

TWO ASPECTS OF SIDM

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SIDM MOTIVATION

DIRECT DETECTION IS TESTING FI

Direct detection is testing Freeze-in

Th. Hambye, M.T., J. Vandecasteele & L. Vanderheyden (2018)

(The Four Basic Ways of Creating Dark Matter Through a Portal)

X. Chu, Th. Hambye & M.T (2012)

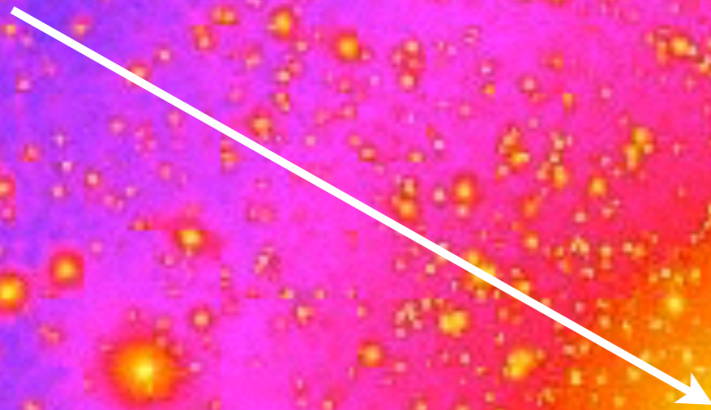
SIDM + NS \rightarrow BH

Non-primordial solar mass black holes

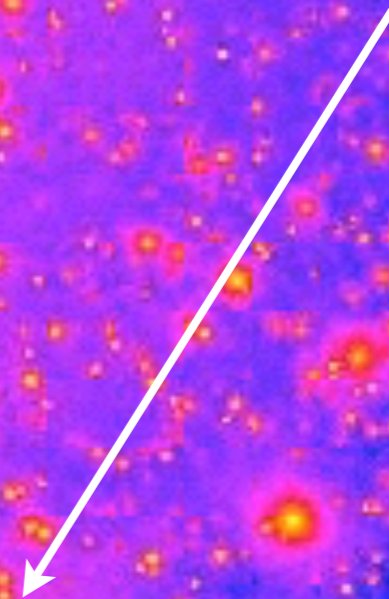
C. Kouvaris, P. Tinyakov, M.T. (2018)

WHY SIDM ?

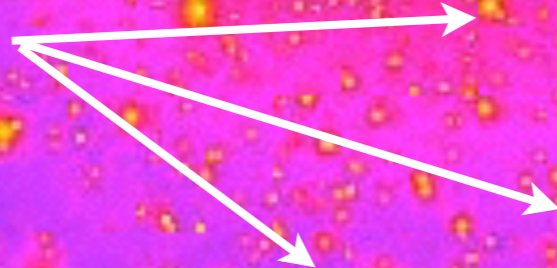
core
or
cusp?



to-big-to-fail ?



missing
satellites ?



CDM only simulation

SIDM may alleviate the small-scale problems

core/cusp

Spergel & Steinhardt (2000),...

too-big-to-fail

Vogelsberger, Zavala & Loeb (2012),...

diversity*

Hamada, Kaplinghat, Pace & Yu (2016),...

collisions \longrightarrow thermalized DM \longrightarrow core instead of cusp



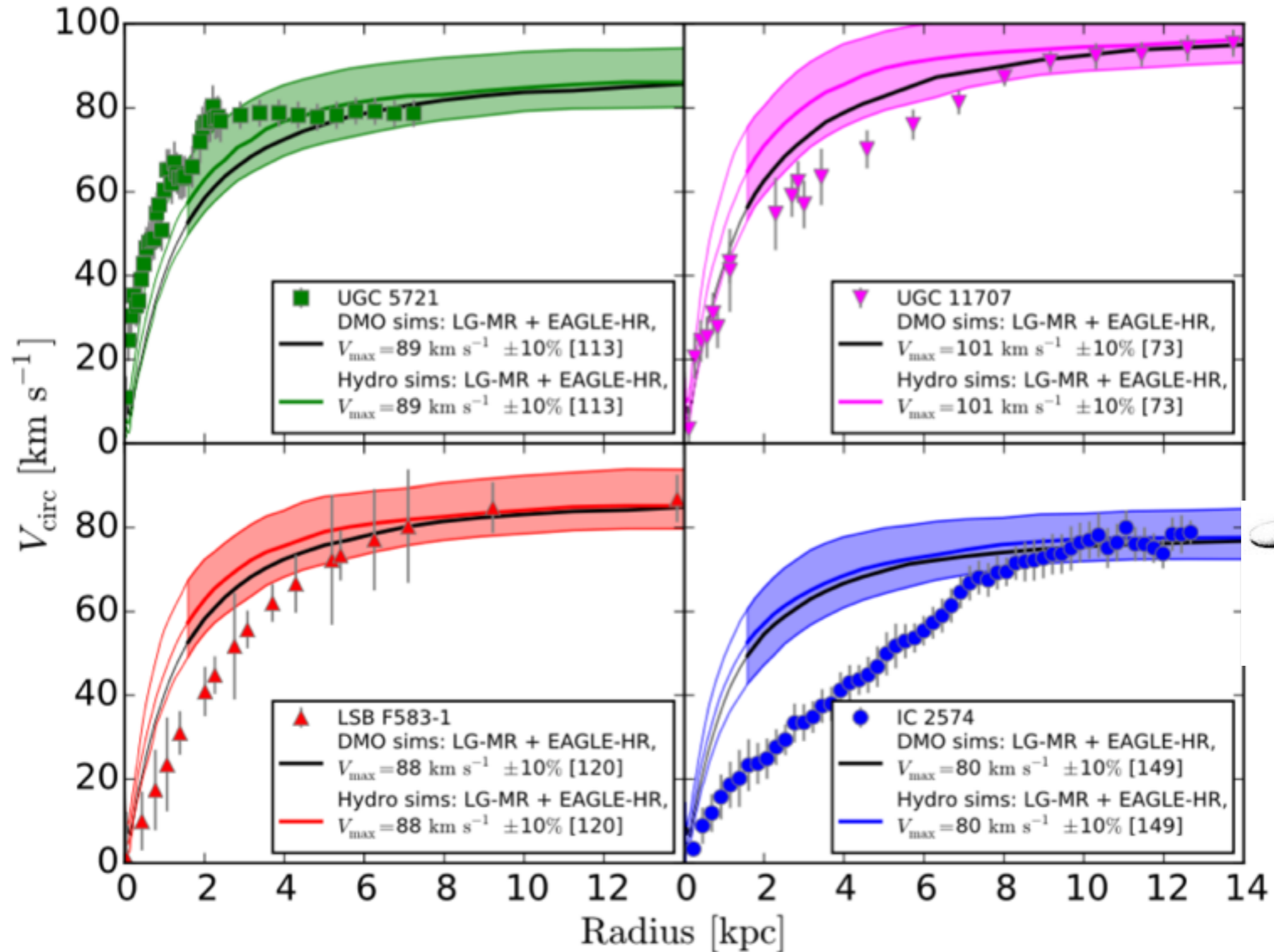
$$\frac{\sigma}{m} \sim \frac{\text{cm}^2}{\text{g}} \equiv \frac{\text{barn}}{\text{GeV}}$$

i.e. seemingly hadronic

but more generally light mediator

DIVERSITY PROBLEM

There is a **diversity problem** unexplained by CDM + BARYONS simulations (mostly dwarf galaxies)

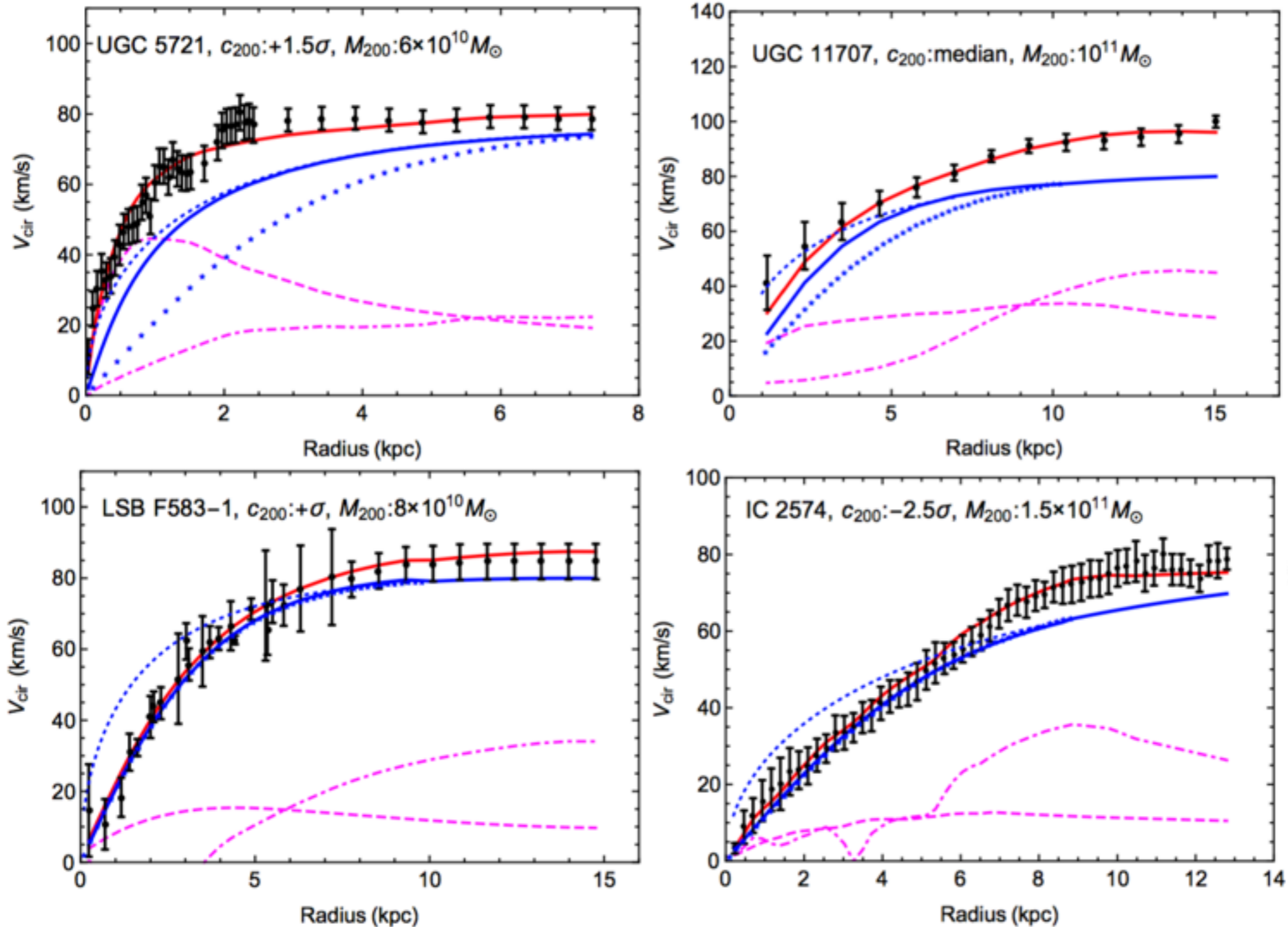


all
same
 V_{max}



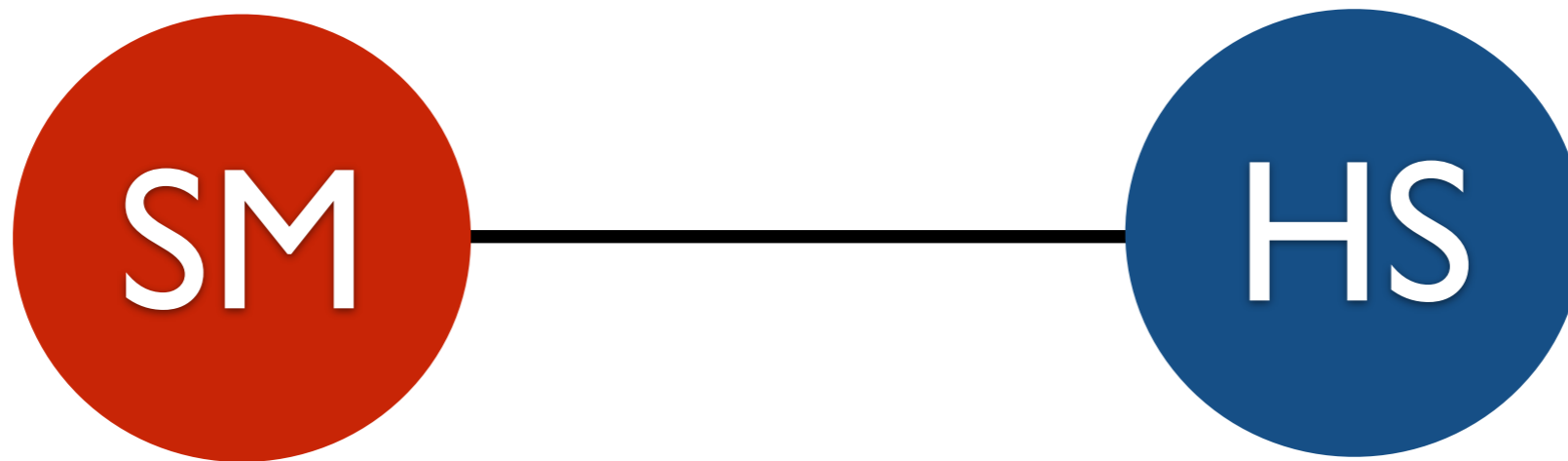
DIVERSITY PROBLEM

Diversity problem solved/alleviated with $\sigma/m \sim \frac{\text{cm}^2}{\text{g}}$



HIDDEN SECTOR

Patt & Wilczek (2006)



THE SM PORTALS

Patt & Wilczek (2006)

SM singlet
operators

renormalizable
interactions
(i.e. dimensionless couplings)

$$\bar{L}\tilde{H}$$

$$\Delta\mathcal{L} \supset y \bar{L}\tilde{H}N$$

Sterile neutrino

Dodelson & Widrow (1994)

...

$$B_{\mu\nu}$$

$$\Delta\mathcal{L} \supset \epsilon B_{\mu\nu} X^{\mu\nu}$$

Kinetic mixing

Holdom (1986)

...

$$H^\dagger H$$

$$\Delta\mathcal{L} \supset \lambda S^2 H^\dagger H$$

Higgs portal

This one is also
Lorentz invariant

Linked to EWSB?

Silveira & Zee (1985)
Veltman & Ynderain (1989)

...

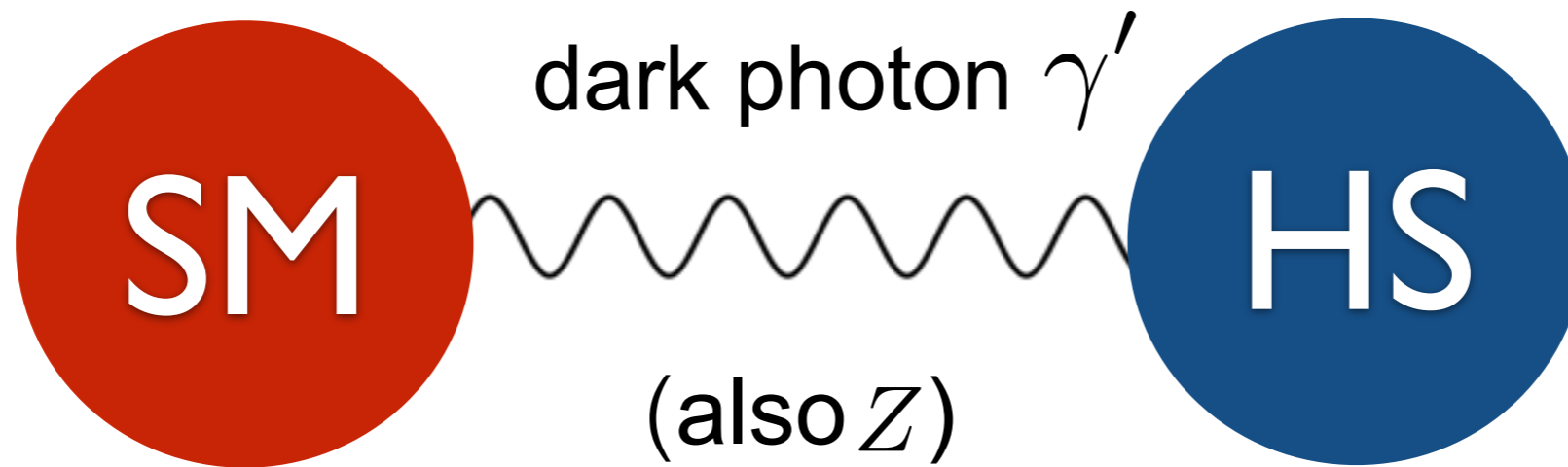
PART I

DIRECT DETECTION IS TESTING FREEZE-IN

Direct detection is testing Freeze-in
Th. Hambye, M.T., J. Vandecasteele & L. Vanderheyden (2018)

KINETIC MIXING

hidden charged χ



χ gauge interaction in HS

α'

χ stable \sim SM electron

m_χ

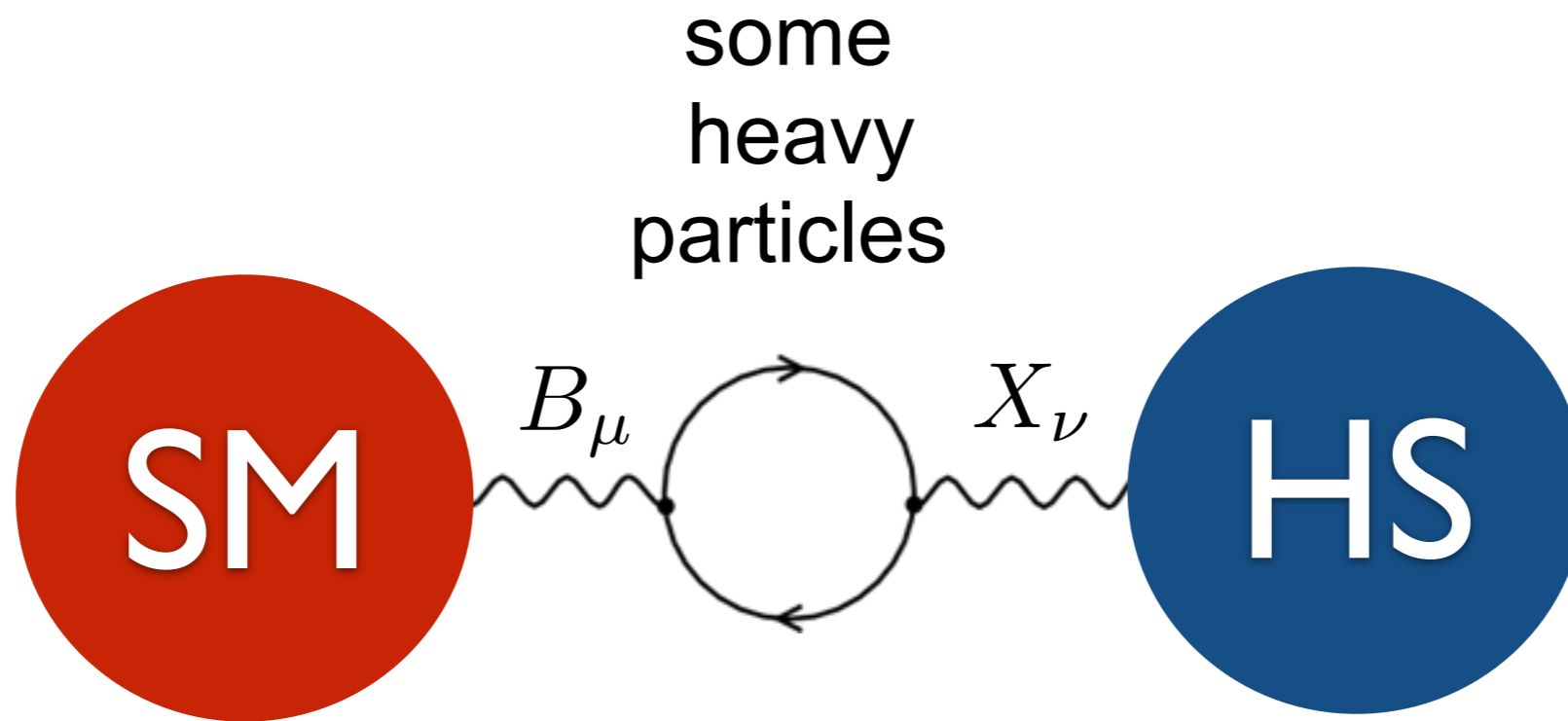
χ suppressed coupling to SM

$$\kappa = \epsilon \sqrt{\alpha' / \alpha}$$

only 4 parameters

$m_{\gamma'}$

FIMP THROUGH KINETIC MIXING



Holdom (1986)

\mathcal{K} is naturally tiny !

DM feebly coupled to SM

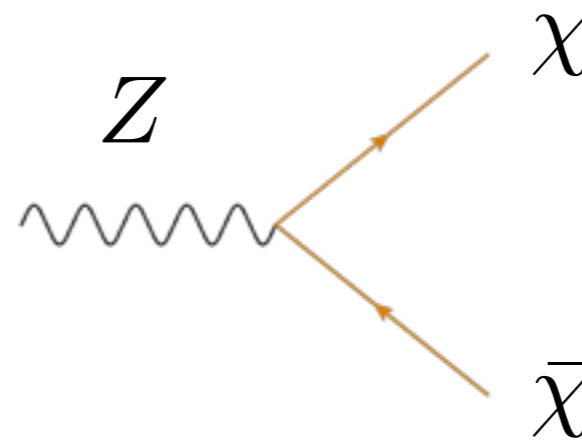
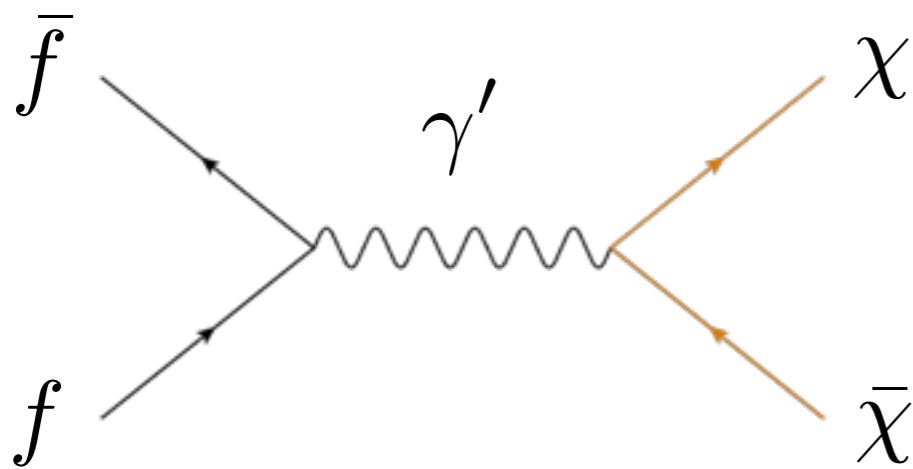
Feebly Interacting Massive Particle
or
FIMP



ABUNDANCE FROM FREEZE-IN

HS feebly coupled ~ no thermal equilibrium

**abundance could built up from
slow processes**

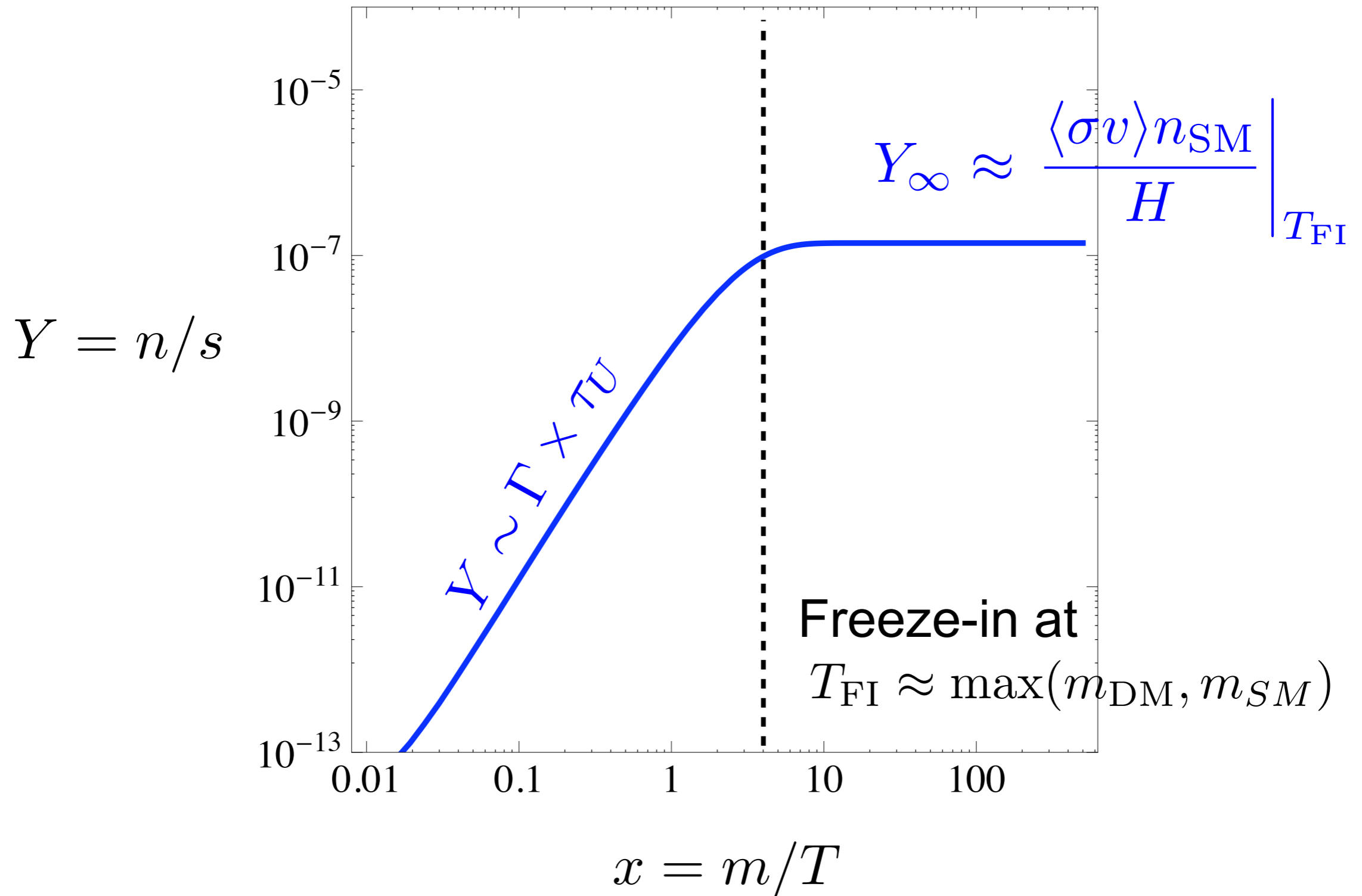


$$\propto \kappa = \epsilon \sqrt{\alpha' / \alpha}$$

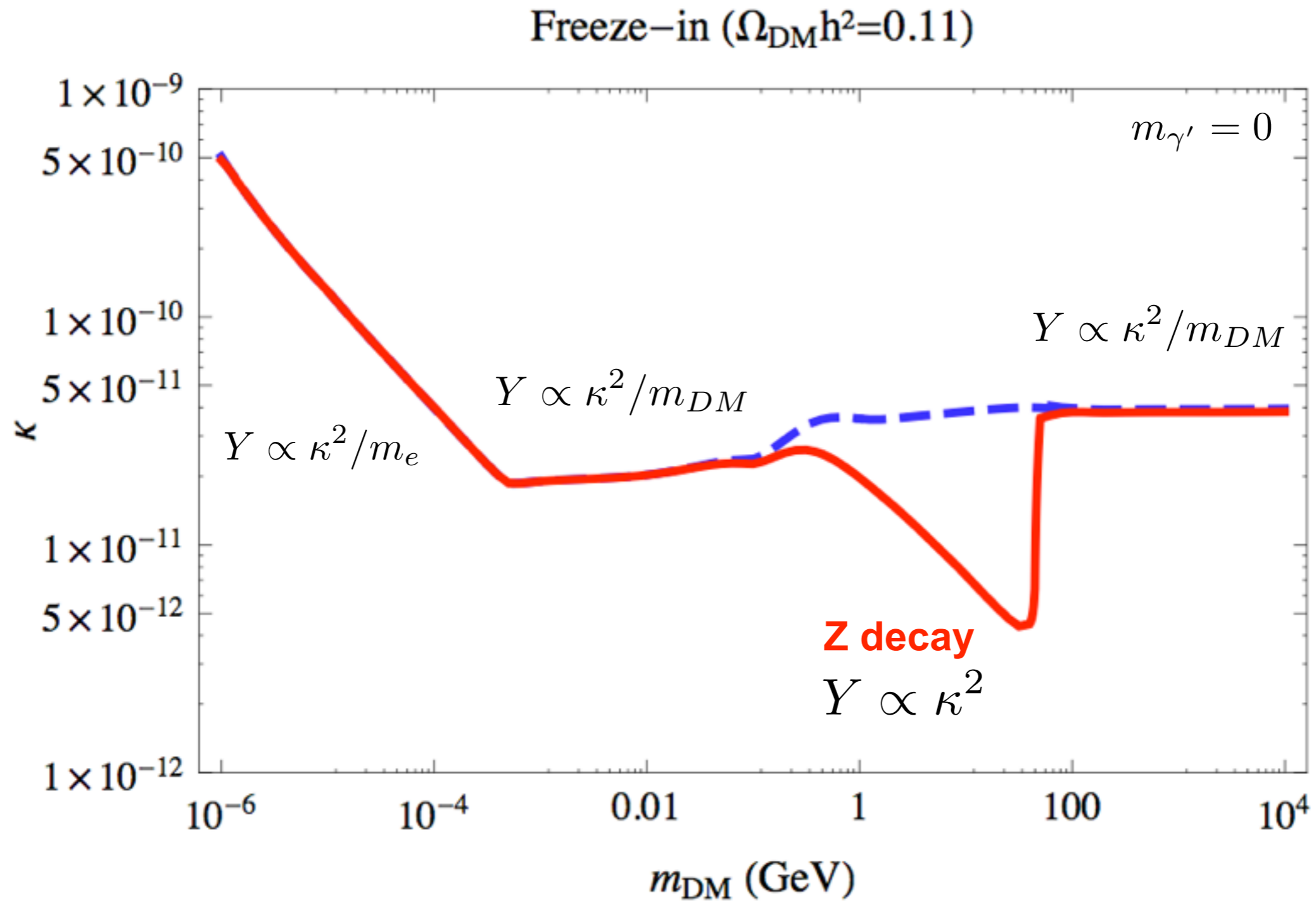


FREEZE-IN

ABUNDANCE FROM FREEZE-IN



ABUNDANCE FROM FREEZE-IN



ABUNDANCE FROM FREEZE-OUT



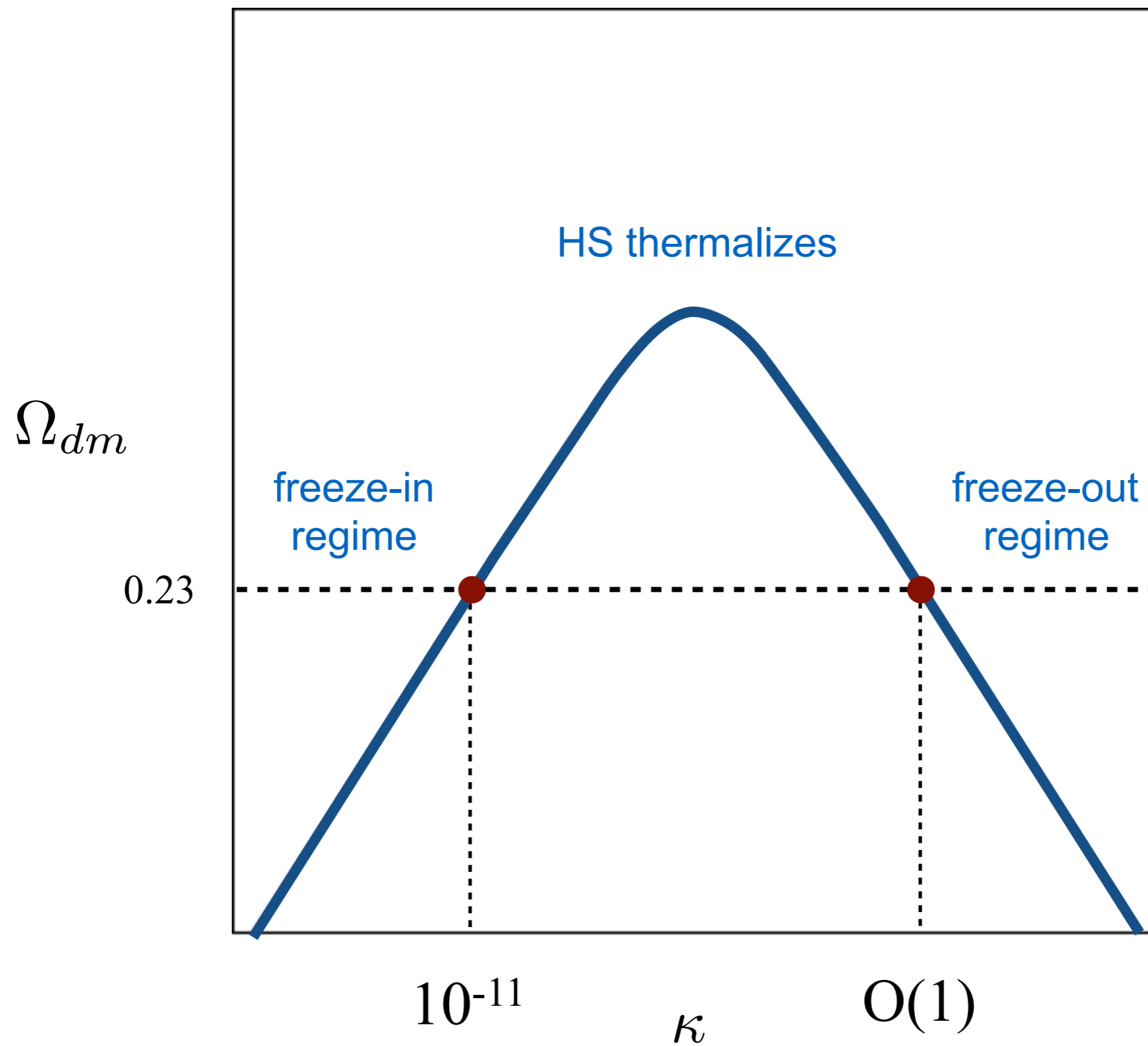
$$\Gamma = \sigma v n_{\text{DM}}$$

ABUNDANCE FROM FREEZE-IN



$$\Gamma = \sigma v n_{\text{SM}}$$

FREEZE-IN vs FREEZE-OUT



4 BASIC WAYS TO CREATE DM THROUGH A PORTAL

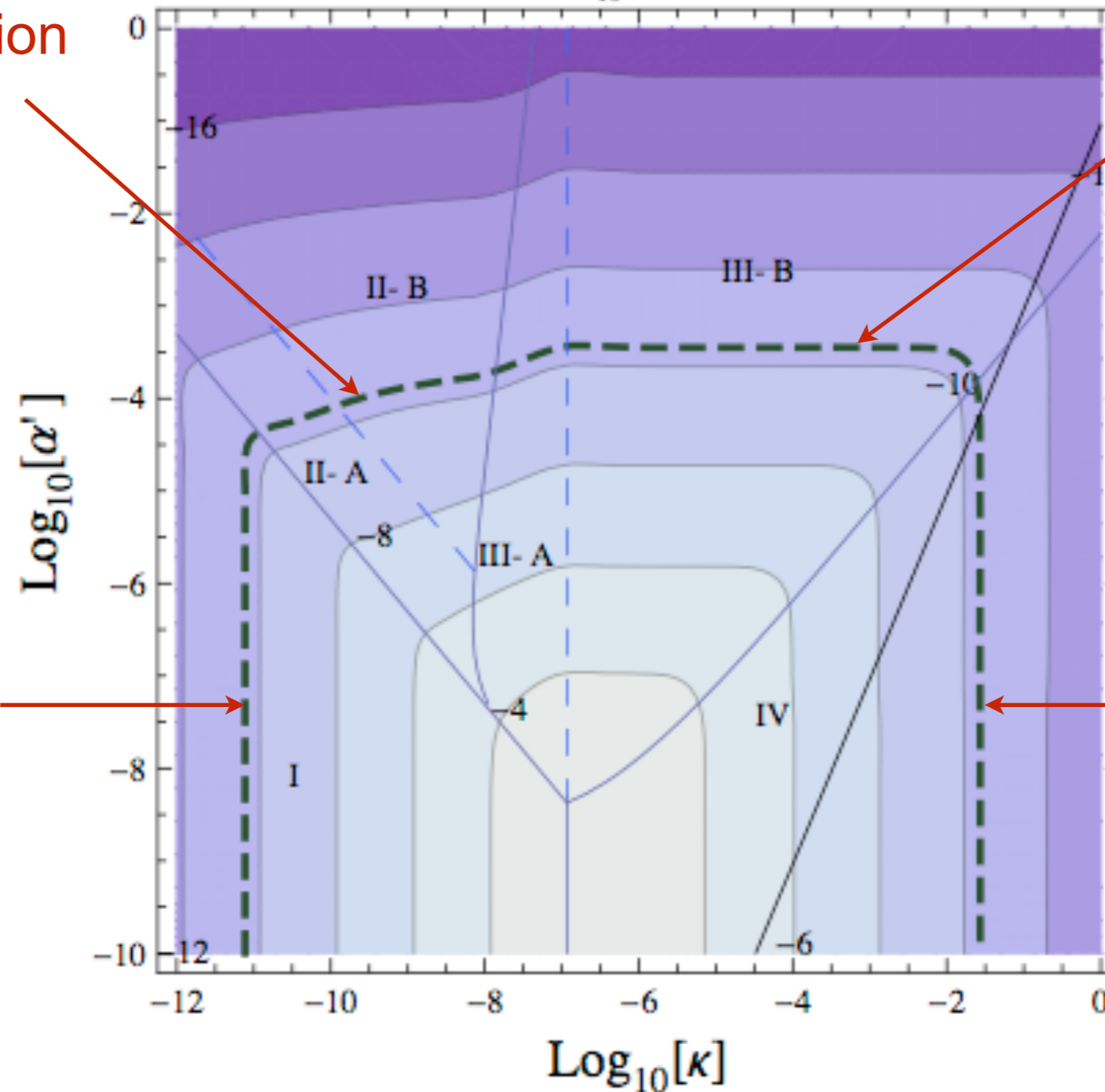
phase diagram $\text{Log}_{10}[Y_{\text{DM}}]$ ($m_{\text{DM}}=10\text{GeV}$)

II : reannihilation

III : freeze-out
in hidden
sector
(secluded DM)

I : freeze-in

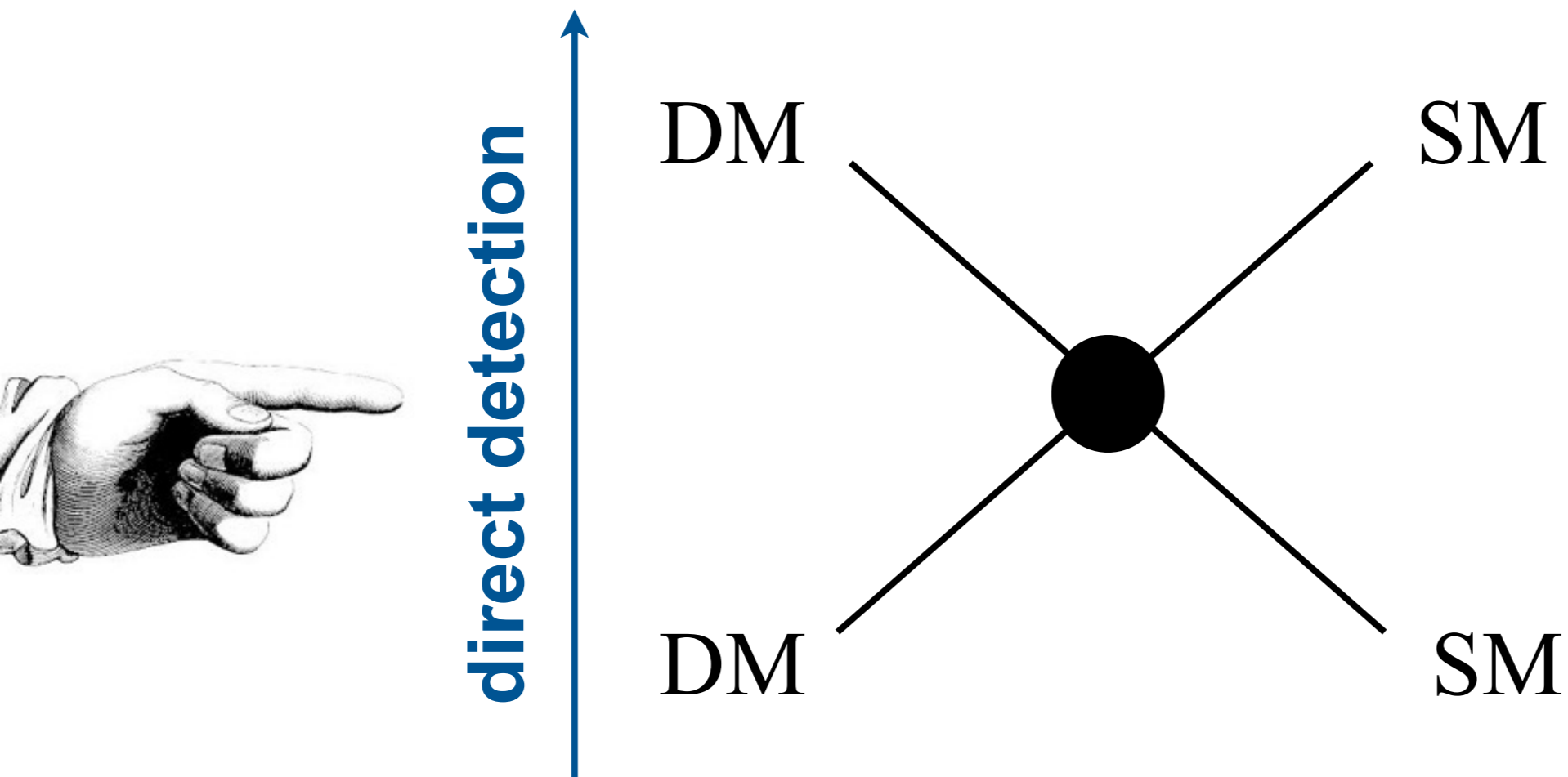
IV : usual
freeze-out



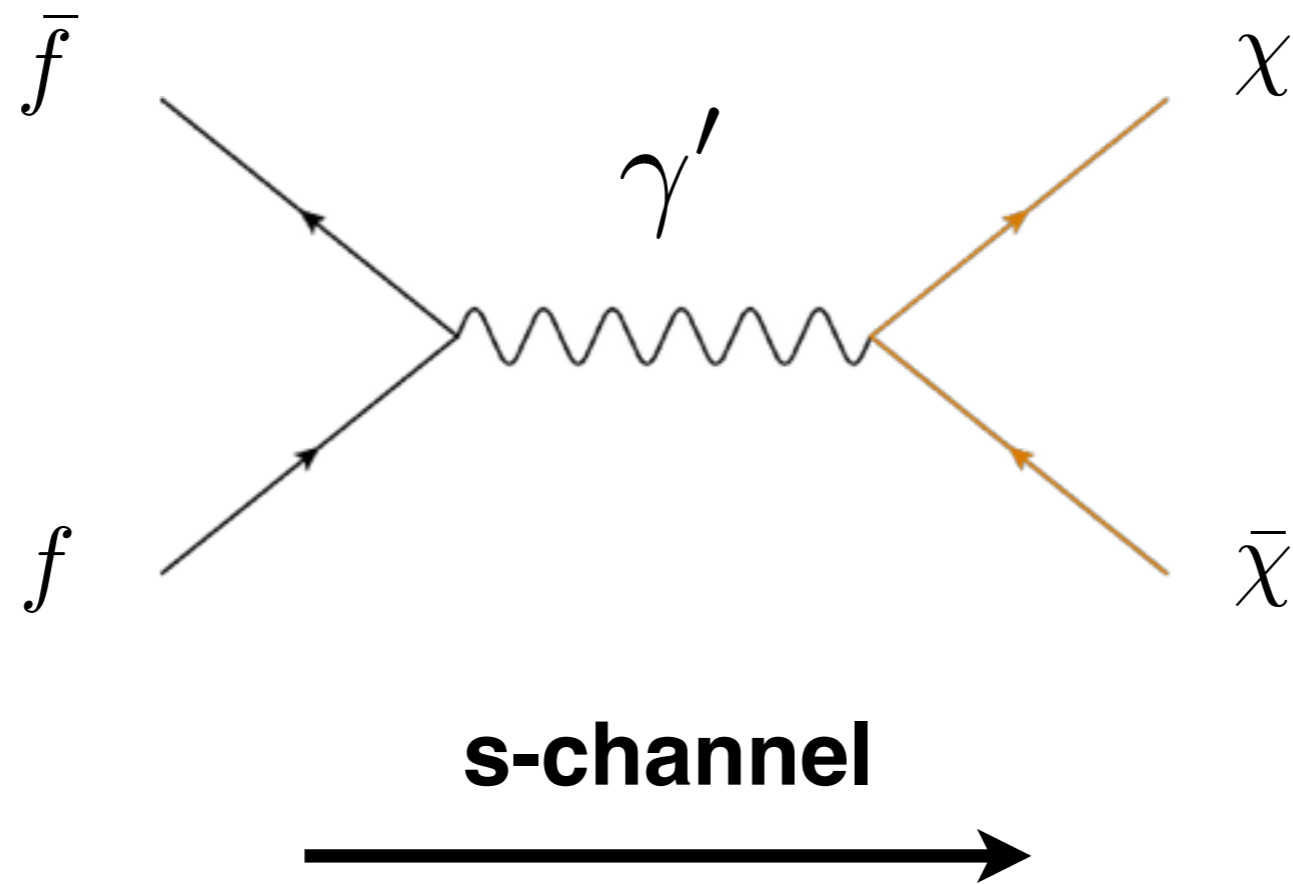
DIRECT DETECTION TEST OF FREEZE-IN

cosmic abundance

$$\kappa = \epsilon \sqrt{\alpha' / \alpha} = (10^{-11})$$



PRODUCTION THROUGH S-CHANNEL



determines relic abundance

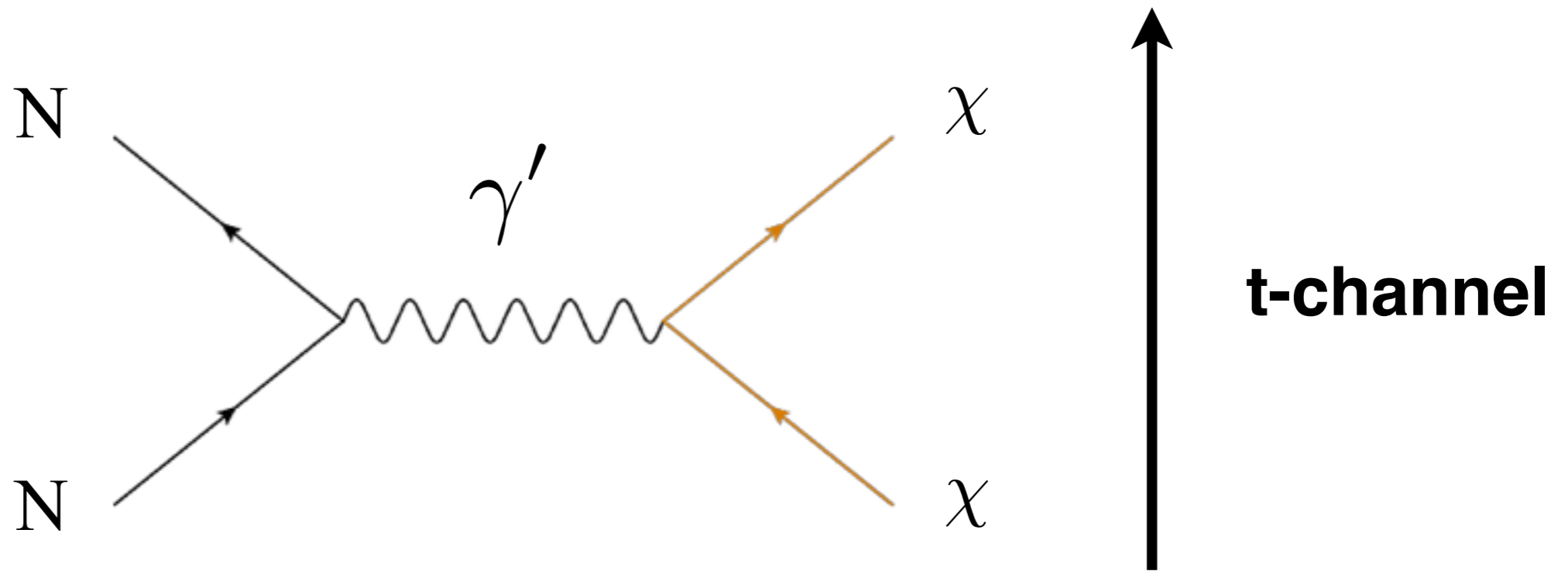
very small cross section

RUTHERFORD SCATTERING - DIRECT DETECTION

recoil energy

$$E_R$$

in keV range

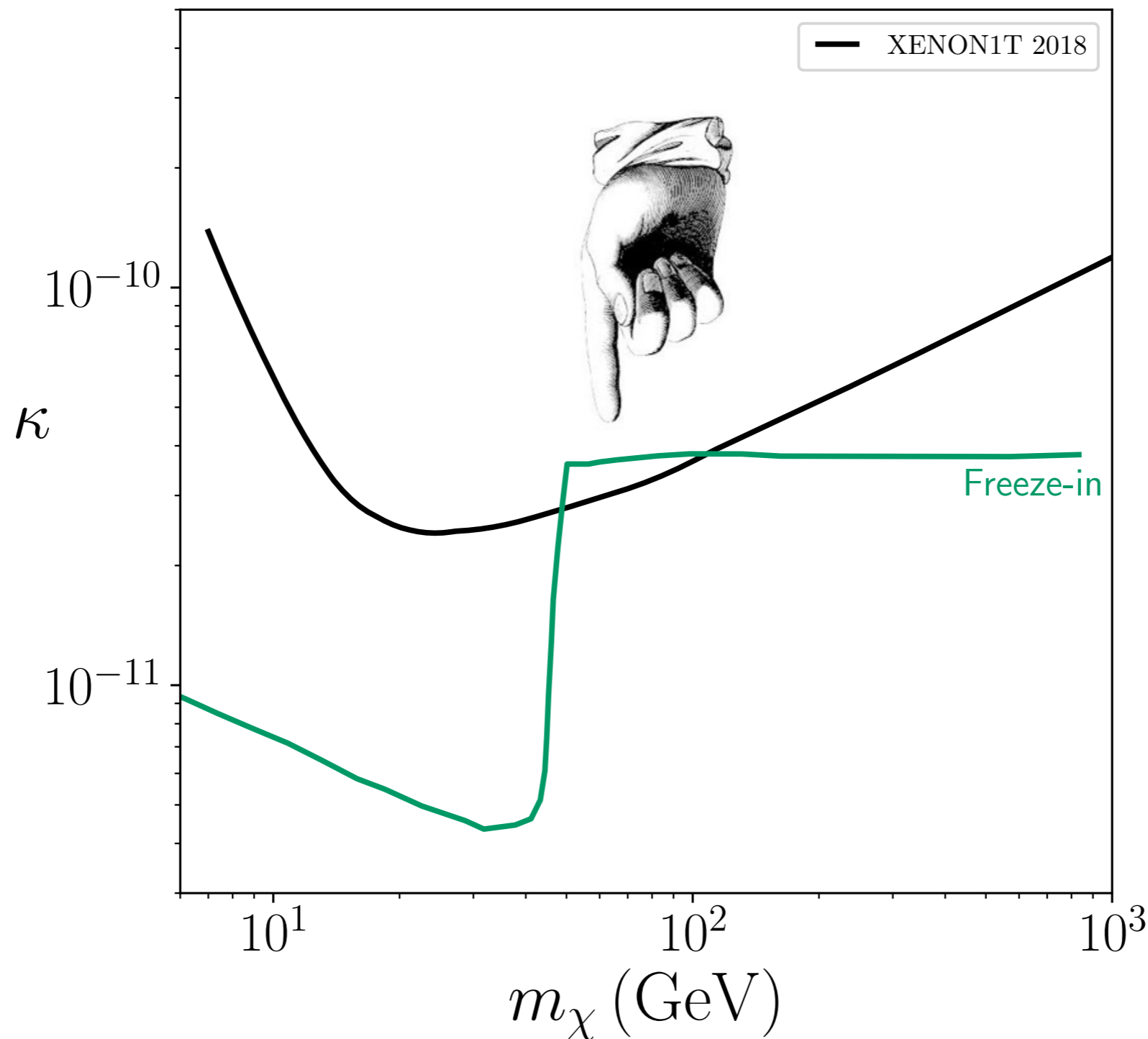


$v \sim 200$ km/s (halo DM)

$$\frac{d\sigma}{dE_R} \propto \frac{m_N \kappa^2 \alpha^2 Z^2}{(2m_N E_R + m_{\gamma'}^2)^2} \sim \frac{1}{E_R^2}$$

large enhancement if $m_{\gamma'} \lesssim 40$ MeV

DIRECT DETECTION IS TESTING FREEZE IN

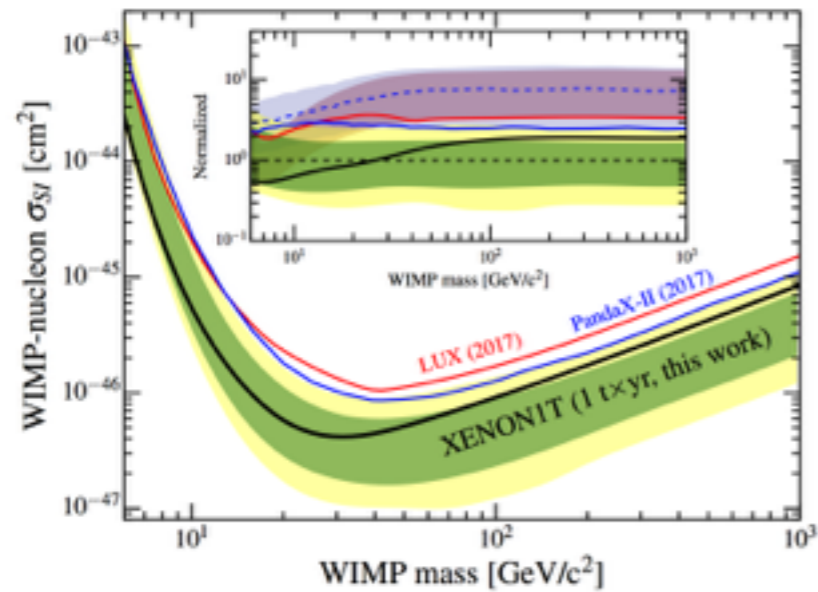


n.b.:
Not the same spectrum as a
WIMP,
Must recast the direct
detection constraints

first direct detection test of a FI scenario

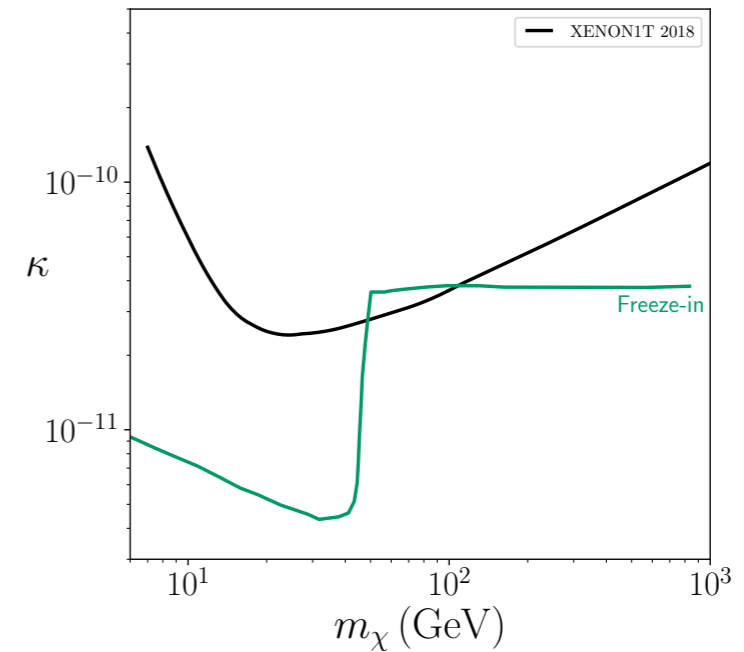
RECASTING DIRECT DETECTION LIMITS

heavy mediator

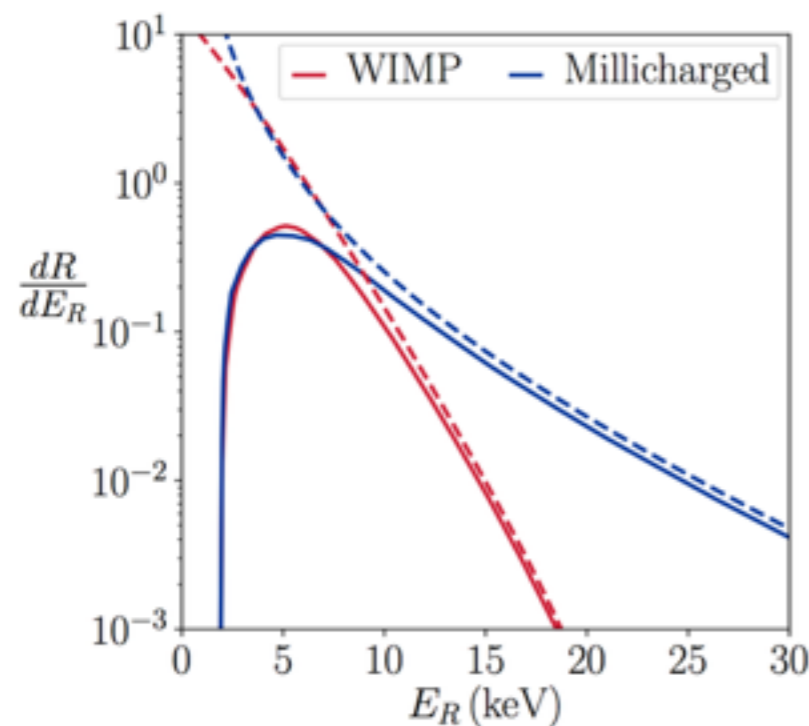


$(m_{\text{dm}}, \sigma_{\text{dm},n})$

light mediator



(m_χ, κ)

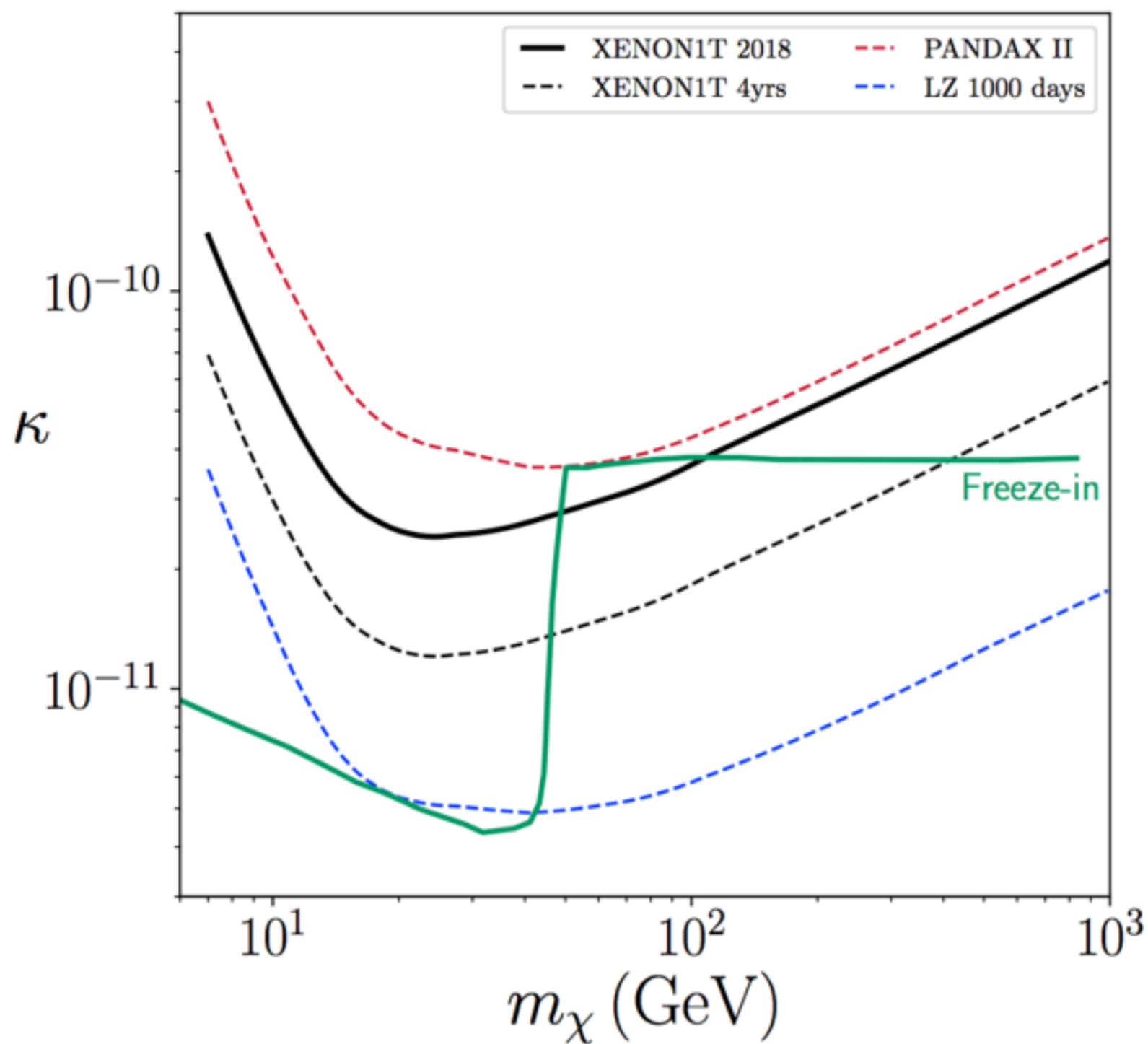


We minimized the differential rate
« quadratic distance »

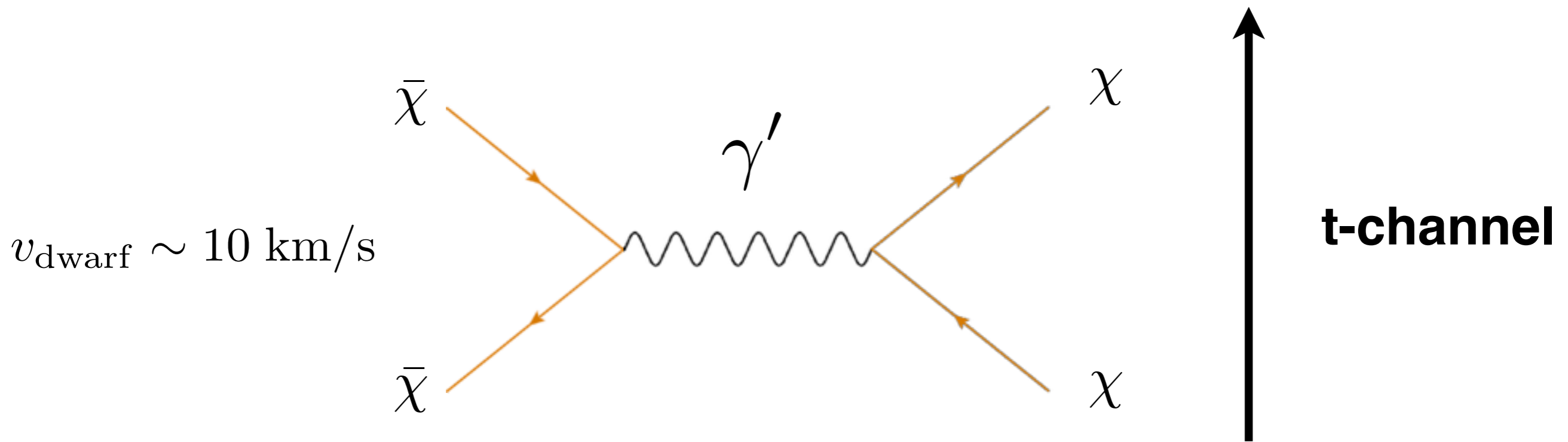
$$\Delta^2 = \frac{1}{R_{\text{exp}}^2} \int dE \epsilon(E)^2 \left(\left(\frac{dR}{dE} \right) - \left(\frac{dR_\chi}{dE} \right) \right)^2$$

XENON1T efficiency

DIRECT DETECTION IS TESTING FREEZE IN

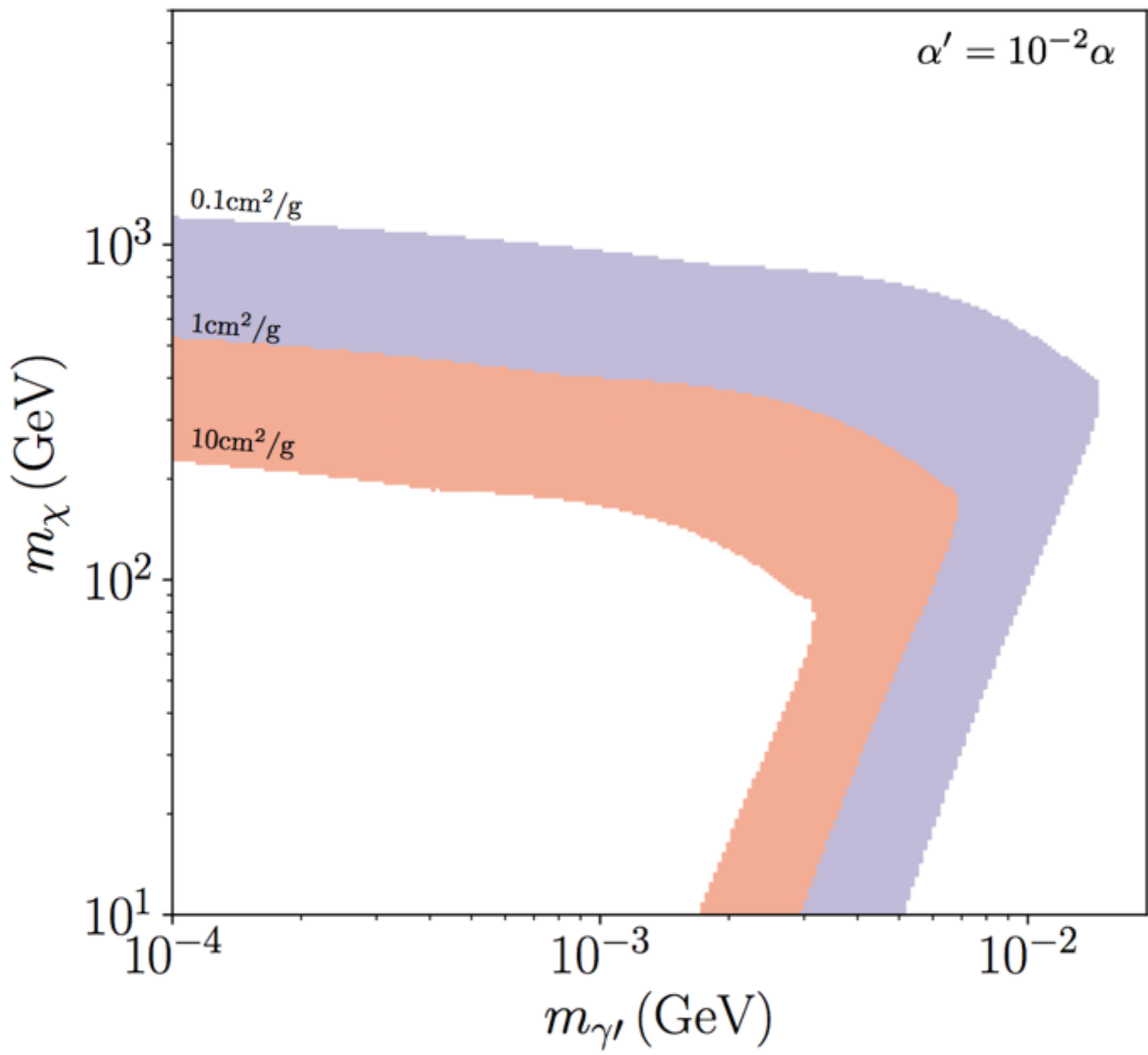


SIDM : RUTHERFORD SCATTERING AGAIN

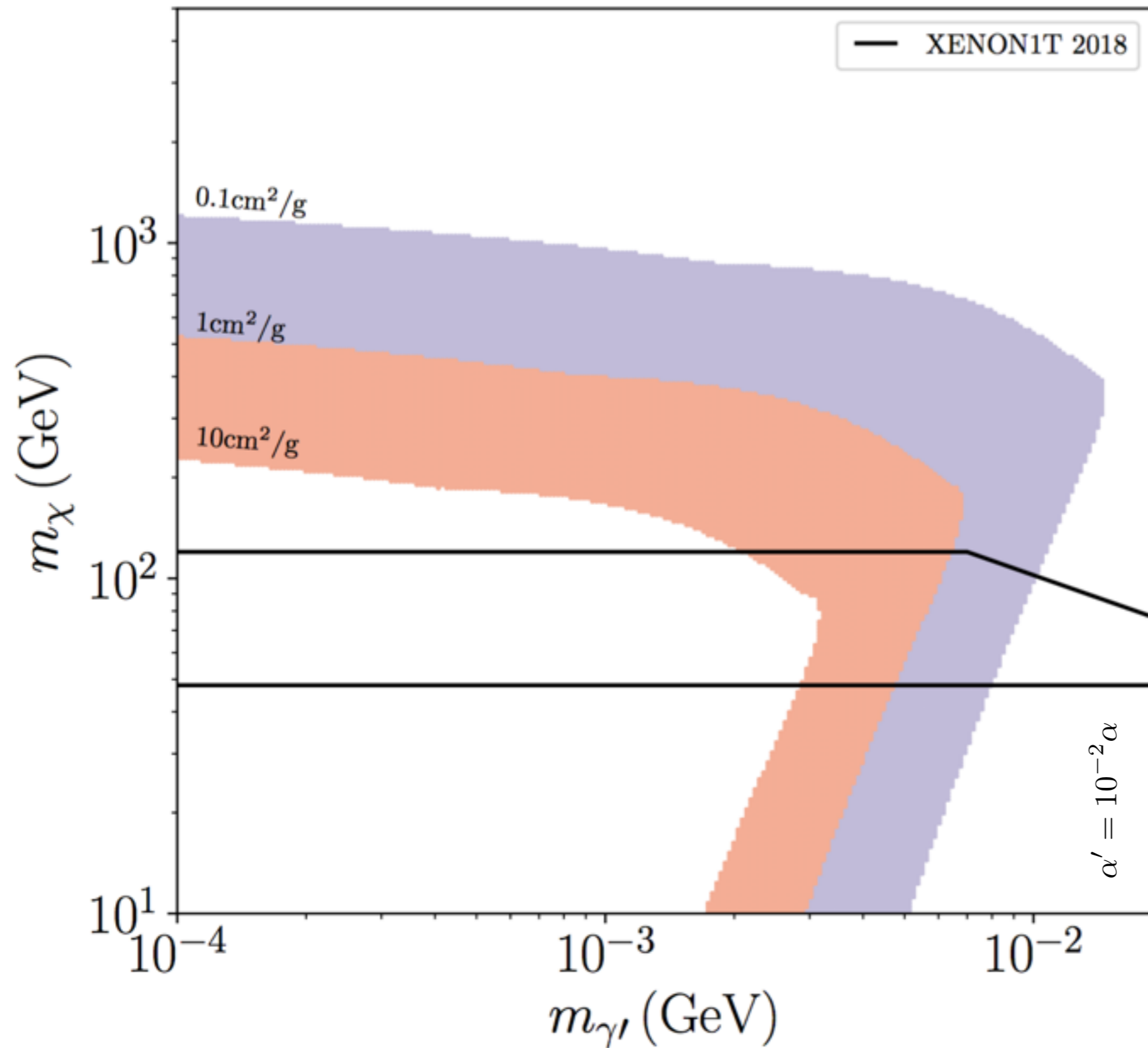


$$\frac{d\sigma}{d\Omega} = \frac{4\alpha'^2 m_\chi^2}{(4m_\chi^2 v^2 \sin^2(\theta/2) + m_{\gamma'}^2)^2} \sim \alpha'^2 \frac{m_\chi^2}{m_{\gamma'}^4}$$

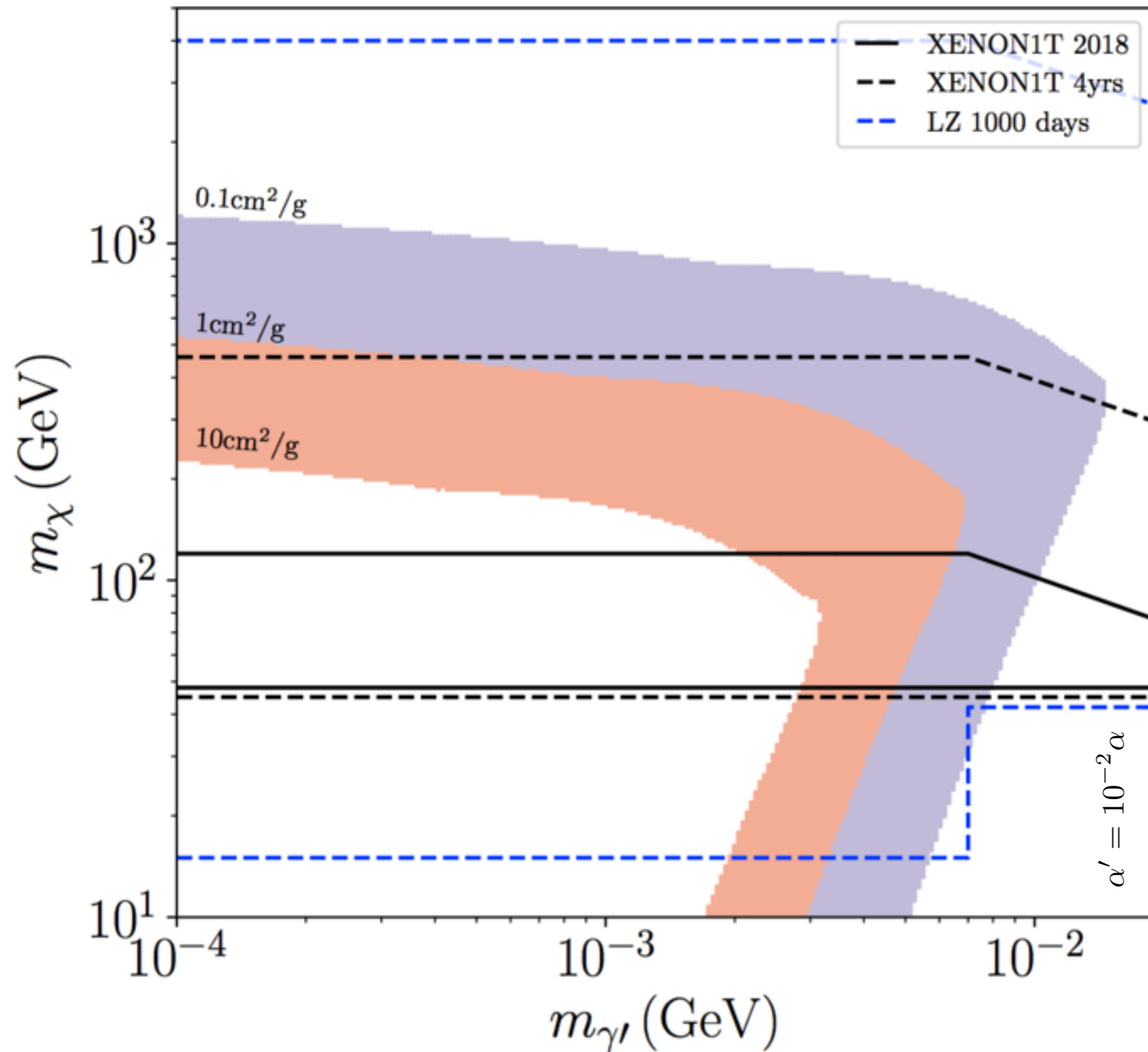
« As big as a barn » for $m_{\gamma'}$ in MeV range



DIRECT DETECTION & SIDM PARAMETER SPACE



DIRECT DETECTION & SIDM PARAMETER SPACE



PART II

SOLAR MASS BH FROM SIDM

$$m_{\text{NS}} \sim m_{\odot}$$
$$N_B \sim 10^{57}$$

$$v_{\text{DM}} \sim 200 \text{ km/s}$$

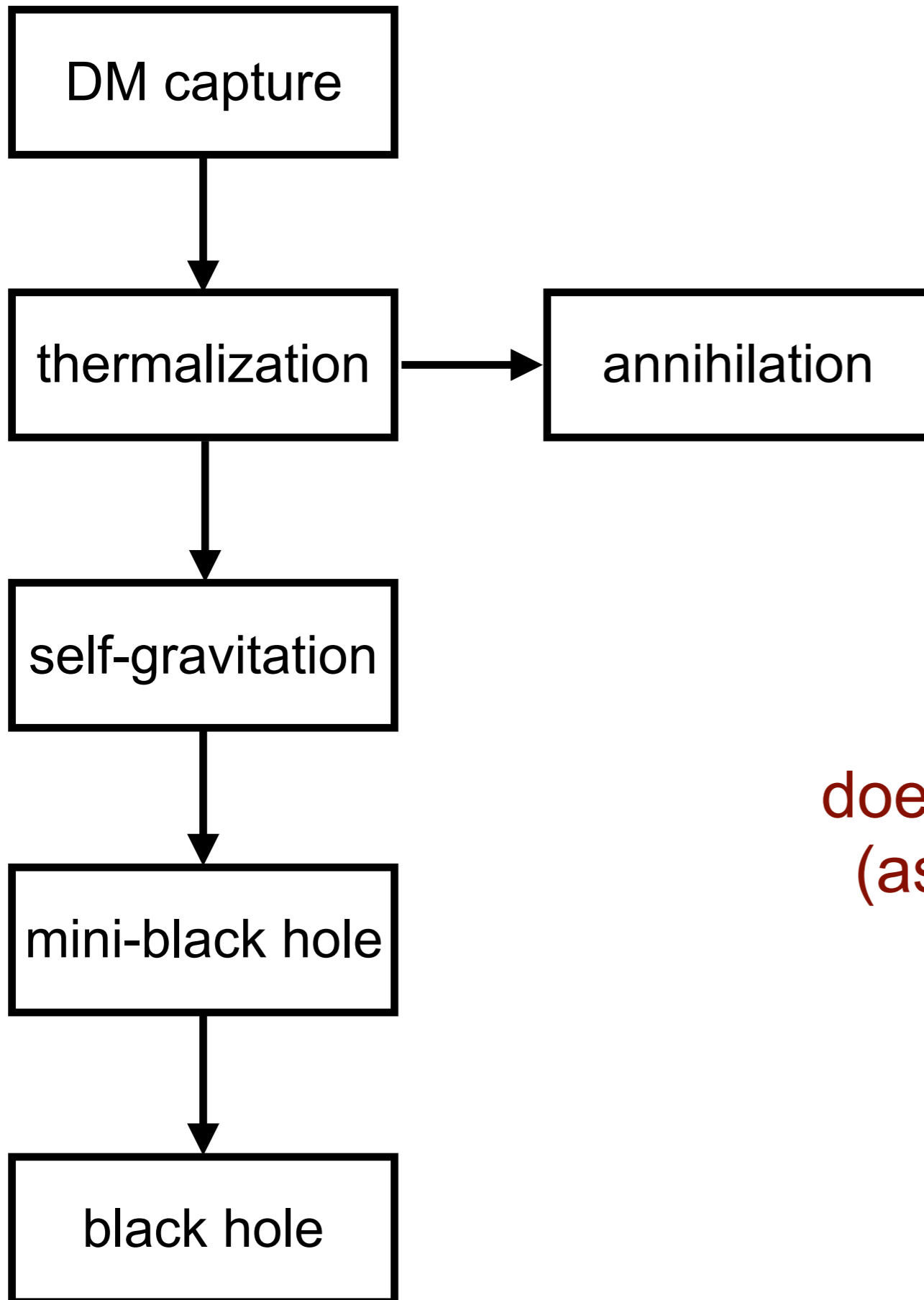


DM

(not to the scale)

Neutron Star
(not to the scale)

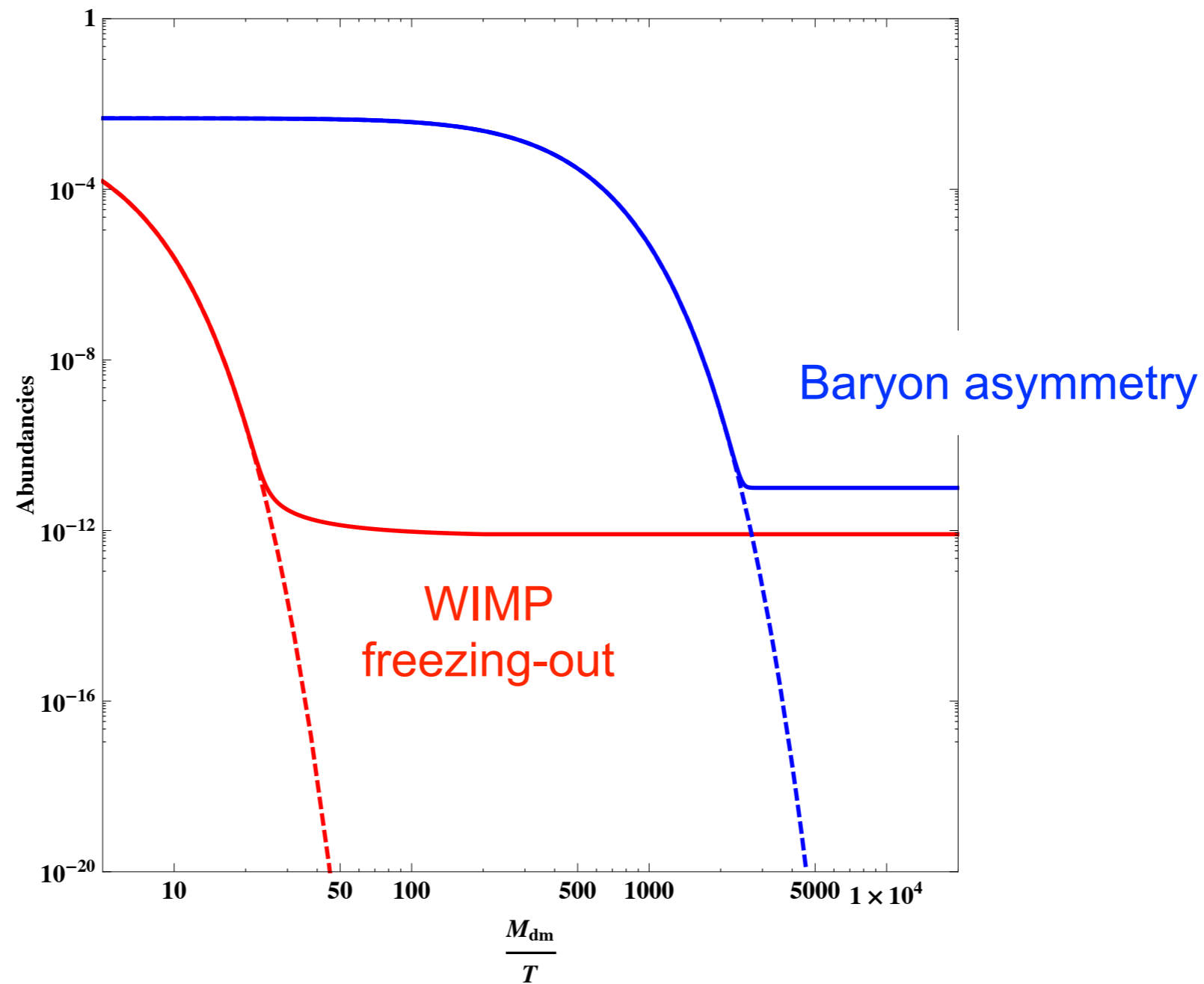
DM + NS \rightarrow BH



assumes DM
does not annihilate
(asymmetric DM)

ASYMMETRIC DM

$$\frac{\Omega_{DM}}{\Omega_B} = \frac{m_{DM}}{m_N} \frac{Y_{DM}}{Y_B} \approx 5 = \mathcal{O}(1)$$



CAPTURE OF DM IN NS

Capture rate

Goldman & Nussinov (1989) - Kouvaris (2008)

$$\frac{dN_{\text{acc}}}{dt} = \sqrt{6\pi} \frac{\rho_{\text{dm}}}{m_{\text{dm}} v_{\text{dm}}} \frac{R_{\star} R_{\text{Sch}}}{1 - R_{\text{Sch}}/R_{\star}} \text{Min} \left(\frac{\sigma}{\sigma_{\text{cr}}}, 1 \right)$$

Critical cross section
(neutron star)

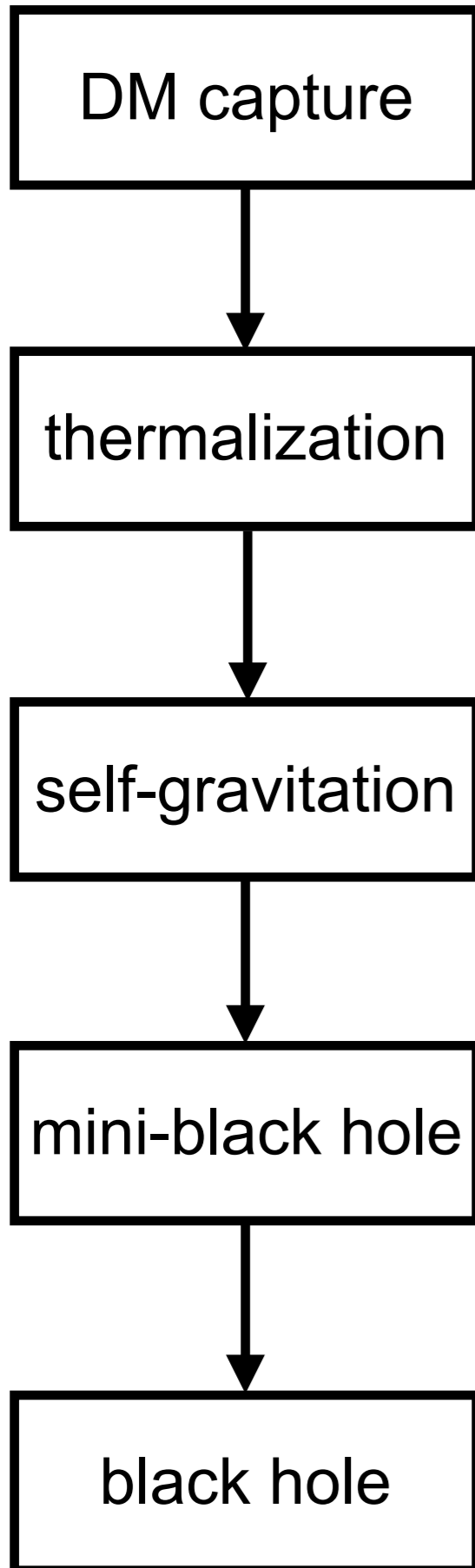
$$\sigma_{\text{cr}} = 0.45 m_n R_{\star}/M_{\star} \approx 1.3 \times 10^{-45} \text{ cm}^2$$

Maximal mass captured

$$M_{\text{acc}} \sim 10^{-15} M_{\odot}$$

$$N_{\text{acc}} \approx 10^{39} (\text{TeV}/m_{\text{dm}})$$

DM + NS \rightarrow BH



$$N_{\text{acc}} \sim 10^{39} (\text{TeV}/m_{\text{dm}})$$

$$t_{\text{th}} \sim 10^2 \text{ yr} \left(\frac{m_{\text{dm}}}{\text{TeV}} \right)^2 \left(\frac{10^{-45} \text{ cm}^2}{\sigma} \right)$$

Goldman & Nussinov (1989)
Kouvaris & Tinyakov (2010)

$$r_{\text{th}} \approx \left(\frac{T_c}{G\rho_c m_{\text{dm}}} \right)^{1/2} \sim 10 \text{ cm} \left(\frac{\text{TeV}}{m_{\text{dm}}} \right)^{1/2} \left(\frac{T_c}{10^5 \text{ K}} \right)^{1/2}$$

$$\frac{N_{\text{acc}} m_{\text{dm}}}{r_{\text{th}}^3} \gtrsim \rho_c \sim 10^{39} \frac{\text{GeV}}{\text{cm}^3} \rightarrow N_{\text{cr}} \gtrsim 10^{38} \left(\frac{\text{TeV}}{m_{\text{dm}}} \right)^{5/2}$$

$$N_{\text{Ch}} \gtrsim \left(\frac{M_{\text{Pl}}}{m} \right)^3 \sim 5 \cdot 10^{48} \left(\frac{\text{TeV}}{m_{\text{dm}}} \right)^3$$

Bondi accretion

$$\frac{dM}{dt} = \frac{4\pi G^2 M^2 \rho_c}{c_s^2} - \frac{\pi^3}{15} R^2 T^4 \rightarrow M_{\text{bh}} \gtrsim 10^{-20} M_{\odot}$$

Hawking radiation

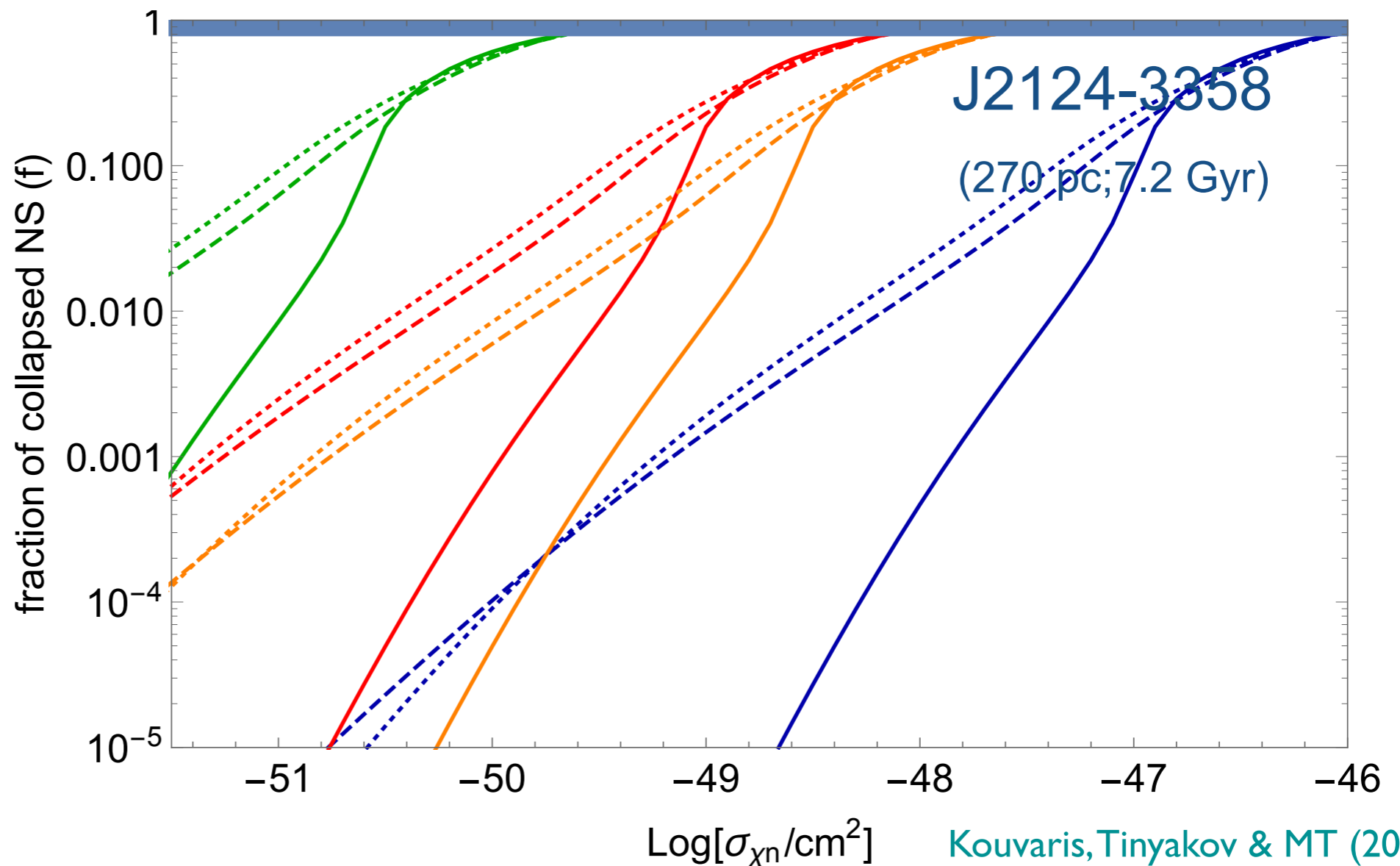
Kouvaris & Tinyakov (2013)

SIDM + NS \rightarrow BH

candidates
that alleviate
CDM issues

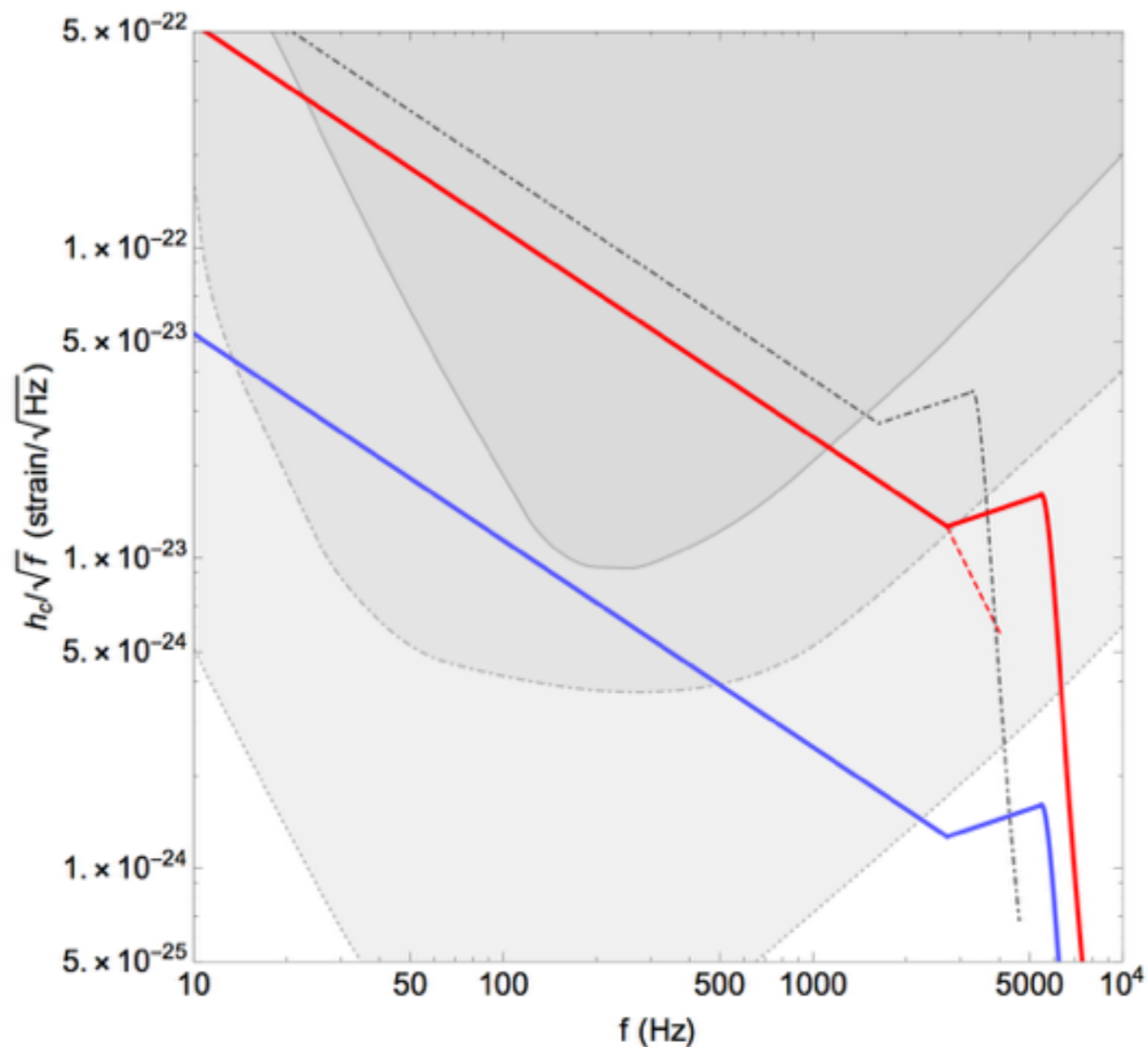
#	α	$\frac{\mu}{\text{MeV}}$	$\frac{m}{\text{TeV}}$	$\frac{\sigma_{10}}{m}$	$\frac{\sigma_{100}}{m}$	N_{cr}	N_{Ch}	$\frac{M_{\text{Ch}}}{M_{\odot}}$
1	10^{-3}	10	1	0.1	$1.6 \cdot 10^{-3}$	$5 \cdot 10^{35}$	$2 \cdot 10^{37}$	10^{-17}
2	10^{-4}	2	0.2	4.5	$4.5 \cdot 10^{-3}$	$2 \cdot 10^{35}$	$7 \cdot 10^{40}$	10^{-14}
3	10^{-4}	1	1	0.2	$5.6 \cdot 10^{-5}$	$3 \cdot 10^{33}$	$6 \cdot 10^{35}$	10^{-18}
4	10^{-3}	3	0.2	6.4	$2.3 \cdot 10^{-1}$	$2 \cdot 10^{35}$	$7 \cdot 10^{39}$	10^{-15}

fraction
of NS
transformed
into BH



SOLAR MASS BINARY MERGERS

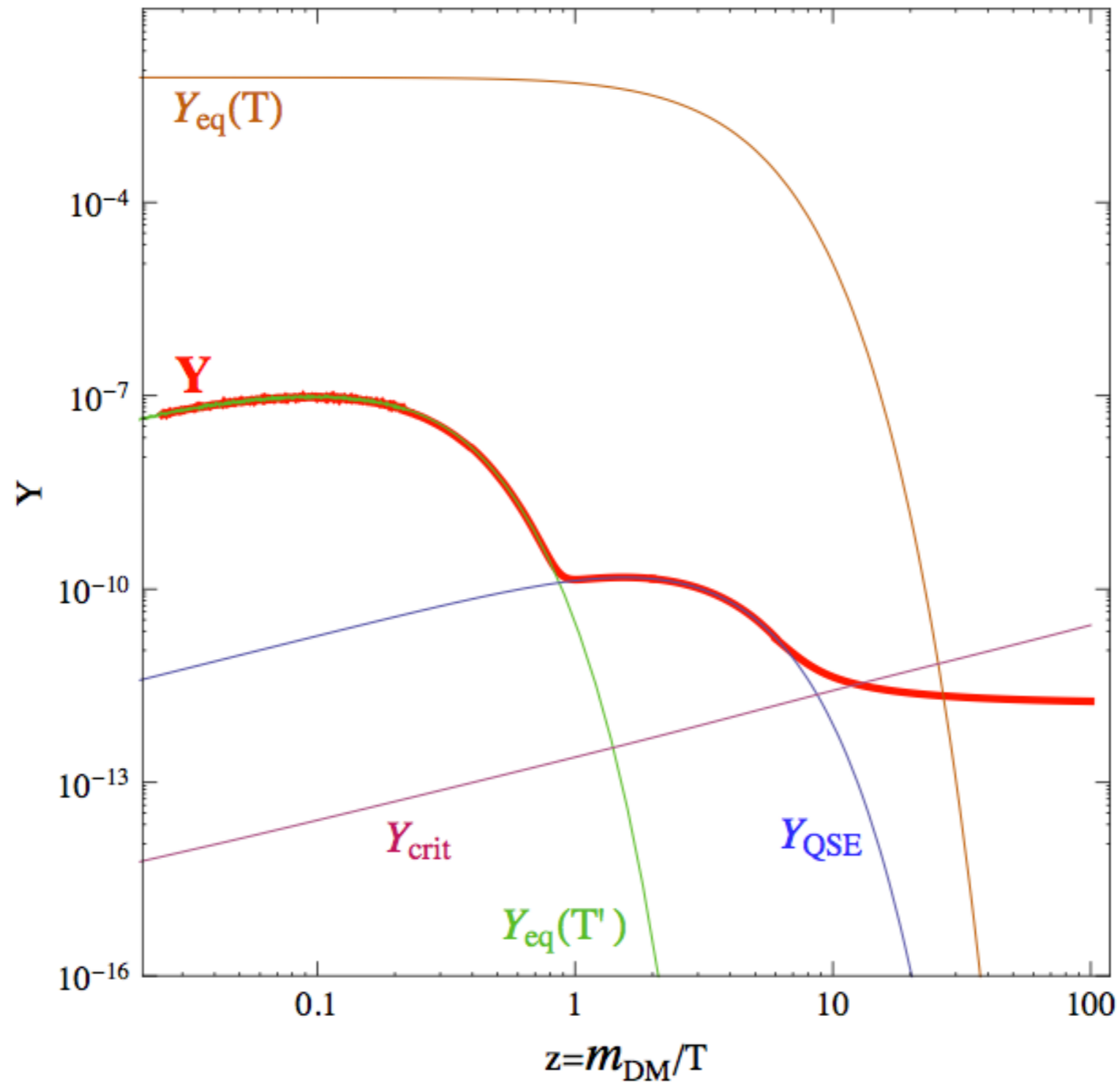
Detectors	BNS range (Mpc)	BNS detections (per year)
LIGO/Virgo	105/80	4 – 80 (2020+)
KAGRA	100	11 – 180 (2024+)
ET	$\sim 5 \cdot 10^3$ ($z \approx 2$)	$\mathcal{O}(10^3 - 10^7)$



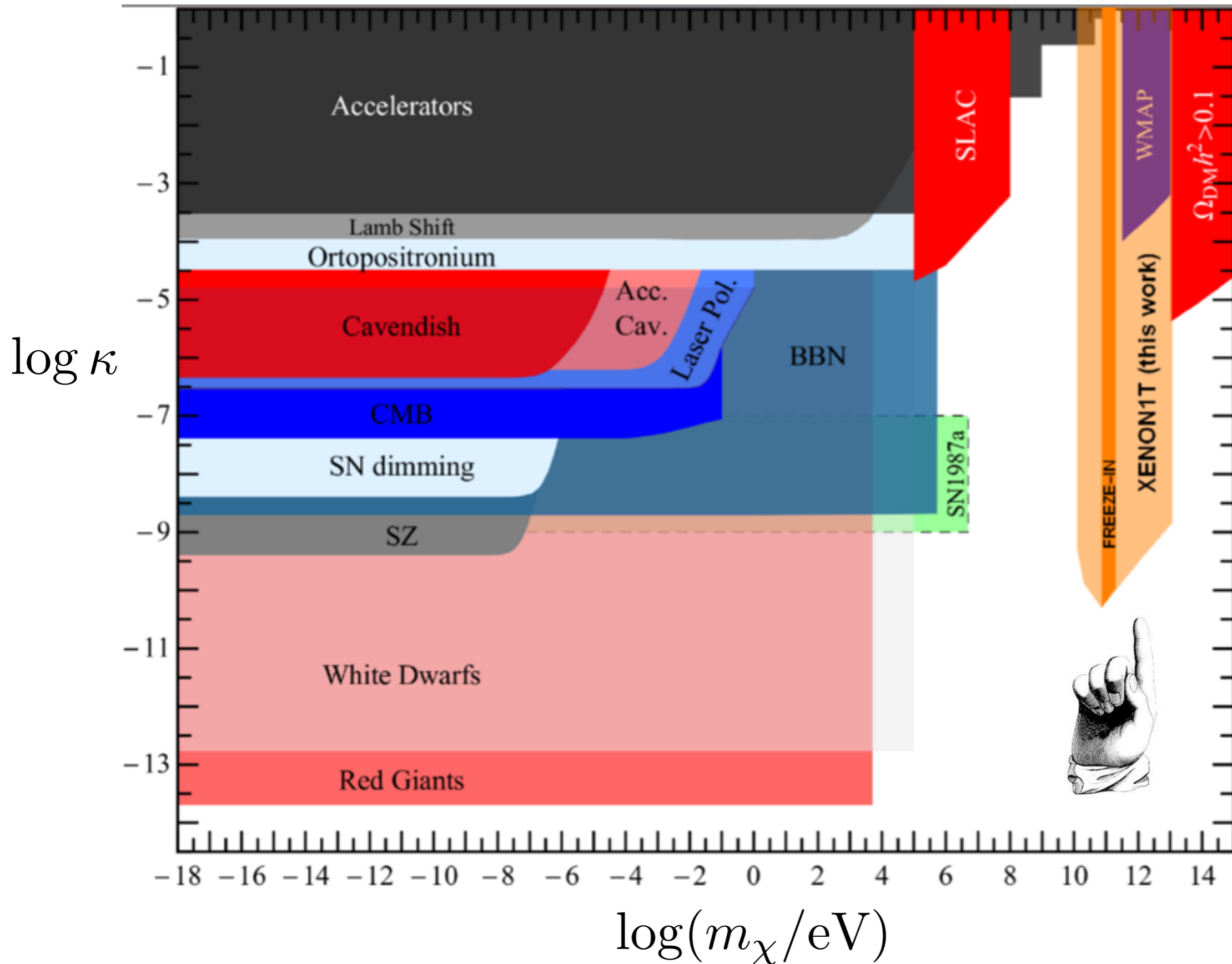
BACKUP

INTERMEDIATE REGIME : RECOMBINATION

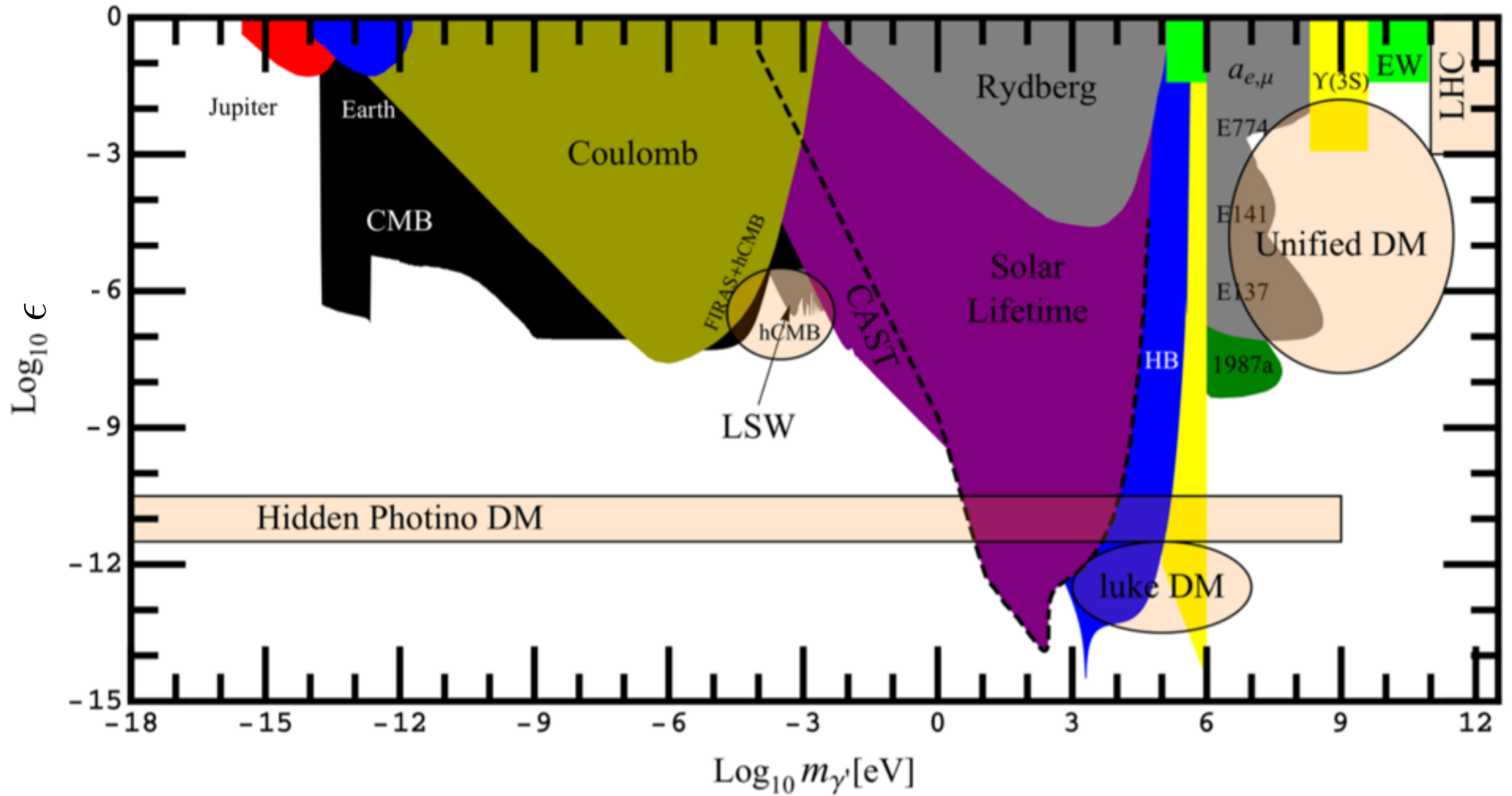
Region.II ($m_{\text{DM}} = 200\text{GeV}$, $\kappa = 10^{-8}$, $\alpha' = 10^{-2.4}$)



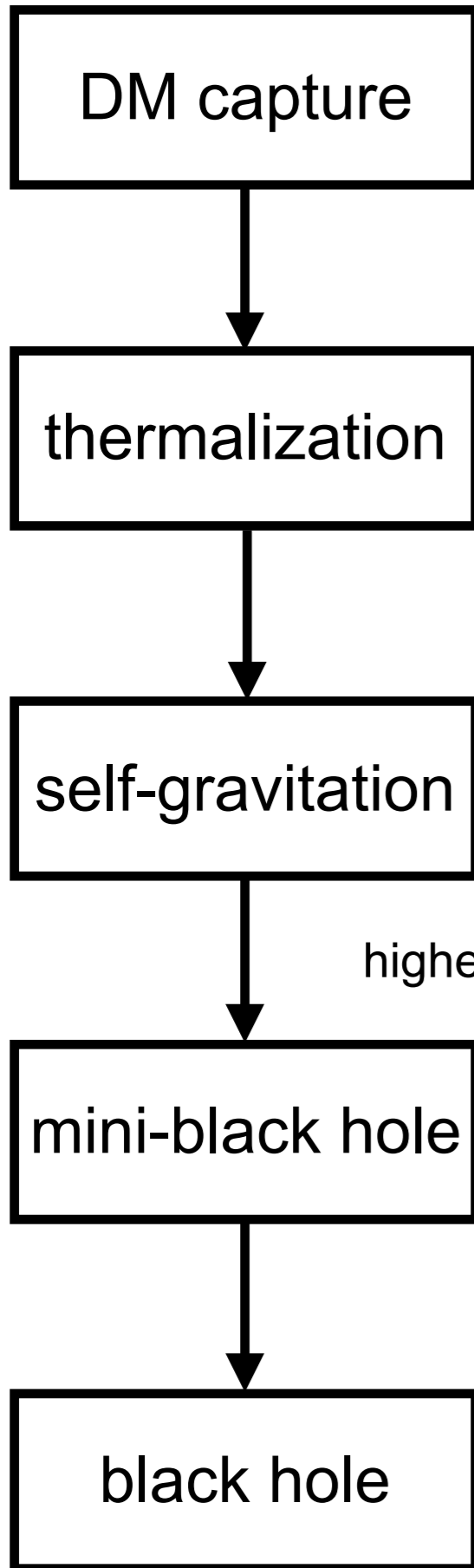
CONSTRAINTS ON MILLICHARGED PARTICLES



CONSTRAINTS ON DARK PHOTON



WIMPS THAT WOULD DEVOUR STARS



$$M_{acc} \sim 10^{39} \text{ TeV} \left(\frac{\rho_{dm}}{\text{GeV}/\text{cm}^3} \right) \left(\frac{t}{\text{Gyr}} \right) \text{Min} \left(\frac{\sigma}{\sigma_{cr}}, 1 \right)$$



$$t_{th} \sim 10^2 \text{ yr} \left(\frac{m_{dm}}{\text{TeV}} \right)^2 \left(\frac{10^{-45} \text{ cm}^2}{\sigma} \right)$$

Goldman & Nussinov (1989)
Kouvaris & Tinyakov (2010)

$$r_{th} \approx \left(\frac{T_c}{G\rho_c m_{dm}} \right)^{1/2} \sim 10 \text{ cm} \left(\frac{\text{TeV}}{m_{dm}} \right)^{1/2} \left(\frac{T_c}{10^5 \text{ K}} \right)^{1/2}$$

$$\frac{N_{acc} m_{dm}}{r_{th}^3} \gtrsim \rho_c \sim 10^{39} \frac{\text{GeV}}{\text{cm}^3} \rightarrow N_{cr} \gtrsim 10^{38} \left(\frac{\text{TeV}}{m_{dm}} \right)^{3/2}$$

higher density, higher T, further cooling, further collapse

$$\frac{GNm_{dm}^2}{R} \gtrsim k_F \sim \frac{N^{1/3}}{R} \rightarrow N_{Ch} \gtrsim \left(\frac{M_{Pl}}{m} \right)^3 \sim 5 \cdot 10^{48} \left(\frac{\text{TeV}}{m_{dm}} \right)^3$$

**Overcoming Fermi pressure
requires $m_{dm} \gtrsim 1000 \text{ TeV}$**



Bramante, Linden & Tsai (2017); Kouvaris, Tinyakov & MT (2018)

PRIMORDIAL BLACK HOLES ?

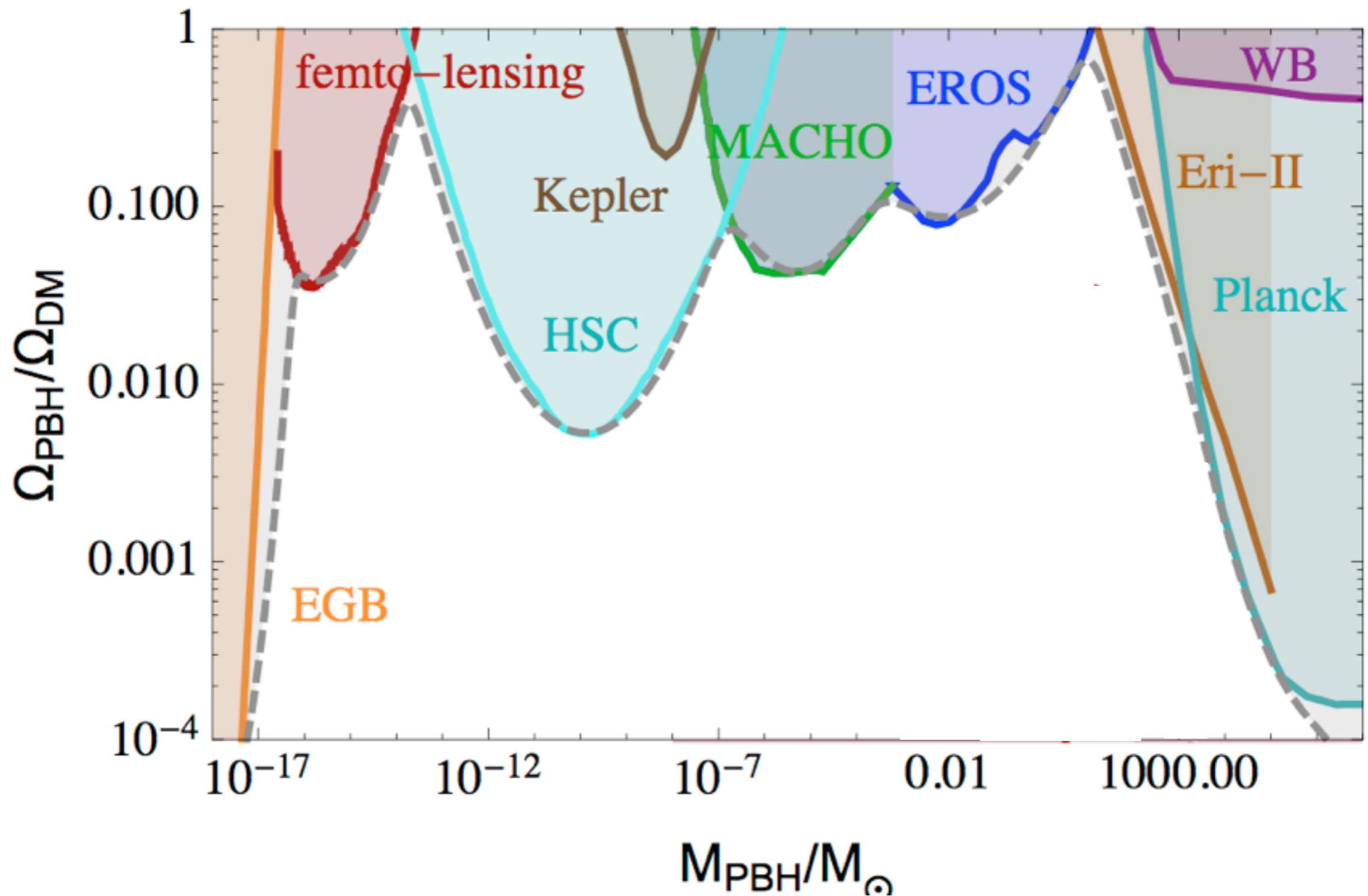


Figure from Garcia-Bellido & Clesse (2018)