}\tilde{\chi}_1^0$	0	2c	Yes	36.1	\tilde{t}_1	0.46 0.85	$m(\tilde{\chi}_1^0) = 0$ GeV	1805.01649
$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow \text{mono-jet}$	0	mono-jet	Yes	36.1	\tilde{t}_1 \tilde{t}_1	0.43	$m(\tilde{t}_1, \tilde{c}) - m(\tilde{\chi}_1^0) = 50$ GeV $m(\tilde{t}_1, \tilde{c}) - m(\tilde{\chi}_1^0) = 5$ GeV	1805.01649 1711.03301	
$\tilde{t}_2\tilde{t}_2, \tilde{t}_2 \rightarrow \tilde{t}_1 + h$	1-2 e, μ	4 b	Yes	36.1	\tilde{t}_2	0.32-0.88	$m(\tilde{\chi}_1^0) = 0$ GeV, $m(\tilde{t}_1) - m(\tilde{\chi}_1^0) = 180$ GeV	1706.03986	
EW direct	$\tilde{\chi}_1^\pm\tilde{\chi}_2^0$ via WZ	2-3 e, μ $ee, \mu\mu$	- ≥ 1	Yes Yes	36.1 36.1	$\tilde{\chi}_1^\pm/\tilde{\chi}_2^0$ $\tilde{\chi}_1^\pm/\tilde{\chi}_2^0$	0.6 0.17	$m(\tilde{\chi}_1^0) = 0$ $m(\tilde{\chi}_1^\pm) - m(\tilde{\chi}_1^0) = 10$ GeV	1403.5294, 1806.02293 1712.08119
	$\tilde{\chi}_1^\pm\tilde{\chi}_2^0$ via Wh	$\ell\ell\ell\gamma\ell b\bar{b}$	-	Yes	20.3	$\tilde{\chi}_1^\pm/\tilde{\chi}_2^0$	0.26	$m(\tilde{\chi}_1^0) = 0$	1501.07110
	$\tilde{\chi}_1^\pm\tilde{\chi}_1^\mp/\tilde{\chi}_2^0, \tilde{\chi}_1^\pm \rightarrow \tilde{\tau}\nu(\tilde{\tau}\bar{\nu}), \tilde{\chi}_2^0 \rightarrow \tilde{\tau}\tau(\tilde{\nu}\bar{\nu})$	2 τ	-	Yes	36.1	$\tilde{\chi}_1^\pm/\tilde{\chi}_2^0$ $\tilde{\chi}_1^\pm/\tilde{\chi}_2^0$	0.76 0.22	$m(\tilde{\chi}_1^0) = 0, m(\tilde{\tau}, \tilde{\nu}) = 0.5(m(\tilde{\chi}_1^\pm) + m(\tilde{\chi}_1^0))$ $m(\tilde{\chi}_1^\pm) - m(\tilde{\chi}_1^0) = 100$ GeV, $m(\tilde{\tau}, \tilde{\nu}) = 0.5(m(\tilde{\chi}_1^\pm) + m(\tilde{\chi}_1^0))$	1708.07875 1708.07875
	$\tilde{\ell}_{L,R}\tilde{\ell}_{L,R}, \tilde{\ell} \rightarrow \ell\tilde{\chi}_1^0$	2 e, μ	0	Yes	36.1	$\tilde{\ell}$	0.5	$m(\tilde{\chi}_1^0) = 0$	1803.02762
	$\tilde{\ell}_{L,R}\tilde{\ell}_{L,R}, \tilde{\ell} \rightarrow \ell\tilde{\chi}_1^0$	2 e, μ	≥ 1	Yes	36.1	$\tilde{\ell}$	0.18	$m(\tilde{\ell}) - m(\tilde{\chi}_1^0) = 5$ GeV	1712.08119
	$\tilde{H}\tilde{H}, \tilde{H} \rightarrow h\tilde{G}/Z\tilde{G}$	0 4 e, μ	$\geq 3b$ 0	Yes Yes	36.1 36.1	\tilde{H} \tilde{H}	0.13-0.23 0.3	BR($\tilde{\chi}_1^0 \rightarrow h\tilde{G}$) = 1 BR($\tilde{\chi}_1^0 \rightarrow Z\tilde{G}$) = 1	1806.04030 1804.03602
Long-lived particles	Direct $\tilde{\chi}_1^\pm\tilde{\chi}_1^\mp$ prod., long-lived $\tilde{\chi}_1^\pm$	Disapp. trk	1 jet	Yes	36.1	$\tilde{\chi}_1^\pm$ $\tilde{\chi}_1^\pm$	0.46 0.15	Pure Wino Pure Higgsino	1712.02118 ATL-PHYS-PUB-2017-019
	Stable \tilde{g} R-hadron	SMP	-	-	3.2	\tilde{g}	1.6		1606.05129
	Metastable \tilde{g} R-hadron, $\tilde{g} \rightarrow q\tilde{q}\tilde{\chi}_1^0$		Multiple		32.8	\tilde{g} [$\tau(\tilde{g}) = 100$ ns, 0.2 ns]	1.6 2.4	$m(\tilde{\chi}_1^0) = 100$ GeV	1710.04901, 1604.04520
	GMSB, $\tilde{\chi}_1^0 \rightarrow \gamma\tilde{G}$, long-lived $\tilde{\chi}_1^0$	2 γ	-	Yes	20.3	$\tilde{\chi}_1^0$	0.44	$1 < c\tau(\tilde{\chi}_1^0) < 3$ ns, SPS8 model	1409.5542
$\tilde{g}\tilde{g}, \tilde{\chi}_1^0 \rightarrow e\bar{e}\nu/\mu\bar{\mu}\nu$	displ. $ee/\mu\mu$	-	-	20.3	\tilde{g}	1.3	$6 < c\tau(\tilde{\chi}_1^0) < 1000$ mm, $m(\tilde{\chi}_1^0) = 1$ TeV	1504.05162	
RPV	LFV $pp \rightarrow \tilde{\nu}_r + X, \tilde{\nu}_r \rightarrow e\mu/\tau\mu/\mu\tau$	$e\mu, e\tau, \mu\tau$	-	-	3.2	$\tilde{\nu}_r$	1.9	$\lambda'_{311} = 0.11, \lambda'_{132/133/233} = 0.07$	1607.08079
	$\tilde{\chi}_1^\pm\tilde{\chi}_1^\mp/\tilde{\chi}_2^0 \rightarrow WW/Z\ell\ell\nu\nu$	4 e, μ	0	Yes	36.1	$\tilde{\chi}_1^\pm/\tilde{\chi}_2^0$ [$\lambda_{133} \neq 0, \lambda_{12k} \neq 0$]	0.82 1.33	$m(\tilde{\chi}_1^0) = 100$ GeV	1804.03602
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}\tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow qq\tilde{q}$	0	4-5 large-R jets	-	36.1 36.1	\tilde{g} [$m(\tilde{\chi}_1^0) = 200$ GeV, 1100 GeV] \tilde{g} [$\lambda'_{112} = 2e-4, 2e-5$]	1.3 1.9 1.05 2.0	Large λ'_{112} $m(\tilde{\chi}_1^0) = 200$ GeV, bino-like	1804.03568 ATLAS-CONF-2018-003
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow t\tilde{b}s/\tilde{g} \rightarrow t\tilde{t}\tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow t\tilde{b}s$		Multiple		36.1	\tilde{g} [$\lambda'_{323} = 1, 1e-2$]	1.8 2.1	$m(\tilde{\chi}_1^0) = 200$ GeV, bino-like	ATLAS-CONF-2018-003
	$\tilde{t}_1, \tilde{t}_1 \rightarrow t\tilde{b}s$		Multiple		36.1	\tilde{g} [$\lambda'_{323} = 2e-4, 1e-2$]	0.55 1.05	$m(\tilde{\chi}_1^0) = 200$ GeV, bino-like	ATLAS-CONF-2018-003
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow b\tilde{s}$	0	2 jets + 2 b	-	36.7	\tilde{t}_1 [$qq, b\bar{s}$]	0.42 0.61		1710.07171
$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow b\tilde{\ell}$	2 e, μ	2 b	-	36.1	\tilde{t}_1	0.4-1.45	BR($\tilde{t}_1 \rightarrow b\ell/h\mu$) > 20%	1710.05544	

*Only a selection of the available mass limits on new states or phenomena is shown. Many of the limits are based on simplified models, c.f. refs. for the assumptions made.

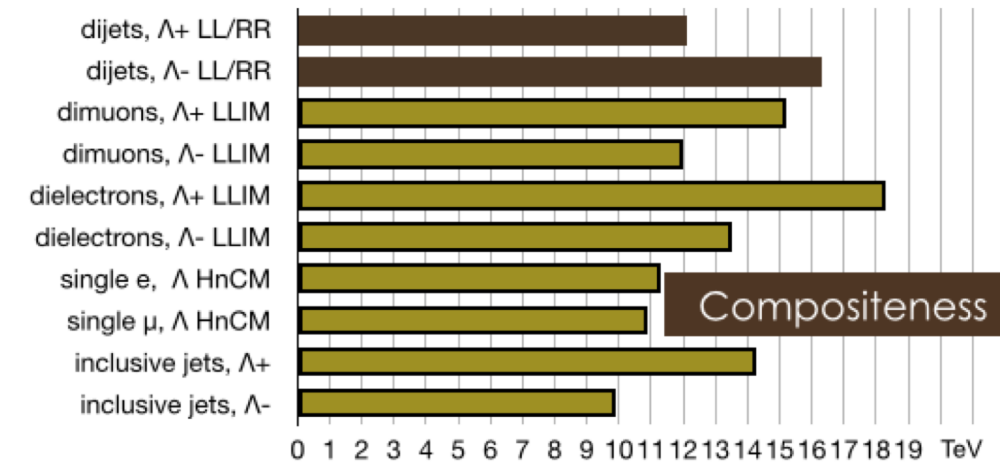
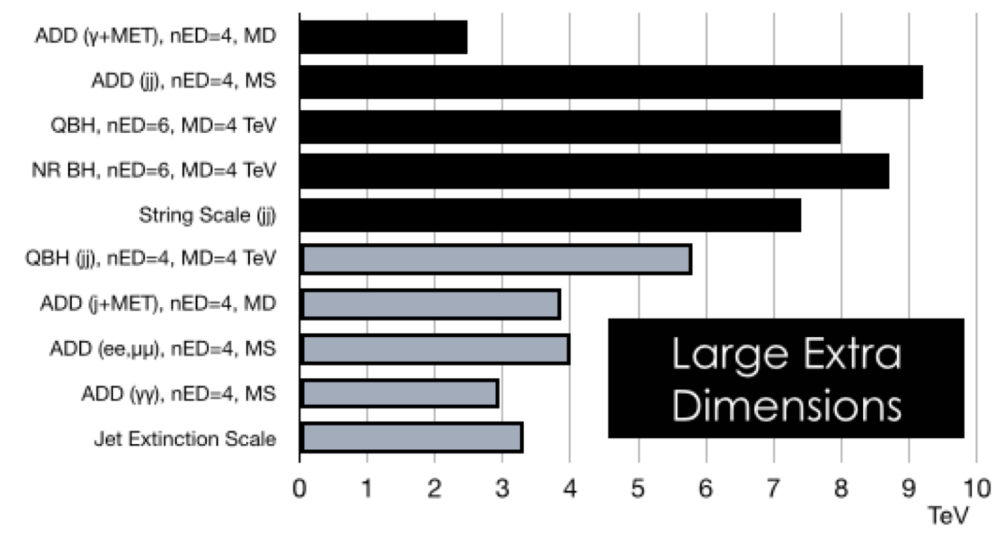
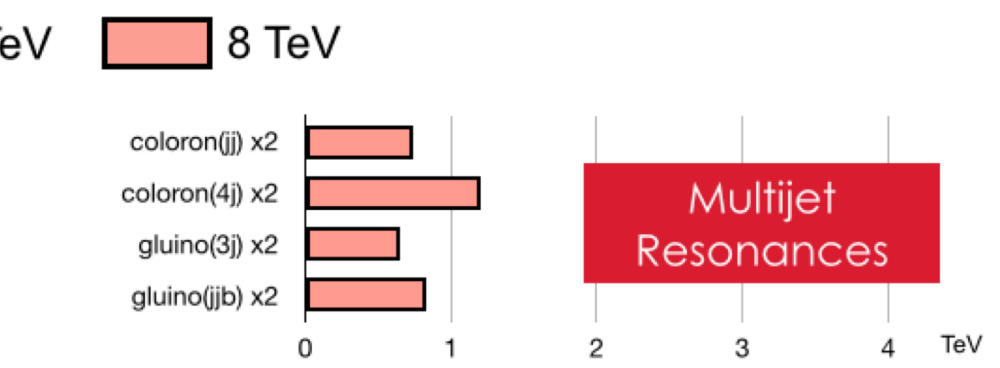
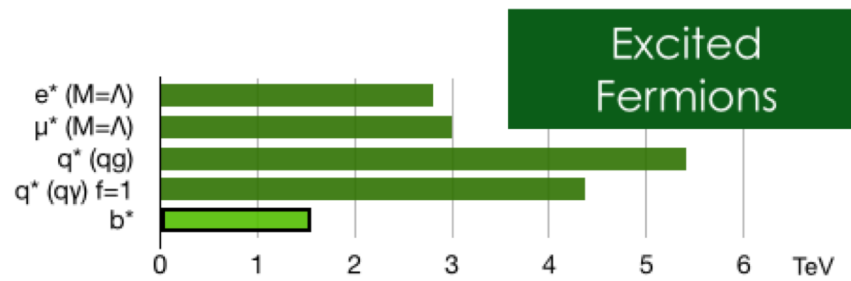
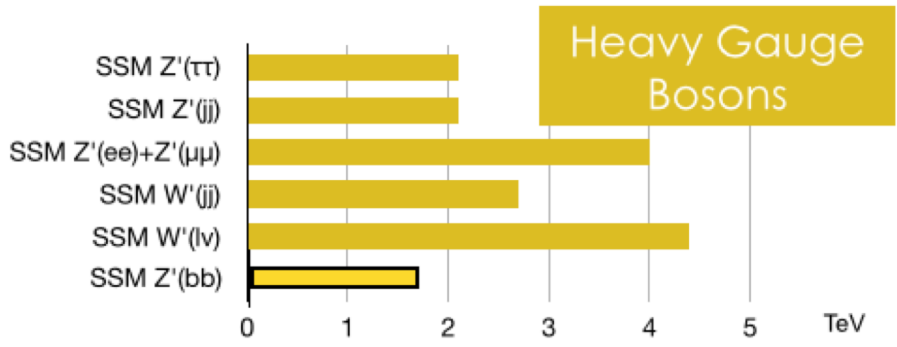
10⁻¹

1

Mass scale [TeV]



CMS Preliminary

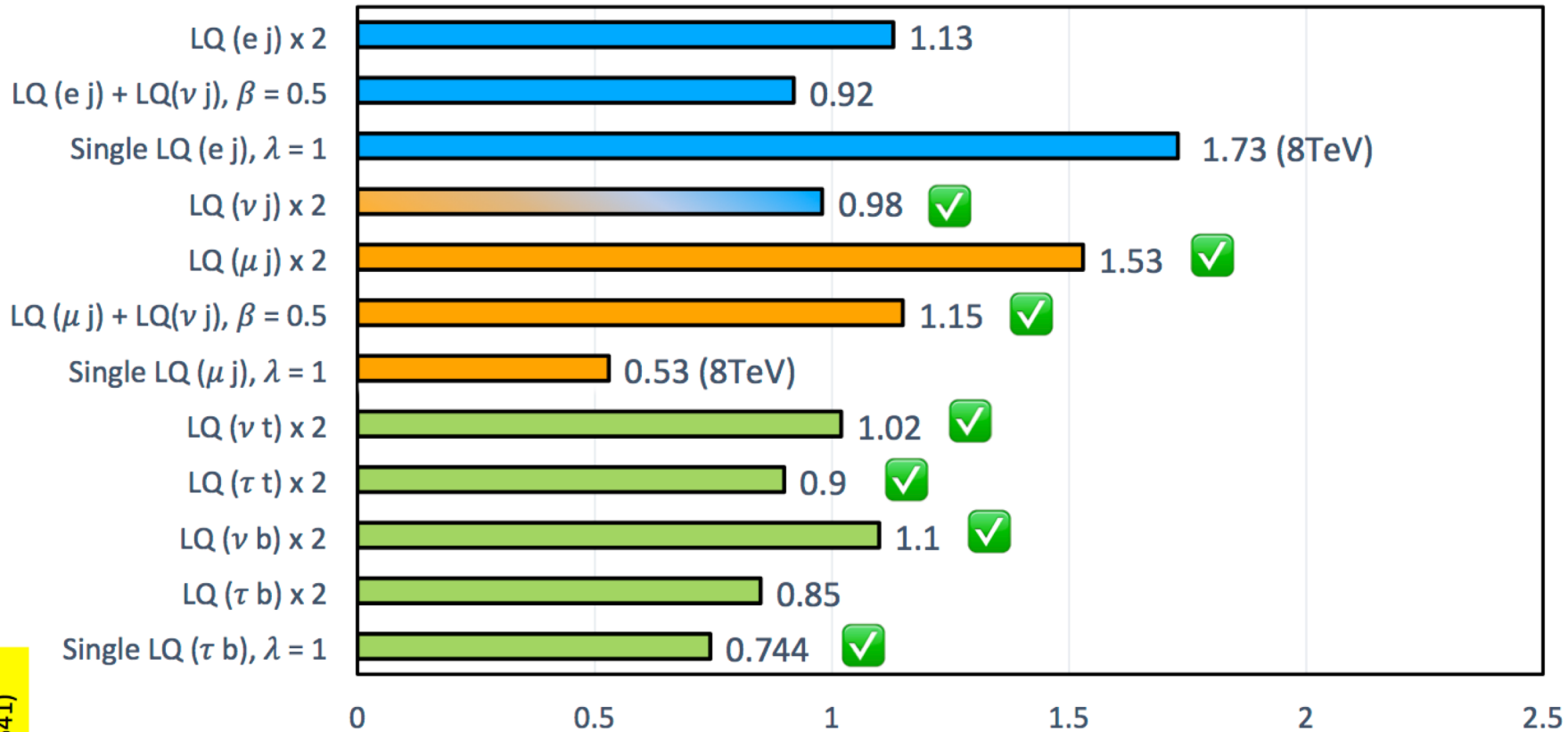


May 2018

LQ → 1st gen. 2nd. gen. 3rd gen.

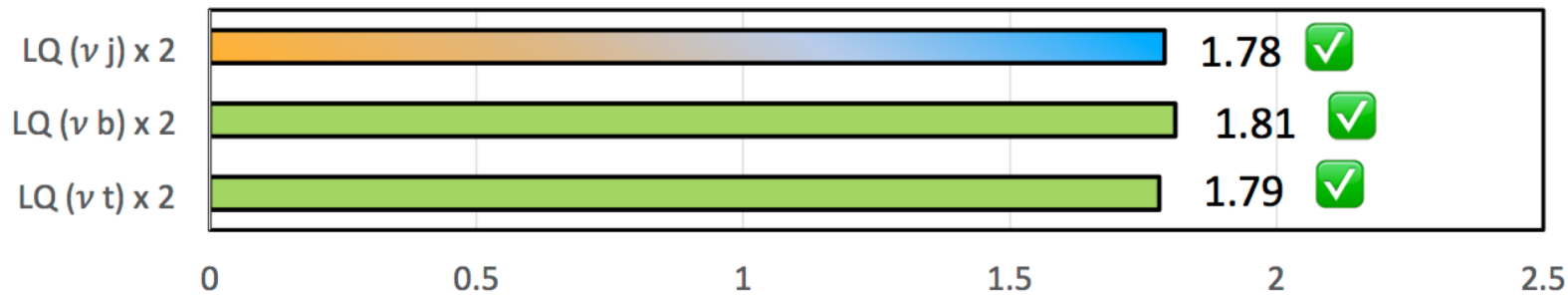
✓ Full 2016 dataset

Scalar LQ



Vector LQ

(LQ model used: 1801.07641)



LeptoQuark mass (TeV)

At what point do you lose interest in extending the new physics searches?

- Keep in mind that after Run 2, you will only have collected 5% of the total luminosity expected during the LHC lifetime.
- If you discover new physics consistent with explanations of the gauge hierarchy problem (why is $m_W/M_{\text{PL}} \sim 10^{-17}$?), the little hierarchy problem becomes much less pressing.

Final thoughts on naturalness

- The announcement of the death of naturalness may be premature.
- There is still room for theoretical innovations.
- However, in evaluating new approaches to naturalness, it is important to consider how one could test these ideas experimentally (i.e. what observable phenomenon would convince you that Nature has employed a natural theory for the dynamics of electroweak symmetry breaking?).

2. Do we really know the particle content of the TeV-scale effective theory?

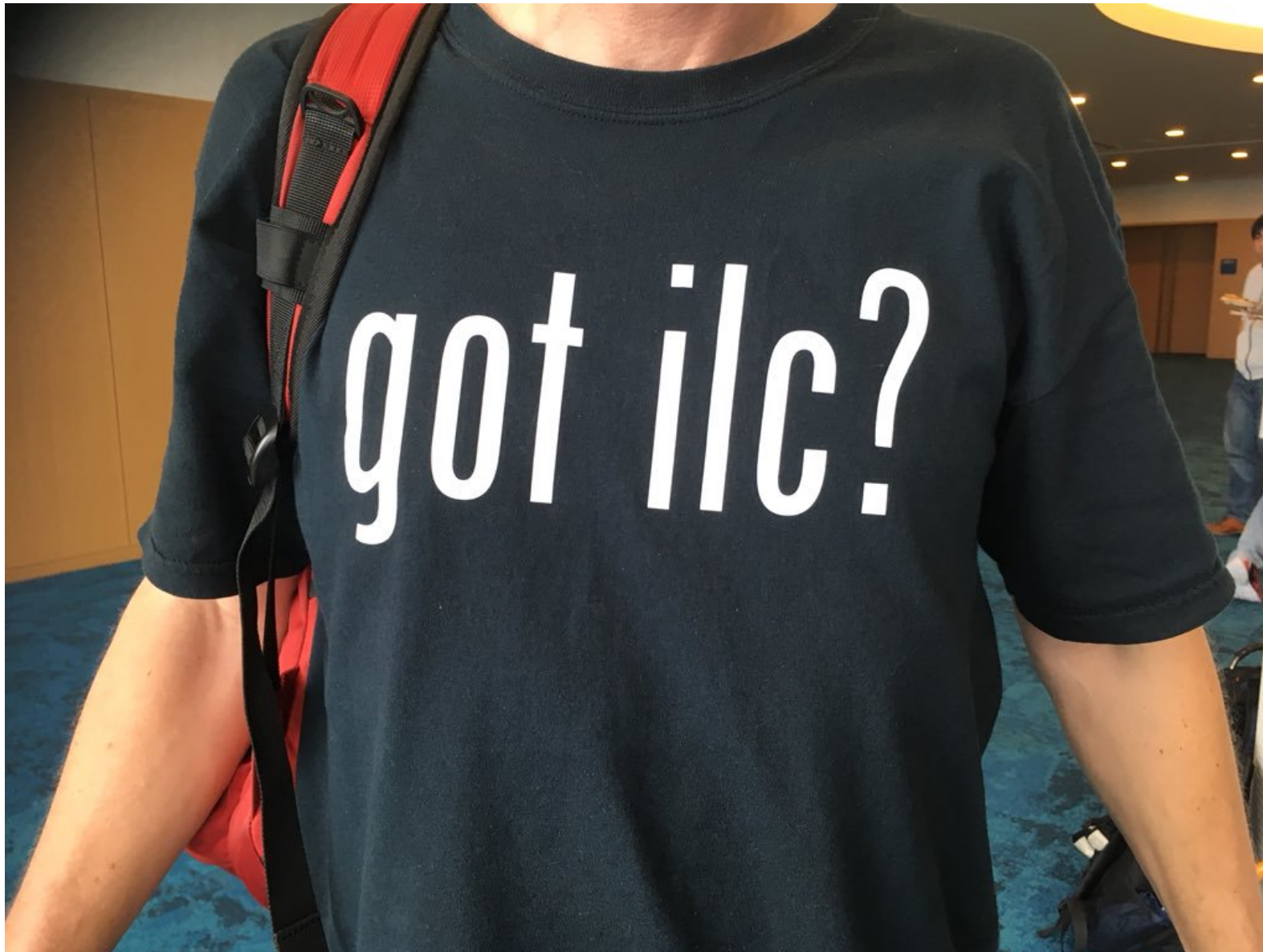
- The fermion sector of the Standard Model (SM) is non-minimal. Three generations—who ordered that?
- The scalar sector of the SM has a single Higgs boson. Why not multiple families of Higgs scalars?
- There are good reasons to think that the number of families of chiral fermions is limited to 3. But what about vector-like quarks and leptons?
- Flavor anomalies have revived interest in leptoquarks.

- Are we really sure that the gauge group of the effective TeV-scale theory is $SU(3) \times SU(2) \times U(1)$? Are there new gauge bosons lurking in the region of 1—10 TeV?

- Of course, don't forget about the dark sector, which I shall define as particles that are neutral with respect to $SU(3) \times SU(2) \times U(1)$. Perhaps motivated by theories of dark matter, but could exist independently. Communications with the SM sector is possible through the various portals.
 - The Higgs portal ($\Phi^\dagger \Phi$ is a SM singlet)
 - U(1) gauge boson mixing ($F^{\mu\nu} F'_{\mu\nu}$)
 - The neutrino portal ($L^\dagger \Phi N$)

3. So, where do we go from here?

- Explore the Higgs sector as thoroughly as possible (since, you have never seen anything like it before).
 - Experimental studies
 - Implications for early universe cosmology
- Precision, precision, precision.
- Don't despair prematurely.
- Search for BSM physics in regions with significant SM backgrounds. (Yes, it is more difficult.)
- Try to expand the area illuminated by the lamppost.



<https://www3.nhk.or.jp/nhkworld/en/vod/scienceview/2015197/>

The popular press has taken notice ...

CANDORVILLE
by DARRIN BELL

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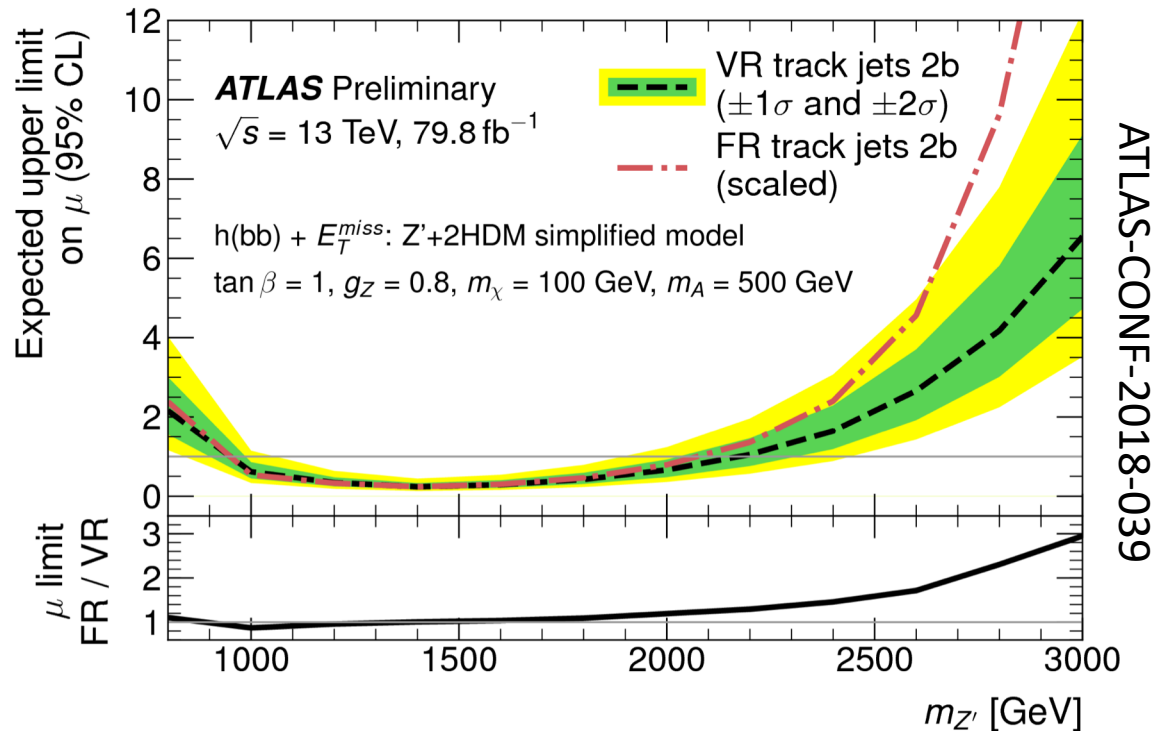
GALAXIES, STARS, PLANETS, YOU, ME, GOLDFISH... ALL OF OUR ATOMS WOULD BE INSTANTLY BLASTED APART.

AND GET THIS: IT MAY HAVE ALREADY HAPPENED. SINCE IT ONLY GROWS AT THE SPEED OF LIGHT, WE'LL HAVE NO WARNING AT ALL.



What should we prioritize right now?

- **Gain on existing analyses: developments in object performance / systematics**
 - This can be helped in some cases by machine learning* (eg top tagging...)
 - This can help us beat the simple increase in the integrated luminosity, there are real gains to be had, eg:



* Machine learning can be great! But one must remember a rule:

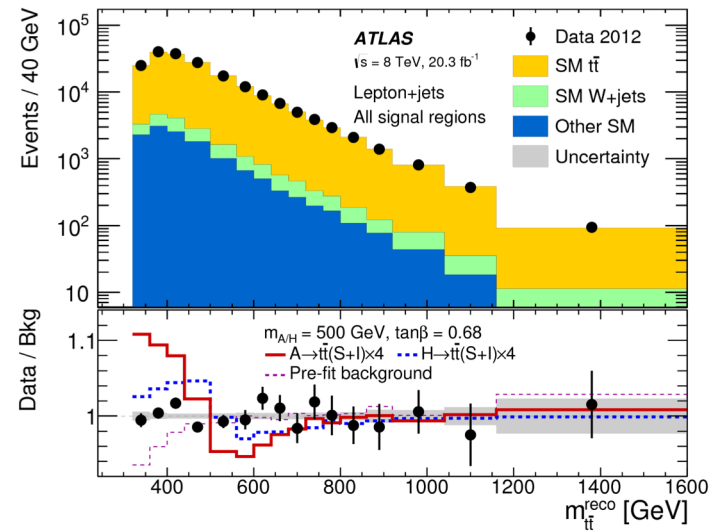


What should we prioritize right now?

- **Uncovered / less covered signatures**

(=> need to assess what has been covered and what not -> improve on recasting?):

- motivating dedicated new searches covering new signatures (eg latest emerging jet paper by CMS)
- More challenging signatures (less covered for a reason!):
 - “Strange” objects & long-lived particles
 - Searches with interference (eg $A \rightarrow t\bar{t}$)



Phys. Rev. Lett. 119 (2017) 191803

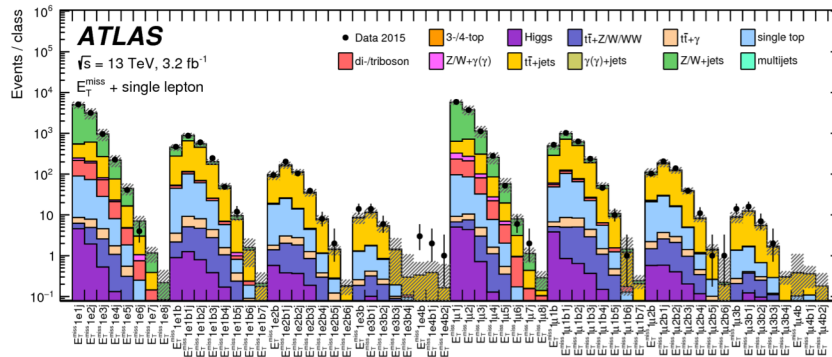
What should we prioritize right now?

- **Measurements :**
 - Higgs couplings obviously
 - Flavour ‘anomalies’
 - Rare processes (e.g. $t\bar{t}t\bar{t}$)
 - Tails sensitive to new physics (through EFTs?)
 - WW scattering
 - Measurement – search a bit blurry: NP searches in large BG regimes (trigger vs killing BG constraints)



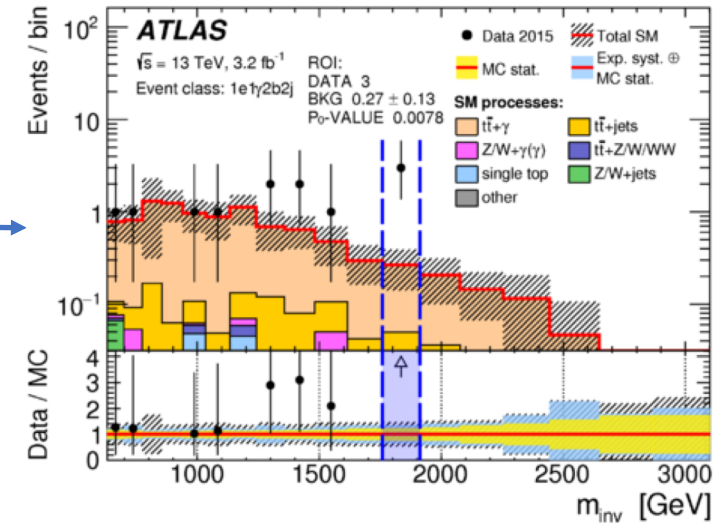
Searching for the unexpected ...?

Already doing something about it in ATLAS (general search) and CMS (MUSIC: <https://cds.cern.ch/record/2256653>)



60 out of 704 channels

Meff and Minv scanned in each category to find the most discrepant range



arXiv:1807.07447

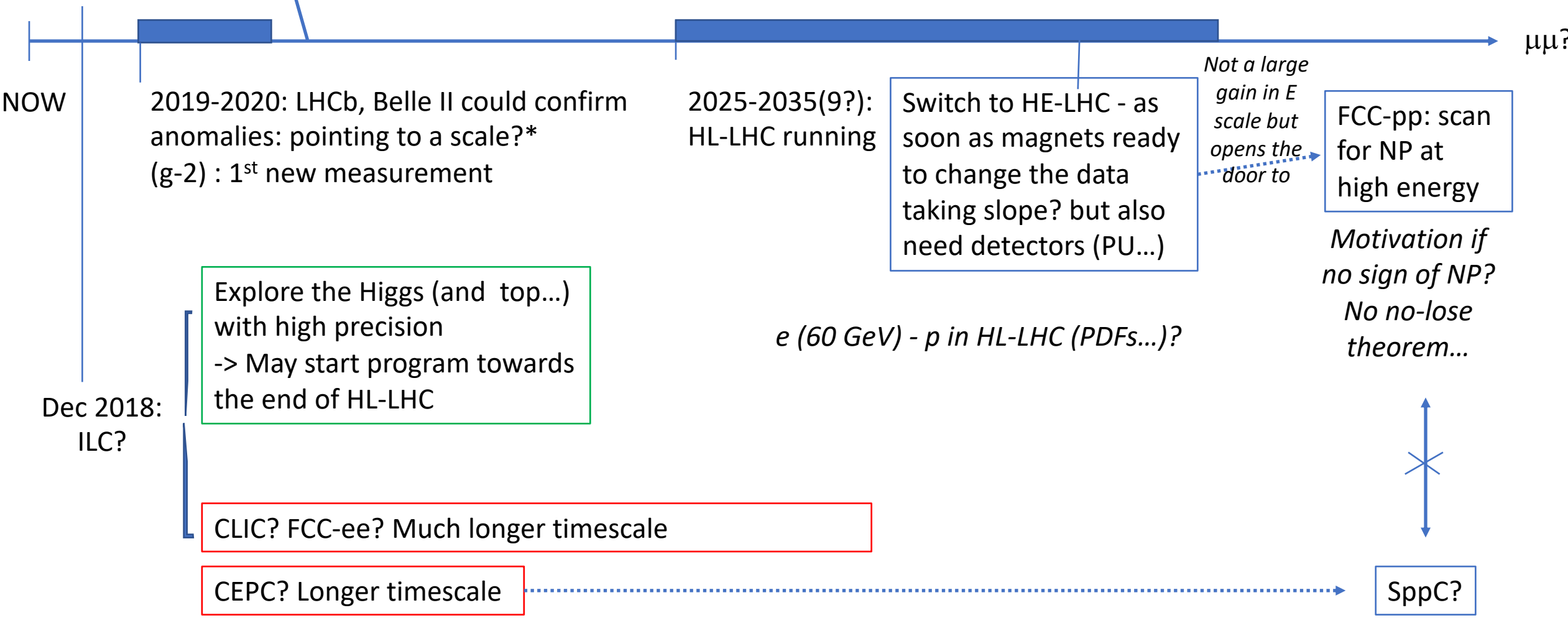
Compare analysis and pseudo-experiment p-values => discrepancies?

Not a discovery tool

=> a tool for discovering potentially interesting channels to be investigated with a dedicated analysis

Can be a 'limit' tool though – if you expect a lot of events with your model in a 'crazy channel' and we saw no data event -> there you go.

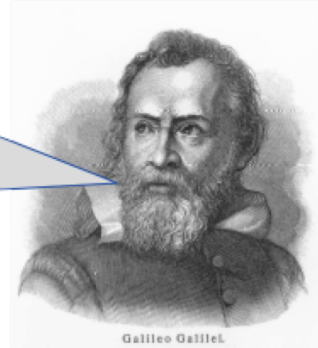
Of course any significant deviation seen in other sectors could have the same impact – there is a **lot** of data left to analyse!



Magnet development needed!

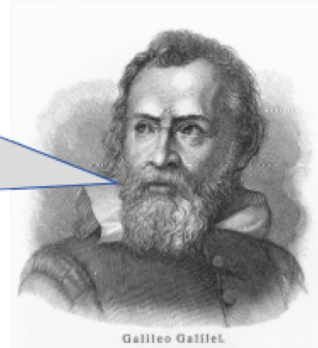
In this data-driven era, one should remind what Galileo himself famously said:

Measure what can be
measured and make
measurable what
cannot be measured.



In this data-driven era, one should remind what Galileo himself famously said:

Measure what can be measured and make measurable what cannot be measured.



And also:

Never fully trust quotes found on internet

