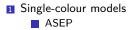
# Transition Probabilities in the Multi-species Asymmetric Exclusion Process

#### William Mead

School of Mathematics and Statistics The University of Melbourne

Based on work with Jan de Gier and Michael Wheeler (arXiv:2109.14232)

Randomness, Integrability and Universality Galileo Galilei Institute for Theoretical Physics, Florence 4 May 2022



- ASEP
- Transition probabilities from vertex models

ASEP

Transition probabilities from vertex models

2 Multi-colour models

ASEP

Transition probabilities from vertex models

2 Multi-colour models

Transition probabilities from higher rank vertex models

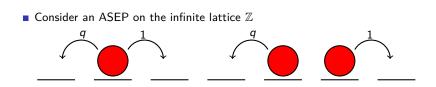
ASEP

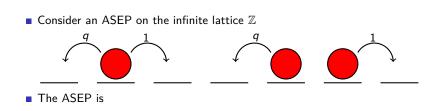
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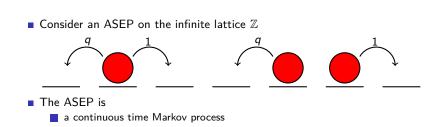
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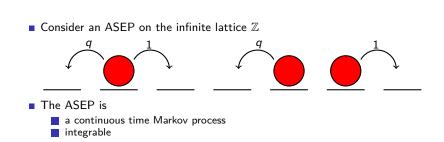
Transition probabilities from higher rank vertex models

3 Progress on two-colours

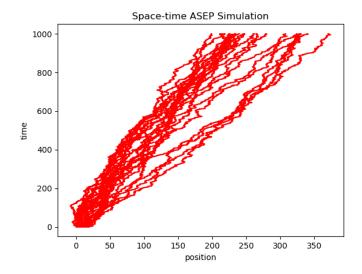








# Asymmetric simple exclusion process



# Time-evolution

The probability satisfies the time-evolution equation

$$rac{\mathsf{d}}{\mathsf{d} t}\mathbb{P}(
u;t) = \sum_{\lambda
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#### Theorem (Schütz (1997), Tracy and Widom (2008))

Given initial  $\mu$  and final conditions  $\nu$  the transition probability on  $\mathbb Z$  is given by

$$\mathbb{P}_{t}^{TASEP}(\mu \to \nu) = \oint_{0} \prod_{i=1}^{n} \frac{\mathrm{d}z_{i}}{2\pi \mathrm{i}} \sum_{\pi \in S_{n}} (-1)^{|\pi|} \prod_{i=1}^{n} \left(\frac{1-z_{i}}{1-z_{\pi_{i}}}\right)^{i} e^{(z_{i}^{-1}-1)t} z_{\pi_{i}}^{\nu_{i}} z_{i}^{-\mu_{i}-1}$$

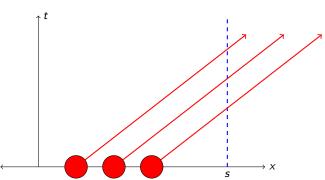
which satisfies the time-evolution with initial condition

$$\mathbb{P}_0^{TASEP}(\mu \to \nu) = \prod_{i=1}^n \delta_{\nu_i,\mu_i}.$$

• We choose the step initial condition  $\mu_i = i$ 

# TASEP crossing probability

- We choose the step initial condition  $\mu_i = i$
- Define probability of *n* particles crossing a wall at position  $s \in \mathbb{N}$  as



$$P_{\text{cross}}(s) = \mathbb{P}(s \leq \nu_1 < \nu_2 < \cdots < \nu_n).$$

# Limiting behaviour of TASEP

• We may find this probability as a Fredholm determinant

$$P_{cross}(s) = \det(1 - K_n(x, y))_{\ell^2(\mathbb{N})},$$

where

$$K_n(x,y) = \sum_{k=0}^{n-1} \phi_k(x)\psi_k(y).$$

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• The functions  $\phi_k, \psi_k$  are defined as contour integrals

$$\phi_k(x) = \oint_1 \frac{\mathrm{d}\eta}{2\pi \mathrm{i}} \frac{\eta^{k-x} e^{-\eta t}}{(\eta-1)^{k+1}}, \qquad \psi_k(y) = \oint_0 \frac{\mathrm{d}\zeta}{2\pi \mathrm{i}} \frac{(\zeta-1)^k e^{\zeta t}}{\zeta^{k-y+2}}.$$

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We change k = vt - κt<sup>1/3</sup> for κ > 0 and through a steepest decent analysis we find

$$\lim_{t \to \infty} f_1(t) \phi_{vt - \kappa t^{1/3}} \left( \xi_1 t^{1/3} \right) = \operatorname{Ai}(\kappa + \xi_1),$$
$$\lim_{t \to \infty} f_2(t) \psi_{vt - \kappa t^{1/3}} \left( \xi_2 t^{1/3} \right) = \operatorname{Ai}(\kappa + \xi_2),$$

which converge uniformly for some unimportant functions  $f_1, f_2$ .

• With the change  $k = vt - \kappa t^{1/3}$  and n = vt as  $t \to \infty$ 

$$\mathcal{K}_n(x,y) = \sum_{k=0}^{n-1} \phi_k(x) \psi_k(y) \sim \int_0^\infty \operatorname{Ai}(\kappa + \xi_1) \operatorname{Ai}(\kappa + \xi_2) \mathrm{d}\kappa =: \mathcal{K}_{\operatorname{Airy}}(\xi_1, \xi_2)$$

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This function satisfies

$$\det(1 - K_{\operatorname{Airy}}(\xi_1, \xi_2))_{L^2(\mathbb{R} \ge \alpha)} = F_2(\alpha)$$

where  $F_2$  is the Tracy-Widom distribution of the largest eigenvalue for the Gaussian unitary ensemble (GUE).

#### Theorem (Johansson, 2000)

For the step initial condition, when setting n = vt, we obtain the limit

$$\lim_{t\to\infty}\mathbb{P}\left(\frac{\nu_1(t)-\mathsf{v}t}{c_0t^{1/3}}\geq\alpha\right)=\mathsf{F}_2(\alpha),$$

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- The TASEP lies within the KPZ universality class for the case of step initial condition.
- There are very few rigorous similar results for models with distinguishable particles. Our work provides a starting point for their asymptotic analysis.

• We start with a square lattice and draw paths between vertices

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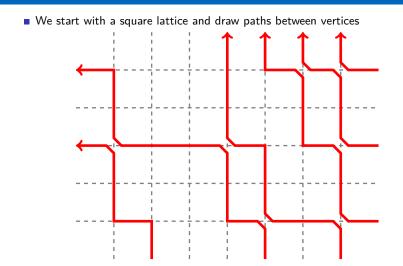
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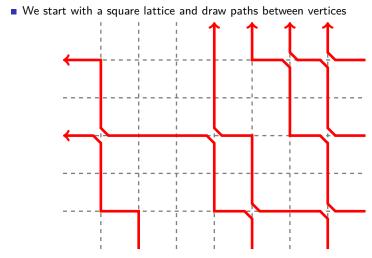
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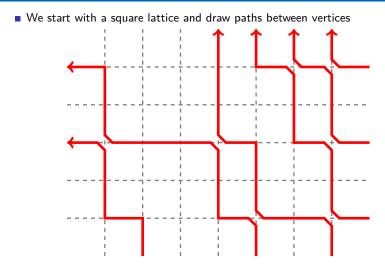
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• We wish to compute the partition function of this model.

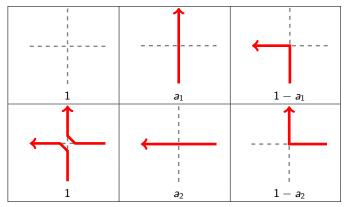


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Schütz's TASEP transition probability can be realised as the partition function of the stochastic six-vertex model.

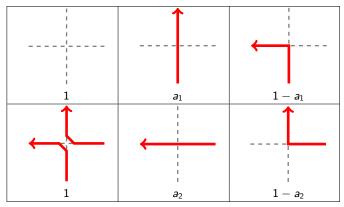
## Six-vertex model weights

• We are allowed to have the following vertex Configurations with Boltzmann weights



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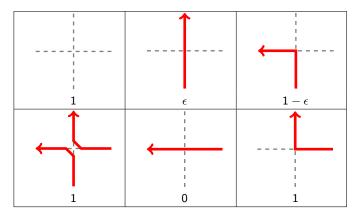
 The classical partition function can be computed by summing over connected path configurations

$$\mathcal{Z} = \sum_{\Omega} a_1^{\#} (1-a_1)^{\#} a_2^{\#} (1-a_2)^{\#}.$$

• Introduce a small parameter  $\epsilon > 0$ 

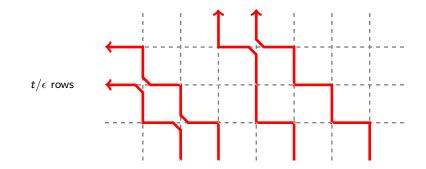
## Six-vertex model weights

- Introduce a small parameter  $\epsilon > 0$
- Set  $a_2 = 0$  and  $a_1 = \epsilon$  which gives the weights as

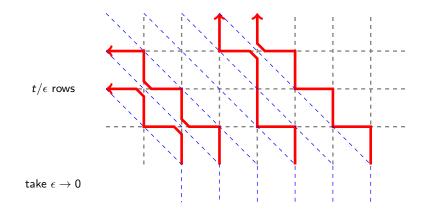


 $t/\epsilon$  rows

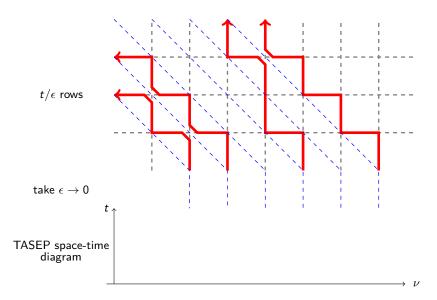
# Reduction to TASEP



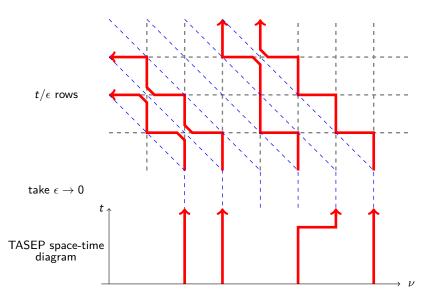
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• The TASEP transition probability can be realised as the partition function of the stochastic six-vertex model.

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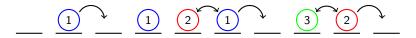
Proposition

$$\lim_{\epsilon \to 0} \mathbb{P}^{\delta VM}[\mu \to \nu - (t/\epsilon)^n] \big|_{\ell = t/\epsilon, \mathsf{a}_1 = \epsilon, \mathsf{a}_2 = 0} = \mathbb{P}_t^{\mathsf{TASEP}}(\mu \to \nu)$$

• We investigate a multi-species version of the TASEP.

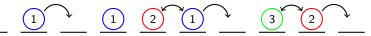
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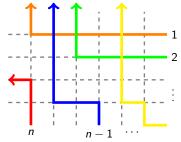
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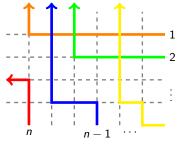
• We aim to recover transition probabilities for the *r*-TASEP from a vertex model.

• We also consider a multi-coloured higher rank version of the stochastic six vertex model with  $U_q\left(\widehat{\mathfrak{sl}}_{n+1}\right)$  symmetry.

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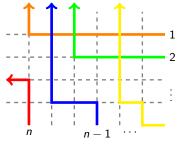


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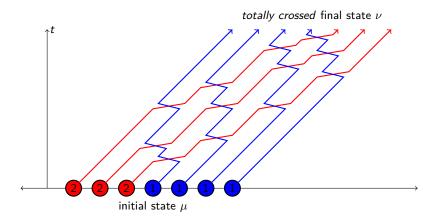
- We can also reduce the multi-coloured partition function to the rainbow TASEP.
- This rainbow TASEP can be partially symmetrized into the form of a general *r*-species TASEP.

$$\mathbb{P}_t^{r\text{-TASEP}}(\mu \to \nu)$$

The simplest multi-species model is the 2-TASEP

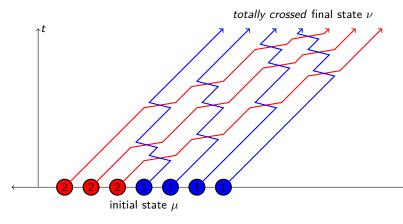
## The Two-species Model

- The simplest multi-species model is the 2-TASEP
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• We consider *n* total particles with *m* of them being type 2.

• The 2-TASEP transition probability simplifies under total crossing events.

### ■ The 2-TASEP transition probability simplifies under total crossing events.

### Proposition

$$\mathbb{P}_{t}^{2\text{-TASEP}}(\mu \to \nu) = \oint \prod_{i=1}^{m} \frac{\mathrm{d}z_{i}}{2\pi \mathrm{i}} \prod_{j=1}^{n-m} \frac{\mathrm{d}w_{j}}{2\pi \mathrm{i}}$$

$$\times \prod_{i=1}^{m} \frac{e^{(z_{i}^{-1}-1)t}}{(1-z_{i})^{n-m}} \prod_{i=1}^{n-m} e^{(w_{i}^{-1}-1)t} \prod_{i=1}^{m} \prod_{j=1}^{n-m} (w_{j}-z_{i})$$

$$\times \det \left( z_{i}^{\nu_{n-m+j}-\mu_{i}-1} (1-z_{i})^{i-j} \right)_{1 \le i,j \le m} \det \left( w_{i}^{\nu_{j}-\mu_{m+i}-1} (1-w_{i})^{i-j} \right)_{1 \le i,j \le n-m}$$

### The 2-TASEP transition probability simplifies under total crossing events.

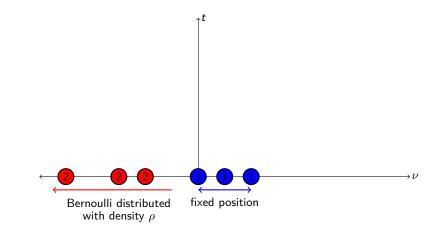
#### Proposition

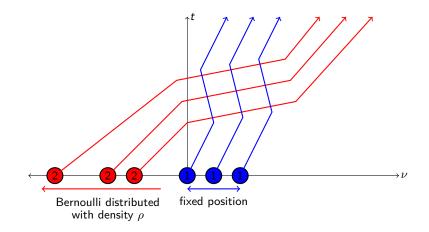
$$\mathbb{P}_{t}^{2\text{-TASEP}}(\mu \to \nu) = \oint \prod_{i=1}^{m} \frac{\mathrm{d}z_{i}}{2\pi \mathrm{i}} \prod_{j=1}^{n-m} \frac{\mathrm{d}w_{j}}{2\pi \mathrm{i}}$$

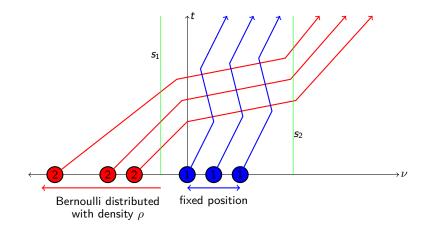
$$\times \prod_{i=1}^{m} \frac{\mathrm{e}^{(z_{i}^{-1}-1)t}}{(1-z_{i})^{n-m}} \prod_{i=1}^{n-m} \mathrm{e}^{(w_{i}^{-1}-1)t} \prod_{i=1}^{m} \prod_{j=1}^{n-m} (w_{j}-z_{i})$$

$$\times \det \left( z_{i}^{\nu_{n-m+j}-\mu_{i}-1} (1-z_{i})^{i-j} \right)_{1 \le i,j \le m} \det \left( w_{i}^{\nu_{j}-\mu_{m+i}-1} (1-w_{i})^{i-j} \right)_{1 \le i,j \le n-m}$$

This result generalises to *r*-species using the vertex model approach.





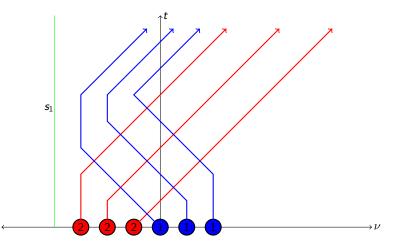


### Proposition

The Bernoulli crossing probability is given by

$$\mathbb{P}_{t}^{\beta\text{-cross}}(s_{1}, s_{2}) = \frac{\rho^{m}}{m!} \oint_{0,1,1-\rho} \prod_{i=1}^{m} \frac{\mathrm{d}z_{i}}{2\pi \mathrm{i}} \prod_{i\neq j} (z_{j} - z_{i}) \prod_{i=1}^{m} \frac{e^{(z_{i}-1)t} z_{i}^{-s_{2}-m+1}}{(z_{i}-1)^{n} (z_{i}-1+\rho)} \\ \times \oint_{0,1} \prod_{i=1}^{n-m} \frac{\mathrm{d}w_{i}}{2\pi \mathrm{i}} \prod_{i=1}^{m} \prod_{j=1}^{n-m} (z_{i} - w_{j}) \prod_{i=1}^{n-m} \frac{-e^{(w_{i}-1)t} w_{i}^{-s_{1}-m}}{(1-w_{i})^{n-m-i+1}} \\ \times \det \left(w_{i}^{j-1} - w_{i}^{n-m+s_{1}-s_{2}-1}\right)_{1 \leq i,j \leq n-m}.$$

What if we take  $s_1 < -m$ ?



Since the type 1 particles move backwards when overtaken, all possible total crossing configurations contribute towards  $\mathbb{P}_t^{\text{B-cross}}(s_1, s_2)$ .

### Proposition

When  $s_1 \leq -m$  the (n - m)-fold integral over type 1 particles collapses into 1 integral

$$\begin{split} \mathbb{P}_{t}^{\beta\text{-cross}}(s_{1},s_{2}) &= \rho^{m} \oint_{0} \frac{\mathrm{d}w}{2\pi \mathrm{i}} \frac{e^{(w-1)t} w^{n-2m-s_{2}-1}}{w-1} \\ &\times \det \left( \oint_{0,1,1-\rho} \frac{\mathrm{d}z}{2\pi \mathrm{i}} \frac{e^{(z-1)t} z^{i+j-s_{2}-m-1}}{(z-1)^{m+1}(z-1+\rho)} (w-z) \right)_{1 \leq i,j \leq m}. \end{split}$$

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  - This has only been investigated in very limited circumstances.
  - Recent work (Nejjar, 2020) investigates the asymptotics with one second-class particle, which we expect to recover.
- These multi-species transition probabilities are useful for constructing higher-rank stochastic dualities and their expectations (work in progress).