Soft Stuff, Heavy Ions and Event Generators III

Leif Lönnblad

Department Physics
Lund University

Frontiers in Nuclear and Hadronic Physics
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Outline of Lectures

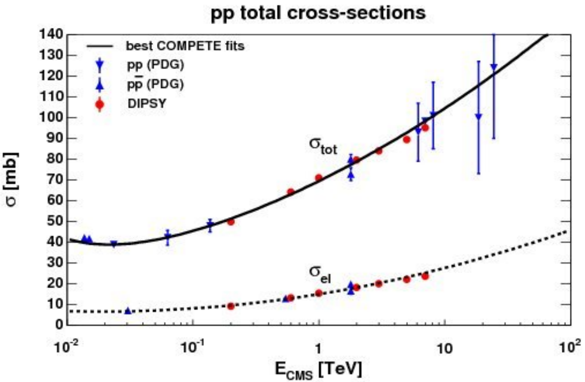
- Lecture I: Basics of Monte Carlo methods, the event generator strategy, matrix elements, LO/NLO, . . .
- Lecture II: Parton showers, initial/final state, (matching/merging), hadronization, decays. . . .
- Lecture III: Minimum bias, multi-parton interactions, pile-up, summary of general purpose event generators, . . .
- Lecture IV: Protons vs. heavy ions, Glauber calculations, initial/final-state interactions, . . .

Buckley et al. (MCnet collaboration), *Phys. Rep.* **504** (2011) 145.

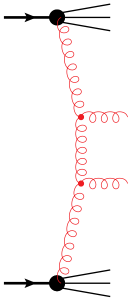
Outline of Lecture III

- Minimum Bias
- Multiple Interactions
  - Interleaved showers
  - Colour connections
- Underlying Events
- General Purpose Event Generators
  - PYTHIA
  - HERWIG
  - SHERPA
- Related Tools

Inclusive cross sections.



Minimum Bias: The typical *pp* collision



soft $gg \to gg$

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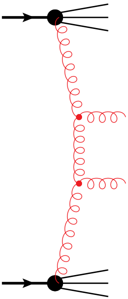

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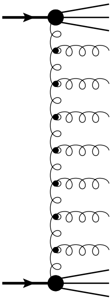

+ISR

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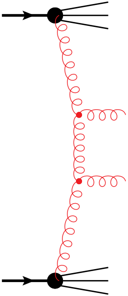

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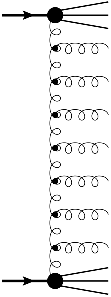

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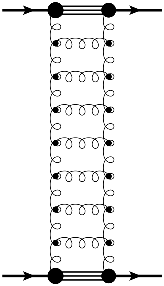
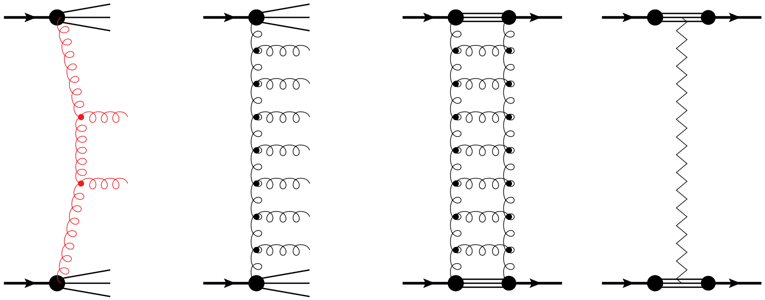


Related to elastic scattering
via optical theorem

Minimum Bias: The typical *pp* collision



soft $gg \to gg$ | +ISR | Related to elastic scattering via optical theorem | Pomeron exchange

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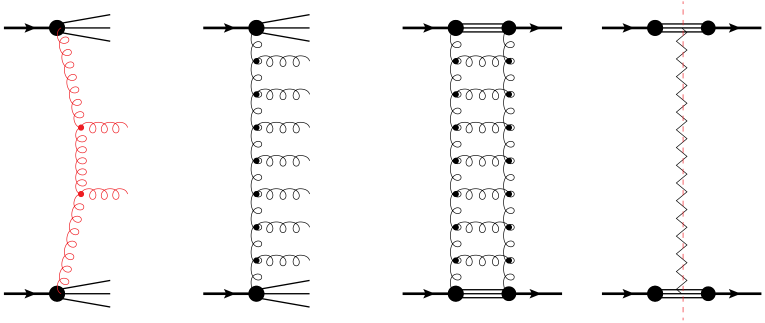

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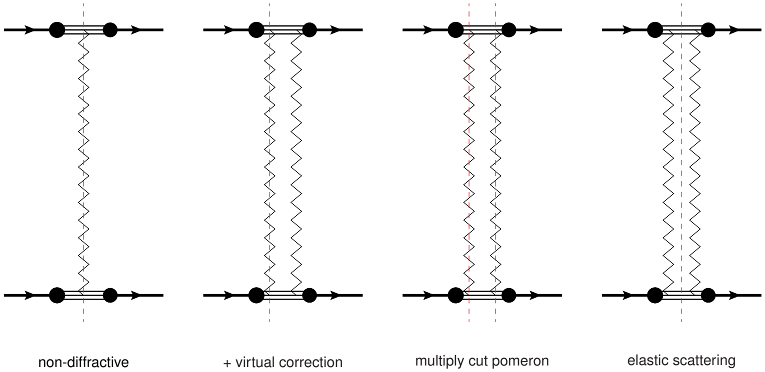
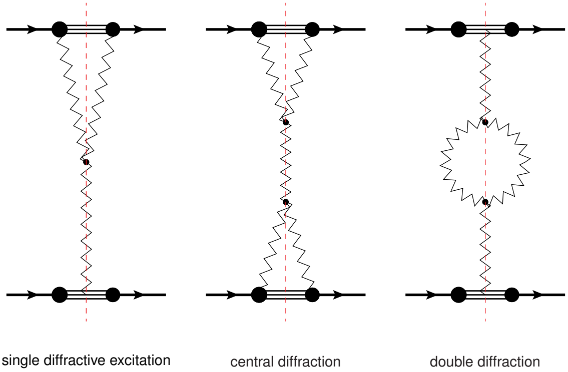
A cut pomeron in an elastic amplitude gives the non-diffractive cross section.
(From Regge theory [Regge. T, *Nuovo Cim.* **14** (1959) 951])

Multi-pomeron diagrams



Each cut pomeron contributes with evenly distributed particle production in the corresponding rapidity interval.
Like two flat strings.

Diffraction and triple-pomeron vertices



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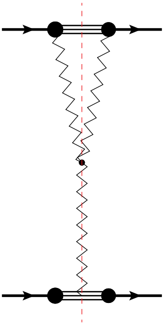
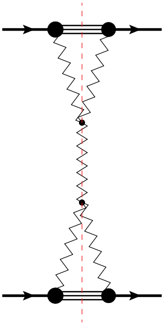
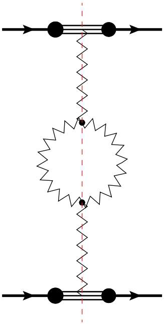


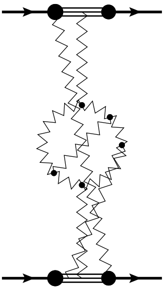
single diffractive excitation



central diffraction



double diffraction



?

Soft multiple interactions

- PHOJET [Engel et al.]
- Shrimps (SHERPA) [Zapp et al.]
- EPOS-LHC (also Heavy ions) [Werner et al.]

Where are the (mini-) jets?

(Semi-) Hard Multiple Interactions

Starting Point in PYTHIA:

$$\frac{d\sigma^H}{dk_\perp^2} = \sum_{ij} \int dx_1 dx_2 f_i(x_1, \mu_F^2) f_j(x_2, \mu_F^2) \frac{d\hat{\sigma}_{ij}^H}{dk_\perp^2}$$

The QCD $2 \to 2$ cross section is divergent $\propto \alpha_S^2(k_\perp^2)/k_\perp^4$

$\int_{k_{\perp c}^2} d\sigma^H$ will exceed the total (non-diffractive) *pp* cross section at the LHC for $k_{\perp c} \lesssim 5$ GeV.

There are more than one partonic interaction per *pp*-collision

$$\langle N_H \rangle (k_{\perp c}) = \frac{\int_{k_{\perp c}^2} d\sigma^H}{\sigma^{\mathrm{ND}}}$$

The trick in PYTHIA is to treat everything as if it is perturbative.

$$\frac{d\hat{\sigma}_{ij}^{H}}{dk_{\perp}^{2}} \rightarrow \frac{d\hat{\sigma}_{ij}^{H}}{dk_{\perp}^{2}} \times \left(\frac{\alpha_S(k_{\perp}^{2} + k_{\perp 0}^{2})}{\alpha_S(k_{\perp}^{2})} \cdot \frac{k_{\perp}^{2}}{k_{\perp}^{2} + k_{\perp 0}^{2}} \right)^{2}$$

Where $k_{\perp 0}^{2}$ is motivated by colour screening (saturation) and is dependent on collision energy.

$$k_{\perp 0}(E_{\mathrm{CM}}) = k_{\perp 0}(E_{\mathrm{CM}}^{\mathrm{ref}}) \times \left(\frac{E_{\mathrm{CM}}}{E_{\mathrm{CM}}^{\mathrm{ref}}} \right)^{\epsilon \sim 0.16}$$

(using handwaving about the the rise of the total cross section)

The total and non-diffractive cross section is put in by hand (or with a Donnachie—Landshoff parameterization).

- Pick a hardest scattering according to

$$\frac{1}{\sigma^{\mathrm{ND}}}\frac{d\sigma^H}{dk_\perp^2} \times \exp\left(-\int_{k_\perp^2} dq_\perp^2 \frac{1}{\sigma^{\mathrm{ND}}}\frac{d\sigma^H}{dq_\perp^2}\right)$$

- Pick an impact parameter, b, from the overlap function (high $k_\perp$ gives bias for small b).
- Generate additional scatterings with decreasing $k_\perp$ using $d\sigma^H(b)/\sigma^{\mathrm{ND}}$

Hadronic matter distributions

We assume that we have factorization

$$\mathcal{L}_{ij}(x_1, x_2, b, \mu_F^2) = \mathcal{O}(b) f_i(x_1, \mu_F^2) f_j(x_2, \mu_F^2)$$

$$\mathcal{O}(b) = \int dt \int dxdydz \rho(x, y, z) \rho(x + b, y, z + t)$$

Where ρ is the matter distribution in the proton
(note: general width determined by σ^{ND})

- A simple Gaussian
- Double Gaussian
- x-dependent Gaussian

x-dependent overlap

Small-x partons are more spread out

$$\rho(r,x) \propto \exp\left(-\frac{r^2}{a^2(x)}\right)$$

with $a(x) = a_0(1 + a_1 \log 1/x)$

Note that high $k_\perp$ generally means higher x and more narrow overlap distribution.

A note on resummation

There are many scales in an event: S, $\hat{s}$, $-\hat{t} \sim k_\perp^2$, Λ_{QCD}^2, m_W^2, $\ldots$

Every time we have two widely separated scales there may be large logarithms that need to be resummed:

- $k_\perp^2 \gg \Lambda_{\mathrm{QCD}}^2$: DGLAP
- $S \gg k_\perp^2$: BFKL (CCFM)
- $\hat{s} \gg k_\perp^2$: FKL

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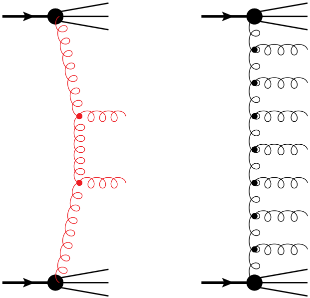
Is it reasonable to use collinear factorization even for very small $k_\perp$?
Soft interactions means very small x, should we not be using $k_\perp$-factorization and BFKL?



For very small x and small $k_\perp$ we also have *saturation*: not only splittings, but also recombinations of gluons.

Energy–momentum conservation

Each scattering consumes momentum from the proton, and eventually we will run out of energy.

- Continue generating MI's with decreasing $k_\perp$, until we run out of energy.
- Or rescale the PDF's after each additional MI. (Taking into account flavour conservation).

Note that also initial-state showers take away momentum from the proton.

Interleaved showers

When do we shower?

- First generate all MI's, then shower each?
- Generate shower after each MI?

Is it reasonable that a low-$k_\perp$ MI prevents a high-$k_\perp$ shower emission? Or vice versa?

- Include MI's in the shower evolution

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After the primary scattering we can have

- Initial-state shower splitting, P_{ISR}
- Final-state shower splitting, P_{FSR}
- Additional scattering, P_{MI}
- Rescattering of final-state partons, P_{RS}

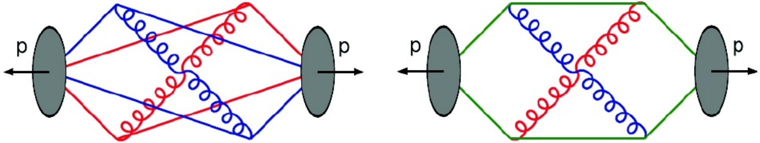
Let them compete

$$\frac{d\mathcal{P}_a}{dk_\perp^2} = \frac{dP_a}{dk_\perp^2} \times \exp - \left(\int_{k_\perp^2} (dP_{\mathrm{ISR}} + dP_{\mathrm{FSR}} + dP_{\mathrm{MI}} + dP_{\mathrm{RS}}) \right)$$

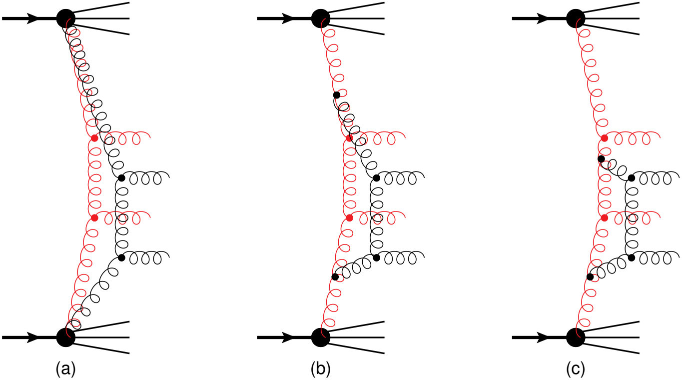
Colour Connections

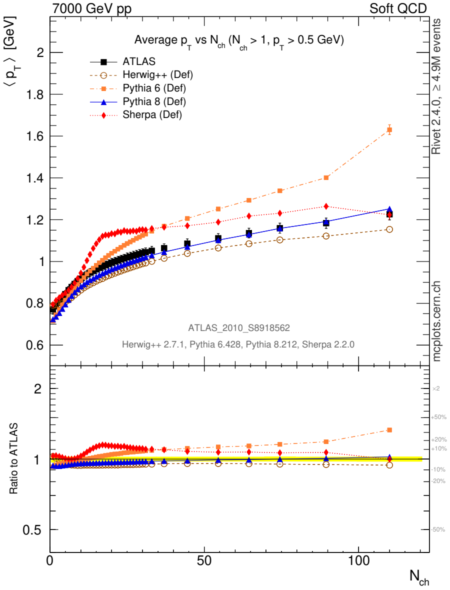
Every MI will stretch out new colour-strings.

Evidently not all of them can stretch all the way back to the proton remnants.



To be able to describe observables such as $\langle p_\perp \rangle (n_{\mathrm{ch}})$ we need (a lot of) colour (re-)connections.


(a)
(b)
(c)


7000 GeV pp
Soft QCD
Average p_T vs N_{ch} (N_{ch} > 1, p_T > 0.5 GeV)
ATLAS
Herwig++ (Def)
Pythia 6 (Def)
Pythia 8 (Def)
Sherpa (Def)
ATLAS_2010_S8918562
Herwig++ 2.7.1, Pythia 6.428, Pythia 8.212, Sherpa 2.2.0
Rivet 2.4.0, ≥ 4.9M events
mcplots.cern.ch
⟨ p_T ⟩ [GeV]
Ratio to ATLAS
N_{ch}

Beyond simple strings

What if we kick out two valens quarks from the same proton?

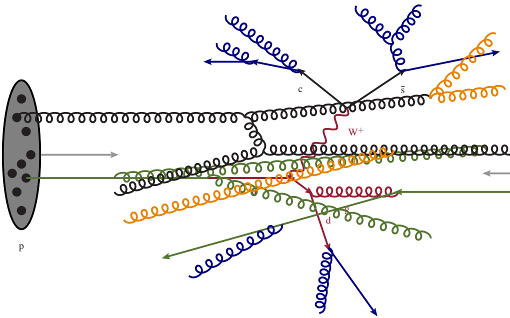
Normally it is assumed that the proton remnant has a di-quark, giving rise to a leading baryon in the target fragmentation.

PYTHIA8 has can hadronize string junctions
(also used for baryon-number violating BSM models)

Non-trivial baryon number distribution in rapidity.

Questions!

What is the Underlying Event?



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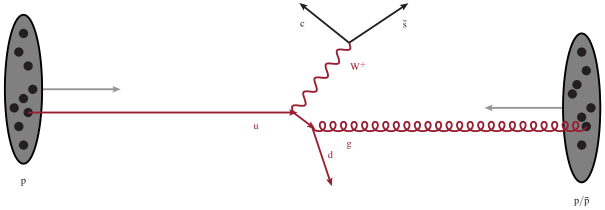

Everything except the hard sub-process?

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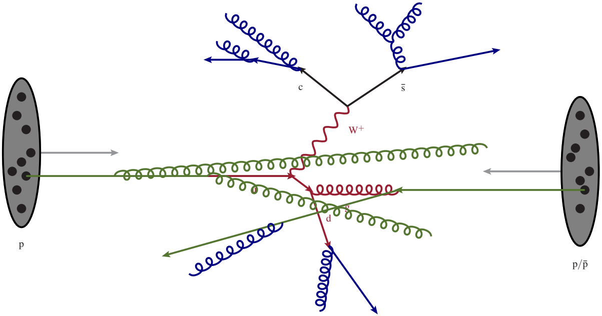

Everything except the hard sub-process
and initial- and final-state showers?

Subtracting underlying events from jets.

- ISR adds energy
- FSR removes energy
- UE adds energy
- Hadronization removes energy

Some of these can be made to cancel eachother by adjusting the size of the jet cone.

But we still need to understand the underlying event.

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UE is not MB

- Harder processes gives a bias towards larger overlap (smaller *b*) giving more UE.
- The UE fluctuates — we can't just subtract a number
- Beware of jet cuts in a steeply falling spectrum

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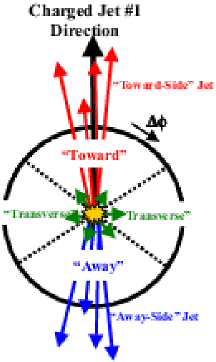
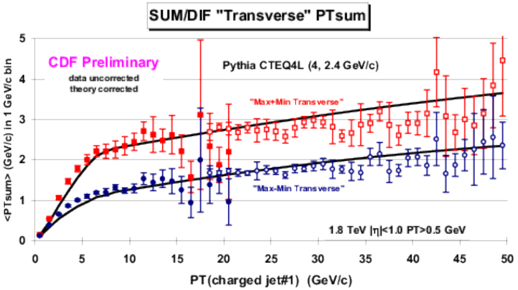
Also relevant for pile-up


Charged Jet #1 Direction
"Toward-Side" Jet
Δφ
"Toward"
"Transverse"
"Transverse"
"Away"
"Away-Side" Jet


SUM/DIF "Transverse" PTsum
CDF Preliminary
data uncorrected
theory corrected
Pythia CTEQ4L (4, 2.4 GeV/c)
"Max+Min Transverse"
"Max-Min Transverse"
1.8 TeV |η|<1.0 PT>0.5 GeV
<PTsum> (GeV/c) in 1 GeV/c bin
PT(charged jet#1) (GeV/c)

A note on Tuning

The Min-bias and UE machineries contains a fair number of parameters that need to be tuned to data. In PYTHIA we have:

- Soft regularisation parameters
- Overlap function parameters
- Cross section parameterisations
- Colour reconnection parameters
- Intrisic transverse momenta
- PDF choices
- ...

Global Tuning

General purpose event generators should describe everything. They should not be tuned to a single observable.

- Hadronization parameters and final-state showers can be tuned to e^+e^- data (LEP).
- Initial-state showers and UE/MPI can be tuned to MB data.
- Anythings else should be fixed by measured Standard Model parameters.
- . . . in principle

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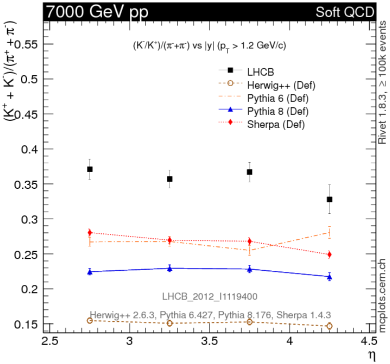
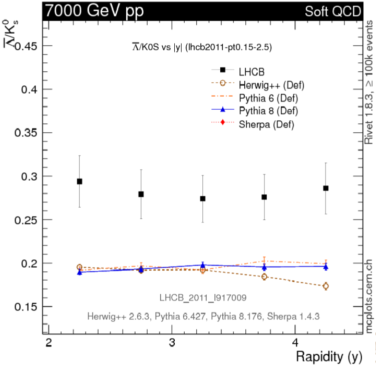
Jet universality





General Purpose Event Generators

There are only a few programs which deals with the whole picture of the event generation

- Hard sub-processes
- Parton showers
- Multiple interactions
- Hadronization
- Decays

Many more programs deal with a specific part of the event generation

- Hard subprocess: AlpGen, MadEvent, . . . can be used with other generators using the Les Houches interface (but be sure to do proper merging)
- Parton Shower: ARIADNE, CASCADE, Vincia, DIRE, . . . need to be integrated with a specific general purpose generator
- Multiple interactions: JIMMY (HERWIG) Shrimps (SHERPA)
- Hadroniziation (?)
- Decays: Tauola, EvtGen, typically called from within other generators.

PYTHIA8

- A few simple MEs, the rest from Les Houches
- $k_\perp$-ordered initial-/final-state DGLAP-based shower
- (N)LO multi-leg matching (not automatic)
- Multiple interactions interleaved with shower
- Lund String Fragmentation
- Particle decays

`https://pythia.org`

HERWIG

- Construction of arbitrary MEs using helicity amplitudes
- Angular ordered and dipole shower
- Different matching schemes via MatchBox
- Soft+hard multiple interactions
- Cluster hadronization
- Particle decays with correlations

`http://projects.hepforge.org/herwig`

SHERPA

- Built-in automated ME generator
- Dipole-based shower
- Semi-automatic (N)LO multi-leg matching
- Multiple interactions ($\sim$ old PYTHIA) with some CKKW features (also Shrimps)
- Cluster hadronization (string fragmentation via old PYTHIA).
- Standard particle decays.

`https://sherpa-team.gitlab.io`

Related Tools

Matrix Element Generators

- **MadGraph5(aMC@NLO)**
- POWHEG
- ALPGEN
- HELAC
- CompHEP
- …

PDF parametrizations

- LHAPDF

Rivet.hepforge.org

(Buckley et al.)

Analyze Event Generator output and compare with published experimental data, using exactly the same cuts, triggers, etc.

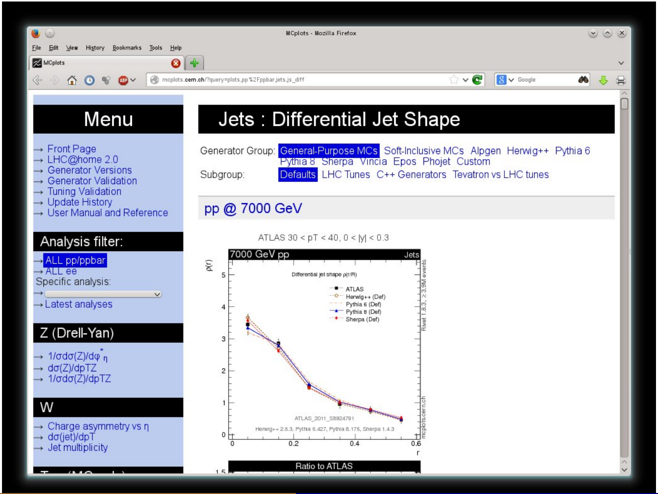
1200+ analyses are already in there.

If you want to make your analyses useful for others —
Publish them in Rivet!

Connected to *Professor* for tuning of parameters

MCplots.cern.ch

(Skands et al.)



Summary III

There are effects which are beyond the formal (leading-twist) precision. We can choose observables that are more or less insensitive to these effects, but they will always be there. We need to understand them better.

General Purpose Event Generators have different solutions to this. The most advanced treatment is found in Pythia8.

Summary III

There are effects which are beyond the formal (leading-twist) precision. We can choose observables that are more or less insensitive to these effects, but they will always be there. We need to understand them better.

General Purpose Event Generators have different solutions to this. The most advanced treatment is found in Pythia8.

And only Pythia8 handles collisions involving Heavy Ions

Questions!

Outline of Lectures

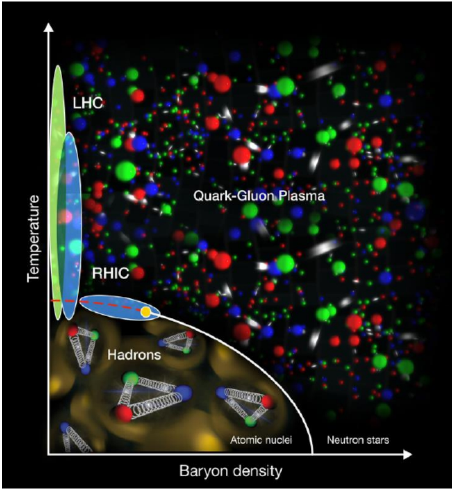
- Lecture I: Basics of Monte Carlo methods, the event generator strategy, matrix elements, LO/NLO, ...
- Lecture II: Parton showers, initial/final state, (matching/merging), hadronization, decays. ...
- Lecture III: Minimum bias, multi-parton interactions, pile-up, summary of general purpose event generators, ...
- Lecture IV: Protons vs. heavy ions, Glauber calculations, initial/final-state interactions, ...

Buckley et al. (MCnet collaboration), *Phys. Rep.* **504** (2011) 145.

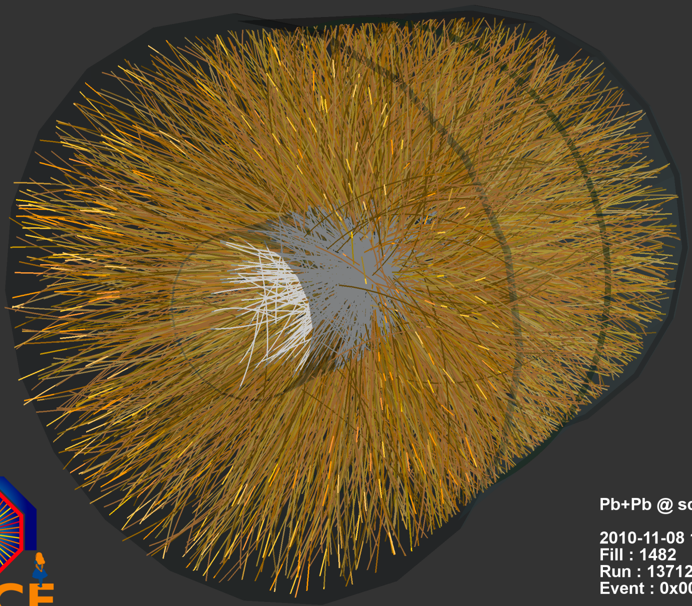
pp vs. *AA* (from the pp point of view)

My immediate reactions when encountering Heavy Ion physics:

- That's just smashing bunches of nucleons together!
- Who is this Glauber guy anyway?
- You do you mean with centrality?
- When is many particles too many?
- I'm from Lund, I want to use string fragmentation!
- You measured what?
- Are you really seeing the Quark–Gluon Plasma?



$$T_C \sim 170 \text{ MeV} \sim 2 \cdot 10^{12} \text{ K}$$

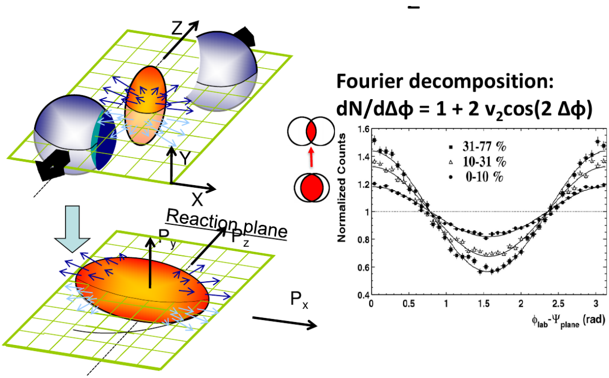


ALICE

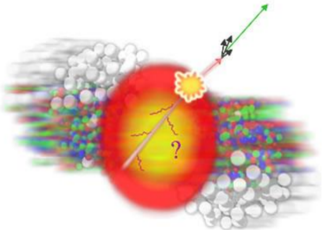
Pb+Pb @ sqrt(s) = 2.76 ATeV

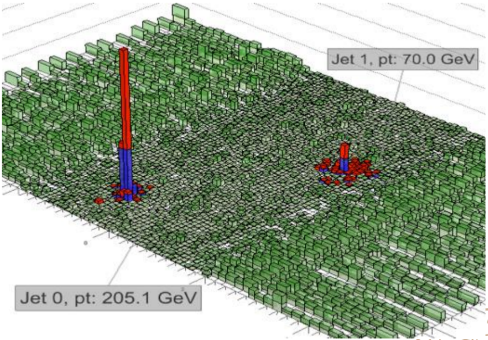
2010-11-08 11:30:46
Fill : 1482
Run : 137124
Event : 0x00000000D3BBE693

Flow

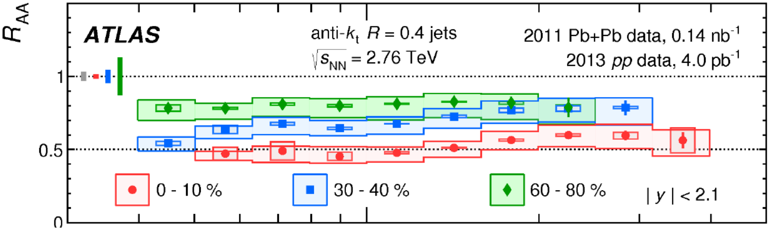

Z
Y
X
Reaction plane
P_y
P_z
P_x
Fourier decomposition:
$dN/d\Delta\phi = 1 + 2\, v_2 \cos(2\, \Delta\phi)$
Normalized Counts
31-77 %
10-31 %
0-10 %
$\phi_{lab}-\Psi_{plane}$ (rad)

Jet quenching





The R_{AA} factor



$$R_{AA} = \frac{d^2N^{AA}/dp_T d\eta}{\langle T_{AA}\rangle d^2\sigma^{pp}/dp_T d\eta}$$

$<T_{AA}>\sigma^{pp} = <N_{coll}>$

N_{coll} is the # of binary collisions

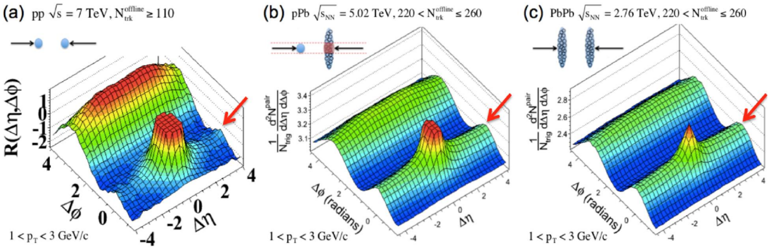
For perturbative QCD processes:

$R_{AA}<1$: suppression

$R_{AA}=1$: no nuclear effects

$R_{AA}>1$: enhancement

The ridge


(a) pp √s = 7 TeV, N_trk^offline ≥ 110
R(Δη,Δϕ)
Δϕ
Δη
1 < p_T < 3 GeV/c
(b) pPb √s_NN = 5.02 TeV, 220 < N_trk^offline ≤ 260
1/N_trig d²N^pair/dΔη dΔϕ
Δϕ (radians)
Δη
1 < p_T < 3 GeV/c
(c) PbPb √s_NN = 2.76 TeV, 220 < N_trk^offline ≤ 260
1/N_trig d²N^pair/dΔη dΔϕ
Δϕ (radians)
Δη
1 < p_T < 3 GeV/c