Dark matter annihilation and non-thermal processes in Galaxy clusters

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The Coma galaxy cluster

D =100 Mpc M_{DM} =1.2 x 10¹⁵ M_{sur} $R_{wir} = 2.7 Mpc$ $B(r) = 4.7 n_{th}(r)^{0.5} \mu G$ $= 2 \mu G$ No cooling flow observed Radio, EUV, X-ray observation

The thermal component (from X-rays)

XMM [0.3 -2] kev





Arnaud et al 2001

Maybe a second thermal component The soft-X-ray excess in the outskirts of Coma (r > 1 Mpc)

Fitted by a second thermal component of T = 0.22 keV

consistent with Warm-Hot Intergalactic Medium filaments found in numerical simulation of cluster formation





Non-thermal components: the radio halo of Coma (r < 1 Mpc)

Diffused over Mpc scale - Requires a population of relativistic non-thermal electrons $\gamma > 10^4$ for B ~ 0.1,1 μ G (Note, Faraday Rotation Measurements suggest $\langle B \rangle \sim 2 \mu$ G - Bonafede et al. 2010)

Primary or reacceleration model: electrons produced by AGN activity (quasars, radio galaxies) or star formation (supernovae, galactic wind, etc). Synchrotron and ICS radiation losses should be balanced by reacceleration (shock waves or magneto-hydrodinamics turbolences) - see Brunetti et al 2004 -

Secondary model: electrons produced in inelastic nuclear collisions between relativistic CR protons and thermal ions of the intracluster medium. B > few μ G is needed + associated photon and neutrino production - see Blasi & Colafrancesco 1999 -

The extreme UV excess in the center of Coma (r < 1 Mpc)



Possibly generated by a secondary population of relativistic electrons produced through inelastic collisions of CRs with cluster plasma (secondary model) - which inverse Compton scatter off CMB photons



Blasi & Colafrancesco 1999 0.001 HEAO-A4 100001 $B=0.1 \ \mu G$ 4.2 σ excess over thermal emission SAX ť, BeppoSAX keV⁻¹ OSSE CH -5 10 $I_{x}(\mathrm{E}_{x})$ [photons Photons cm⁻² 10 10 10^{-6} 10^{-7} ٩ 10 Fusco-Femiano et al 2004 15 20 30 40 60 70 50 80 100 1000 Energy (keV) E_x (keV)

The hard X-ray excess in the center of Coma (r < 1 Mpc)

Possibly generated by ICS off CMB of the same electrons responsible for the radio halo (primary or secondary). Warning: in case of secondary model, the magnetic field needed may overproduce γ -rays (Blasi&Colafrancesco 1999)

Alternative: maybe due to a supra-thermal electron tail developed in the thermal electron distribution due to stochastic acceleration in the turbulent intra-cluster medium (Ensslin et al 1998)

Thermal gas at T=8.2 kev

Thermal gas in the filaments at T=0.22 kev

Non-thermal electrons

 Produced by astrophysical source and continuously reaccelerated by cluster turbolences or merger shock waves
 Produced by interaction of CRs with thermal ions

An Alternative Non Thermal Hypothesis: (although non asked for...) Relativistic electrons are produced by DM annihilation



MILLENNIUM Simulation CDM universe Springel et al. 2005

Simulates halos on cosmological scales, then resimulates a smaller patch with higher mass resolution down to cluster scale.

Tracks the formation of galaxies and quasars in the simulation, by implementing a semianalytic model to follow gas, star and supermassive black hole processes within the merger history trees of dark matter halos and their substructures



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DM Density profiles can be inferred from astronomical measurements or derived from numerical simulations



DM best fit to the radio halo spectrum of Coma



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1st step:

Magnetic field and particle physics are left as free parameters to fit the radio halo of Coma

Fig. 2. The radio halo spectrum of Coma and the best fits of three DM models with $M_{\chi} = 9$ GeV (upper panel) for bb (solid with $\sigma V = 1.7 \cdot 10^{-25}$ cm³ s⁻¹ and $B_0 = 20 \ \mu$ G) and τ^{\pm} (dashed with $\sigma V = 5.6 \cdot 10^{-25}$ cm³ s⁻¹ and $B_0 = 2.2 \ \mu$ G), $M_{\chi} = 60$ GeV (mid panel) for bb (solid with $\sigma V = 8.5 \cdot 10^{-23}$ cm³ s⁻¹ and $B_0 = 0.5 \ \mu$ G) and τ^{\pm} (dashed with $\sigma V = 9.0 \cdot 10^{-22}$ cm³ s⁻¹ and $B_0 = 0.1 \ \mu$ G), and with $M_{\chi} = 500$ GeV (lower panel) with bb (solid with $\sigma V = 2.0 \cdot 10^{-18}$ cm³ s⁻¹ and $B_0 = 0.01 \ \mu$ G), τ^{\pm} (dashed with $\sigma V = 1.9 \cdot 10^{-17}$ cm³ s⁻¹ and $B_0 = 0.002 \ \mu$ G) and W[±] (dotted with $\sigma V = 1.3 \cdot 10^{-16}$ cm³ s⁻¹ and $B_0 = 0.001 \ \mu$ G). Data from Thierbach et al. (2003). The results are obtained using $B(r) \propto n_{th}(r)$ and NFW smooth DM profiles, without considering the effect of substructures.

Colafrancesco, Lieu, Marchegiani, Pato & LP 2010

Multiwavelenght DM interpretation or exclusion?



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Compatibility with the cluster heating rate



Fig. 5. The heating rate induced by DM secondary particles in Coma as predicted by the three DM models with $M_{\chi} = 9$ GeV, $M_{\chi} = 60$ GeV and $M_{\chi} = 500$ GeV for $b\bar{b}$ (solid), τ^{\pm} (dashed), W[±] (dot-dashed) compositions. We consider an NFW DM profile (green, blue and red peaking curves for 9, 60 and 500 GeV respectively) and a cored DM profile (green, blue and red flattening curves for 9, 60 and 500 GeV respectively). The values of σV and B_0 for the different models are the same used in Fig.2. The solid black curve shows the bremsstrahlung cooling rate of the intra-cluster gas at a temperature of 8.2 keV. The results are obtained using $B(r) \propto n_{th}(r)$ and smooth DM profiles, without considering the effect of substructures.

3rd step:

The heating rate of the inctracluster gas due to Coulomb collision of low-energy non-thermal electrons should not exceed the bremmstrahlung cooling rate of thermal electrons

otherwise the heated gas would get a temperature higher than the one observed in the cluster, which is related to the cooling rate of thermal gas. We would observe a fast gas heating and expansion, while the cluster is thermally stable.

$$\frac{dE}{dtdV} = \int dE \frac{dn_e}{dE} \cdot \left(\frac{dE}{dt}\right)_{Coul}$$

Colafrancesco, Lieu, Marchegiani, Pato & LP 2010

Compatibility with multimessenger constraints



Cross-sections are compared with available constraints from GC γs , diffuse γs , antimatter, CMB, radio ... which excludes ANY dark matter interpretation for smooth profiles

Look at the upper curves: smooth cluster halo All DM explanation are excluded





Adding subhalos: modeling the structure of dark matter halos

 $\begin{array}{l} \mbox{Halos form through a hierarchical process of successive mergers.} \\ \mbox{The halo of our Galaxy will be self-similarly composed by:} \\ \mbox{-a smoothly distributed component } (\rho^2_{\mbox{DM(h)}} \mbox{ single halo }) \\ \mbox{-a number of virialized substructures } (\rho^2_{\mbox{DM(subh)}} \mbox{ all halos}) \end{array}$



Make use of simulations on galactic scale and use self-similarity arguments to infer cluster properties. Note: self-similarity proven from cluster to galactic scale Adding subhalos: modeling the structure of dark matter halos

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N-body simulations study the smooth halo and the larger halos (M> $10^5 M_{sun}$).

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Microphysics and theory of structure formation sets the mass of the smallest halo because there is no enough cpu power to simulate small halos from collapse till today.

Modeling the structure of dark matter halos from theory of structure formation (M< $10^5 M_{sun}$)

Theory: Damping of the primordial power spectrum due to CDM free streaming or acoustic oscillations after kinetic decoupling

Typical M_{min} for a WIMP = 10⁻⁶ M_{sun}







MW-like halos at z=0

Via Lactea 2, Diemand et al

Aquarius, Springel et al

*Note σ_8 =0.8 (WMAP 7yr)



Halo and subhalo profile shape and concentration



Halo and subhalo profile shape and concentration

Aq-A-1

Concentration parameter differ (because of σ_8)

Concentration parameter (R_{vir}/r_s) has radial dependence higher concentration -> higher flux!

The higher concentration parameter at small radial distance from the GC reflects:

1) The halo had to be more concentrated not to be disrupted by tides, encounters, etc.

2) The closer to the center, the larger is the subhalo permanence in the parent halo, i.e. the older is the subhalo. Older subhalos are the ones that formed at higher σ -peak of the fluctuation density field, i.e. more concentrated than halos of same mass which formed later - at $M_h = M^*(z)^{-1}$

Subhalo abundance and density distribution



Note the different subhalo definition (v_{max} VS mass) Slope -1.95 is consistent with both simulations within the fit errors

Compatibility with multimessenger constraints adding subhalos



Subhalo population

In presence of a population of substructures with M_{min} =10⁻⁶ Msun and radial dependence of the concentration parameter, a boost of ~ 35 still let some models allowed, providing a favourable environment (MW DM structure and propagation model)

Note that subhalos are also needed to explain the surface brigthness profile of the radio halo



The multiwavelength/multimessenger/multitarget approach

 Φ = ParticlePhysics x Cosmology/Astrophysics x Transport



Slide: courtesy of M. Pato

The γ -ray sky Φ_{γ} = $\Phi_{\text{particle physics}} \times \Phi_{\text{cosmology}}$

MW smooth and single subhalo contribution

$$\Phi_{COSMO}^{halo}(M,R,r) \propto \int_{l.o.s.} d\lambda d\Omega \left[\frac{\rho_{DM}^2(M,c(M,R),r,\psi))}{d^2} \right]$$

Integrated contribution of all the GALACTIC halos along the LOS

$$\Phi^{\text{allhalos}} cosmo(\psi, \Delta \Omega) \propto \int_{M} dM \int_{C} dc \iint_{\Delta \Omega} d\vartheta d\phi \int_{\text{los}} d\lambda \rho_{\text{sh}}(M, R) \cdot P(c) \Phi^{\text{halo}}_{COSMO}$$

Integrated contribution of EXTRAGALACTIC halos and subhalos

Computing the cosmological γ -ray flux due to DM annihilation in halos and subhalos

$$\frac{d\phi_{\gamma}}{dE_{0}} = \frac{\sigma v}{8\pi} \frac{c}{H_{0}} \frac{\overline{\rho}_{0}^{2}}{m_{\chi}^{2}} \int dz (1+z)^{3} \frac{\Delta^{2}(z)}{h(z)} \frac{dN_{\gamma}(E_{0}(1+z))}{dE} e^{-\tau(z,E_{0})} e^{-\tau(z,E_{0})} \frac{dI_{0}}{dI_{0}} \frac{M_{\gamma}(E_{0}(1+z))}{h(z)} e^{-\tau(z,E_{0})} \frac{dI_{0}}{dI_{0}} \frac{M_{\gamma}(E_{0}(1+z))}{dI_{0}} \frac{dE_{0}}{dI_{0}} \frac{dI_{0}}{M_{\gamma}} \frac{M_{\gamma}(E_{0}(1+z))}{dI_{0}} \frac{dE_{0}}{dI_{0}} \frac{M_{\gamma}(E_{0}(1+z))}{dI_{0}} \frac{dI_{0}}{M_{\gamma}} \frac{M_{\gamma}(E_{0}(1+z))}{dI_{0}} \frac{dE_{0}}{dI_{0}} \frac{M_{\gamma}(E_{0}(1+z))}{dI_{0}} \frac{M_{\gamma}(E_{0}(1+z))}{M_{\gamma}(E_{0}(1+z))} \frac{M_{\gamma}(E_{0}(1+z))}{M_{\gamma}(E_{0}(1+z))}} \frac{M_{\gamma}(E_{0}(1+z))}{M_{\gamma}(E_{0}(1+z))}} \frac{M_{\gamma}(E_{0}(1+z))}{M_{\gamma}(E_{0}(1+z))}} \frac{M_{\gamma}(E_{0}(1+z))}{M_{\gamma}(E_{0}(1+z))}} \frac{M_{\gamma}(E_{0}(1+z))}{M_{\gamma}(E_{0}(1+z))}} \frac{M_{\gamma}(E_{0}(1+z))}{M_{\gamma}(E_{0}(1+z))}} \frac{M_{\gamma}(E_{0}(1+z))}{M_{\gamma}(E_{0}(1+z))}}$$

The $\gamma\text{-ray sky}$ Galactic and extragalactic: smooth + subhalos



LP, Lavalle, Bertone & Branchini 2009



The antimatter sky - coherent halo description wrt γ -rays

 $n_{CR}(t, \bar{x}, E_{CR}) = \frac{d^2 N_{CR}}{dV dF_{cR}}$

Compute the number density à la Delahaye et al. 2008

electrons and positrons



Compute fluxes and boosts à la Lavalle et al. 2008

$$\phi_{CR,sm}(E_{CR}) \propto < \sigma v > \int_{E_{CR}}^{\infty} dE \frac{dN_{CR}}{dE} \int_{diff.zone} d^{3}\bar{x} \left(\frac{\rho_{sm}(\bar{x})}{\rho_{sun}}\right)^{2} G_{sun}^{CR}(\bar{x}, \lambda_{D})$$

$$< \phi_{CR,cl} > (E_{CR}) \propto < \sigma v > N_{cl} \int_{E_{CR}}^{\infty} dE \frac{dN_{CR}}{dE} \int_{diff.zone} d^{3}\bar{x} < \xi >_{M} (R) \frac{dP_{V}}{dV}(R) G_{sun}^{CR}(\bar{x}, \lambda_{D}) = N_{tot}^{sub} < \phi_{sub} > 0$$

The antimatter sky - coherent halo description wrt γ -rays



Compute fluxes and boosts à la Lavalle et al. 2008



$$\phi_{CR,sm}(E_{CR}) \propto <\sigma v > \int_{E_{CR}}^{\infty} dE \frac{dN_{CR}}{dE} \int_{diff.zone}^{d^{3}} \bar{x} \left(\frac{\rho_{sm}(\bar{x})}{\rho_{sun}}\right)^{2} G_{sun}^{CR}(\bar{x},\lambda_{D})$$

$$<\phi_{CR,cl}>(E_{CR}) \propto <\sigma_{V}>N_{cl}\int_{E_{CR}}^{\infty} dE \frac{dN_{CR}}{dE} \int_{diff.zone} d^{3}\vec{x} <\xi>_{M} (R)\frac{dP_{V}}{dV}(R)G_{sun}^{CR}(\vec{x},\lambda_{D}) = N_{tot}^{sub} <\phi_{sub}>$$

The radio sky GC modeled coherently with γ-rays and antimatter

Compute synchtrotron power



Constraints from CMB - no structure dependence -

Injection of DM annihilation around z=1000 would affect recombination and hence modify the CMB



Multi³ constraints on annihilation cross-section



Different messenger play different roles for different channels Yet the amount of exclusion is almost the same..

Multi³ constraints on annihilation cross-section



Catena, Fornengo, Pato, LP & Masiero 2010

In order to get bands of exclusion we change profile (Via Lactea II or Aquarius with subhalos, isocored without subhalos) and propagation parameters (inside the MIN-MED-MAX propagation model)



Pato, LP, Bertone 2009



Minimal Dark Matter models are excluded



Leptonic models are excluded



Nomura&Thaler models are excluded

Arkani-Hamed et al. model is the only one surviving



Conclusions

Multi³ analysis must be applied to any candidate claimed to explain any excess

In the case of Coma, it proved that the DM explanation of the radio halo is possible only under favourable environment (profile and propagation)