# Global asymptotics of particle systems at high temperature

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Based on work w/ Florent Benaych-Georges & Vadim E. Gorin

Eigenvalues of Gaussian beta ensemble (G $\beta$ E)

General theorems and proof ideas

Further questions  $\gamma$ -Semifree Probability Discrete Ensembles

# Plan of the talk

#### Eigenvalues of Gaussian beta ensemble (G $\beta$ E)

General theorems and proof ideas

Further questions γ-Semifree Probability Discrete Ensembles

# Gaussian Unitary Ensemble (GUE)

 $GUE_N$  is the probability measure on  $\mathcal{H}_N := \{A \in \mathbb{C}^{N \times N} \mid A^* = A\}$ on Hermitian matrices with density

$$\frac{1}{\mathcal{Z}_{N}} \cdot \exp\left\{-\frac{\operatorname{Trace}(A^{2})}{2}\right\} \prod_{i=1}^{N} da_{ii} \prod_{1 \le i < j \le N} d\Re a_{ij} d\Im a_{ij}.$$

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If 
$$A \in \mathcal{H}_N$$
 is  $\mathsf{GUE}_N$ -distributed, its real eigenvalues  
 $x_1 \leq x_2 \leq \cdots \leq x_{N-1} \leq x_N$ 

are random and their distribution is:

$$\mathsf{Eigen}_{N}^{(2)}(x_{1},\cdots,x_{N}) = \frac{1}{Z_{N}^{(2)}} \prod_{1 \le i < j \le N} (x_{i} - x_{j})^{2} \prod_{k=1}^{N} e^{-\frac{1}{2}x_{k}^{2}}$$

We call this the GUE (eigenvalue) density.

# Global asymptotics of Hermite N-particle ensemble

# Consider the empirical measures $\mu_N := \frac{1}{N} \sum_{i=1}^N \delta_{\frac{x_i}{\sqrt{N}}}, \quad \text{where } x_1 \leq \cdots \leq x_N \text{ is } \mathsf{Eigen}_N^{(2)} - \mathsf{distributed}.$

# Theorem (Wigner '55)

The (random) probability measures  $\mu_N$  converge weakly, in probability, to the semicircle distribution — with density

$$s(t) := \mathbf{1}_{\{-2 \le t \le 2\}} \cdot rac{\sqrt{4-t^2}}{2\pi},$$

i.e. for any  $f \in C_b(\mathbb{R})$ :  $\lim_{N \to \infty} \mathbb{E}_{x_1 \le \dots \le x_N} \left[ \int_{\mathbb{R}} f(t) \mu_N(dt) \right] = \int_{-2}^2 f(t) s(t) \, dt.$ 

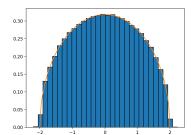
# Global asymptotics of GUE eigenvalues

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# Eigenvalues of Gaussian Beta Ensemble $(G\beta E)$

For general  $\beta \ge 0$ , we study the random *N*-tuple  $x_1 \le x_2 \le \cdots \le x_{N-1} \le x_N$ 

determined by the probability measure

$$\mathsf{Eigen}_{N}^{(\beta)}(x_{1}, \cdots, x_{N}) = \frac{1}{Z_{N}^{(\beta)}} \prod_{1 \le i < j \le N} (x_{i} - x_{j})^{\beta} \prod_{k=1}^{N} e^{-\frac{1}{2}x_{k}^{2}}$$

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#### Why?

1. For  $\beta = 1 \& 4$ , it's the eigenvalue density of Gaussian Orthogonal Ensemble (GOE) & Gaussian Symplectic Ensemble (GSE).

- 2. Relation with particle systems in physics (log-gas);
- $\beta$  is called the inverse temperature.

# Global asymptotics of $G\beta E$ eigenvalues

Nothing changes if  $\beta > 0$  is fixed: as  $N \to \infty$ , then  $\mu_N := \frac{1}{N} \sum_{i=1}^N \delta_{\frac{x_i}{\sqrt{N}}}, \quad \text{where } x_1 \leq \cdots \leq x_N \text{ is Eigen}_N^{(\beta)} - \text{distributed},$ 

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Eigen<sub>N</sub><sup>(\beta=0)</sup>(x<sub>1</sub>,...,x<sub>N</sub>) = 
$$\frac{1}{(2\pi)^{N/2}} \prod_{k=1}^{N} e^{-\frac{1}{2}x_k^2}$$
.

Then  $x_1, \dots, x_N$  are i.i.d. standard Gaussian r.v.'s. Hence if  $\beta = 0$ ,  $N \to \infty \Longrightarrow \frac{1}{N} \sum_{i=1}^N \delta_{x_i} \longrightarrow$  Gaussian distribution. Global asymptotics of  $G\beta E$  eigenvalues at high temp

Theorem (Duy, Shirai '15 & Benaych-Georges, C, Gorin '22) Consider the empirical measures  $\mu_{N,\beta} := \frac{1}{N} \sum_{i=1}^{N} \delta_{\frac{x_i}{N}}, \quad \text{where } x_1 \leq \cdots \leq x_N \text{ is } Eigen_N^{(\beta)} - distributed.$ 

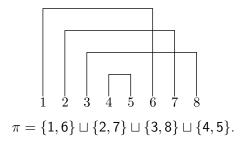
In the limit

$$N o \infty, \quad \beta o 0^+, \quad \frac{N\beta}{2} o \gamma \in (0,\infty),$$

the measures  $\mu_{N,\beta}$  converge weakly, in probability, to certain probability measure  $\mu_{\gamma}$  which can be completely described.

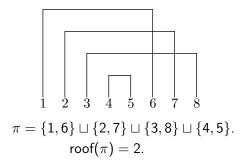
# Global asymptotics of Hermite N-particle $\beta$ -ensemble at high temperature

For a *perfect matching*  $\pi = \{B_1, \dots, B_n\}$  of  $\{1, \dots, 2n\}$ , draw the *arc diagram*. Define  $roof(\pi) := \#$  roofs with no intersections.



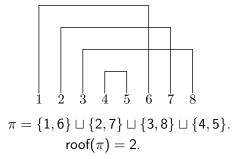
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Theorem (Benaych-Georges, Cuenca, Gorin '22) The limiting measure  $\mu_{\gamma}$  is uniquely determined by its moments:

$$\int_{-\infty}^{\infty} x^k \, \mu_{\gamma}(dx) = \sum_{\text{perfect matchings } \pi \text{ of } \{1, \cdots, k\}} (\gamma + 1)^{\text{roof}(\pi)}$$

Limits as  $\gamma \rightarrow \mathrm{0^+}$  and  $\gamma \rightarrow \infty$ 

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#### Comments:

1. If k is odd, the k-th moment of  $\mu_{\gamma}$  is zero.

2. If 
$$k = 2n$$
 and  $\gamma \to 0^+$ , then  
RHS = number of perfect matchings of  $\{1, 2, \dots, 2n\}$   
=  $(2n-1)(2n-3)\cdots 3\cdot 1$ .

3. If k = 2n and  $\gamma \to \infty$  (need to divide by  $\gamma^n$  first), then RHS = number of <u>noncrossing</u> perfect matchings of  $\{1, 2, \dots, 2n\}$ 

= Catalan number 
$$C_n = \frac{(2n)!}{(n+1)!n!}$$
.

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# Recall: Levy's continuity theorem

Let  $\{\mu_N\}_{N\geq 1}$ ,  $\mu$  be probability measures on  $\mathbb{R}^d$ . The *Fourier transform* of  $\mu_N$  is

$$\phi_N(\vec{\mathbf{x}}) := \int_{\mathbb{R}^d} \mathcal{K}(\vec{\mathbf{a}}, \vec{\mathbf{x}}) \, \mu_N(\vec{\mathbf{a}}).$$
where  $\vec{\mathbf{x}} := (x_1, \cdots, x_d)$ ,  $\vec{\mathbf{a}} := (a_1, \cdots, a_d)$ ,  
 $\mathcal{K}(\vec{\mathbf{a}}, \vec{\mathbf{x}}) := e^{\mathbf{i}(a_1 x_1 + \cdots + a_d x_d)}.$ 

Similarly, let  $\phi(\vec{\mathbf{x}})$  be the Fourier transform of  $\mu$ .

# Recall: Levy's continuity theorem

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#### Theorem

$$\mu_N \rightarrow \mu$$
 weakly  $\iff \phi_N(\vec{\mathbf{x}}) \rightarrow \phi(\vec{\mathbf{x}})$  pointwise.

Intuition: At least when all measures are compactly supported, use

$$\mathbb{E}_{\mu}\left[\boldsymbol{a}_{1}^{k_{1}}\cdots\boldsymbol{a}_{d}^{k_{d}}\right] = \left.\frac{\partial^{k_{1}+\cdots+k_{d}}}{\partial \boldsymbol{x}_{1}^{k_{1}}\cdots\boldsymbol{x}_{d}^{k_{d}}}\,\phi(\vec{\mathbf{x}})\right|_{\boldsymbol{x}_{1}=\cdots=\boldsymbol{x}_{d}=0}$$

# Multivariate Bessel functions

Theorem (Benaych-Georges, Cuenca, Gorin '22)

(Abbreviated) LLN for empirical measures of  $x_1 \leq \cdots \leq x_N \iff$ Taylor coeffs of the logarithm of  $\beta$ -Fourier transforms converge.

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Our  $\beta$ -Fourier transform = Dunkl transform relies on a kernel  $K(\vec{a}, \vec{x})$  that depends on  $\beta$ , the multivariate Bessel function:

$$K(ec{\mathbf{a}},ec{\mathbf{x}}) = B_N^{(eta)}(ec{\mathbf{a}},ec{\mathbf{x}}), \quad eta \geq 0$$

defined from the (differential, symmetrized) Dunkl operators  $P_k^{(\beta)}$ :  $B_N^{(\beta)}(\vec{\mathbf{a}}, \vec{\mathbf{x}})$  is symmetric in the variables  $x_1, \dots, x_N$ ,  $B_N^{(\beta)}(\vec{\mathbf{a}}, \vec{\mathbf{0}}) = 1$ ,  $P_k^{(\beta)} B_N^{(\beta)}(\vec{\mathbf{a}}, \vec{\mathbf{x}}) = \left(\sum_{i=1}^N a_i^k\right) \cdot B_N^{(\beta)}(\vec{\mathbf{a}}, \vec{\mathbf{x}}), \ \forall k = 1, 2, \dots$ .

# How to think of the Bessel generating function?

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When 
$$\beta = 0$$
:  
 $P_k^{(\beta=0)} = \left(\frac{\partial}{\partial x_1}\right)^k + \dots + \left(\frac{\partial}{\partial x_N}\right)^k$ ,  
 $B_N^{(\beta=0)}(\vec{\mathbf{a}}, \vec{\mathbf{x}}) = \frac{1}{N!} \sum_{\sigma \in S(N)} e^{\mathbf{a}_1 x_{\sigma(1)} + \dots + \mathbf{a}_N x_{\sigma(N)}}$ .

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• When  $\beta = 2$ : they are the HCIZ, spherical integral  $B_N^{(\beta=2)}(\vec{\mathbf{a}}, \vec{\mathbf{x}}) := \int_{U(N)} e^{\operatorname{Trace}(UD(\vec{\mathbf{a}})U^*D(\vec{\mathbf{x}}))} \operatorname{Haar}(dU),$ 

where  $D(\vec{\mathbf{a}}) := \text{diag}(\vec{\mathbf{a}}), D(\vec{\mathbf{x}}) := \text{diag}(\vec{\mathbf{x}})$ ; the integral is over the *Haar probability measure* on U(N).

When β = 1, 4: B<sub>N</sub><sup>(β=1)</sup>(**a**, **x**), B<sub>N</sub><sup>(β=4)</sup>(**a**, **x**) are spherical integrals over orthogonal and symplectic compact groups.

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- They are limits of Macdonald polynomials.

# Our general approach

The idea is to apply moment / operator method to the  $\beta\mbox{-}{\rm Fourier}$  transform

$$G_N^{(eta)}(ec{\mathbf{x}}) := \int_{\mathbb{R}^N} B_N^{(eta)}(ec{\mathbf{a}}, ec{\mathbf{x}}) \, \mathrm{d} \mu_N(ec{\mathbf{a}}),$$

by analogy with classical theory

$$e^{\mathbf{i}(a_1x_1+\dots+a_Nx_N)} \longrightarrow B_N^{(eta)}(\mathbf{\vec{a}},\mathbf{\vec{x}})$$
  
 $rac{\partial^k}{\partial x_1^k} + \dots + rac{\partial^k}{\partial x_N^k} \longrightarrow P_k^{(eta)}$ 

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by analogy with classical theory

$$e^{\mathbf{i}(a_1 \times 1 + \dots + a_N \times N)} \longrightarrow B_N^{(\beta)}(\mathbf{\vec{a}}, \mathbf{\vec{x}})$$
$$\frac{\partial^k}{\partial x_1^k} + \dots + \frac{\partial^k}{\partial x_N^k} \longrightarrow P_k^{(\beta)}$$

$$\left.\prod_{i=1}^{s} P_{k_i}^{(\beta)}\left(G_N^{(\beta)}\right)\right|_{x_1=\cdots=x_N=0} = \mathbb{E}_{\mu_N}\left[\prod_{i=1}^{s} \left(a_1^{k_i}+\cdots+a_N^{k_i}\right)\right]$$

# Our general approach

$$\left.\prod_{i=1}^{s} \mathcal{P}_{k_{i}}^{(\beta)}\left(e^{\ln(G_{N}^{\beta})}\right)\right|_{x_{1}=\cdots=x_{N}=0}=\mathbb{E}_{\mu_{N}}\left[\prod_{i=1}^{s}\left(a_{1}^{k_{i}}+\cdots+a_{N}^{k_{i}}\right)\right]$$

These equations link: **analytic info of**  $G_N^{(\beta)} \leftrightarrow$  **probabilistic info of**  $\mu_N$ . In the high temperature limit, they link: **limits of Taylor coeffs of**  $\ln(G_N^{(\beta)}) \leftrightarrow$  **limits of moments of**  $\mu_N$ .

# The first main theorem

$$\left.\prod_{i=1}^{s} \mathcal{P}_{k_{i}}^{\left(\beta\right)}\left(e^{\ln(G_{N}^{\beta})}\right)\right|_{x_{1}=\cdots=x_{N}=0}=\mathbb{E}_{\mu_{N}}\left[\prod_{i=1}^{s}\left(a_{1}^{k_{i}}+\cdots+a_{N}^{k_{i}}\right)\right]$$

Theorem (Benaych-Georges, Cuenca, Gorin '22) LLN  $\iff$  limits of Taylor coeffs of  $\ln(G_N^{(\beta)})$ , i.e. TFAE: (1) There exist  $m_1, m_2, \cdots$  such that  $\lim_{N\to\infty}\mathbb{E}_{\mu_N}\left[N^{-s}\prod_{i=1}^s\left(a_1^{k_i}+\cdots+a_N^{k_i}\right)\right]=\prod_{i=1}^s m_{k_i}.$ (2) There exist  $\kappa_1, \kappa_2, \cdots$  such that  $\lim_{N\to\infty,\,\beta\to 0^+} \frac{1}{\ell!} \cdot \frac{\partial^{\ell}}{\partial x_1^{\ell}} \ln \left( G_N^{(\beta)} \right) \bigg|_{\chi_1 = \dots = \chi_{\ell} = 0} = \kappa_{\ell}/\ell, \quad \forall \, \ell \in \mathbb{Z}_{\geq 1},$  $\lim_{N\to\infty,\,\beta\to0^+}\frac{\partial^r}{\partial x_{i_1}\cdots\partial x_{i_r}}\ln\left(G_N^{(\beta)}\right)\Big|_{x_1=\cdots=x_r=0}=0,\ if\ |\{i_1,\cdots,i_r\}|\geq 2.$ 

The moments of the limiting measure in the LLN are  $m_1, m_2, \cdots$ .

# Moments of the limiting measure

In example of eigenvalues of  $G\beta E_N$ , the  $\beta$ -Fourier Transform is:

$$G_N^{(\beta)}(x_1,\cdots,x_N)=\exp\left(\frac{x_1^2+\cdots+x_N^2}{2}\right),$$

so  $\kappa_2 = 1$ ;  $\kappa_1 = \kappa_3 = \kappa_4 = \cdots = 0$ .

What are the corresponding moments  $m_k$ 's?

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What are the corresponding moments  $m_k$ 's?

The first few relations  $m_k$ 's  $\leftrightarrow \kappa_\ell$ 's are:

$$m_{1} = \kappa_{1}$$

$$m_{2} = (\gamma + 1)\kappa_{2} + \kappa_{1}^{2}$$

$$m_{3} = (\gamma + 1)(\gamma + 2)\kappa_{3} + 3(\gamma + 1)\kappa_{2}\kappa_{1} + \kappa_{1}^{3}$$

$$m_{4} = (\gamma + 1)(\gamma + 2)(\gamma + 3)\kappa_{4} + (\gamma + 1)(2\gamma + 3)\kappa_{2}^{2} + \cdots$$
...

The second main theorem of [BG - C - G] is an explicit formula.

The second main theorem: moments of limiting measure

For a <u>set partition</u>  $\pi = \{B_1, \dots, B_m\}$  of  $\{1, 2, \dots, k\}$ , draw the arc diagram of  $\pi$  and define the weight

$$W_{\gamma}(\pi) := \prod_{i=1}^{m} rac{p(i)! (\gamma + |B_i| - 1)!}{(\gamma + p(i))!}.$$

 $\{1, \dots, 8\} = \{1, 3, 5, 7\} \sqcup \{2, 4, 8\} \sqcup \{6\}.$  $p(i) := \# \text{ roofs of } B_i \text{ with some intersection.}$  $|B_i| := \text{ size of the block } B_i.$  The second main theorem: moments of limiting measure

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 $p(1) = p(3) = 0, \ p(2) = 2 \Rightarrow W_{\gamma}(\pi) = 2(\gamma + 1)(\gamma + 2)(\gamma + 3).$ 

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Theorem (Benaych-Georges, Cuenca, Gorin '22)

$$m_k = \sum_{\textit{set partitions $\pi$ of $\{1,\cdots,k$\}}} W_\gamma(\pi) \prod_{B \in \pi} \kappa_{|B|}, \quad \forall \, k \geq 1.$$

Eigenvalues of Gaussian beta ensemble  $(G\beta E)$ 

General theorems and proof ideas

Further questions  $\gamma$ -Semifree Probability Discrete Ensembles

# $\gamma$ -cumulants and $\gamma$ -semifree probability

The relation  $m_k$ 's  $\leftrightarrow \kappa_\ell$ 's generalizes the relation between moments  $\leftrightarrow$  cumulants of a probability measure (at  $\gamma = 0$ ), and between moments  $\leftrightarrow$  free cumulants (at  $\gamma = \infty$ ).

We call  $\kappa_{\ell}$ 's the  $\gamma$ -semifree cumulants.

Problem

Study the  $\gamma$ -Semifree Probability.

# $\gamma\text{-}\mathsf{cumulants}$ and $\gamma\text{-}\mathsf{semifree}$ probability

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For example,

1. **Conjecture:** Given probability measures  $\mu$ ,  $\nu$  of compact support, and  $\gamma$ -semifree cumulants  $\{\kappa_{\ell}^{\mu}\}_{\ell \geq 1}$ ,  $\{\kappa_{\ell}^{\nu}\}_{\ell \geq 1}$ , there exists a unique probability measure  $\mu \boxplus_{\gamma} \nu$  of compact support such that

$$\kappa_{\ell}^{\mu\boxplus_{\gamma}\nu} = \kappa_{\ell}^{\mu} + \kappa_{\ell}^{\mu}, \quad \ell \ge 1.$$

This would be the  $\gamma\text{-}\mathsf{Semifree}$  Convolution of  $\mu$  and  $\nu.$ 

 $\gamma\text{-}\mathsf{cumulants}$  and  $\gamma\text{-}\mathsf{semifree}$  probability

2. Assuming the conjecture, classify the *infinitely divisible laws* with respect to the operation of  $\gamma$ -semifree convolution.

3. Theorem (BG – C – G '22). (LLN and connection to RMT) If the empirical measures of  $(a_1 \leq \cdots \leq a_N)$ ,  $(b_1 \leq \cdots \leq b_N)$ converge weakly to  $\mu$ ,  $\nu$ , then the eigenvalues  $(c_1 \leq \cdots \leq c_N)$  of the  $\beta$ -sum  $A_N +_{\beta} B_N$  converge in the high temperature regime to the  $\gamma$ -semifree convolution  $\mu \boxplus_{\gamma} \nu$ .

etc, etc.

Eigenvalues of Gaussian beta ensemble  $(G\beta E)$ 

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### Discrete ensembles: The Jack-Plancherel measure

Begin with the Burnside identity for symmetric group  $S_N$ :

$$\sum_{\lambda \in \mathbb{Y}_N} \dim(\lambda)^2 = N!,$$
  
where  $\lambda$  ranges over partitions of size  $N$ :  
 $\mathbb{Y}_N := \{\lambda = (\lambda_1 \ge \lambda_2 \ge \cdots \ge 0) \text{ s.t. } |\lambda| = \sum_{i \ge 1} \lambda_i = N\},$   
 $\dim(\lambda) = \frac{N!}{\prod_{s \in \lambda} (a(s) + l(s) + 1)} = \dim. \text{ of an irred. } S_N\text{-module.}$ 

#### Discrete ensembles: The Jack-Plancherel measure

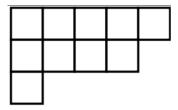
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#### Definition

The Plancherel measure is the measure on  $\mathbb{Y}_N$  is:  $Planch^{(1)}(\lambda) := \frac{\dim(\lambda)^2}{N!} = \frac{N!}{\prod_{s \in \lambda} (a(s) + l(s) + 1)^2}.$ 

# Limits of the Plancherel measure

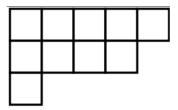


Young diagram of partition (5, 4, 1)



Large  $\frac{1}{\sqrt{N}}$ -normalized partition

# Limits of the Plancherel measure





Young diagram of partition (5, 4, 1)

Large  $\frac{1}{\sqrt{N}}$ -normalized partition

Theorem (Vershik-Kerov '77 & Logan-Shepp '77)

The  $\frac{1}{\sqrt{N}}$ -normalized profiles of Planch<sup>(1)</sup>-distributed partitions  $\lambda \in \mathbb{Y}_N$  converge as  $N \to \infty$  to

$$\omega(u) := \begin{cases} \frac{2}{\pi} (u \operatorname{arcsin}(u/2) + \sqrt{4 - u^2}), & \text{if } |u| \le 2, \\ |u|, & \text{if } |u| \ge 2. \end{cases}$$

# The $\alpha$ -Jack-Plancherel measure

Definition

The  $\alpha$ -Jack-Plancherel measure is the measure on  $\mathbb{Y}_N$  is:  $Planch^{(\alpha)}(\lambda) := \frac{\alpha^N \cdot N!}{\prod_{s \in \lambda} (\alpha a(s) + l(s) + 1)(\alpha a(s) + l(s) + \alpha)}.$ 

This is a " $\beta$ -ensemble-type measure" with  $\beta = 2/\alpha$ . Use  $\alpha$ -anisotropic partitions with boxes  $\sqrt{\alpha} \times (1/\sqrt{\alpha})$ .

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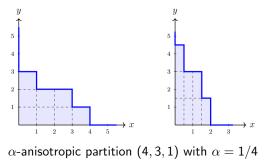


Image taken from [Dolega-Sniady, Gaussian fluctuations of Jack-deformed random Young diagrams, PTRF '19].

#### Theorem (Dolega-Sniady '19)

The  $\frac{1}{\sqrt{N}}$ -normalized profiles of  $\alpha$ -anisotropic Planch<sup>( $\alpha$ )</sup>-distributed partitions  $\lambda \in \mathbb{Y}_N$  converge in the high temperature regime

 $N \to \infty$ ,  $\alpha \to \infty$ ,  $N/\alpha \to g \in (0,\infty)$ ,

to certain limit shape  $\omega_g : \mathbb{R} \to \mathbb{R}_{\geq 0}$  such that  $\omega_g(u) = |u|$ , whenever  $|u| \gg 0$ .

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But what is the limit shape  $\omega_g$ ? Where are the beautiful formulas for the moments? [Cuenca, Dolega, Moll '22+] to the rescue!

The process is a bit indirect: consider the Markov-Krein correspondence  $\omega \rightarrow \nu_{\omega}$ , which makes an association:

shapes  $\omega : \mathbb{R} \to \mathbb{R}_{\geq 0} \longrightarrow$  probability measures  $K[\omega]$  on  $\mathbb{R}$ .

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#### Theorem (Cuenca-Dolega-Moll '22+)

The limit shapes  $\omega_g$  are uniquely determined by the moments of  $\nu_{\omega_g}$ , which have nice combinatorial formulas:

$$\int_{\mathbb{R}} x^m \, \mathcal{K}[\omega_g](dx) = \sum_{Motzkin \text{ paths } P \text{ of length } m} g^{-|H^{\rightarrow}(P)|/2} \cdot \prod_{j \ge 0} j^{|H^{\rightarrow}(P;j)|}.$$

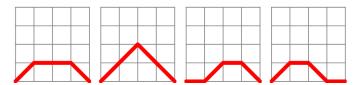


Image taken from Wolfram MathWorld - Motzkin number

Thank you for your attention)

