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# Jets and differential operators in noncommutative geometry

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Jets and differential operators in noncommutative geometry

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## Setting and algebraic preliminaries

- To generalize the notion of jet and differential operator from classical differential geometry to the noncommutative setting all that is required is the following.
  - A first order differential calculus
  - An exterior algebra

### Definition

A first order calculus  $(\Omega^1_d, d)$  for a unital associative  $\Bbbk$ -algebra A is

- An A-bimodule  $\Omega_d^1$ ,
- **2** a map  $d: A \to \Omega^1_d$  satisfying d(ab) = adb + (da)b,

**3** such that 
$$AdA = \Omega_d^1$$
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#### Definition

An exterior algebra Ω<sup>a</sup><sub>d</sub> over a k-algebra A, is an associative graded algebra (Ω<sup>a</sup><sub>d</sub> = ⊕<sub>n≥0</sub> Ω<sup>n</sup><sub>d</sub>, ∧) equipped with a map d such that
Ω<sup>a</sup><sub>d</sub> = A;
d is a differential, that is a k-linear map d: Ω<sup>a</sup><sub>d</sub> → Ω<sup>a</sup><sub>d</sub> such that d(Ω<sup>n</sup><sub>d</sub>) ⊆ Ω<sup>n+1</sup><sub>d</sub> for all n ≥ 0, which satisfies d<sup>2</sup> = 0 and d(α ∧ β) = dα ∧ β + (-1)<sup>n</sup>α ∧ dβ, ∀α ∈ Ω<sup>n</sup><sub>d</sub>, β ∈ Ω<sup>h</sup><sub>d</sub>.
A generates Ω<sup>a</sup><sub>d</sub> via the d and the ∧.

### Remark

Given an exterior algebra  $\Omega_d^{\bullet}$ , the first grade and  $d: \Omega_d^0 = A \to \Omega_d^1$  form a first order calculus for A. Vice versa, given a first order differential calculus,  $\Omega_d^1$ , a maximal exterior algebra,  $\Omega_{d,\max}^{\bullet}$ , is given by quotienting the tensor algebra by the minimal relations for  $d^2 = 0$  to hold.

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Quantum symmetric forms

Given an exterior algebra  $\Omega^{\bullet}_d$  over A, we can define the tensor functors

$$\Omega_d^n \colon {}_A \mathrm{Mod} \longrightarrow {}_A \mathrm{Mod} \qquad \qquad E \longmapsto \Omega_d^n \otimes_A E.$$

We define the functors

$$S^0_d = \Omega^0_d = \mathrm{id}_{A \operatorname{Mod}}, \qquad \qquad S^1_d = \Omega^1_d \coloneqq \Omega^1_d \otimes_A -$$

For  $n \ge 0$ , the functor of quantum symmetric forms  $S_d^n$  is defined by induction as the kernel of the following composition

$$\begin{split} \Omega^1_d \circ S^{n-1}_d & \xrightarrow{\Omega^1_d(\iota^{n-1}_{\wedge})} \Omega^1_d \circ \Omega^1_d \circ S^{n-2}_d \xrightarrow{\wedge_{S^{n-2}_d}} \Omega^2_d \circ S^{n-2}_d \\ \text{and} \ \iota^n_{\wedge} \colon S^n_d & \longrightarrow \Omega^1_d \circ S^{n-1}_d \text{ is the inclusion.} \end{split}$$

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## Spencer cohomology

For all  $k,h\geq 0,$  consider the functor  $\Omega^k_d\circ S^h_d,$  and define  $\delta^{h,k}$  as the following composition

$$\Omega_d^k \circ S_d^h \xrightarrow{\Omega_d^k(\iota_{\wedge}^n)} \Omega_d^k \circ \Omega_d^1 \circ S_d^{h-1} \xrightarrow{(-1)^k \wedge_{S_d^{h-1}}^{k,1}} \Omega_d^{k+1} \circ S_d^{h-1}$$

We thus obtain a complex in the category of endofunctors on  $_AMod$ .

$$0 \longrightarrow S_d^n \xrightarrow{\delta^{n,0}} \Omega_d^1 \circ S_d^{n-1} \xrightarrow{\delta^{n-1,1}} \Omega_d^2 \circ S_d^{n-2} \xrightarrow{\delta^{n-2,2}} \Omega_d^3 \circ S_d^{n-3} \xrightarrow{\delta^{n-3,3}} \cdots$$

### Definition (Spencer cohomology)

We call this the Spencer  $\delta$ -complex, its cohomology the Spencer cohomology, and we denote the cohomology at  $\Omega_d^k \circ S_d^h$  by  $H^{h,k}$ .

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## Universal calculus and 1-jets

#### Definition

The universal calculus  $\Omega_u^1$  for A is given by the kernel of the multiplication map,

$$\Omega^1_u = \ker(\cdot) \subset A \otimes A.$$

The corresponding universal differential is

 $d_u : a \mapsto 1 \otimes a - a \otimes 1$ 

### Proposition

We have

$$0 \to \Omega^1_u \to A \otimes A \to A \to 0$$

For the bimodule map  $\pi_u^{1,0}: a \otimes b \mapsto ab$ . Moreover, the universal prolongation  $j_u^1: a \mapsto 1 \otimes a$  splits the sequence in  $Mod_A$ .

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## Universal property of the universal calculus

Recall the following result.

### Proposition

For any first order calculus  $\Omega_d^1$ , there is an epimorphism  $p_d: \Omega_u^1 \to \Omega_d^1$  which takes  $d_u$  to d, and we have

$$0 \to N_d \to \Omega^1_u \to \Omega^1_d \to 0$$

For a sub-bimodule  $N_d \subset \Omega^1_u$ .

We call  $N_d$  the space of first order differential relations. In the case of classical differential geometry, it contains things such as

$$d_u f - \sum_i f_{x^i} d_u x^i.$$

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For any left A-module E, we define

$$N_d(E) = \ker(p_{d,E}) = \left\{ \sum_i a_i \otimes e_i | \sum_i a_i e_i = 0, \sum_i da_i \otimes_A e_i = 0 \right\} \subset A \otimes E$$

#### Definition

The 1-jet module for a left A-module E is

$$J_d^1 E := J_u^1 E / N_d(E) = A \otimes E / N_d(E).$$

The prolongation operator is  $j_{d,E}^1: E \to J_d^1 E$ , given by  $j_{d,E}^1(e) = [1 \otimes e]$ . The projection map is  $\pi_{d,E}^{1,0}: [\sum_i a_i \otimes e_i] \mapsto \sum_i a_i e_i$ .

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## A plethora of jet functors

We construct the following three families of functors:

- $\bullet$  The nonholonomic jet functors  $J^{(n)}_d\colon {}_A\mathrm{Mod}\to {}_A\mathrm{Mod}$
- The semiholonomic jet functors  $J_d^{[n]} \colon {}_A\mathrm{Mod} \to {}_A\mathrm{Mod}$
- The holonomic jet functors  $J_d^n \colon {}_A \mathrm{Mod} \to {}_A \mathrm{Mod}$

In particular we have  $J_d^{(0)} = J_d^{[0]} = J_d^0 = \mathrm{id}_{A\mathrm{Mod}}$ , and  $J_d^{(1)} = J_d^{[1]} = J_d^1$ . These functors come equipped with natural transformations

$$j_d^{(n)} \colon \mathrm{id}_{{}^A\mathrm{Mod}} \longrightarrow J_d^{(n)} \quad \ j_d^{[n]} \colon \mathrm{id}_{{}^A\mathrm{Mod}} \longrightarrow J_d^{[n]} \quad \ j_d^n \colon \mathrm{id}_{{}^A\mathrm{Mod}} \longrightarrow J_d^n,$$

which are respectively called the nonholonomic, semiholonomic, and holonomic jet prolongation maps. We also have the natural transformations,

$$\pi_d^{(n,n-1;m)} \colon J_d^{(n)} \longrightarrow J_d^{(n-1)}, \quad \pi_d^{[n,n-1]} \colon J_d^{[n]} \longrightarrow J_d^{[n-1]}, \quad \pi_d^{n,n-1} \colon J_d^n \longrightarrow J_d^{n-1},$$

respectively called the nonholonomic, semiholonomic, and holonomic jet projections.

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We can describe  $J^2_d E$  implicitly as the kernel of a left-linear (bilinear for  $E=A)\ \mathrm{map}$ 

$$\widetilde{\mathbf{D}}_E \colon J_d^{(2)} E \longrightarrow (\Omega^1_d \ltimes \Omega^2_d)(E),$$

where  $(\Omega_d^1 \ltimes \Omega_d^2)(E) \cong (\Omega_d^1 \ltimes \Omega_d^2) \otimes_A E$ . As a right A-module,  $\Omega_d^1 \ltimes \Omega_d^2 \cong \Omega_d^1 \oplus \Omega_d^2$ , but as an A-bimodule, it comes equipped with a non-trivial left action

$$f \star (\alpha + \omega) = f\alpha + df \wedge \alpha + f\omega, \qquad \forall f \in A, \ \alpha \in \Omega^1_d, \ \omega \in \Omega^2_d.$$

Explicitly, we have

$$\widetilde{\mathbf{D}}_E \colon J_d^{(2)} E \longrightarrow (\Omega_d^1 \ltimes \Omega_d^2)(E)$$
$$[a \otimes b] \otimes_A [c \otimes e] \longmapsto (ad(bc) \otimes_A e, da \wedge d(bc) \otimes_A e).$$

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## Definition (Holonomic *n*-jet functor)

Let A be a k-algebra endowed with an exterior algebra  $\Omega_d^{\bullet}$ . We define  $J_d^n$  as the kernel of the natural transformation

$$\widetilde{\mathrm{D}}_{J_d^{n-2}} \circ J_d^1(l_d^{n-1}) \colon J_d^1 \circ J_d^{n-1} \longrightarrow (\Omega_d^1 \ltimes \Omega_d^2) \circ J_d^{n-2},$$

where we denote the natural inclusion by  $l_d^n : J_d^n \longrightarrow J_d^1 \circ J_d^{n-1}$ . We call  $J_d^n$  the (holonomic) *n*-jet functor.

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## Higher jet exact sequence

### Theorem (Holonomic jet exact sequence)

Let A be a k-algebra endowed with an exterior algebra  $\Omega_d^{\bullet}$  such that  $\Omega_d^1$ ,  $\Omega_d^2$ , and  $\Omega_d^3$  are flat in  $\operatorname{Mod}_A$ . For  $n \ge 1$ , if the Spencer cohomology  $H^{m,2}$ vanishes, for all  $1 \le m < n-2$ , then the following sequence is exact,

$$0 \longrightarrow S_d^n \xrightarrow{\iota_d^n} J_d^n \xrightarrow{\pi_d^{n,n-1}} J_d^{n-1} \longrightarrow H^{n-2,2}$$

Therefore, if  $H^{n-2,2} = 0$  we obtain a short exact sequence

$$0 \longrightarrow S_d^n \stackrel{\iota_d^n}{\longleftrightarrow} J_d^n \stackrel{\pi_d^{n,n-1}}{\longrightarrow} J_d^{n-1} \longrightarrow 0.$$

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#### Definition

Let  $E, F \in {}_{A}Mod$ . A k-linear map  $\Delta : E \to F$  is called a linear differential operator of order at most n with respect to the exterior algebra  $\Omega_{d}^{\bullet}$ , if it factors through the holonomic prolongation operator  $j_{d}^{n}$ , i.e. there exists an A-module map  $\widetilde{\Delta} \in {}_{A}Hom(J_{d}^{n}E, F)$  such that the following diagram commutes:



If n is minimal, we say that  $\Delta$  is a linear differential operator of order n.

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## Examples of differential operators

- The differential d.
- Suppose  $\Omega_d^1$  is free and finitely-generated as a left *A*-module, i.e. parallelizable. Given a basis  $\theta_1, \ldots, \theta_n$ , we can define the partial derivative operator  $\partial_i \colon A \to A$ , by  $da = \sum_i \partial_i(a)\theta_i$ .
- A (left) connection with respect to the first order differential calculus  $\Omega_d^1$  on a left A-module E is a k-linear map

$$E \longrightarrow \Omega^1_d \otimes_A E,$$

satisfying the identity

$$\nabla(fe) = df \otimes_A e + f \nabla e.$$

- The operator  $\mathfrak{D}_E \colon J^1_d(E) \to (\Omega^1_d \ltimes \Omega^2_d)(E)$ , whose lift  $\widetilde{\mathfrak{D}}_E$  defines second order jets.
- Vector fields:  $\mathfrak{X}_d := \operatorname{Diff}_d^1(A, A) \cap \operatorname{Ann}(\ker(d))$ . Note that  $\mathfrak{X}_d \simeq {}_d\operatorname{Hom}(\Omega^1_d, A)$ .

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#### Proposition

Let  $\Delta_1: E \to F$  and  $\Delta_2: F \to G$  be differential operators of order at most nand m, respectively. Then the composition  $\Delta_2 \circ \Delta_1: E \to G$  is a differential operator of order at most n + m.

#### Proof.



#### Corollary

There is a category  $\operatorname{Diff}_d^{\operatorname{fin}}$  with the same objects as  ${}_A\operatorname{Mod}$  and with morphisms given by the finite order differential operators.

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## Representability

### Definition

We say that the functor  $\operatorname{Diff}_d^n(E, -)$  is representable if there is an object Q such that  $\operatorname{Diff}_d^n(E, -) \simeq {}_A\operatorname{Hom}(Q, -)$ .

### Question

When is  $\operatorname{Diff}_{d}^{n}(E, -)$  represented by  $J_{d}^{n}E$ ?

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## Higher order differential relations

There is a natural transformation

$$\hat{p}_d^n: J_u^1 = A \otimes - \to J_d^n$$

Explicitly given by

$$a \otimes e \to a j_d^n(e)$$

We term its kernel the differential relations of order n, written  $\ker \hat{p}_d^n(E) = N_d^n(E)$ .

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#### Proposition

Let the *n*-jet sequence be left exact. We have  $\hat{p}_d^n|_{N^{n-1}} \to S_d^n$ . The following are equivalent:

- $\hat{p}_d^n|_{N^{n-1}}(E) \to S_d^n E$  is surjective.
- $S^n_d(E) \subset Aj^n_d(E).$

When the latter condition is satisfied, we say symmetric forms of degree n are generated by differential relations.

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#### Theorem

Suppose the *n*-jet sequence is left exact, and that  $J_d^{n-1}E$  represents differential operators. Then  $J_d^n E$  represents differential operators if and only if symmetric forms of degree *n* are generated by differential relations. In that case, we have

 $S_d^n E = (N^{n-1}/N^n)(E)$ 

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## Skeletal jet functors

Let us introduce one more "jet functor".

### Definition

The skeletal n-jet functor is defined as the image subfunctor

 $\mathbb{J}_d^n = \operatorname{Im} \hat{p}_d^n$ 

#### Theorem

We have the following:

- $\mathbb{J}_d^0 = J_d^0$  and  $\mathbb{J}_d^1 = J_d^1$ .
- Let  $\operatorname{Tor}(\Omega^1_d, E) = 0$ . Then  $\mathbb{J}^2_d = J^2_d$  if and only if  $\Omega^2_d = \Omega^2_{d,\max}$ , the maximal prolongation
- $\operatorname{Diff}_d^n(E, -)$  is representable if and only if  $\mathbb{J}_d^n = J_d^n$ , and then the representing object is  $J_d^n$ .

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Consider the following diagram in the abelian category of endofunctors on  ${}_AMod$ , constructed using the jet projections.

$$\cdots J_d^n \xrightarrow{\pi_d^{n,n-1}} J_d^{n-1} \cdots \xrightarrow{\pi_d^{3,2}} J_d^2 \xrightarrow{\pi_d^{2,1}} J_d^1 \xrightarrow{\pi_d^{1,0}} \mathrm{id}_{A\mathrm{Mod}}.$$

We call the above diagram the *(holonomic) jet tower*, and its limit in the category of endofunctors on  $_AMod$  the (holonomic)  $\infty$ -jet functor, denoted

$$J_d^{\infty} := \lim_{n \in \mathbb{N}} J_d^n.$$

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The $\infty$ -jet			

## Jet comonad

Given an object E in  ${}_A\mathrm{Mod}$ , denote an element of  $J_d^{\infty}E$  by

$$\mathbf{e}_0 := (\ldots, \hat{e}_n, \ldots, \hat{e}_1, \hat{e}_0)$$

Similarly, we will denote a truncated element as follows.

$$\mathbf{e}_k := (\dots, \hat{e}_{k+n}, \dots, \hat{e}_{k+1}, \hat{e}_k)$$

#### Proposition

 $J_d^{\infty}$  is a comonad with comultiplication  $\iota: J_d^{\infty} \longrightarrow J_d^{\infty} J_d^{\infty}$  and counit  $\epsilon: J_d^{\infty} \longrightarrow \mathrm{id}_{A \operatorname{Mod}}$  given, respectively, by

 $\mathbf{e}_0 \mapsto (\dots, \mathbf{e}_n, \dots, \mathbf{e}_1, \mathbf{e}_0)$  and  $\mathbf{e}_0 \mapsto \hat{e}_0$ 

#### Proof.

Show that  $(\ldots, \mathbf{e}_n, \ldots, \mathbf{e}_1, \mathbf{e}_0) \in J_d^{\infty} J_d^{\infty} E$ , then check counitality and coassociativity.

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## Noncommutative differential equations

This gives us several ways to proceed in developing a noncommutative theory of (linear) differential equations:

- Via  $\mathcal{D}$ -modules and  $\mathcal{D}$ -algebras give notions of linear and (polynomial) nonlinear differential equations.
- Via a noncommutative analogue of the geometric theory of differential equations developed by Spencer, Quillen, Goldschmidt, etc., where one considers monomorphisms  $\mathcal{E} \hookrightarrow J_d^k E$ .
- Via the comandic analogue of Vinogradov's category of differential equations. The Eilenberg-Moore category for the jet comonad (cf. Marvan, Schreiber, Khavkine) where the objects are diffieties (i.e. infinitely prolonged differential equations) and the morphisms are Cartan distribution preserving maps.