## Modifications of gravity

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#### Tensor-Vector-Scalar theory

Ingredients

(Sanders 1997, Bekenstein 2004)

- Tensor field (metric)  $\tilde{g}_{ab}$
- Unit-timelike Vector field  $A_a$
- Scalar field  $\phi$

Physical metric  $g_{ab} = e^{-2\phi} (\tilde{g}_{ab} + A_a A_b) - e^{2\phi} A_a A_b$ 

- toy-theory (phenomenological)
- gives MOND in the non-relativistic limit
- good platform for studying alternatives to  $\Lambda \text{CDM}$

See topical review in CQG (C.S. on arXiv next week)

#### Growth of structure in CDM and TeVeS



(C.S., D. Mota, P. Ferreira, C.Boehm, 2005)

(S.Dodelson & M.Liguori, 2006)

#### Lessons from TeVeS

- gravity may depend on additional fields
- these fields may mimic dark matter
- however, some other effect different from dark matter may appear

example :

Both  $\Lambda$ CDM and TeVeS give same P(k)But  $\Phi - \Psi$  is very different

Can be distinguished from combinations of other observables

#### How special is General Relativity?

- A principle theory (SEP, diffeo-invariance)
- Lovelock-Grigore theorem : In any dimension, the only local diffeomorphism invariant action leading to 2nd order field equations and which depends only on a metric is a linear combination of the E-H action with a cosmological constant up to a total derivative.

#### Any other theory must : (at least one applies)

- be non-local
- Have absolute elements
- Depend on other fields
- have higher than 2nd order field equations

- (Sousa-Woodard, Dvali-Gabadaze-Porrati, etc)
  - (stratified theories)
    - (JFBD, TeVeS, etc)
    - (Weyl gravity)

### What is/not gravity?

- "f(R) and scalar-tensor theories are not modified gravity because they are equivalent to GR + scalar"
- "Scalar-tensor and TeVeS theories are not modified gravity because they depend on extra fields"
- TeVeS is not modified gravity because it can be written in a single-metric frame without coupling to matter.

$$S = \frac{1}{16\pi G} \int d^4x \sqrt{-g} R - \int d^4x \sqrt{-g} \ K^{abcd}(g, B_c) \nabla_a B_b \nabla_c B_d + S_m[g]$$

• "If  $\Phi - \Psi 
eq 0$  then we have modified gravity"

these are certainly all incorrect statements

#### Cosmological tests of gravity

How do we treat complicated theories of gravity?  $f_{ab} \left[g_{cd}, Riem, Ric, R, \phi^A, \ldots\right] = 8\pi Gh_{ab} \left[T_{cd}, g_{cd}, Ric, R, \phi^A, \ldots\right]$ metric curvature extra fields matter

Trick : add  $G_{ab}$  to both sides add and subtract  $8\pi GT_{ab}$  on RHS

Collect terms 
$$G_{ab} = 8\pi G T_{ab}^{(known)} + U_{ab}$$

where  $U_{ab} = 8\pi G [h_{ab} - T_{ab}] + G_{ab} - f_{ab}$ 

Bianchi identity :  $\nabla_a U^a_{\ b} = 0 \rightarrow$  Field equations for  $\phi^A$ 

C.S. (arXiv:0806.1238)



To distinguish gravity from fluids or other forces we must specify the field content

# Can we distinguish modified gravity from matter at the FRW level?

FRW 
$$ds^2 = -dt^2 + a^2 d\ell_K^2$$
  
 $G^0_0: \quad 3H^2 + \frac{3K}{a^2} = 8\pi G \sum_i \rho_i + X$   
 $G^i_i: \quad -2\frac{\ddot{a}}{a} - H^2 = 8\pi G \sum_i P_i + Y$ 

Bianchi identity gives  $\dot{X} + 3H(X + Y) = 0$ 

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the answer is therefore : NO

#### Linear perturbation level

C.S. (arXiv:0806.1238)

Metric has 4 scalar dof :  $\Psi$ ,  $\zeta$ ,  $\Phi$ ,  $\nu$ Given a vector field  $\xi^a$  s.t.  $\xi_\mu = a(-\xi, \vec{\nabla}_i \psi)$ reduced to 2 by gauge transformations

Distinguishing gravity from fluids : field content

• Linearized equations must be gauge form-invariant.

$$\delta G^{\mu}{}_{\nu} = \sum_{i} \mathcal{O}_{i} \Delta_{i} \to \sum_{i} \mathcal{O}_{i} \Delta_{i} + [FRWeq.]\xi$$

- Holds iff background equations satisfied.
- Fixes all gauge non-invariant terms
- Bianchi identity holds (local energy conservation)

#### Parameterizing field equations

$$2(\vec{\nabla}^{2} + 3K)(\Phi - \vec{\nabla}^{2}\nu) - 6\frac{\dot{a}}{a}(\dot{\Phi} - \frac{1}{3}\vec{\nabla}^{2}\zeta) - 6\frac{\dot{a}^{2}}{a^{2}}\Psi = 8\pi Ga^{2}\rho\delta$$
gauge
 $transform$ 

$$2(\vec{\nabla}^{2} + 3K)(\Phi' - \vec{\nabla}^{2}\nu') - 6\frac{\dot{a}}{a}(\dot{\Phi}' - \frac{1}{3}\vec{\nabla}^{2}\zeta') - 6\frac{\dot{a}^{2}}{a^{2}}\Psi' = 8\pi Ga^{2}\rho\delta' + [FRWeq.]\xi$$

- Parameterizations in a fixed gauge are inconsistent
- Parameterizations using gauge invariant combinations are consistent but may be too arbitrary (from Stewart-Walker lemma)
- All parameterizations must take into account the field content.

# Distinguishing gravity from fluids at the linear level

- Must specify field content
- Specify the parameterization
- Determine the force between 2 well separated masses in vacuum



This requires writing an action leading to the parameterized equations

#### The extended $~~\Lambda \text{CDM}$ model

C.S. (arXiv:0806.1238)

Background : <u>ACDM</u> No additional fields No higher than 2 time derivatives

→ No Gauge Non-Invariant terms allowed

 $\delta U^a{}_b$  contains  $\Phi_{GI}$   $\Psi_{GI}$  and derivatives

$$\Psi_{GI} = \Psi - \ddot{
u} - \dot{\zeta} - rac{a}{a}(\dot{
u} + \zeta)$$
 contains 2nd derivatives

$$\Phi_{GI} = \Phi + \frac{1}{3} \vec{\nabla}^2 \nu + \frac{\dot{a}}{a} (\dot{\nu} + \zeta) \quad \text{contains Ist derivatives}$$

#### Constructing the U-tensor

$$G_{ab} = 8\pi G T_{ab}^{(known)} + U_{ab}$$

#### Constraints : Ist derivatives

 $\delta U^0_{\ 0} = \mathcal{A}\Phi_{GI}$ 

$$\delta U^0_{\ i} = \mathcal{B}\Phi_{GI}$$

## Propagation : 2nd derivatives $\delta U^{i}{}_{i} = C_{1}\Phi_{GI} + C_{2}\dot{\Phi}_{GI} + C_{3}\Psi_{GI}$ $\delta U^{i}{}_{j} - \frac{1}{3}\delta U^{k}{}_{k}\delta^{i}{}_{j} = D_{1}\Phi_{GI} + D_{2}\dot{\Phi}_{GI} + D_{3}\Psi_{GI}$

#### Bianchi identity gives

 $\mathcal{C}_3 = \mathcal{D}_3 = 0 \qquad \qquad \mathcal{A} = -\frac{\dot{a}}{a}\mathcal{C}_2 \qquad \qquad \mathcal{B} = \frac{1}{3}\mathcal{C}_2 + \frac{2}{3}\left(\vec{\nabla^2} + 3K\right)\mathcal{D}_2$  $\dot{\mathcal{A}} + \frac{\dot{a}}{a}\mathcal{A} - \vec{\nabla}^2\mathcal{B} + \frac{\dot{a}}{a}\mathcal{C}_1 = 0 \qquad \qquad \dot{\mathcal{B}} + 2\frac{\dot{a}}{a}\mathcal{B} - \frac{1}{3}\mathcal{C}_1 - \frac{2}{3}\left(\vec{\nabla}^2 + 3K\right)\mathcal{D}_1 = 0$ 

Corollary I: If  $\mathcal{D}_1 = \mathcal{D}_2 = 0$ Then  $U^a_{\ b} = 0$  no shear : no modification

Corollary II: If  $\mathcal{A} = \mathcal{B} = 0$ Then  $U^a_{\ b} = 0$  no constraints : no modification

Corollary III: If  $\mathcal{D}_2 = \mathcal{B} = 0$  i.e.  $\Phi_{GI} - \Phi_{GI} = \mathcal{D}_1 \Phi_{GI}$ Then  $U^a_{\ b} = 0$  no modification

#### **A simple model :** $\Phi_{GI} - \Psi_{GI} = \mathcal{D}_2 \dot{\Phi}_{GI}$



#### Further consistency requirements

(in progress)

- Action for parameterized perturbed cosmological equations
- Quantize on de Sitter
- Eliminate ghosts : Should impose further constraints on the allowed terms
- Initial conditions : e.g. Inflation
- Modified gravity at the non-linear level (non-linear completion)

#### The end

- Detecting  $\Phi \Psi \neq 0$  not enough
- Parametrizing only in terms of the potentials can ignore important physics. Gravity may depend on additional fields.
- Constraints depend on the field content
- Field content important for distinguishing gravity from fluids.