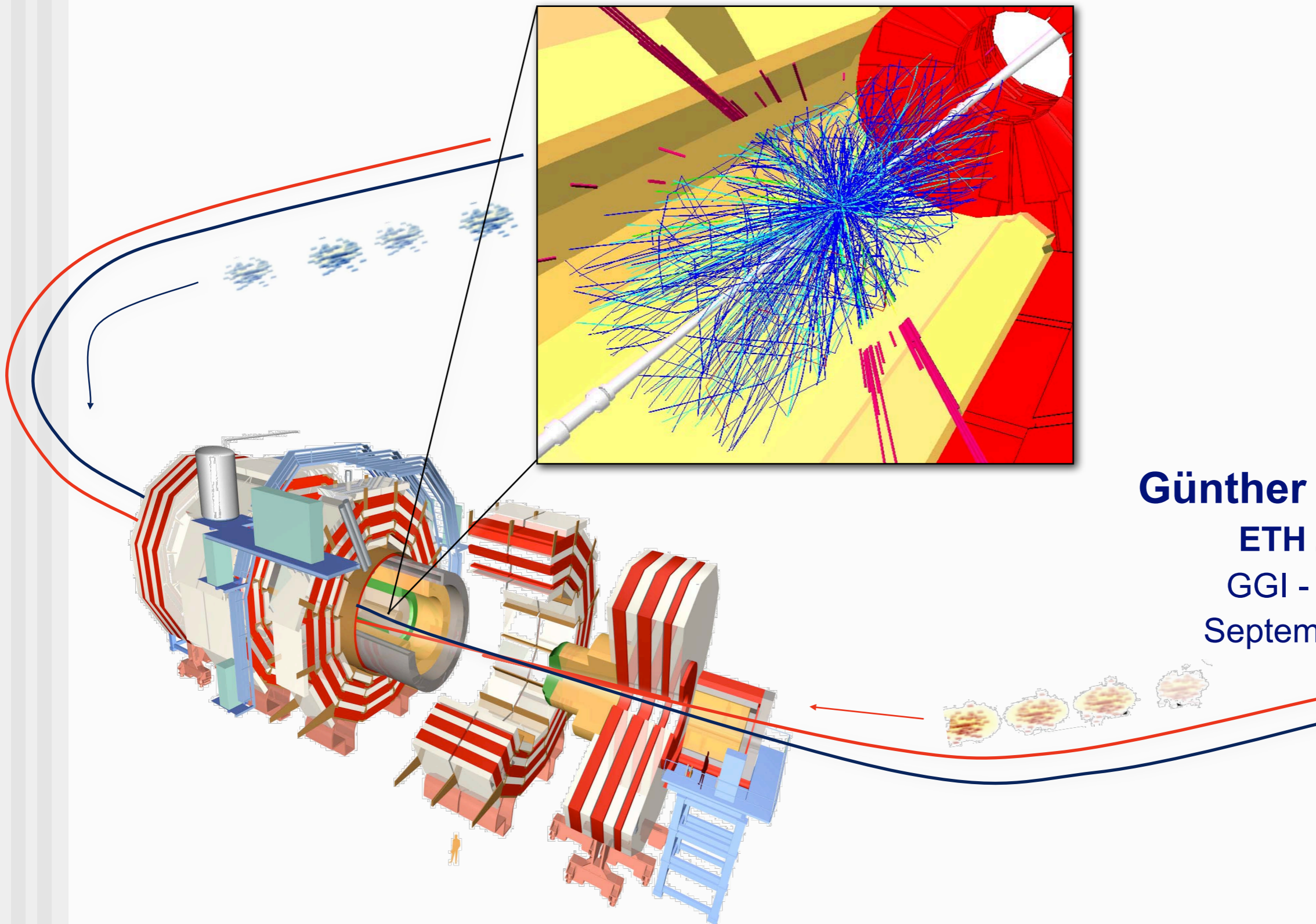




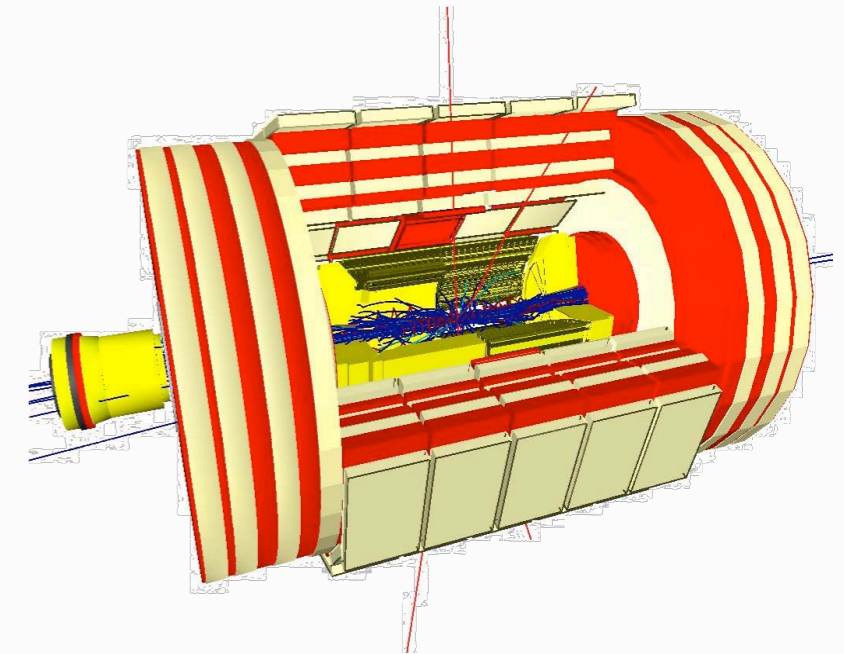
# LHC Detectors : Part 3



**Günther Dissertori**  
ETH Zürich  
GGI - Firenze  
September 2007

## What is measured, how and why?

- Basic processes, rates
- Resulting difficulties and requirements
- Basic detector layout

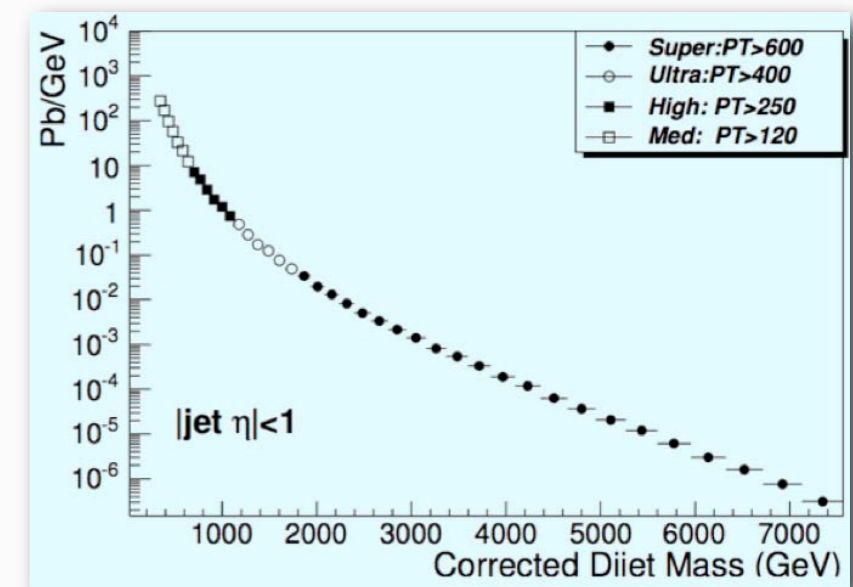


## ATLAS and CMS

- Overview
- Construction status
- Comparison

## Experimental issues

- Some examples of experimental issues to be addressed
- such as Jet Energy Calibration
- and background estimations



Disclaimer 1 : I concentrate on multi-purpose detector ATLAS and CMS

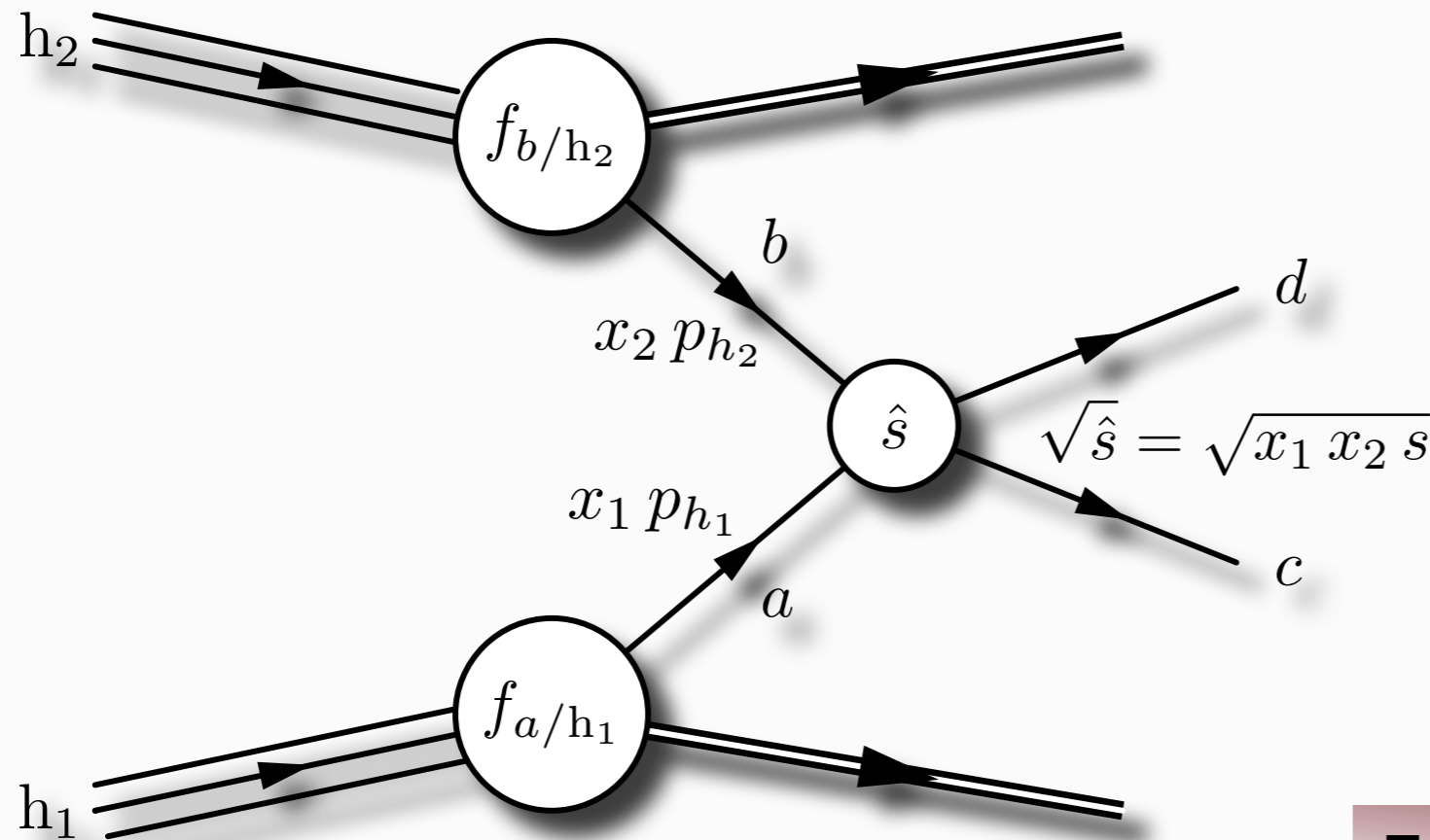
Disclaimer 2 : Some slides or slide content taken from seminars/lectures of other LHC colleagues, eg. K. Jakobs, O. Buchmüller, L. Dixon, M. Dittmar, D. Froidevaux, F. Gianotti



# Introduction : Measurements of hard processes



# The hard scattering



To produce (at central rapidity, ie.  $x_1 \sim x_2$ ) a mass of

	LHC	TEVATRON
100 GeV	$x \sim 0.007$	0.05
5 TeV	$x \sim 0.36$	--

From where do we know these?

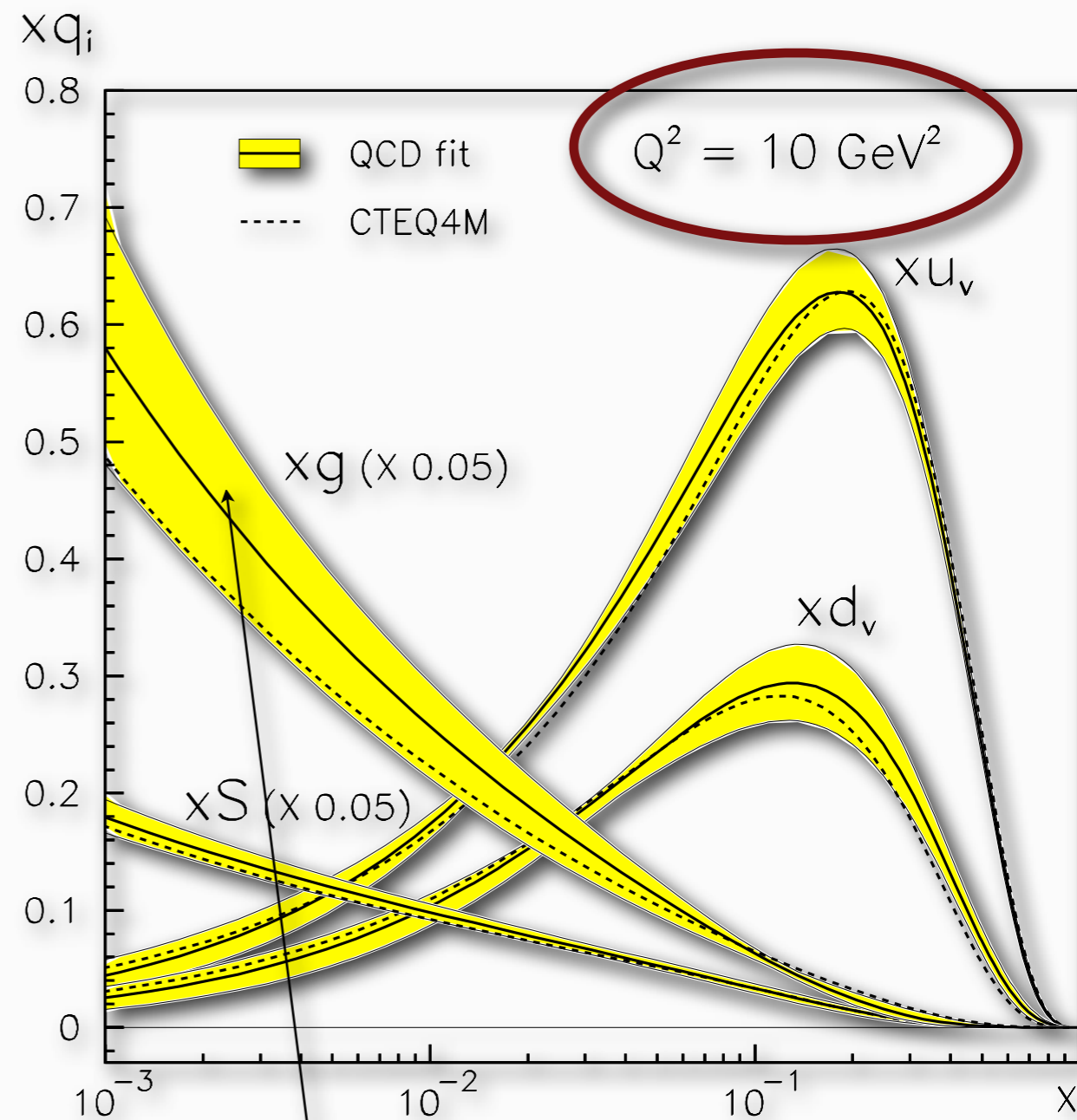
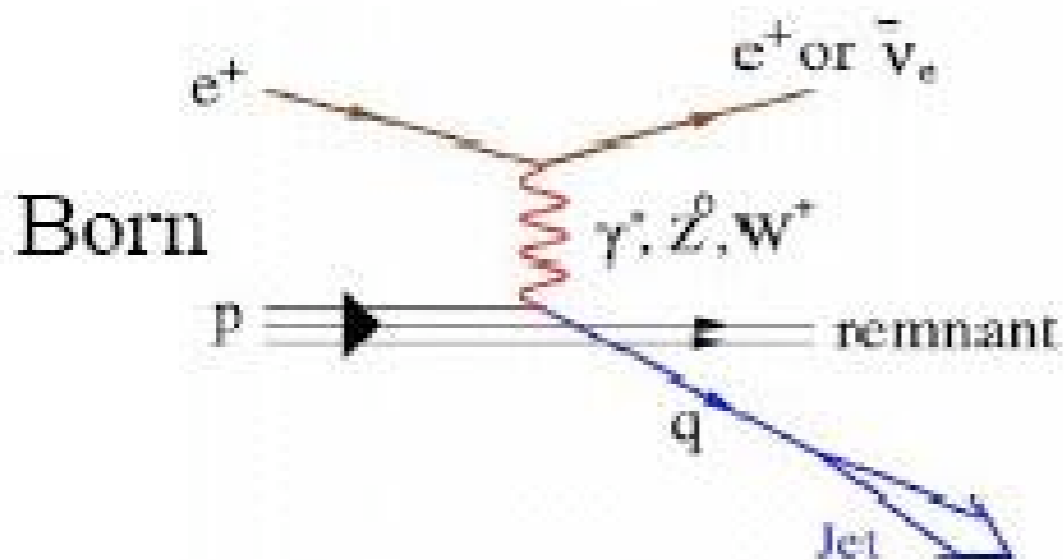
$$d\sigma(h_1 h_2 \rightarrow cd) = \int_0^1 dx_1 dx_2 \sum_{a,b} f_{a/h_1}(x_1, \mu_F^2) f_{b/h_2}(x_2, \mu_F^2) d\hat{\sigma}^{(ab \rightarrow cd)}(Q^2, \mu_F^2)$$

**Hard Scattering** = processes with large momentum transfer ( $Q^2$ )

Represent only a tiny fraction of the total inelastic pp cross section ( $\sim 70$  mb)

eg.  $\sigma(pp \rightarrow W+X) \sim 150$  nb  $\sim 2 \cdot 10^{-6} \sigma_{\text{tot}}(pp)$

HERA ep accelerator, 6.3 km circumference

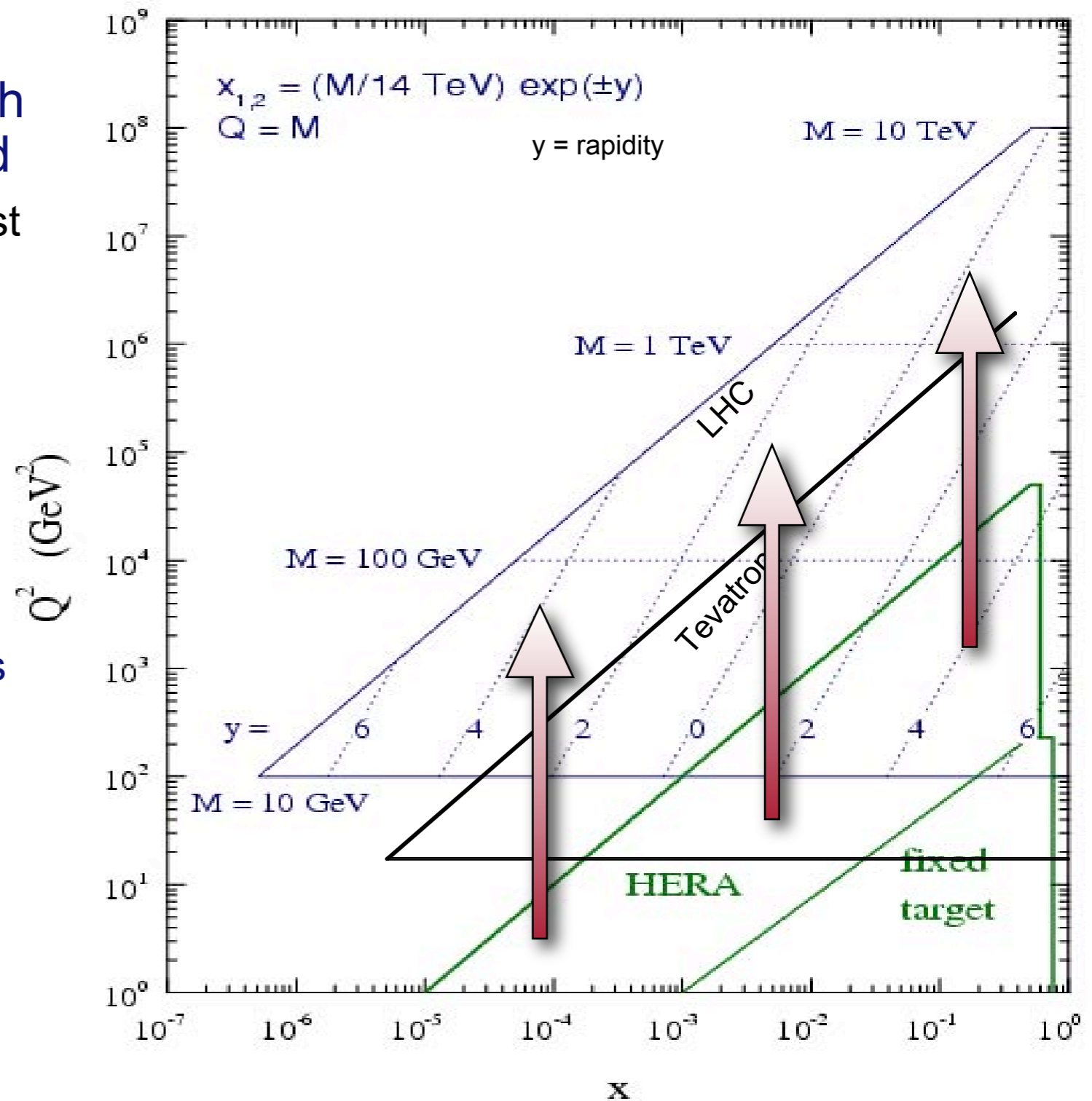


gluons dominate at low x !

Scattering of 30 GeV electrons on 900 GeV protons:  
 → Test of proton structure down to  $10^{-18}$  m

→ the LHC is a gluon-gluon collider !

- Kinematic regime for LHC much broader than currently explored
  - for example, HERA covers most of the relevant  $x$  range, but at much smaller values of  $Q^2$
- Is NLO DGLAP evolution sufficient for LHC?
- Have to propagate correctly the uncertainties of PDF determinations into predictions of LHC processes
  - important when comparing to data
- Have to determine / constrain the pdfs at LHC itself
  - what are useful processes for this?





# Event rates

Event production rates at  $L=10^{33} \text{ cm}^{-2} \text{ s}^{-1}$  and statistics to tape

Process	Events/s	Evts on tape, $10 \text{ fb}^{-1}$
$W \rightarrow e\nu$	15	$10^8$
$Z \rightarrow ee$	1	$10^7$
$t\bar{t}$	1	$10^6$
Minimum bias	$10^8$	$10^7$
QCD jets $p_T > 150 \text{ GeV}/c$	$10^2$	$10^7$
$b\bar{b} \rightarrow \mu X$	$10^3$	$10^7$
gluinos, $m=1 \text{ TeV}$	0.001	$10^3$
Higgs, $m=130 \text{ GeV}$	0.02	$10^4$

assuming 1% of trigger bandwidth

$10^7$  events to tape every 3 days, assuming 30% data taking efficiency, 1 PB/year/exp

statistical error negligible after few days (in most cases) !  
dominated by systematic errors (detector understanding, luminosity, theory)



# First Physics runs ( 2008, 2009, ... )

## After first “good” $10 \text{ pb}^{-1}$

- many jets...
- $\sim 20000$  W, decaying to lepton + neutrinos
- $\sim 2500$  Z, decaying into two leptons
- $\sim 200$  semi-leptonic top-pair events
  - Measure rates, align and calibrate better

## After first “good” $100 \text{ pb}^{-1}$

- W(Z)+jets rates well measurable
  - Jet calibration, MET calibration (for SUSY)
- Inclusive leptons, di-leptons, photons, di-photon triggers (for Higgs)

## From $100 \text{ pb}^{-1}$ to $1 \text{ fb}^{-1}$

- Standard model candles
  - Top pair prod., W/Z cross sections, PDF studies, QCD studies, b-jet production
  - Do extensive MC tuning
- Early Higgs boson search
  - $H \rightarrow \gamma\gamma, WW, ZZ$
- Early SUSY-BSM searches
  - MET + anything, di-jet, di-leptons, di-photon, resonances....





# Why SM physics?

- Interesting in its own right
  - measure (calculable) event rates, cross sections
  - establish (dis)agreement with SM, constrain SM
  - challenge theoretical calculations at high  $Q^2$
  - demonstrate “working” experiment with well known processes
- Understanding, commissioning, calibration of detector and the software. eg. well suited :  $Z \rightarrow e^+e^-$
- Backgrounds to many searches : check MC simulations
  - eg. W/Z+jets, Multi-Jets, top-pair events
- Constrain (relative) PDFs
- Alternative measurements of luminosity



# Our Master Equation

Event rates (absolute, relative, differential)  
Stat vs syst errors, backgrounds from data or MC?  
Resolution, Energy Scale, Signal Significance

$$\sigma_{\text{meas}} = \frac{N_{\text{obs}} - N_{\text{bkg}}}{\epsilon L}$$

Proton-Proton Luminosity  
uncertainty ~ 5% ? Do better?

**Experimental issues** : Triggers, reconstruction, isolation cuts, low- $p_T$  jets (jet veto)

**Theoretical issues** :  $p_T$  distributions at NLO + resummation;  
differential calculations for detectable acceptance.

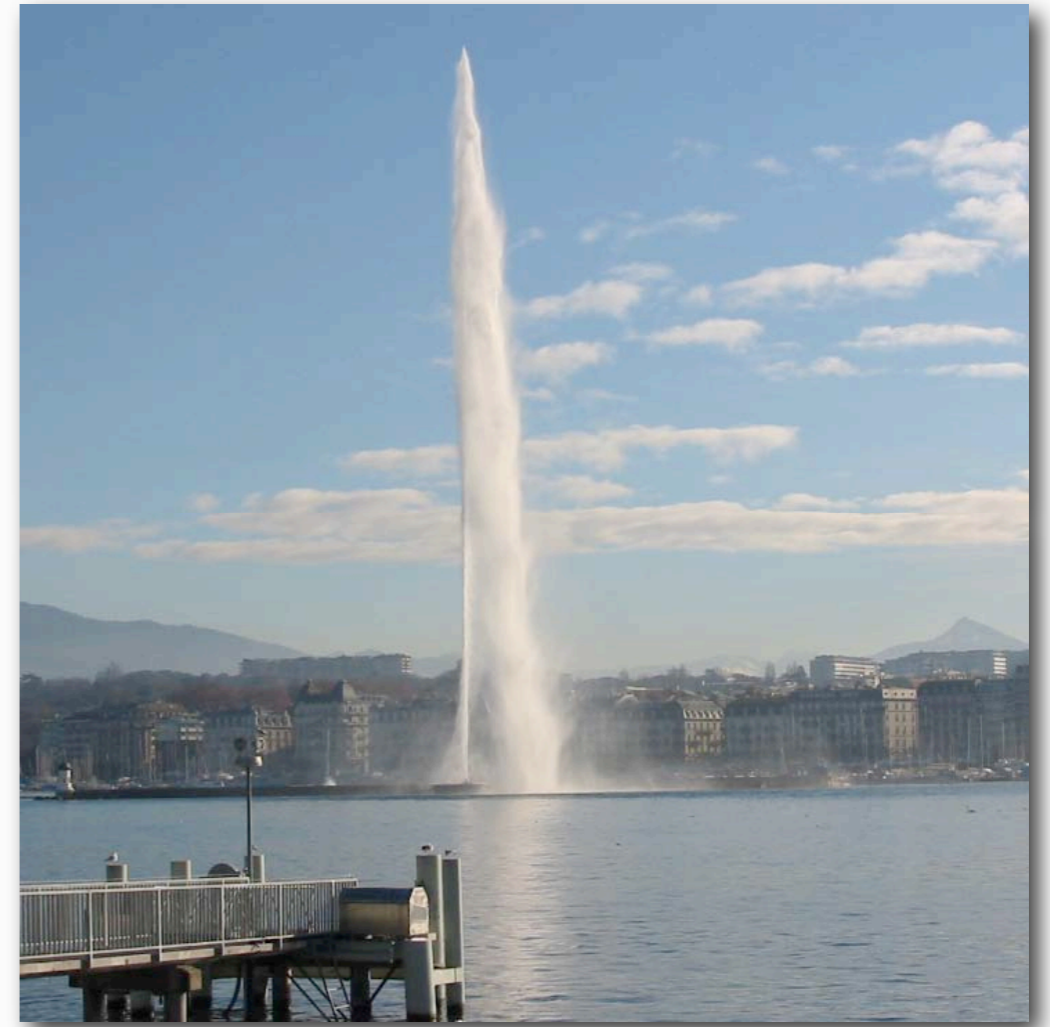
$$\sigma_{\text{theo}} = PDF(x_1, x_2, Q^2) \otimes \hat{\sigma}_{\text{hard}}$$

constrain, define uncertainties

HO calculations,  
implement in MC

Goal : test SM (in)consistency :  $\sigma_{\text{exp}} \pm \Delta_{\text{exp}} \stackrel{?}{=} \sigma_{\text{SM}} \pm \Delta_{\text{th}}$

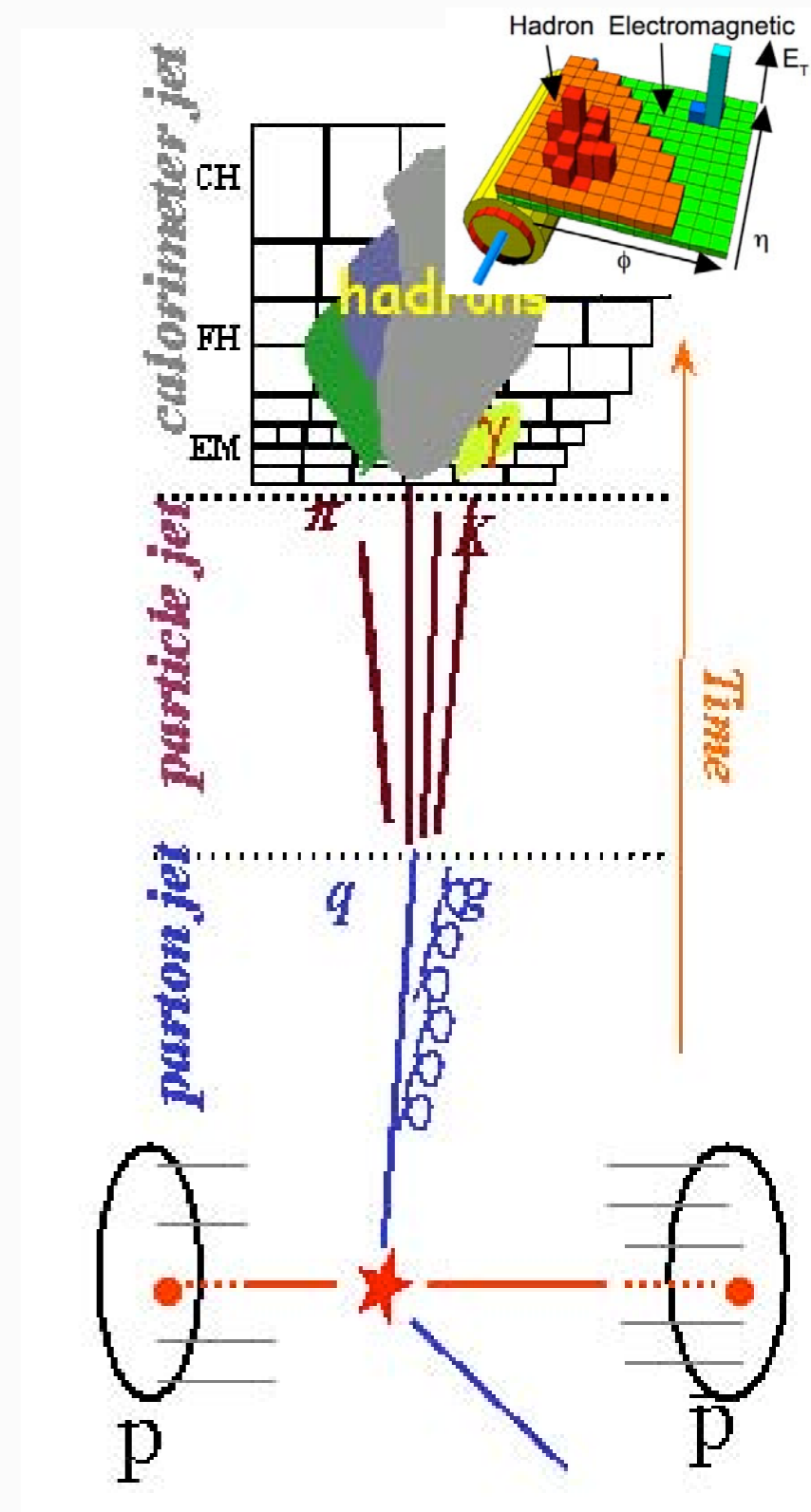
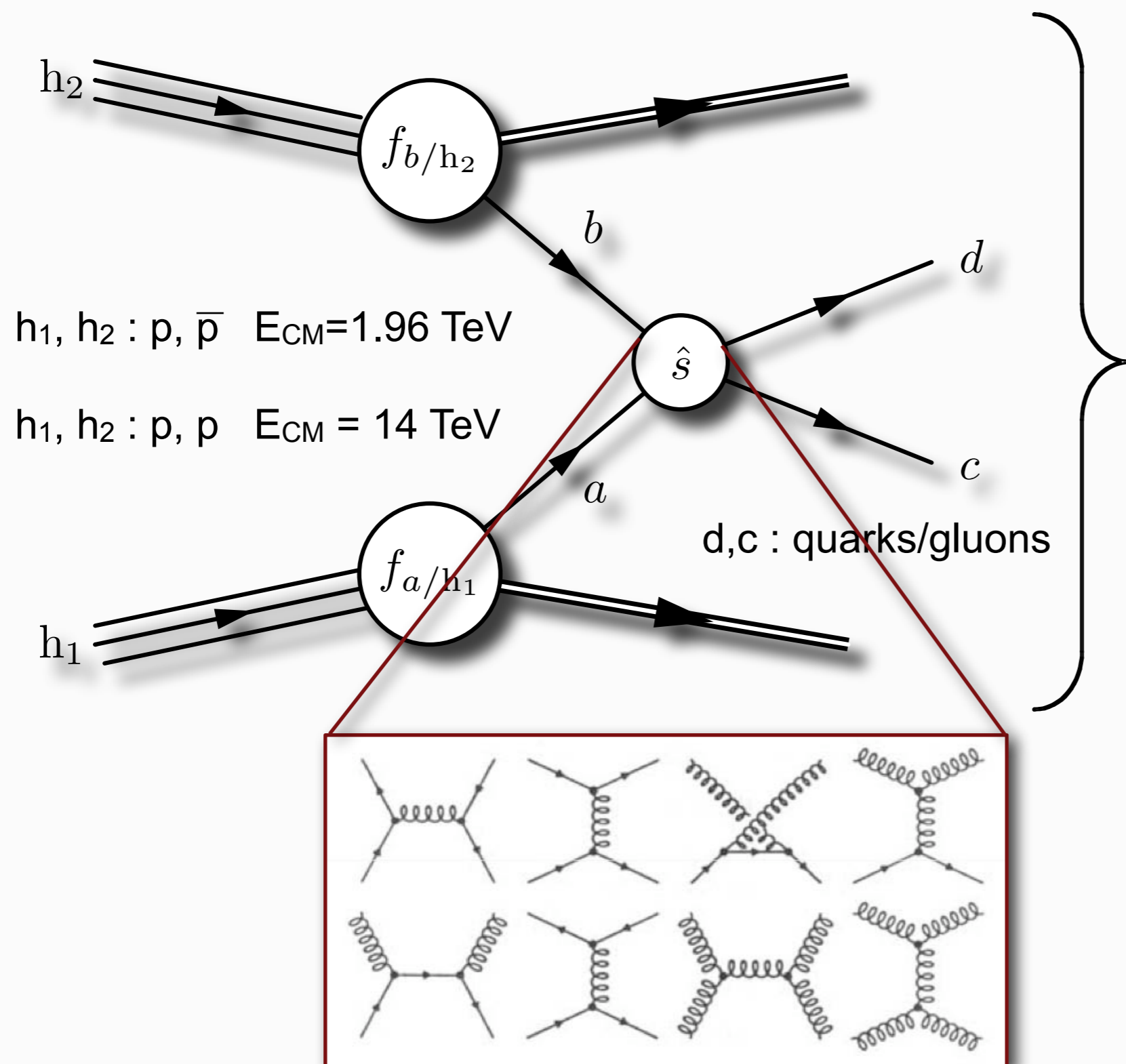
# Jets

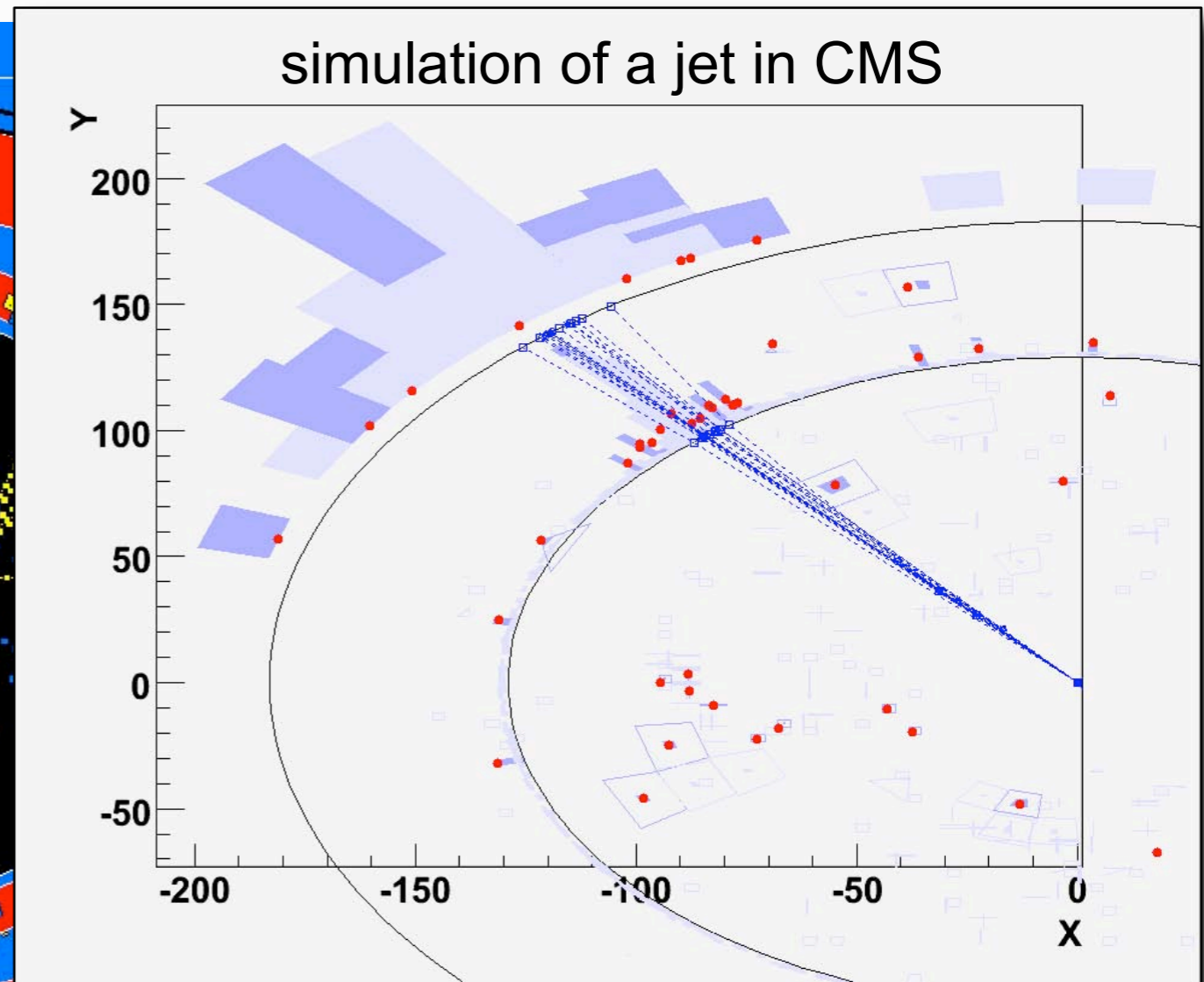
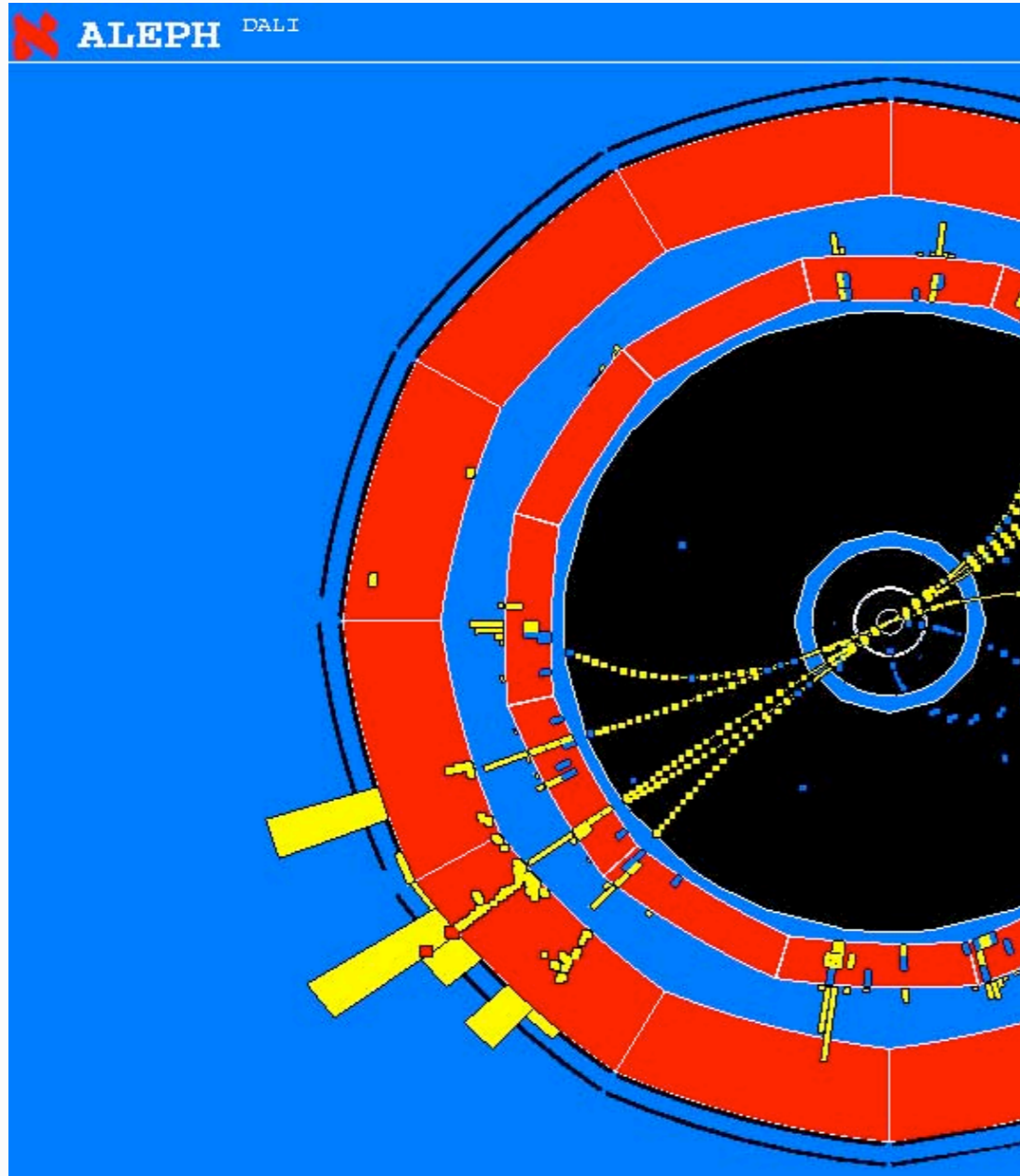




# JET production at hadron colliders

at the Tevatron, or in the future at the LHC





*“cluster/spray of particles (tracks, calorimeter deposits) or flow of energy in a restricted angular region”*

- clear : need some algorithmic definition



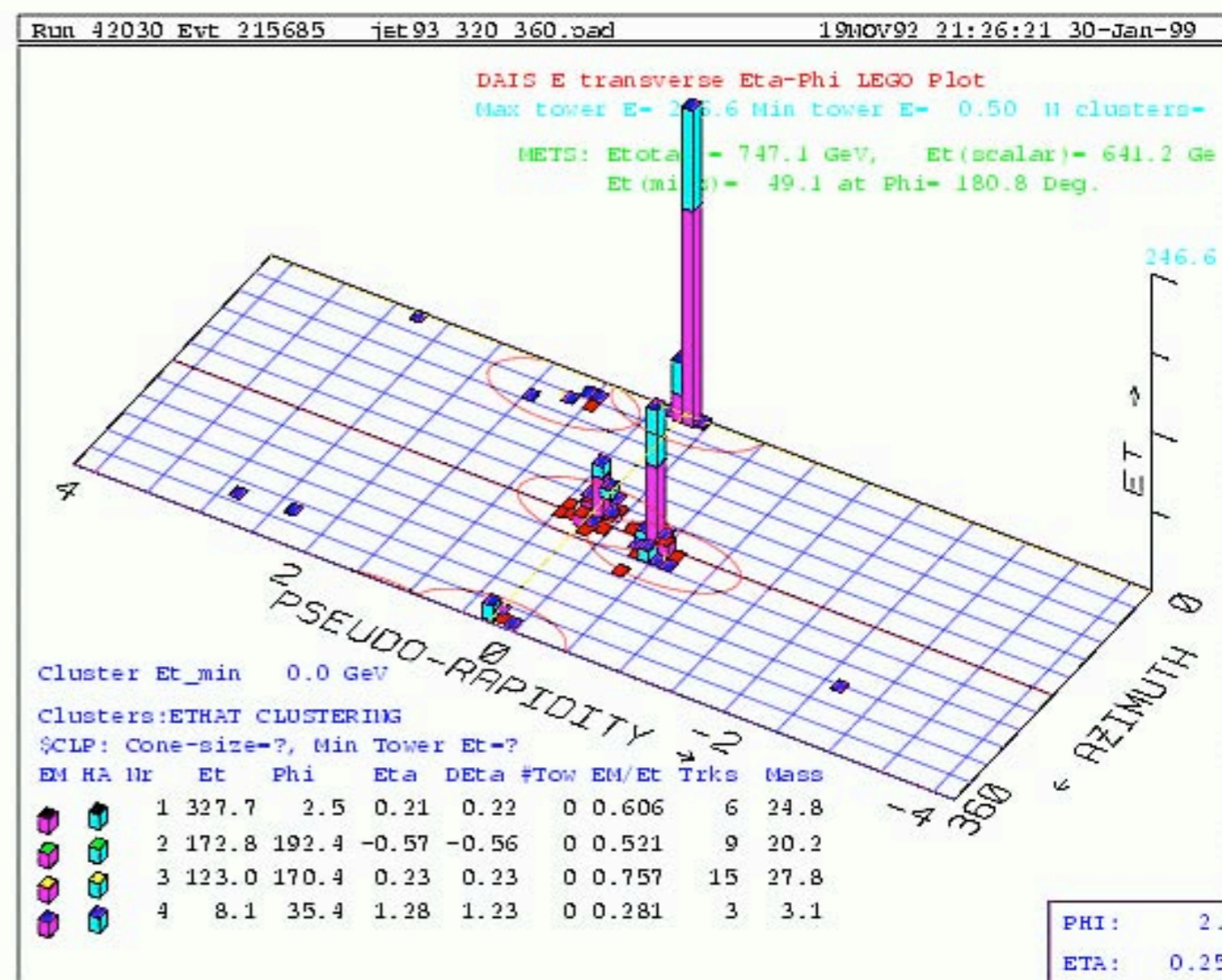
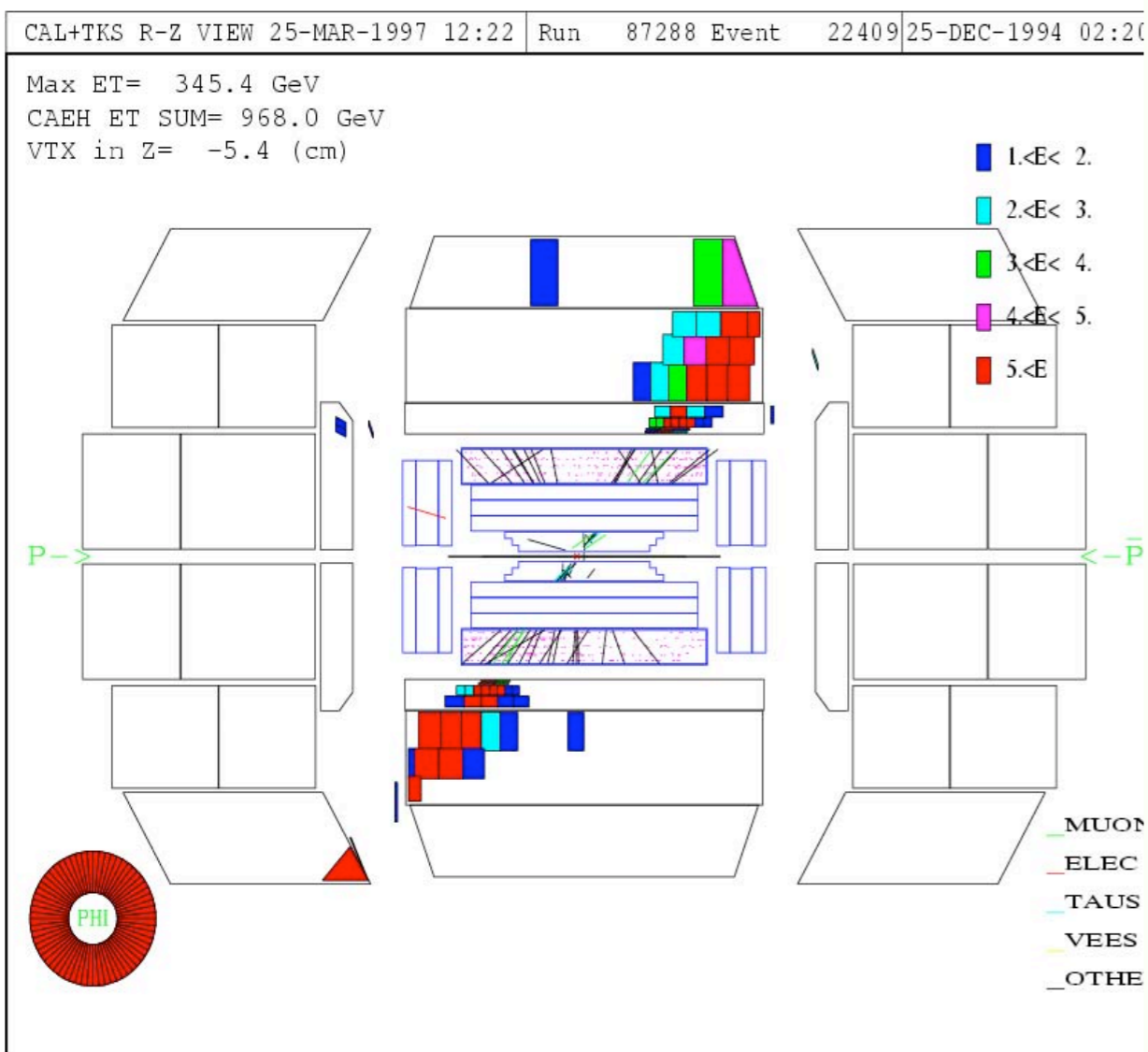
# A short digression: Jet Algorithms



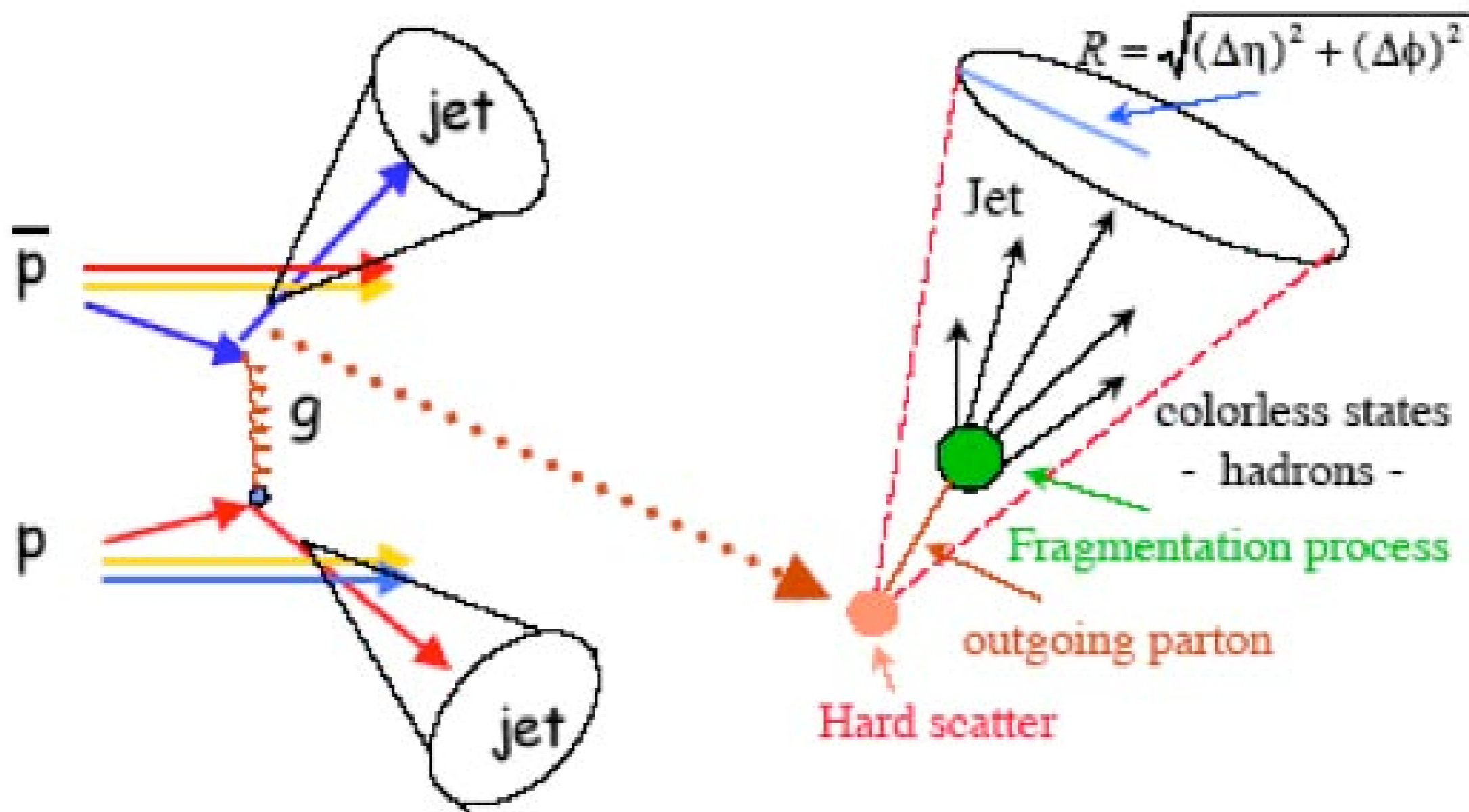
# Jets in Hadron Collider Detectors

## Jets in DØ

## CDF



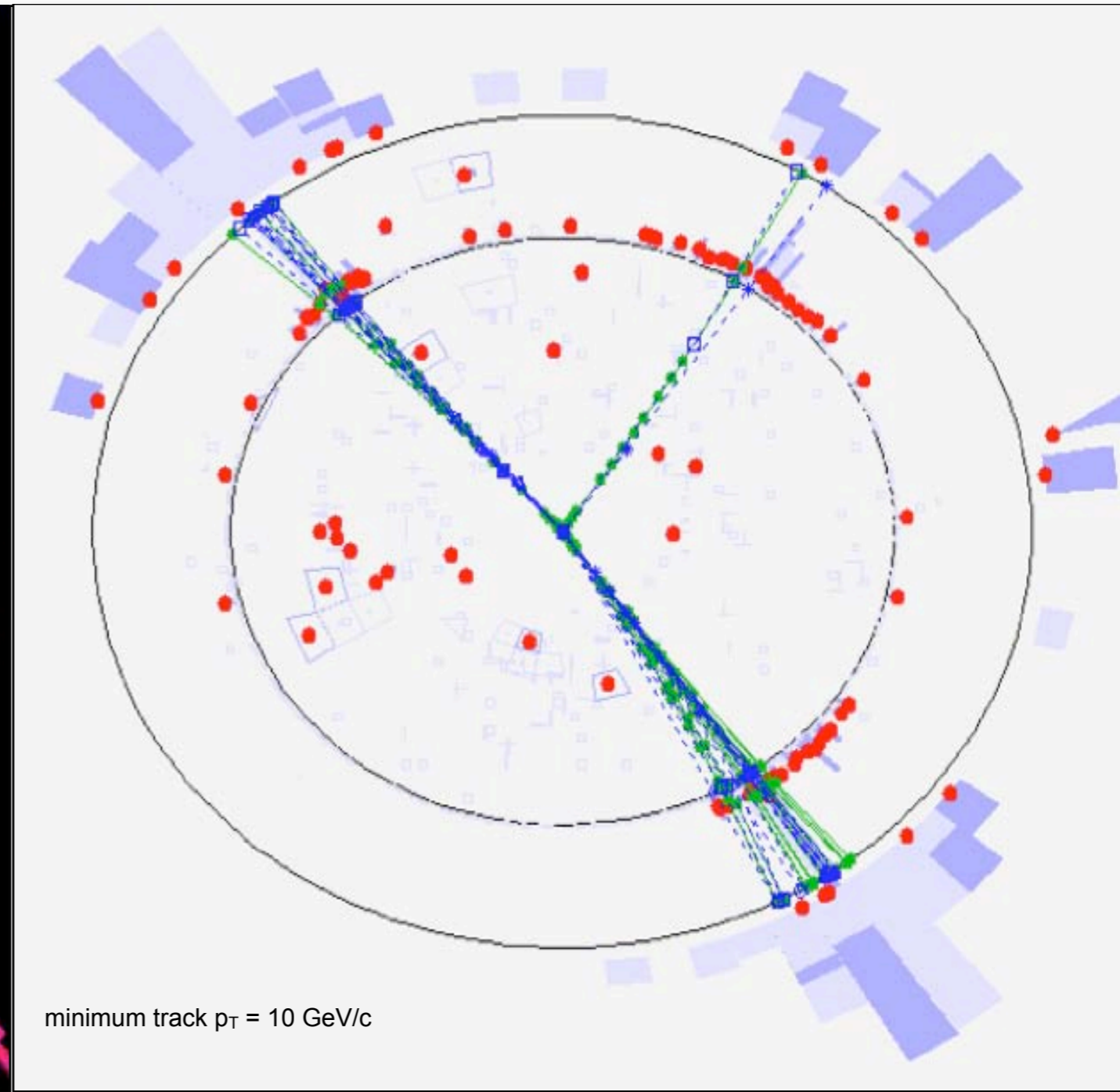
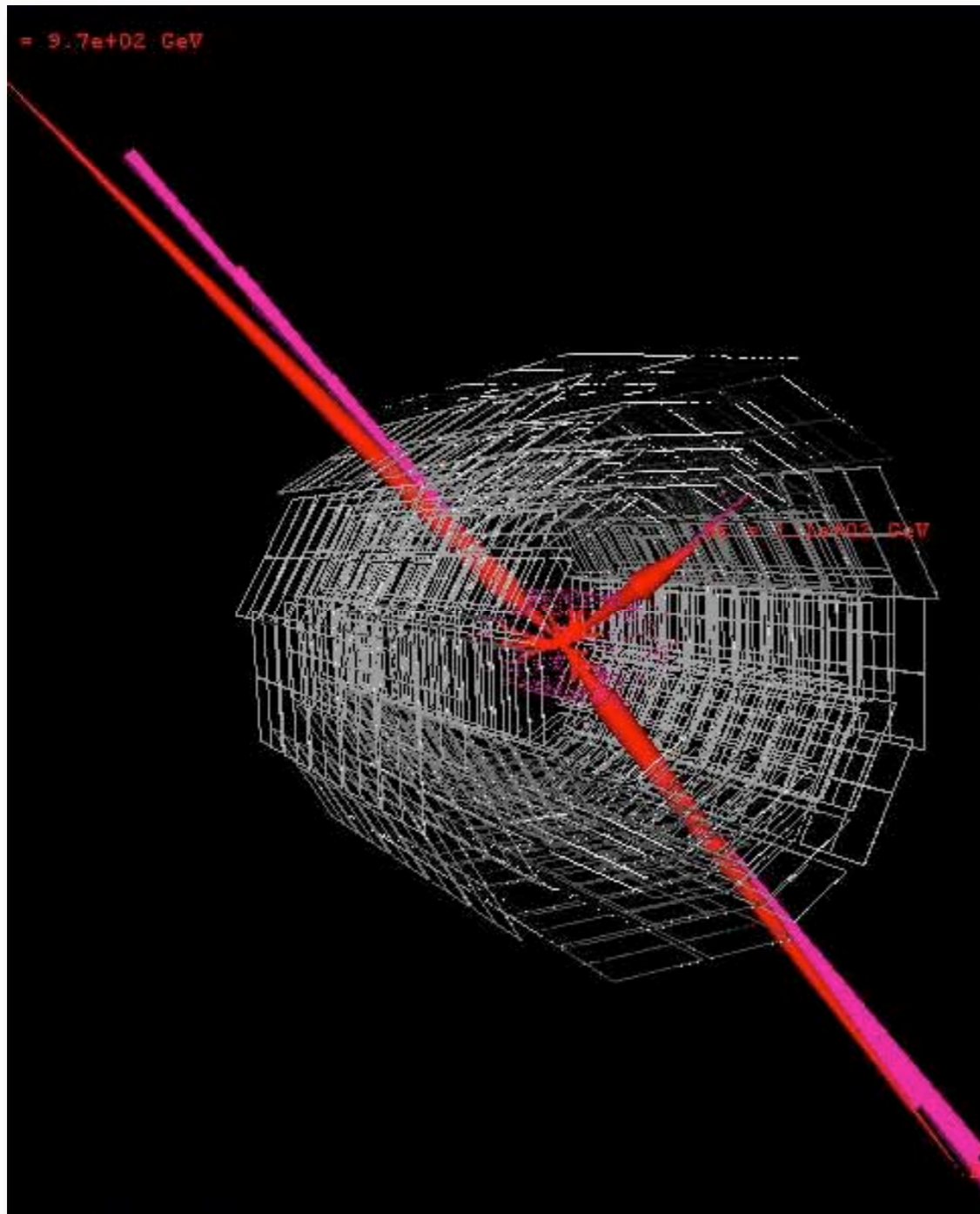
- Introducing a cone prescription seems “natural”...
- But how to make it more quantitative?
  - don't want people “guessing” at whether there are 2,3, ... jets



The natural (?) definition of a jet in a hadron collider environment



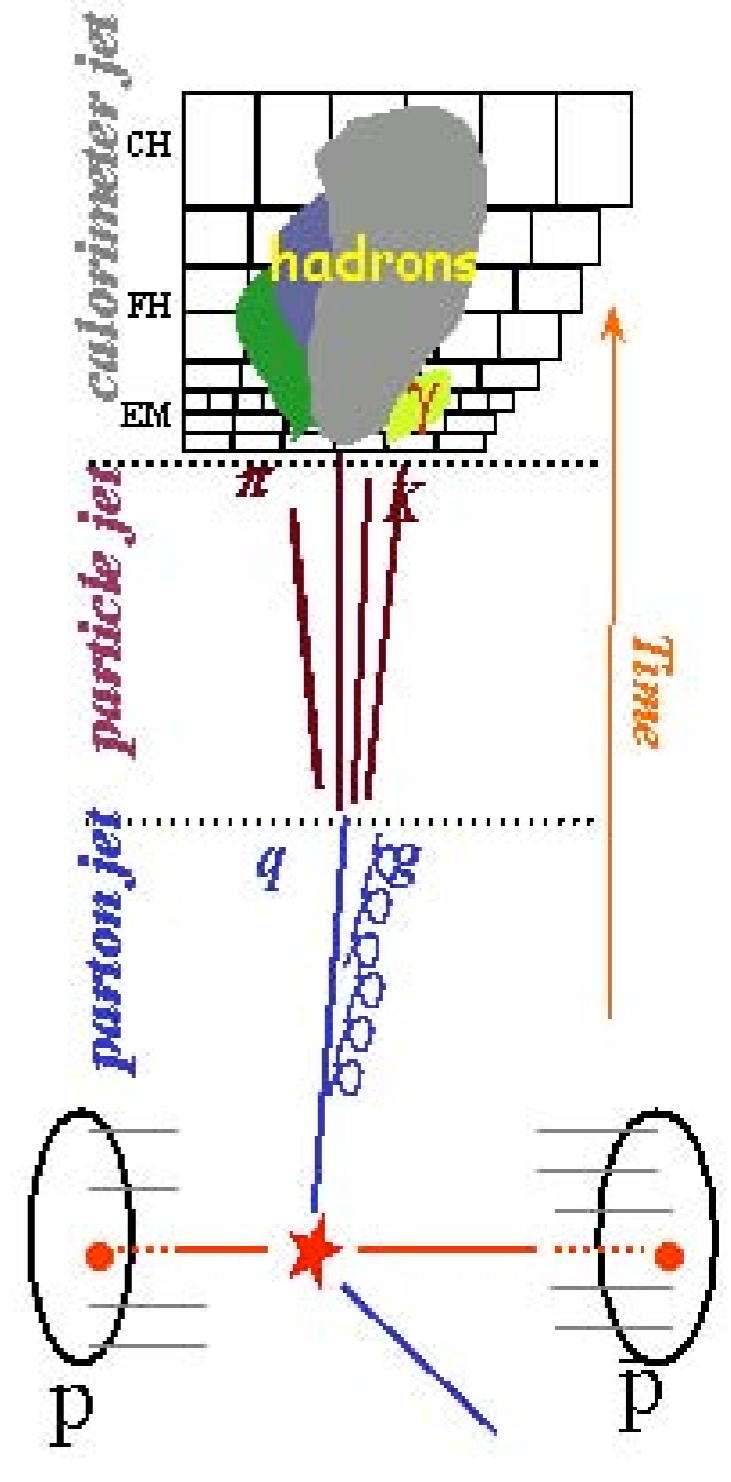
## (simulated) Jets in CMS



- Applicable at all levels
  - partons, stable particles
    - ◆ for theoretical calculations
  - measured objects (calorimeter objects, tracks, etc)
  - and always **find the same jet**
  
- Independent of the very details of the detector
  - example : granularity of the calorimeter, energy response,...
  
- Easy to implement !
  
- Close correspondence between

$$P_{\text{parton}} \longleftrightarrow P_{\text{jet}}$$

Energy  
Momentum  
angle

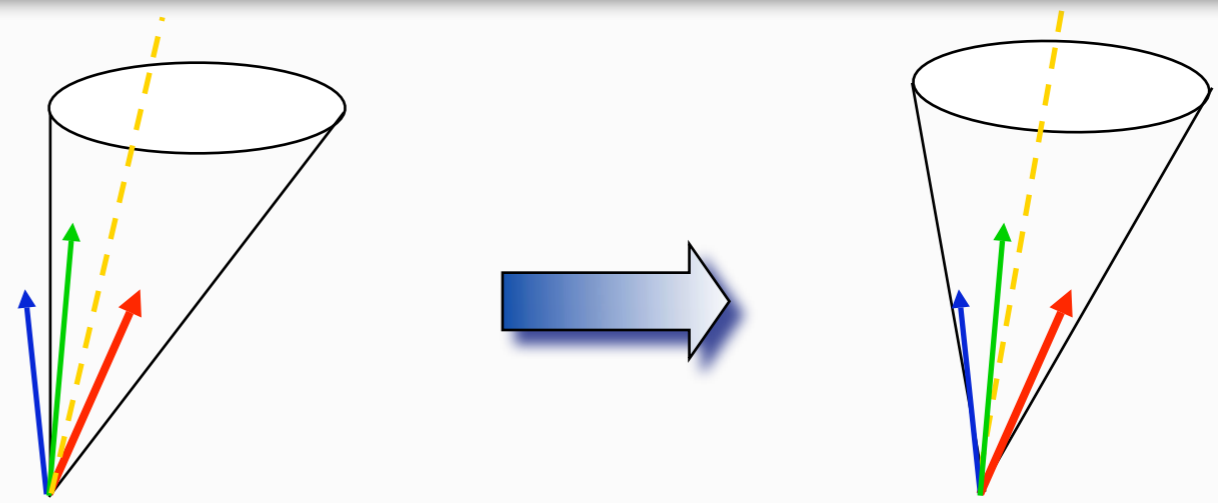
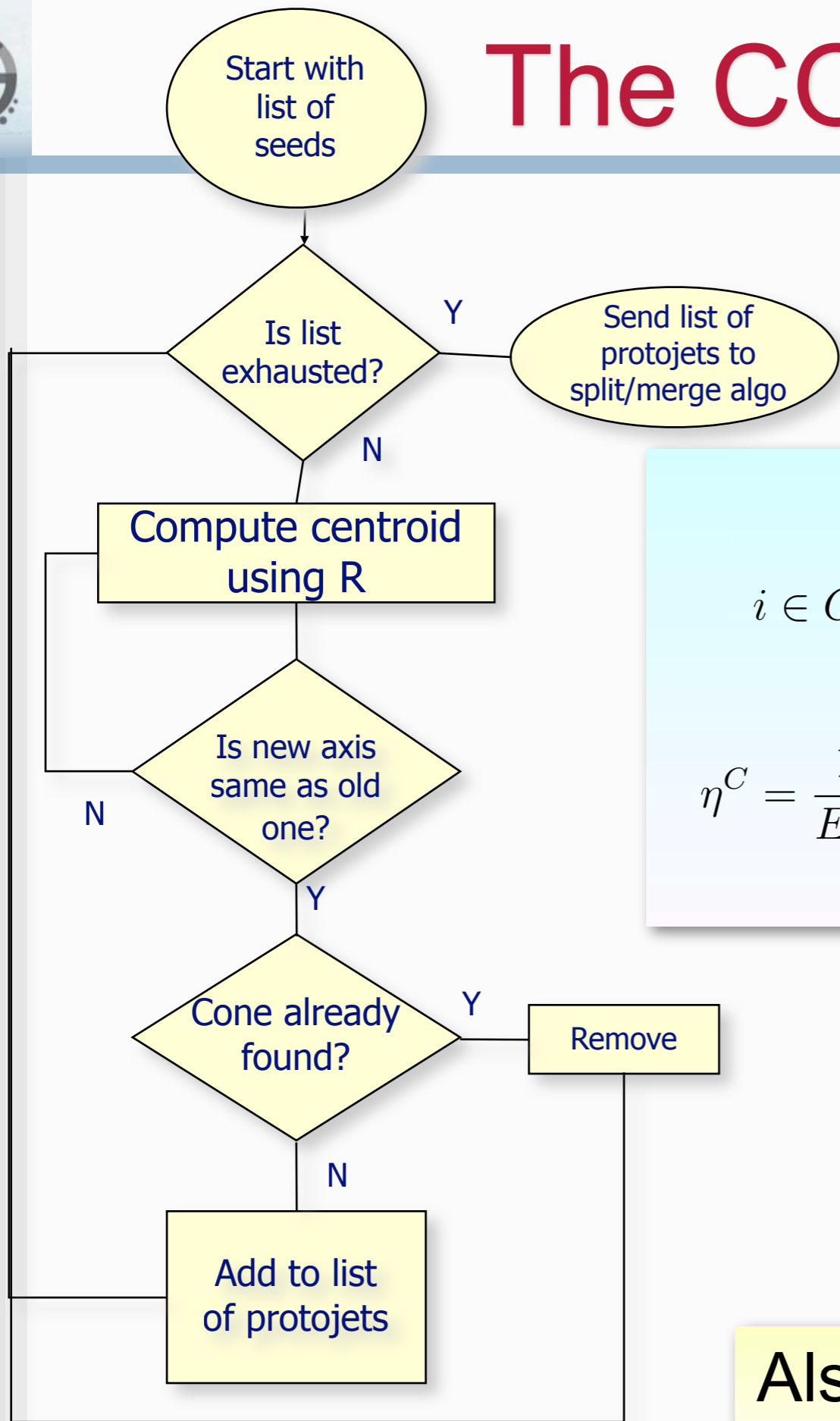


**Seeds:** for example, energy deposits with transverse energy ( $E_T = E \sin\theta$ ) > 2 GeV in a tower of the calorimeter

Centroid (one possible def) :

$$i \in C : \sqrt{(\eta^i - \eta^C)^2 + (\Phi^i - \Phi^C)^2} \leq R \quad \text{cone radius}$$

$$\eta^C = \frac{1}{E_T^C} \sum_{i \in C} E_T^i \eta^i \quad ; \quad \Phi^C = \frac{1}{E_T^C} \sum_{i \in C} E_T^i \Phi^i \quad ; \quad E_T^C = \sum_{i \in C} E_T^i$$



**Also : new seedless cone algorithm and  $k_T$  algorithm (not cone based) !**

# Further difficulties

- Pile Up** : many additional soft proton-proton interactions

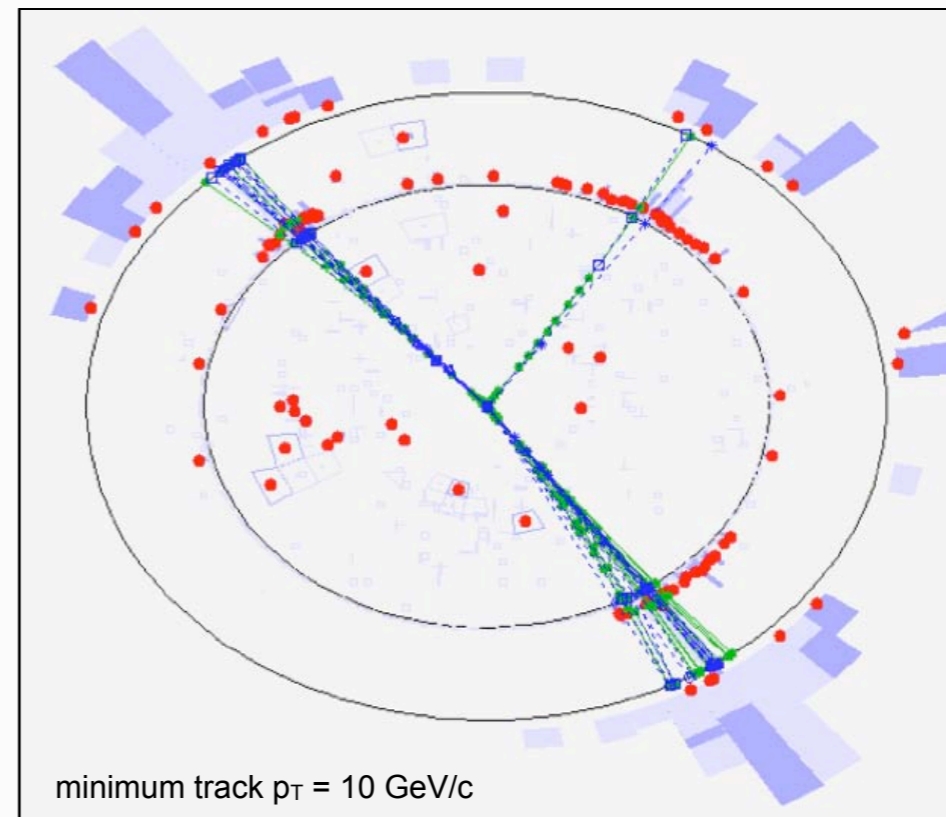
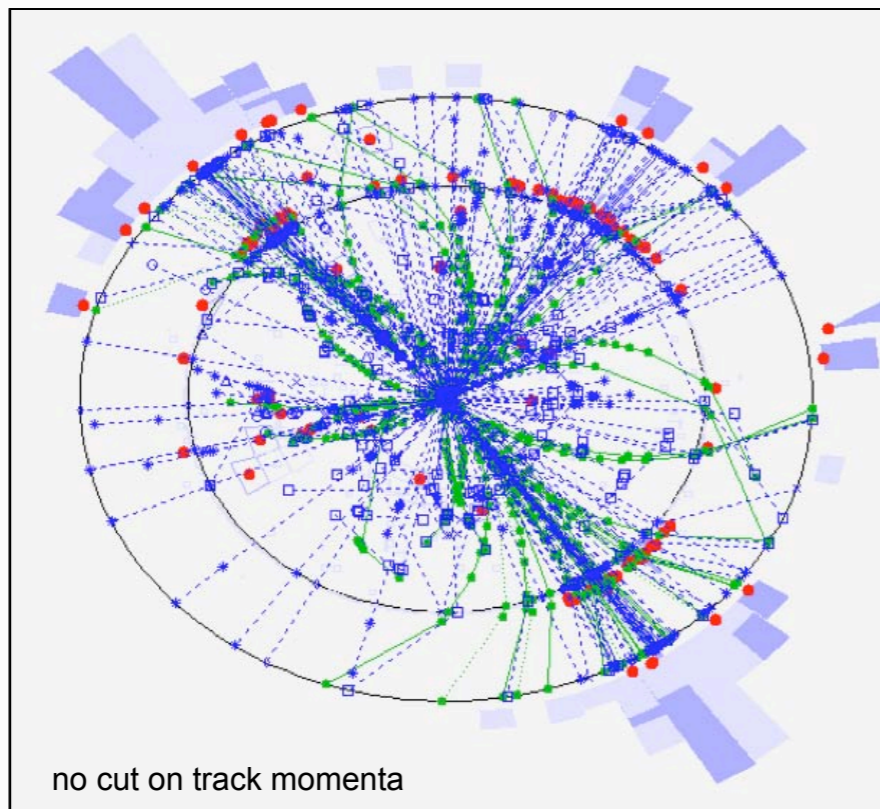
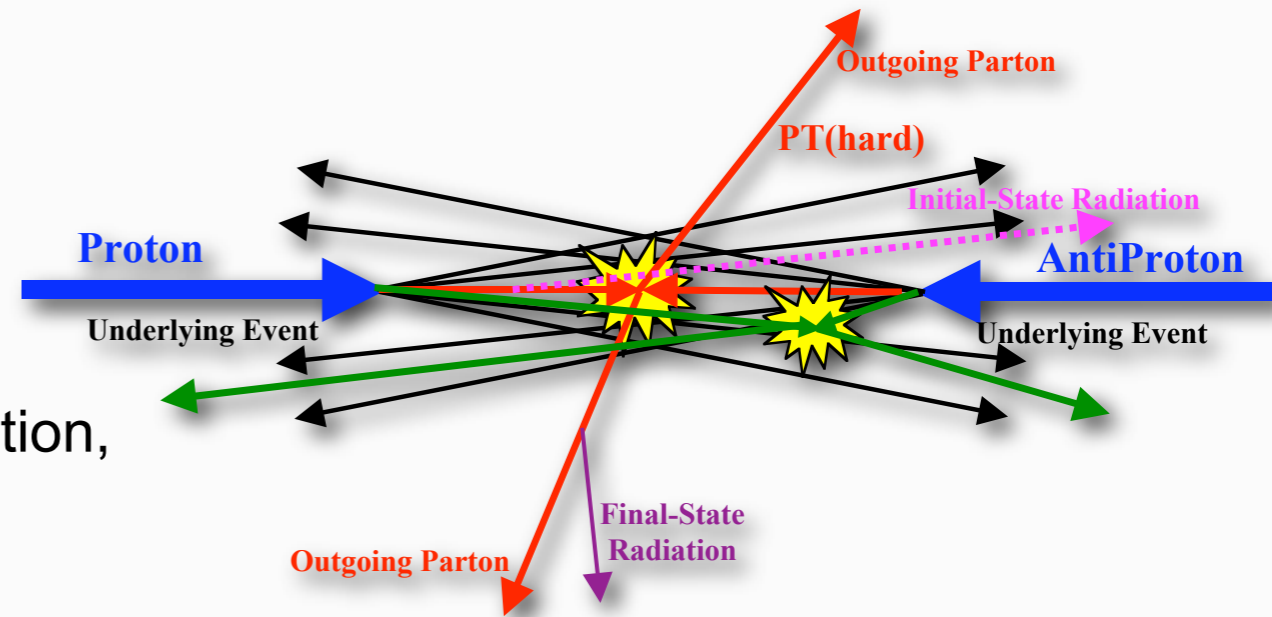
- up to 20 at highest LHC luminosity

- Underlying event**

- beam-beam remnants, initial state radiation, multiple parton interactions
  - gives additional energy in the event

- All this additional energy has nothing to do with jet energies**

- have to subtract it**

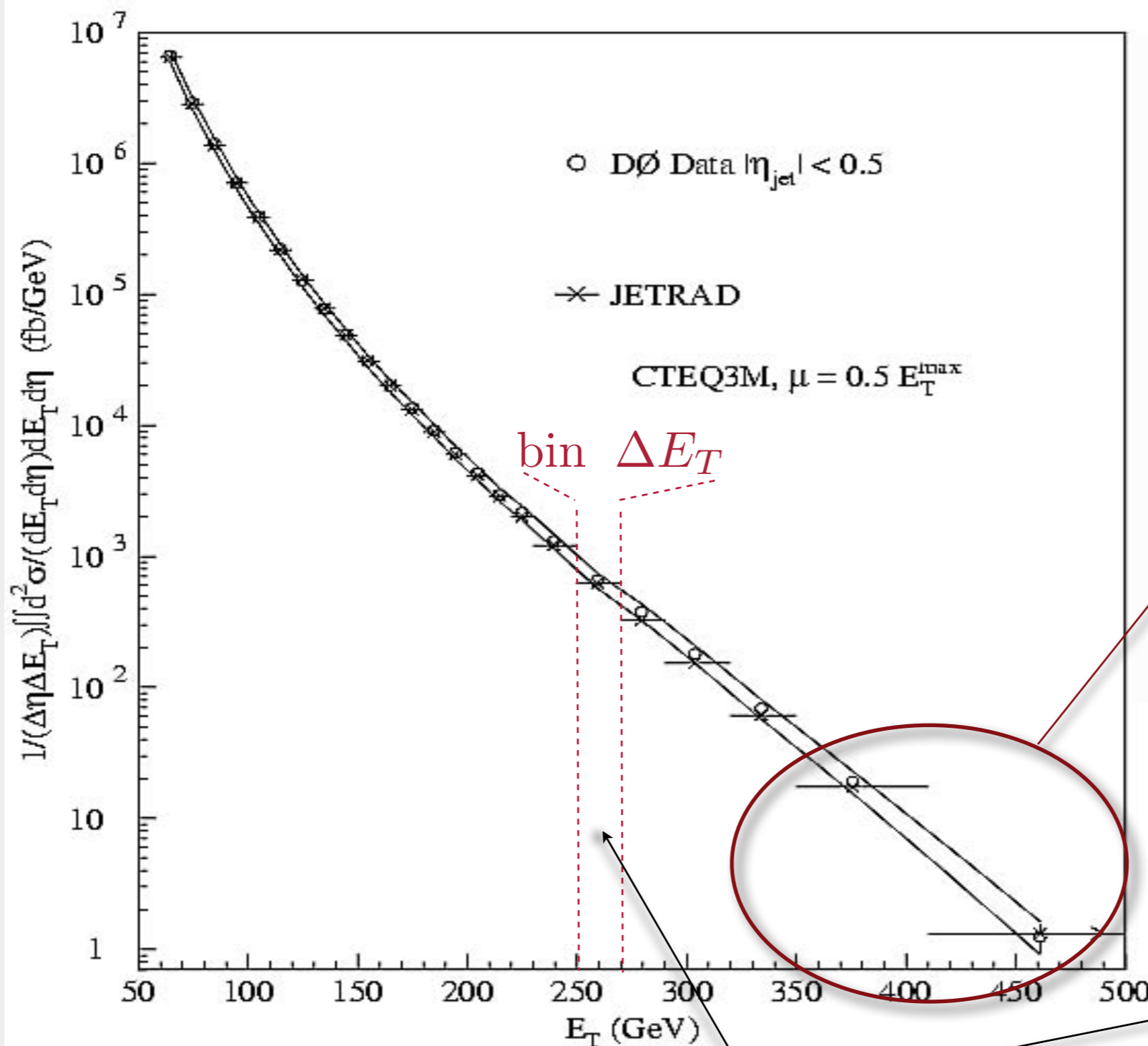




# End of the digression



# What do we have to measure?



## Goal

- measure **cross section** inclusively for **jets** that are produced with a certain **transverse energy**  $E_T$ , within a certain **rapidity range**
- Test of perturbative QCD**, over many orders of magnitude!
- Look at **very high energy tail**, **new physics** could show up there in form of excess (eg. sub-structure of quarks?)

can be calculated in pert. QCD

$$\left\langle \frac{d^2\sigma}{dE_T d\eta} \right\rangle = \frac{N}{\Delta E_T \Delta \eta \epsilon \mathcal{L}}$$

- count** number of events,  $N$ , in this bin
- for a certain range in **rapidity** (angle)  $\Delta\eta$

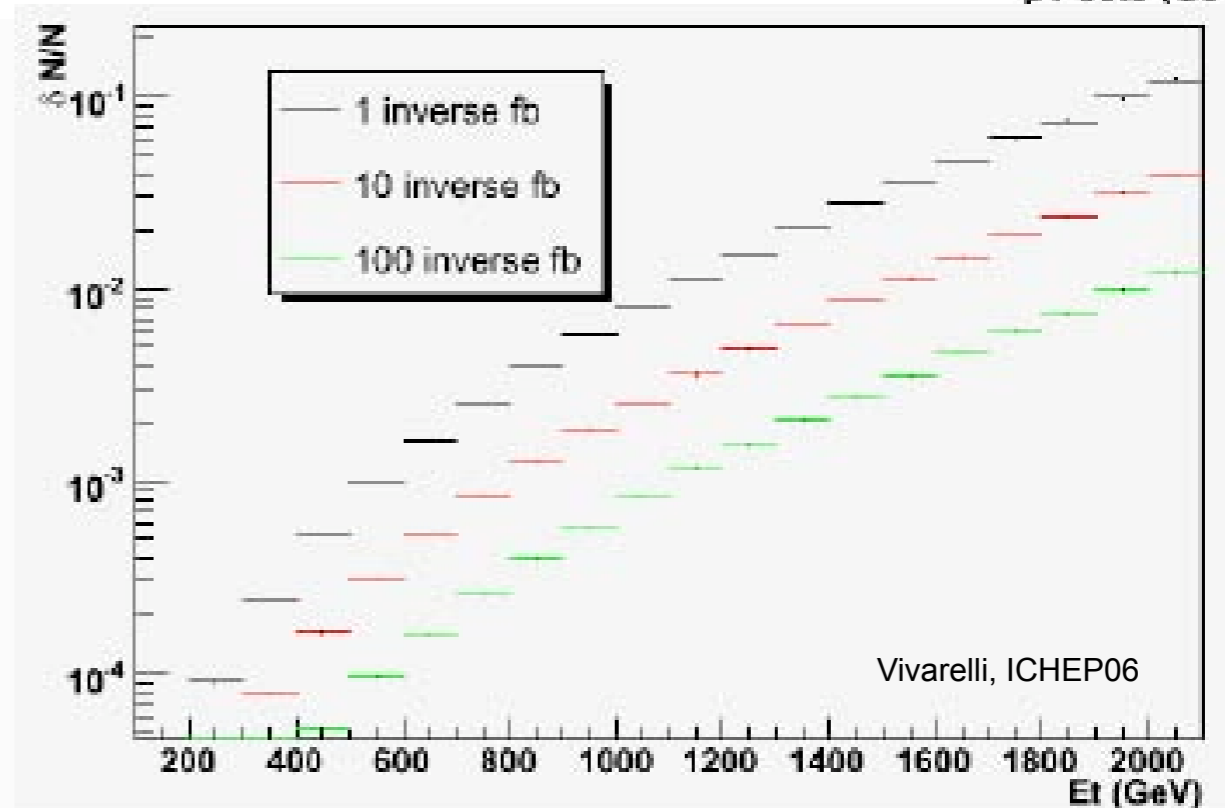
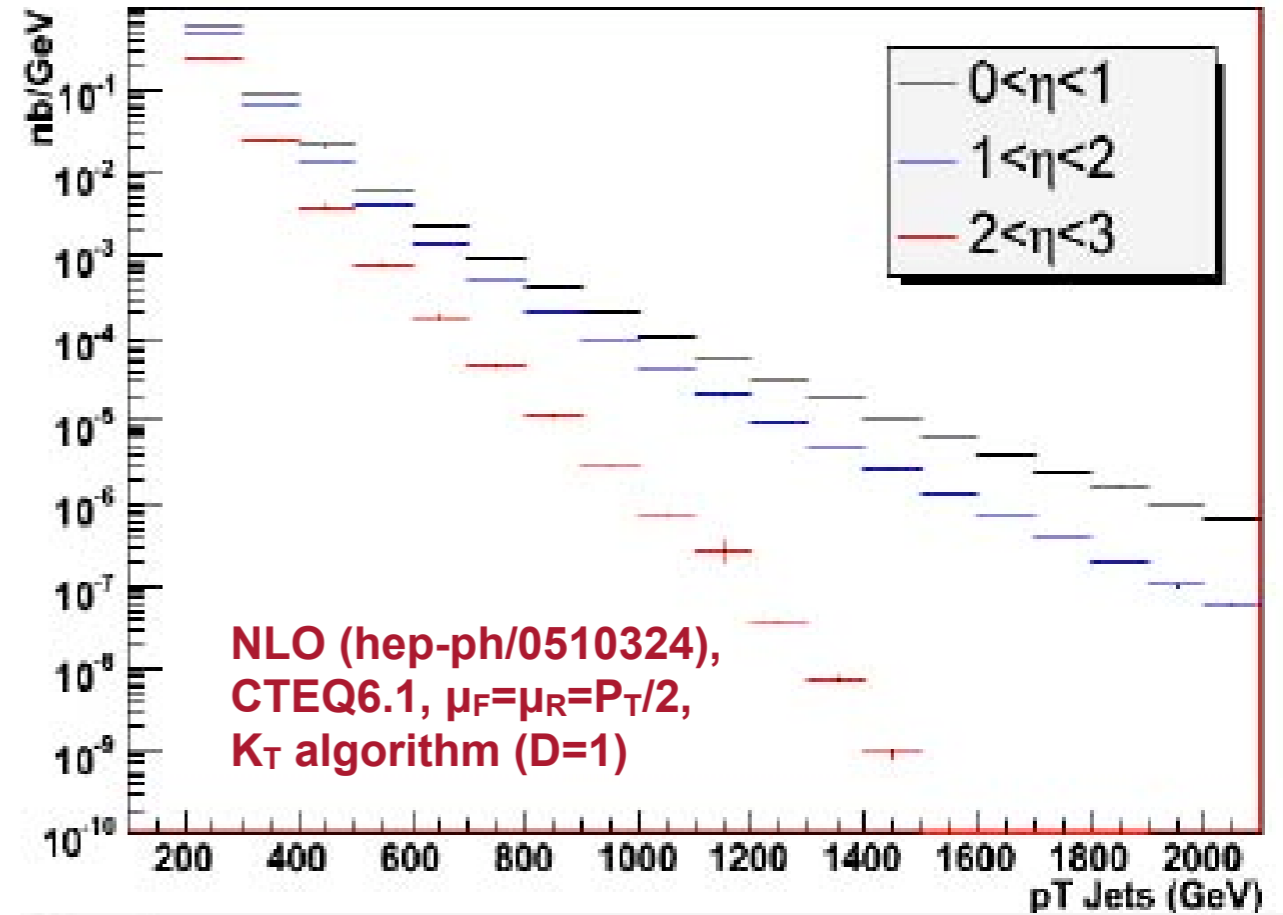
**efficiency** to reconstruct jets

integrated **luminosity**

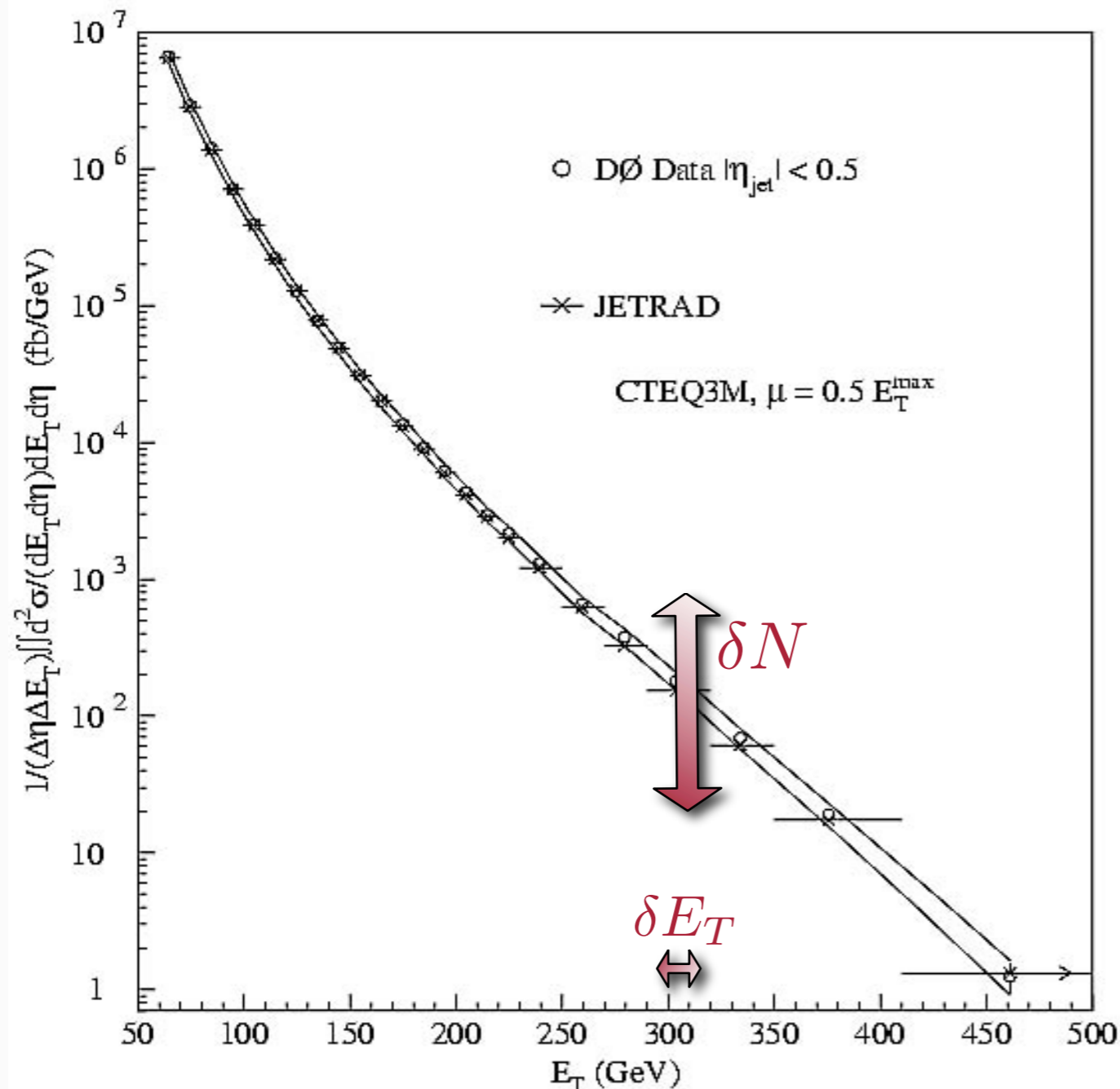


# Inclusive Jet cross section at the LHC

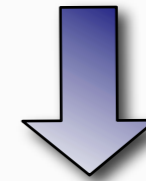
- After MB studies, jets will be the first objects seen and measured
- Enormous cross section, so statistical errors quickly negligible
  - 1% at  $p_T=1$  TeV for  $1 \text{ fb}^{-1}$  (central)
  - 10% for  $3 < \eta < 5$
- Steeply falling cross section : energy scale knowledge most relevant



- Question : how well do we know the **energy calibration**?
- Critical because of very steeply falling spectrum!



$$\frac{d^2\sigma}{dE_T d\eta} \approx \text{const} \cdot E_T^{-6}$$



**relative uncertainties**

$$\frac{\delta N}{N} \approx 6 \cdot \frac{\delta E_T}{E_T}$$

**so beware:**

eg. an uncertainty of **5%** on absolute energy scale (calibration)

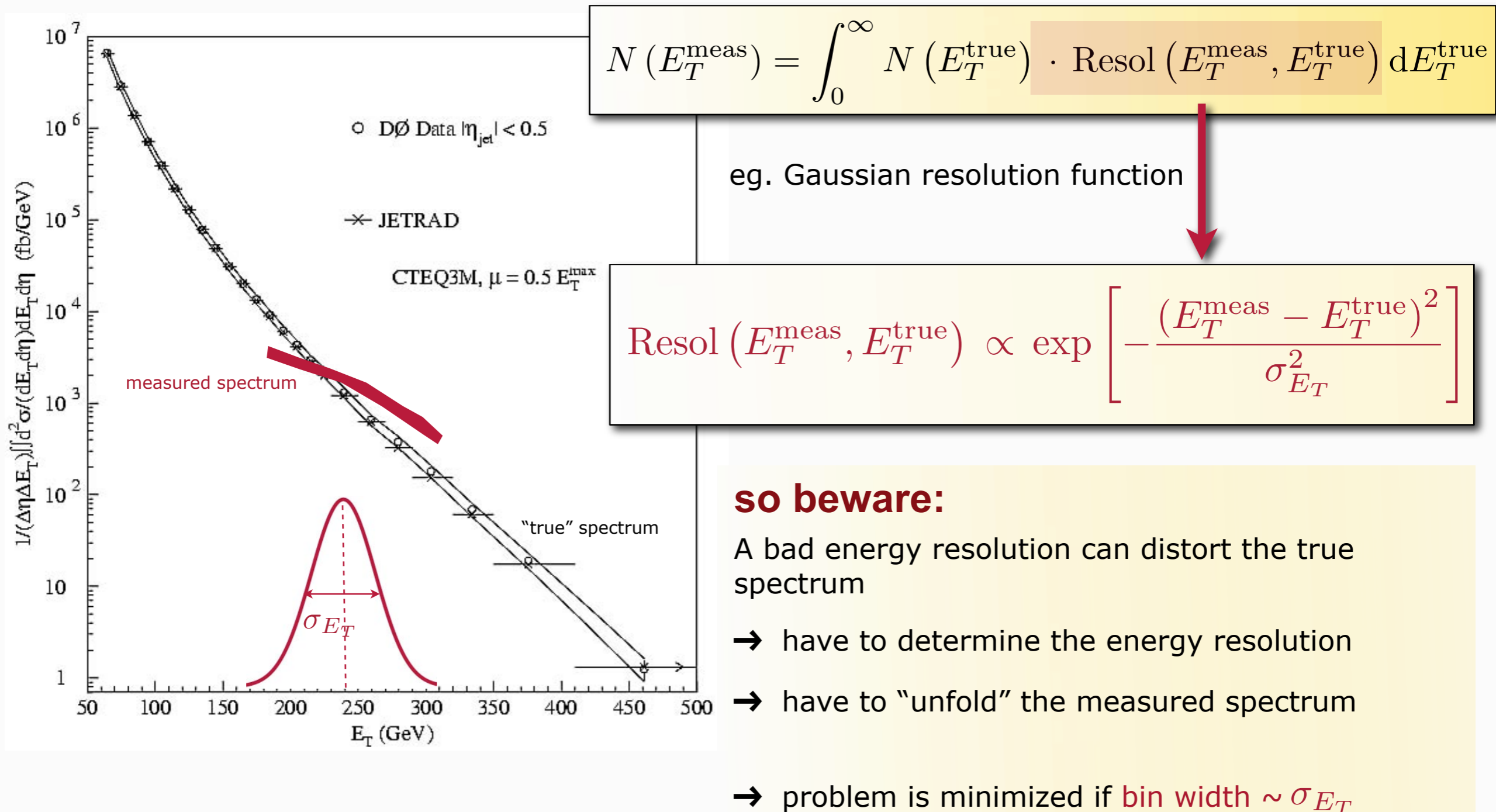
→ an uncertainty of **30%** (!) on the measured cross section

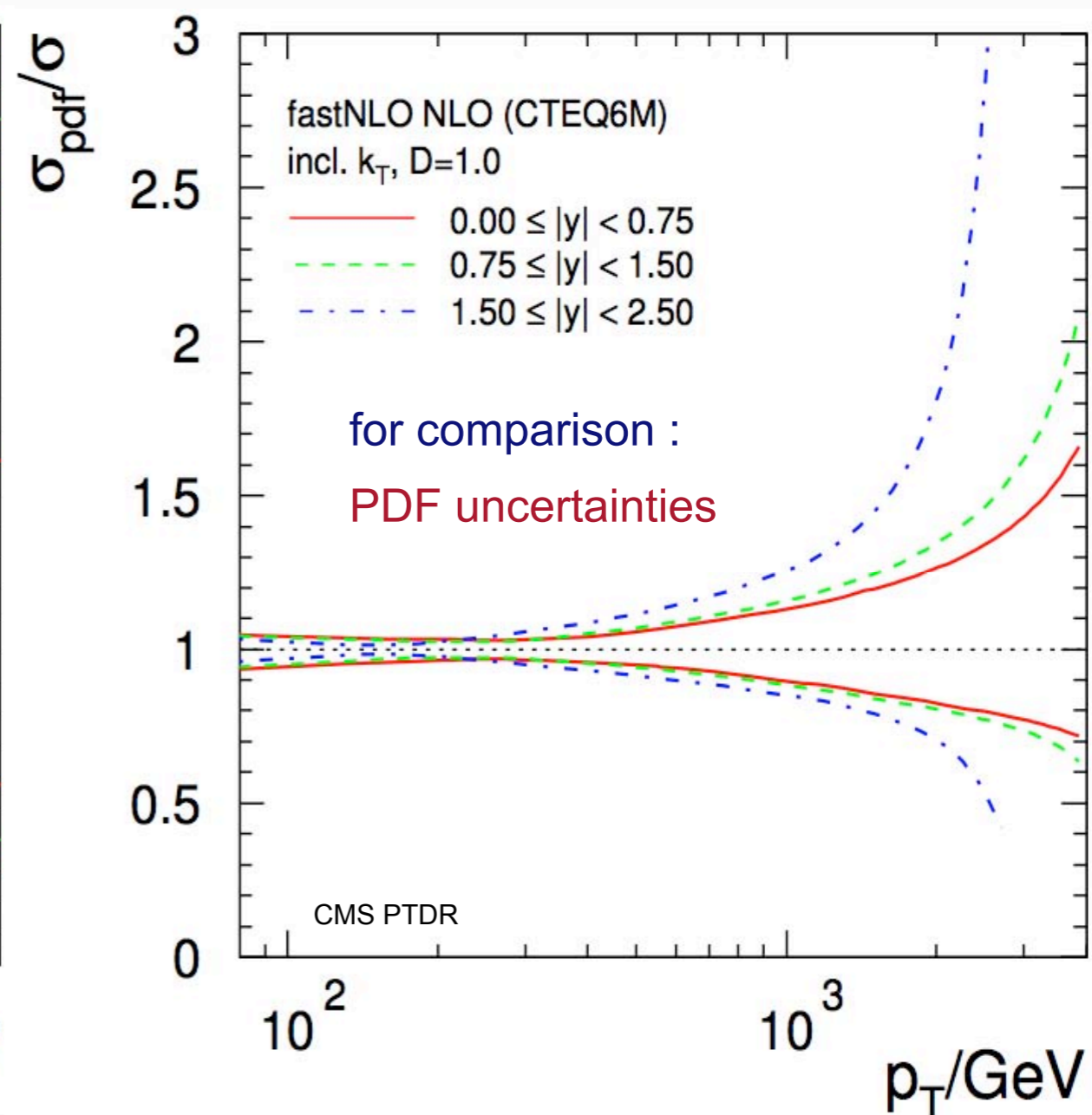
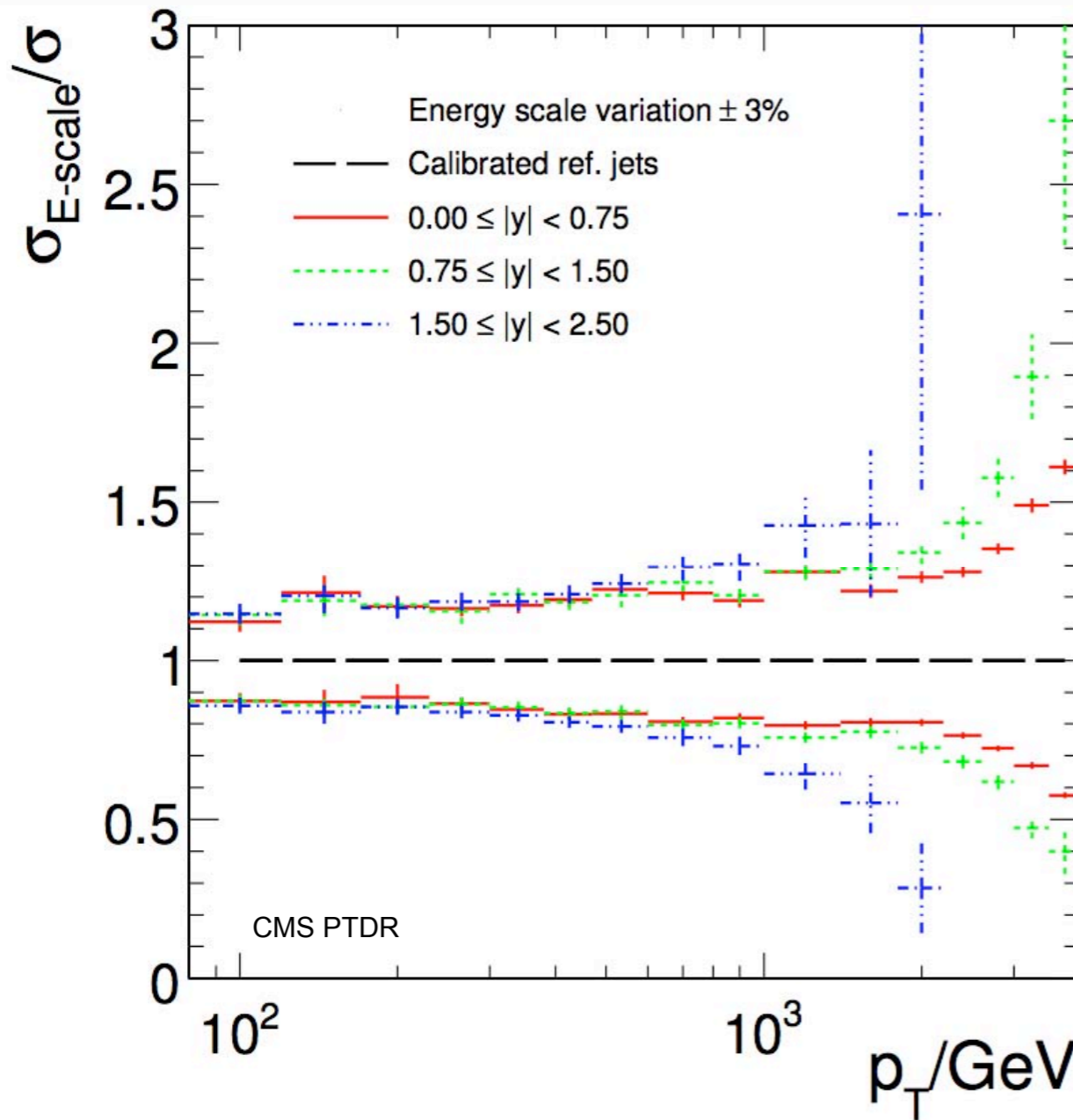




# Problem 2 : Energy resolution

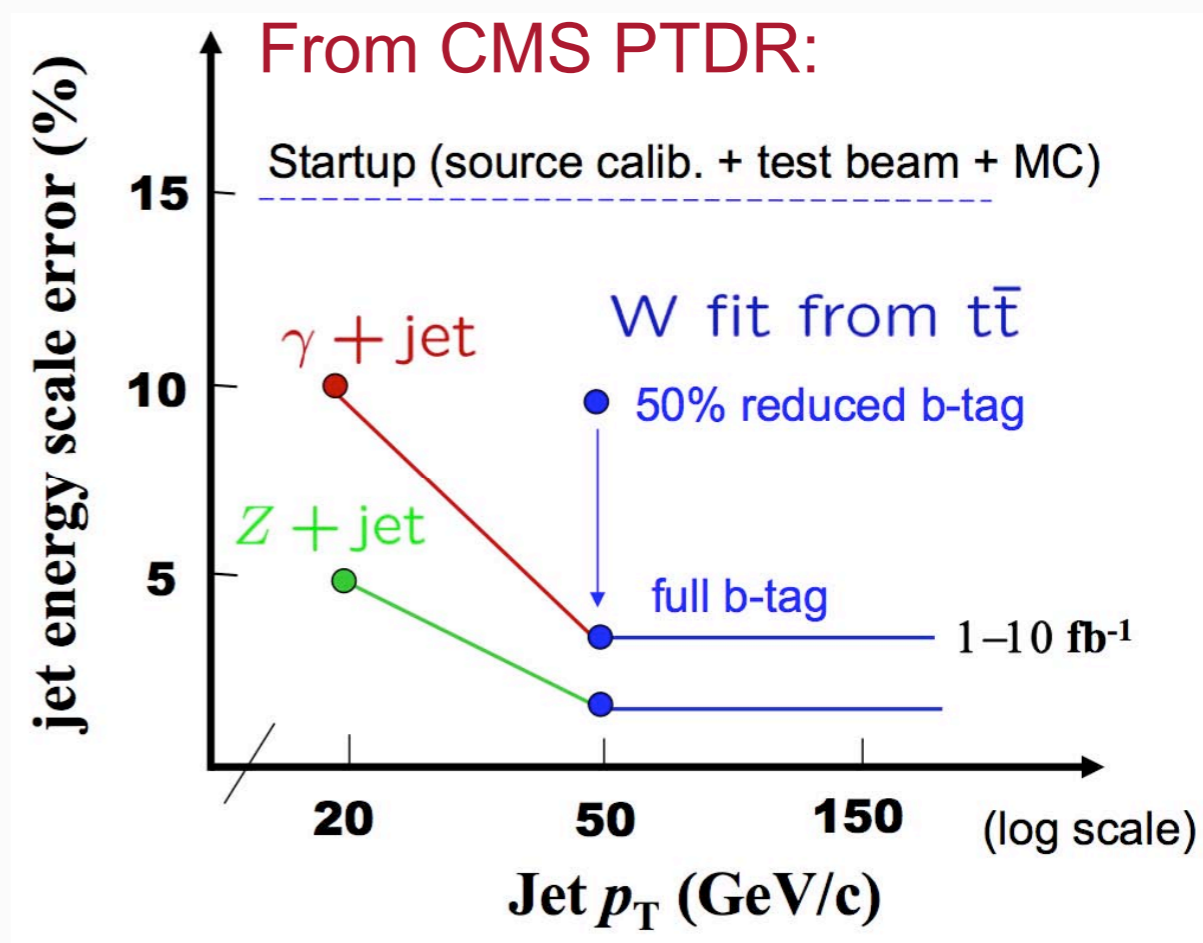
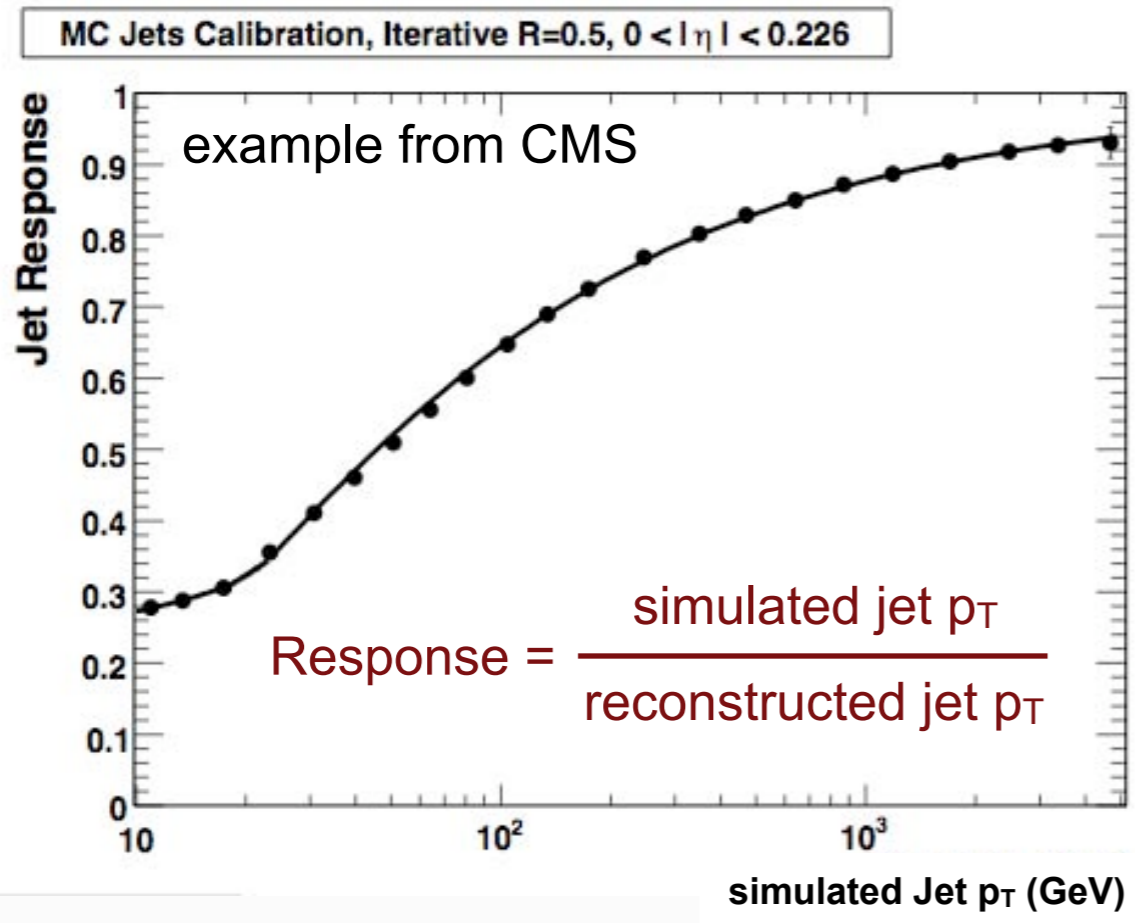
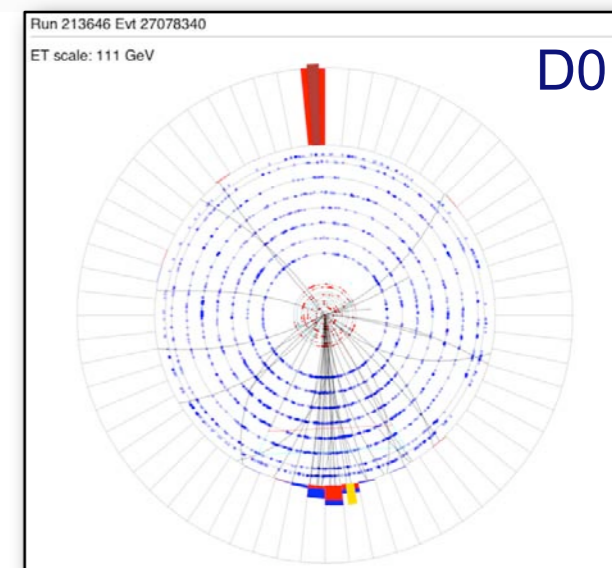
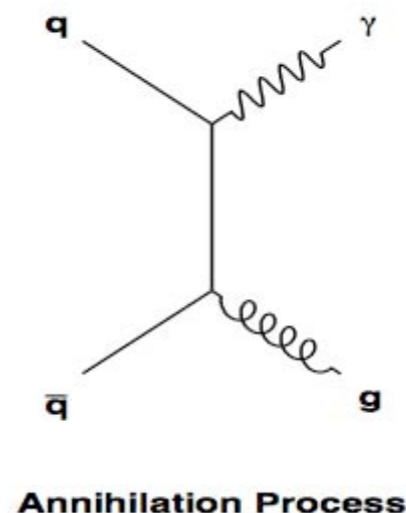
- The energy resolution can distort the spectrum
- Again : Critical because of very steeply falling spectrum!





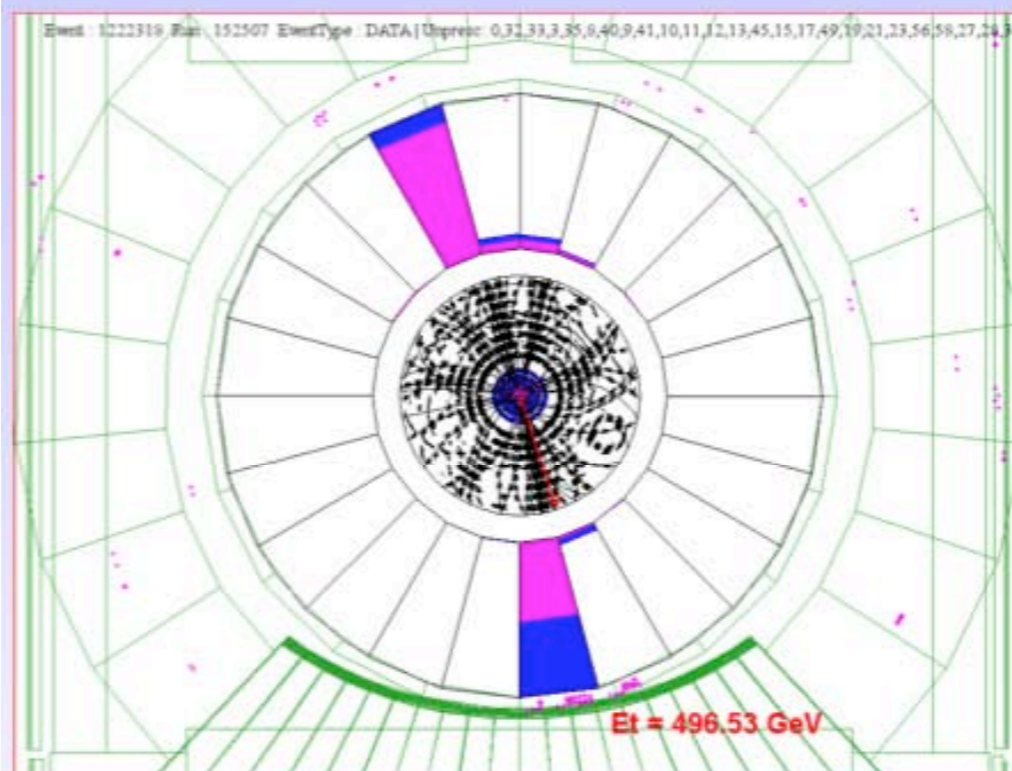
- a **5% jet energy scale** uncertainty (which is more realistic at start-up) gives a **30% error** on the cross section!
- Control in-situ with : **photon/Z+jets** and  **$W \rightarrow JJ$**  in top decays
- Other sources : **jet corrections** (det  $\rightarrow$  had  $\rightarrow$  part), **UE subtraction**

- Jet calibration using  $p_T$  balance in Jet+Photon events
  - Selection : isolated photons, no high- $p_T$  secondary jet, photon and jet well separated in transverse plane
  - Statistical error well below 1% after  $10 \text{ fb}^{-1}$

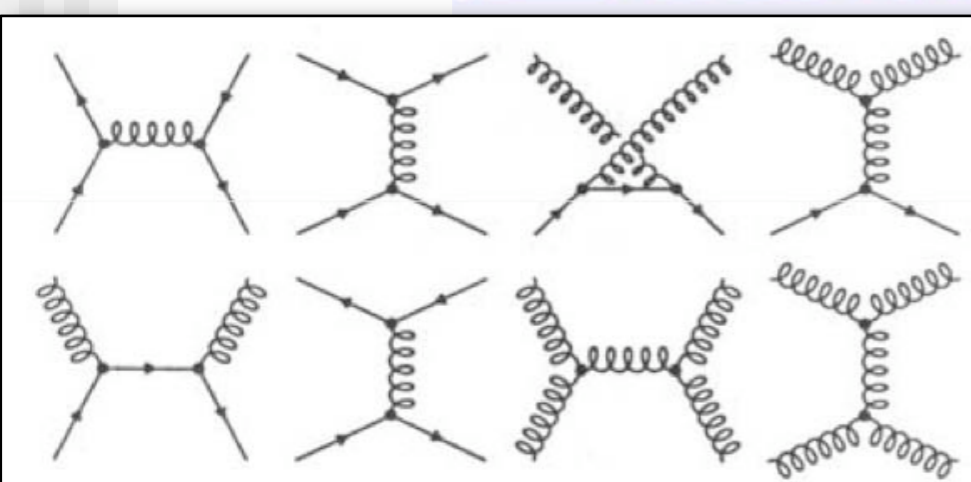
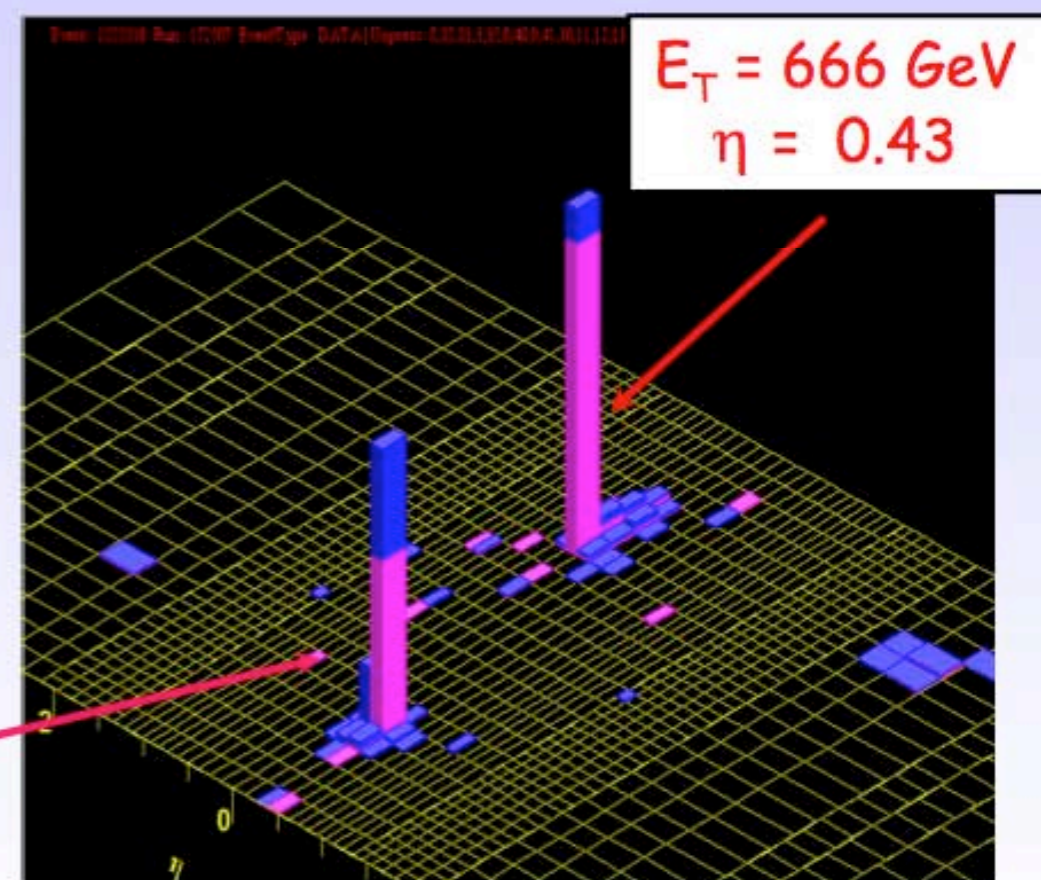


**Currently:**  
 “monolithic” MC-corrections, i.e.  
 one-step correction from calorimeter-level to particle level  
 (inverse of response function)

## A two jet event at the Tevatron (CDF)



Dijet Mass = 1364 GeV/c<sup>2</sup>



CDF ( $\phi$ -r view)

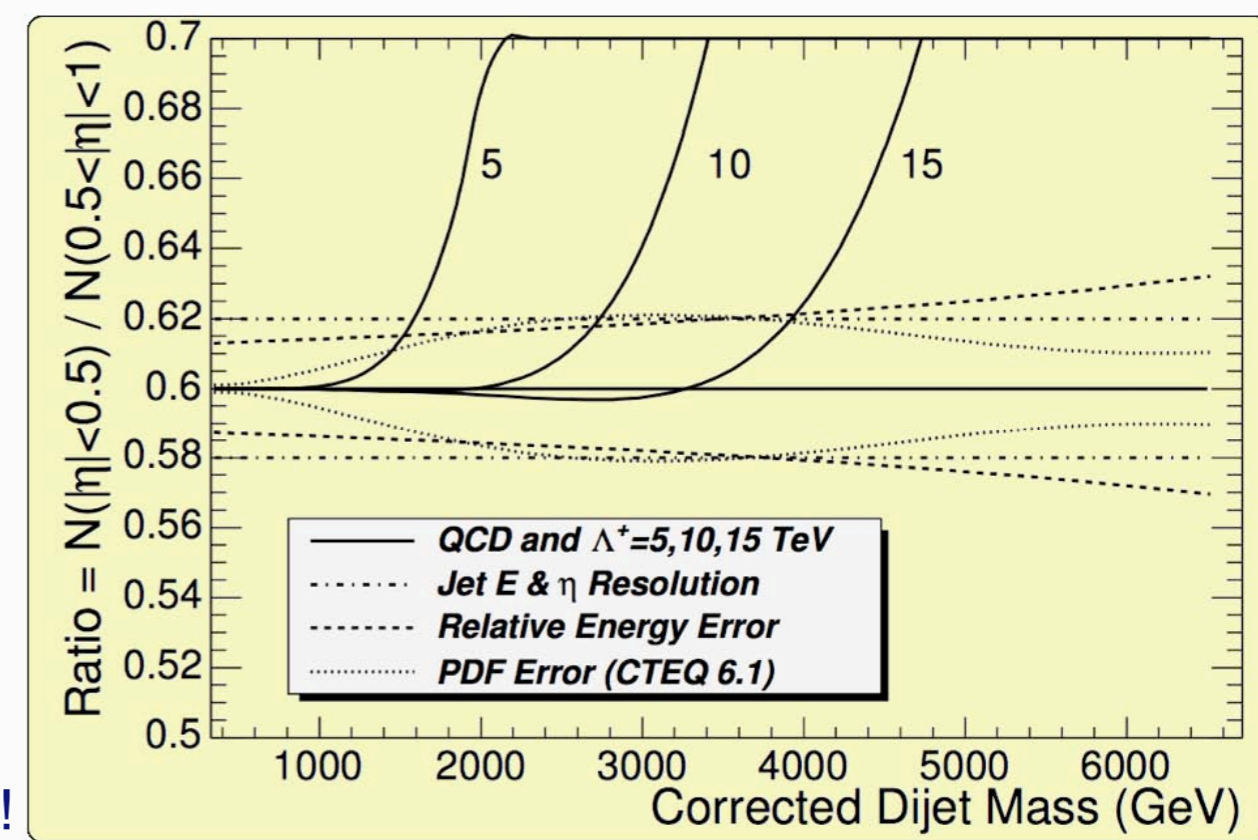
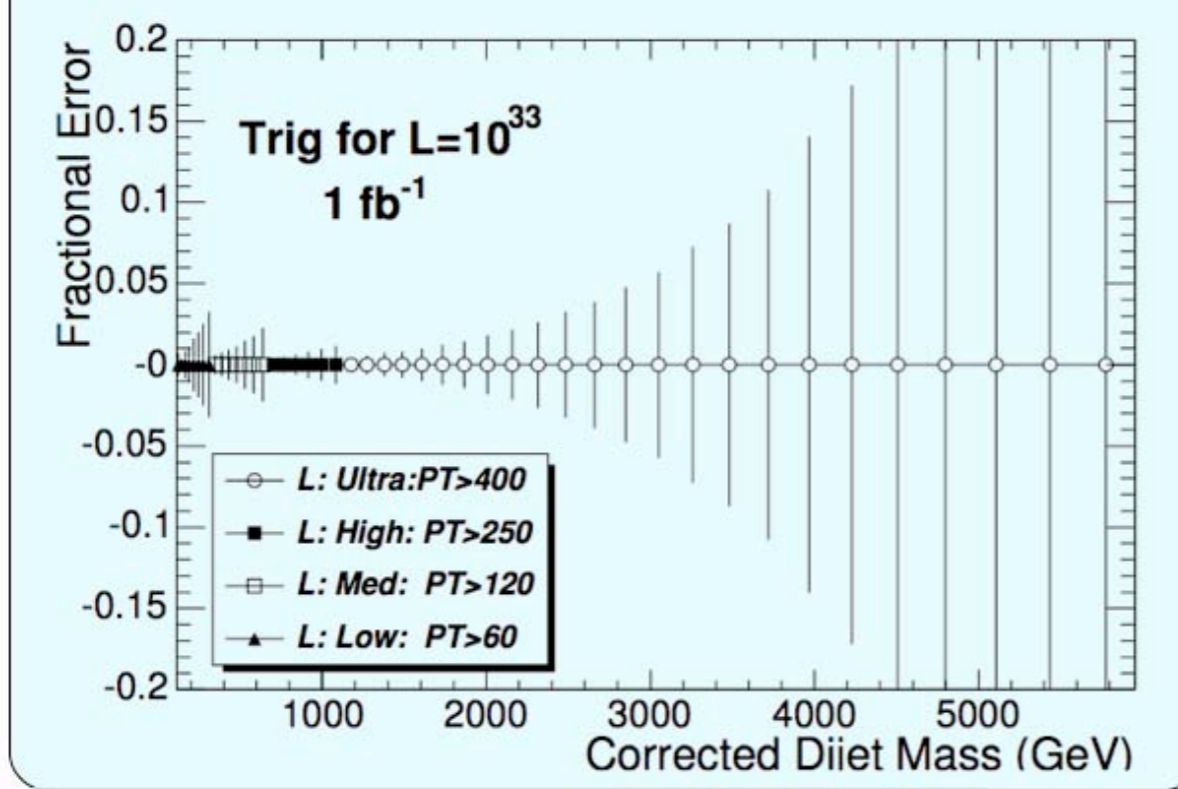
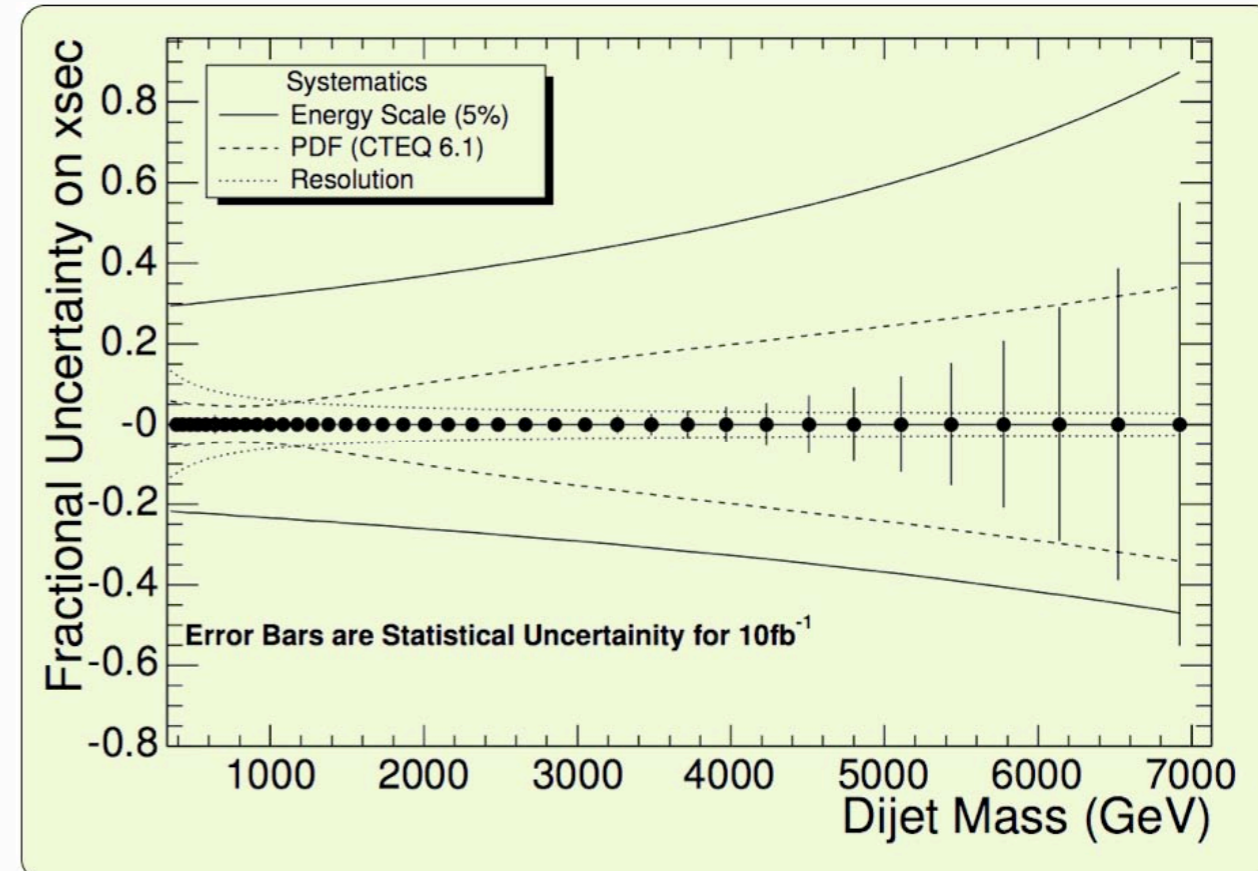
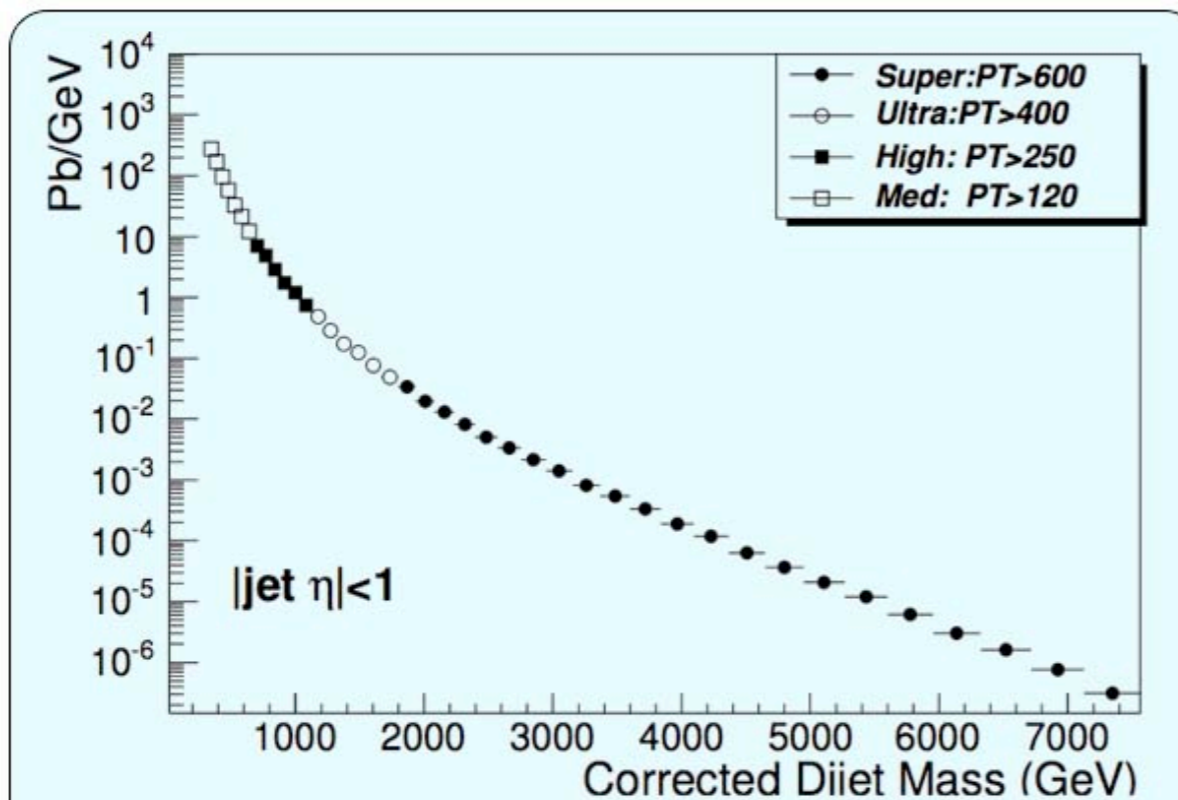
ET = 633 GeV  
eta = -0.19

K. Jakobs, CSS07

**Goal** : Measure cross section as function of invariant mass of the two jets.  
Test QCD predictions and look for resonances at high invariant mass.



# Di-Jets (CMS PDTR)



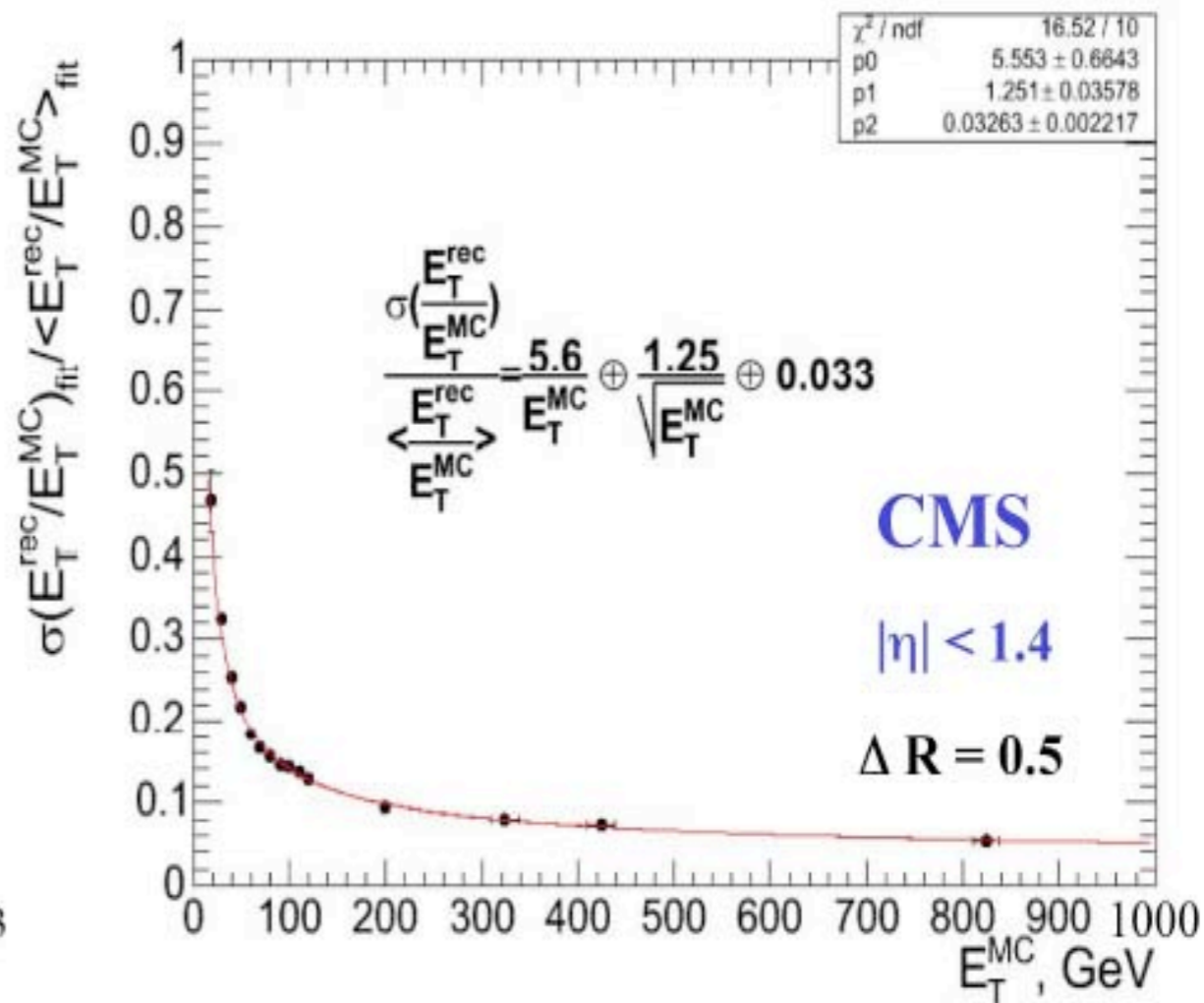
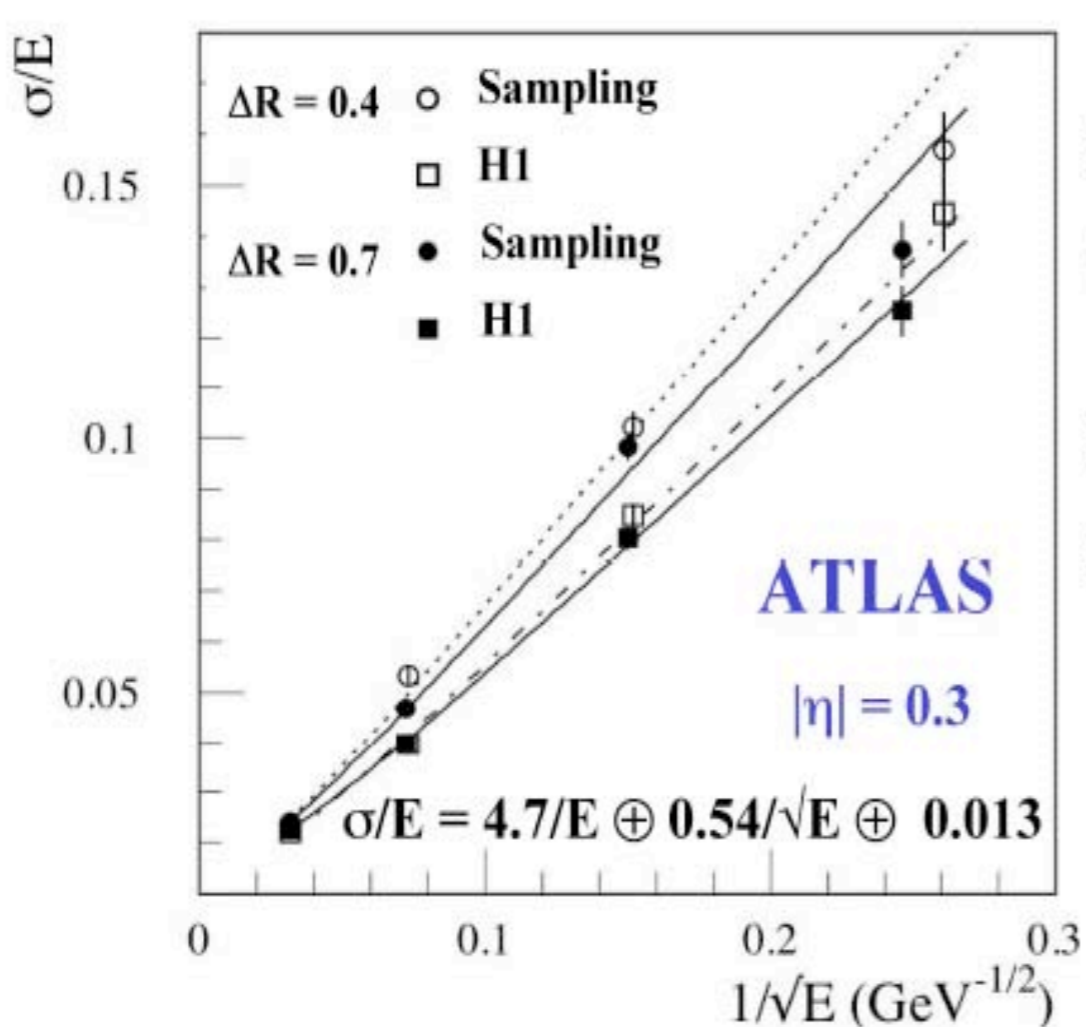
Note : Ratio of cross-section central / forward !



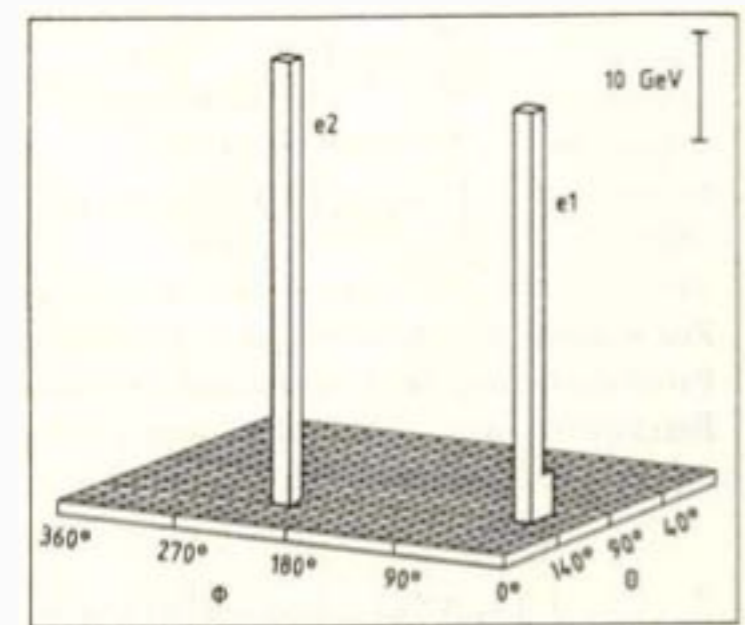
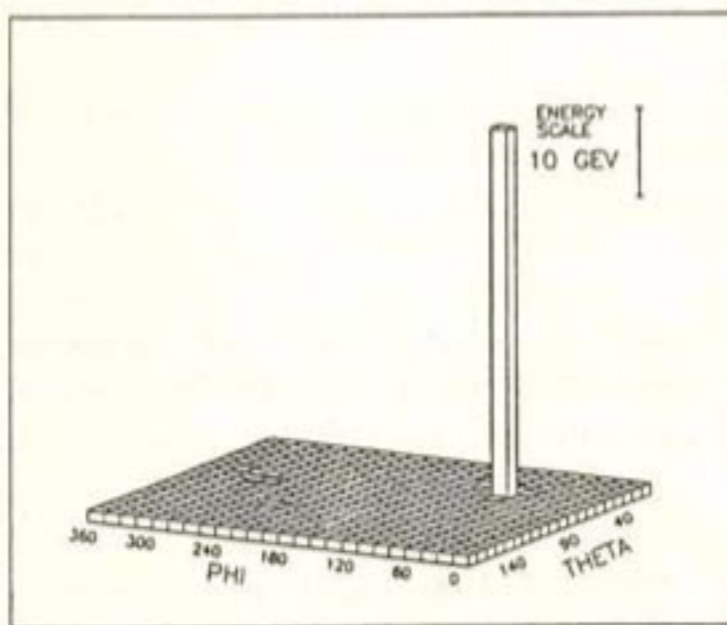
# Comparison ATLAS - CMS

Biggest difference in performance perhaps for hadronic calo

**Jets at 1000 GeV** ATLAS ~ 2% energy resolution CMS ~ 5% energy resolution, but expect sizeable improvement using tracks (especially at lower E)

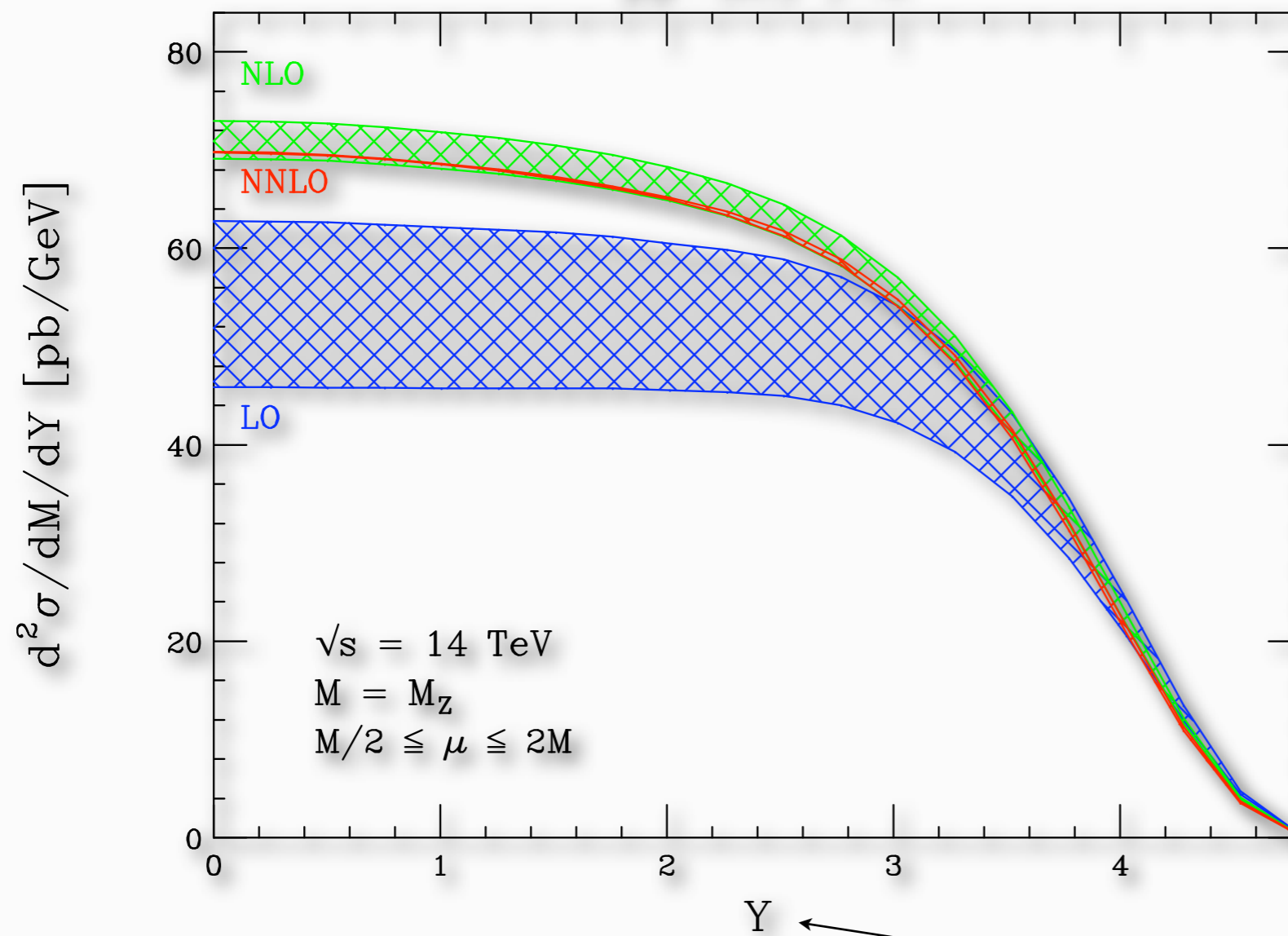
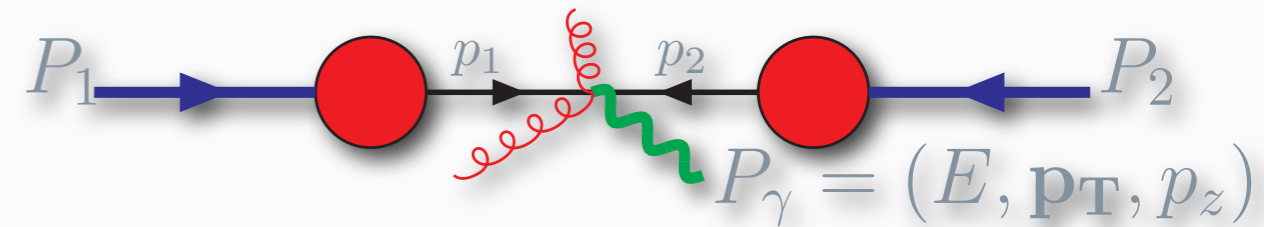


# W and Z production



... one of the first W and Z's in UA1/2

- Probably best known cross section at LHC, NNLO, differentially
- a well suited normalization process



small NNLO scale uncertainty:

LO :	25% - 30%
NLO :	6%
NNLO :	0.1 % (Y=0) -
	1% (Y<3) -
	3% (Y ~ 4)

shape stabilizes at NNLO

Anastasiou, Dixon, Petriello, Melnikov : differential in W/Z rapidity  
 Petriello, Melnikov : fully differential in lepton momenta

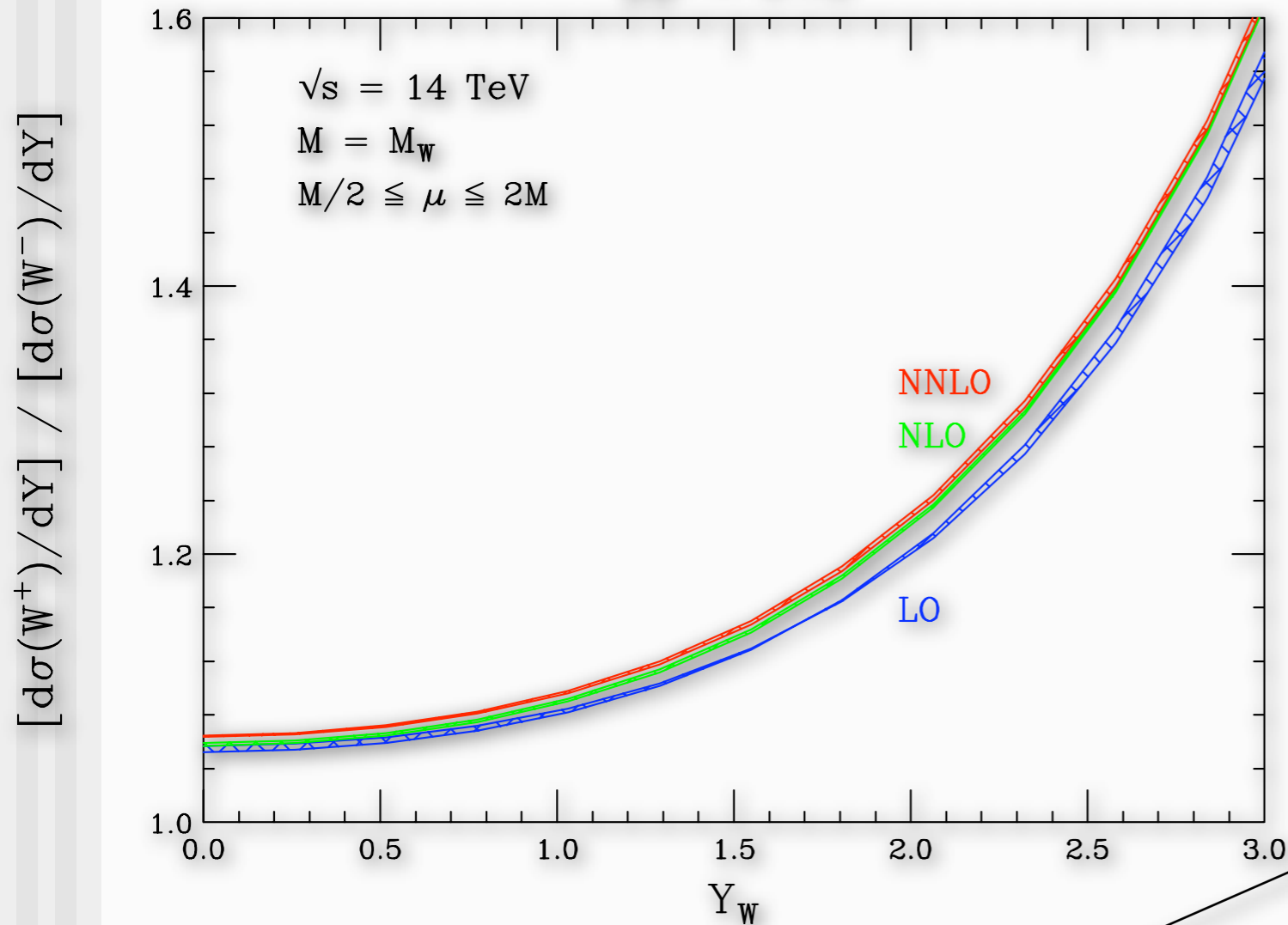




# Predictions

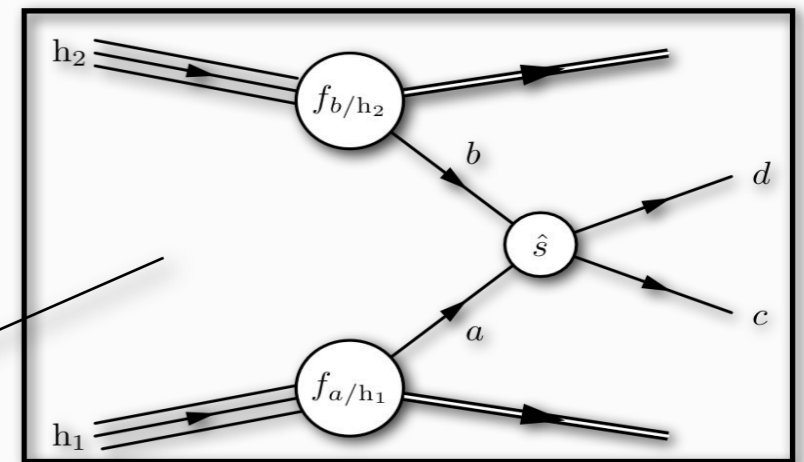
similarly for  $W^+ / W^-$  : very small  
NNLO scale uncertainty: **0.5 - 0.7 %**

$pp \rightarrow W+X$



However :  
 dominant  
 uncertainty for  
**absolute W and Z**  
 cross section:  
**PDFs : 4 - 5 %**

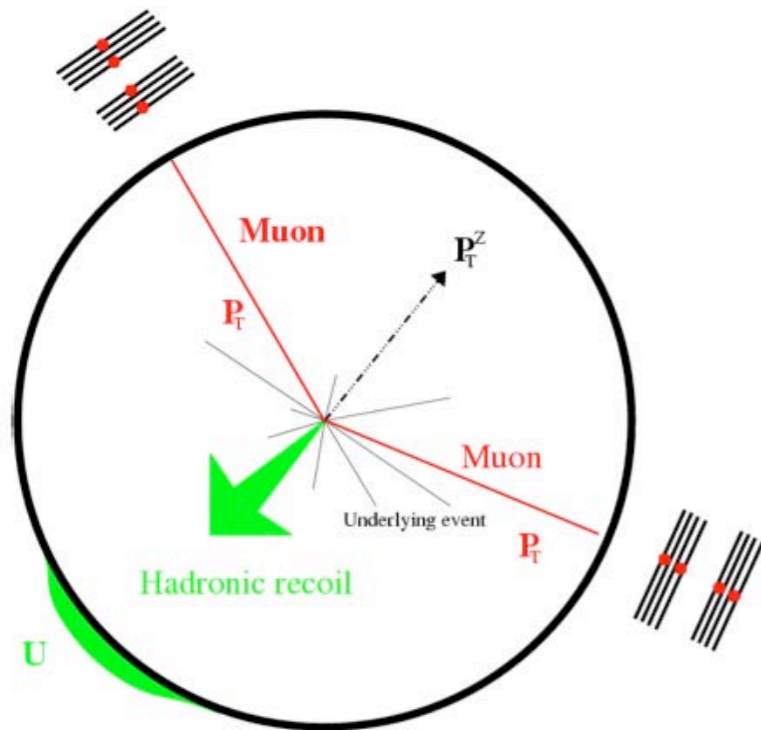
remember:



$$d\sigma(h_1 h_2 \rightarrow cd) = \int_0^1 dx_1 dx_2 \sum_{a,b} f_{a/h_1}(x_1, \mu_F^2) f_{b/h_2}(x_2, \mu_F^2) d\hat{\sigma}^{(ab \rightarrow cd)}(Q^2, \mu_F^2)$$

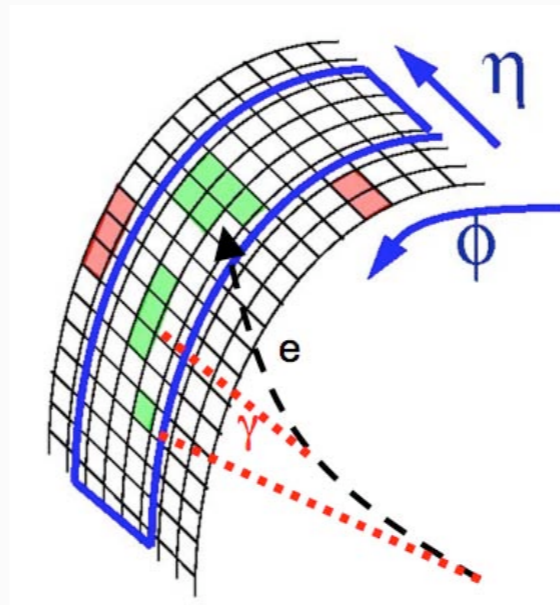
so : if we measure precisely this, can we constrain this?

known precisely up to NNLO



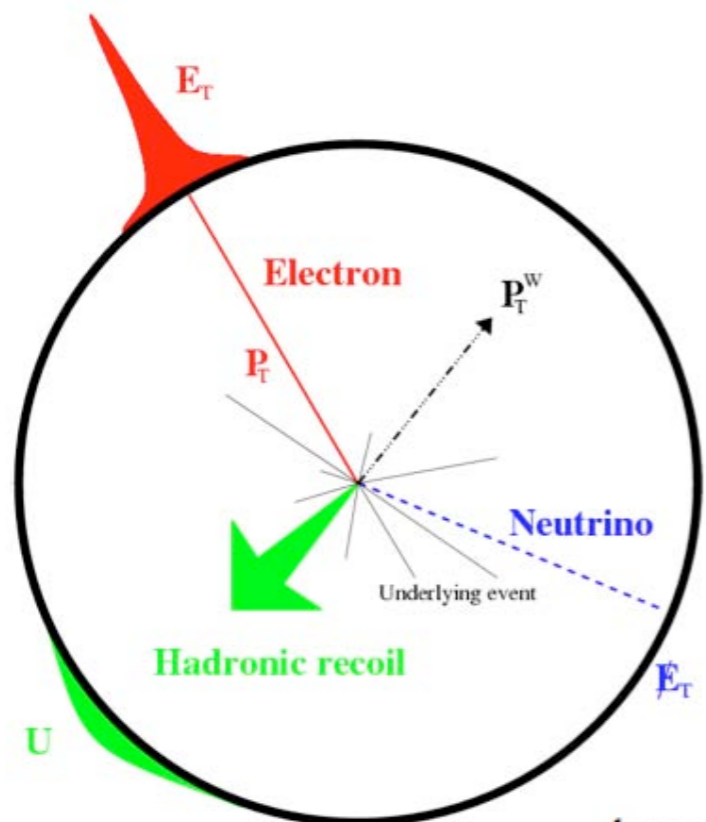
**Z: pair of charged leptons**

- high- $p_T$
- isolated
- opposite charge
- $\sim 70 < m_{ll} < \sim 110$  GeV



**Example: electron reconstruction**

- isolated cluster in EM calorimeter
- $p_T > 20$  GeV
- shower shape consistent with expectation from electrons
- matching charged track

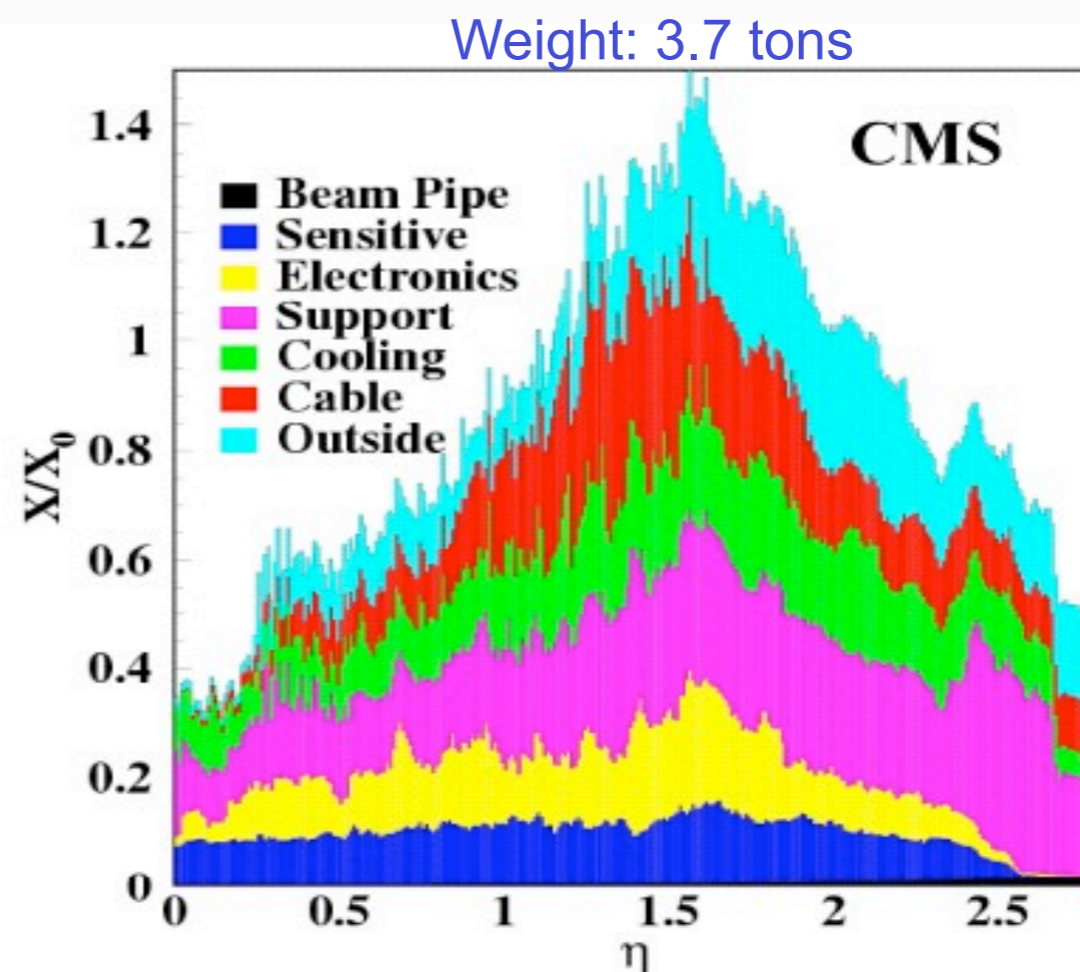
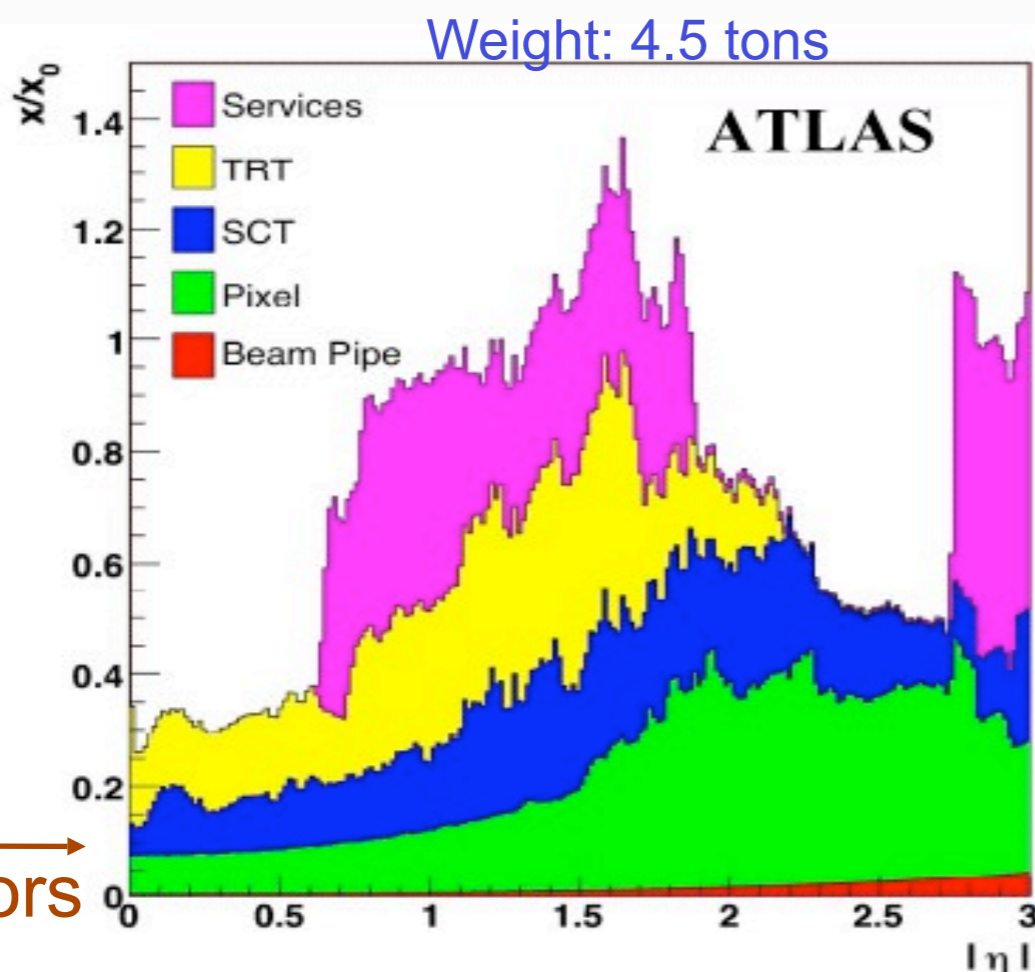


**W: single charged leptons**

- high- $p_T$
- isolated
- $E_{T,miss}$  (from neutrino)

transverse mass:  $M_W^T = \sqrt{2 \cdot P_T^l \cdot P_T^\nu \cdot (1 - \cos \Delta\phi^{l,\nu})}$

## Amount of material in ATLAS and CMS inner trackers



LEP  
→  
detectors

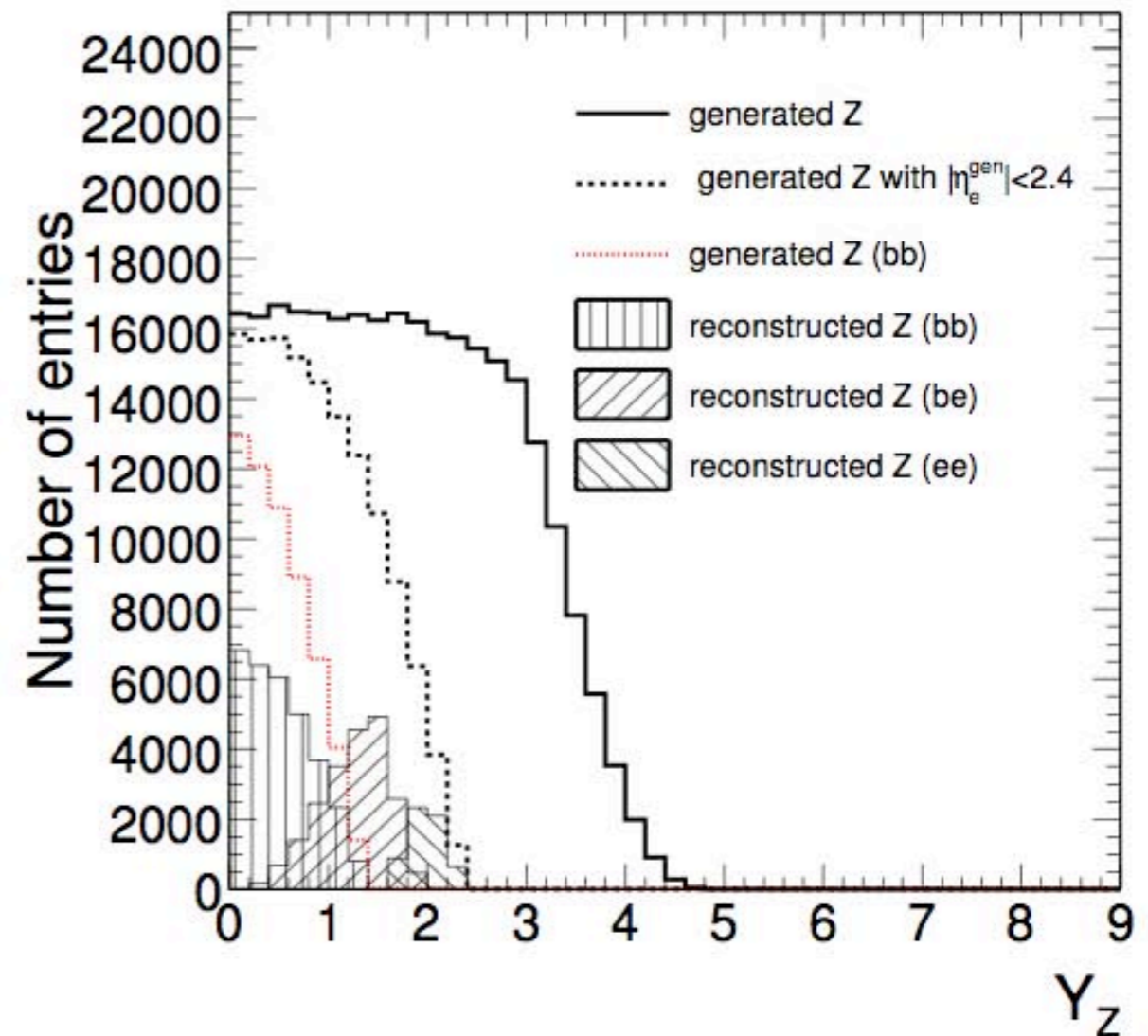
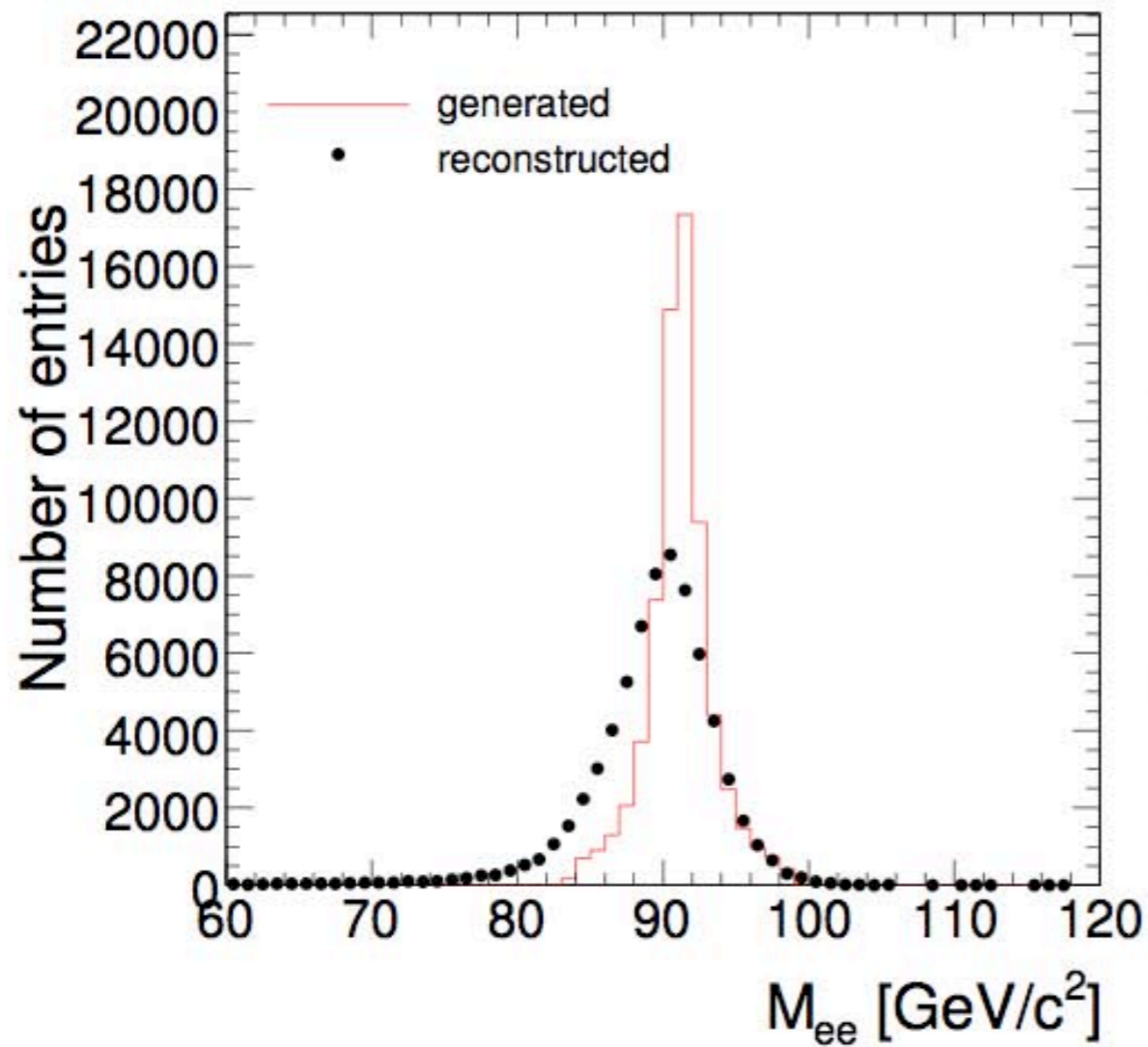
- Active sensors and mechanics account each only for  $\sim 10\%$  of material budget
- Need to bring 70 kW power into tracker and to remove similar amount of heat
- Very distributed set of heat sources and power-hungry electronics inside volume: this has led to complex layout of services, most of which not at all understood at the time of the TDRs

- Material increased by  $\sim$  factor 2 from 1994 (approval) to now (end constr.)
- Electrons lose between 25% and 70% of their energy before reaching EM calo
- Between 20% and 65% of photons convert into  $e^+e^-$  pair before EM calo
- Need to know material to  $\sim 1\%$   $X_0$  for precision measurement of  $m_W$  ( $< 10$  MeV)!

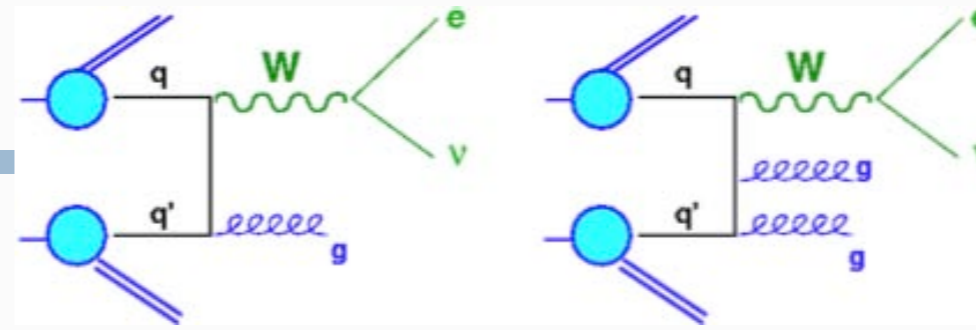
## Example of selection from CMS (simulation studies)

- Z : 2 isolated leptons,  $p_T > 20$  GeV,  $|\eta| < 2.5$ , W : 1 isolated lepton + MET
- Studied : electrons, muons.
- Difficult issue : MET (=neutrino reconstruction)

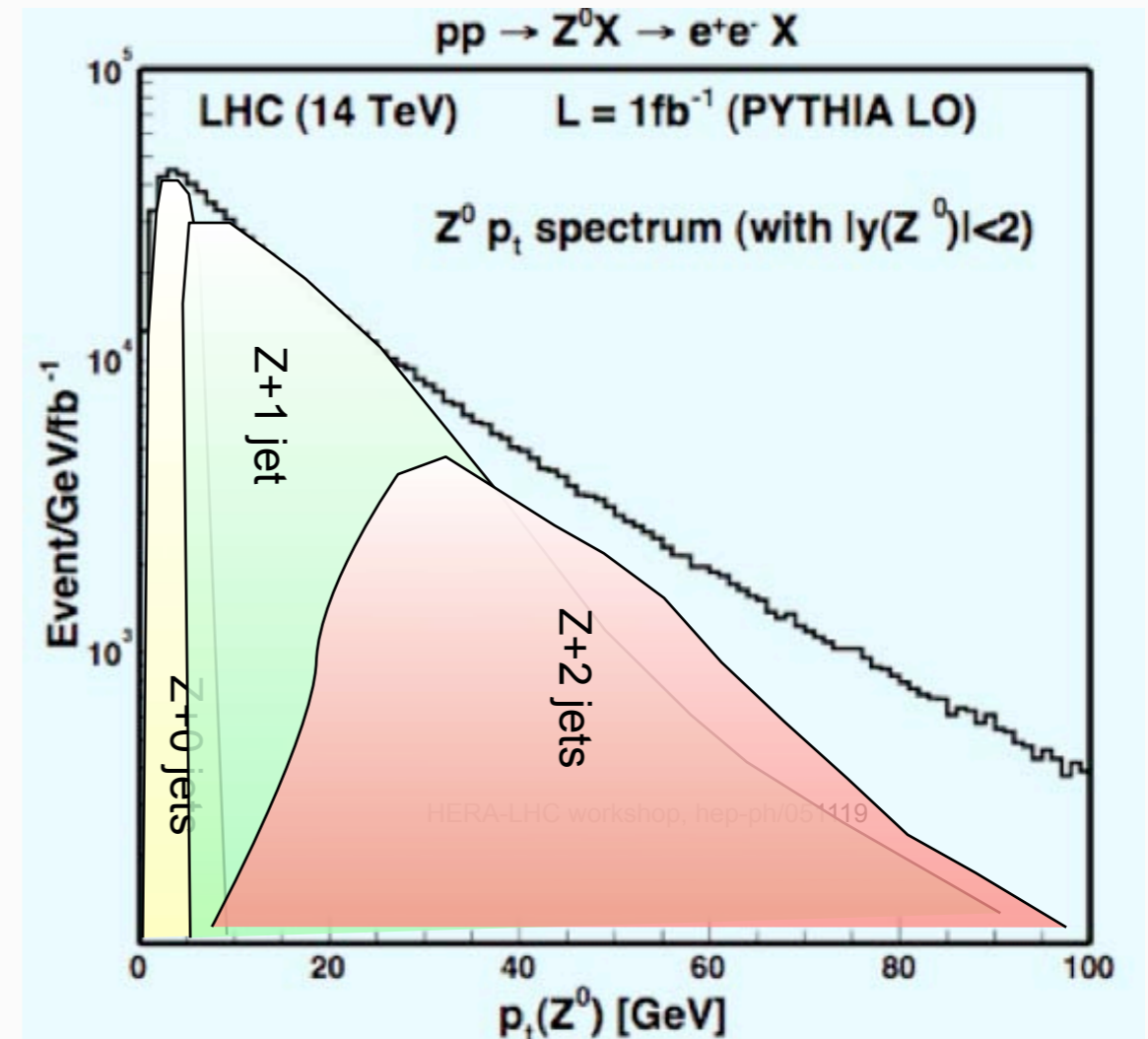
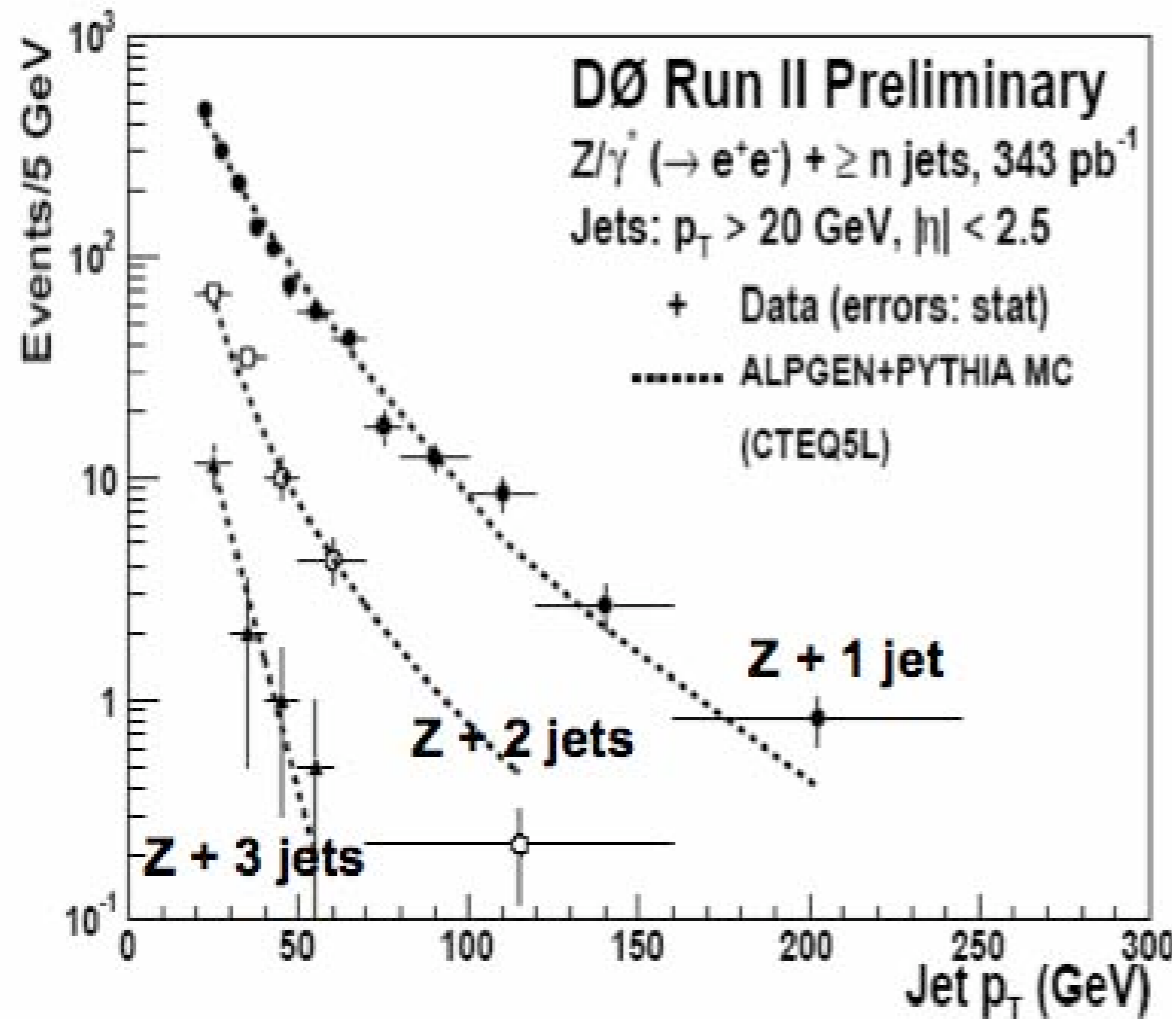
CMS PTDR, G.D., Dittmar, Ehlers, Holzner



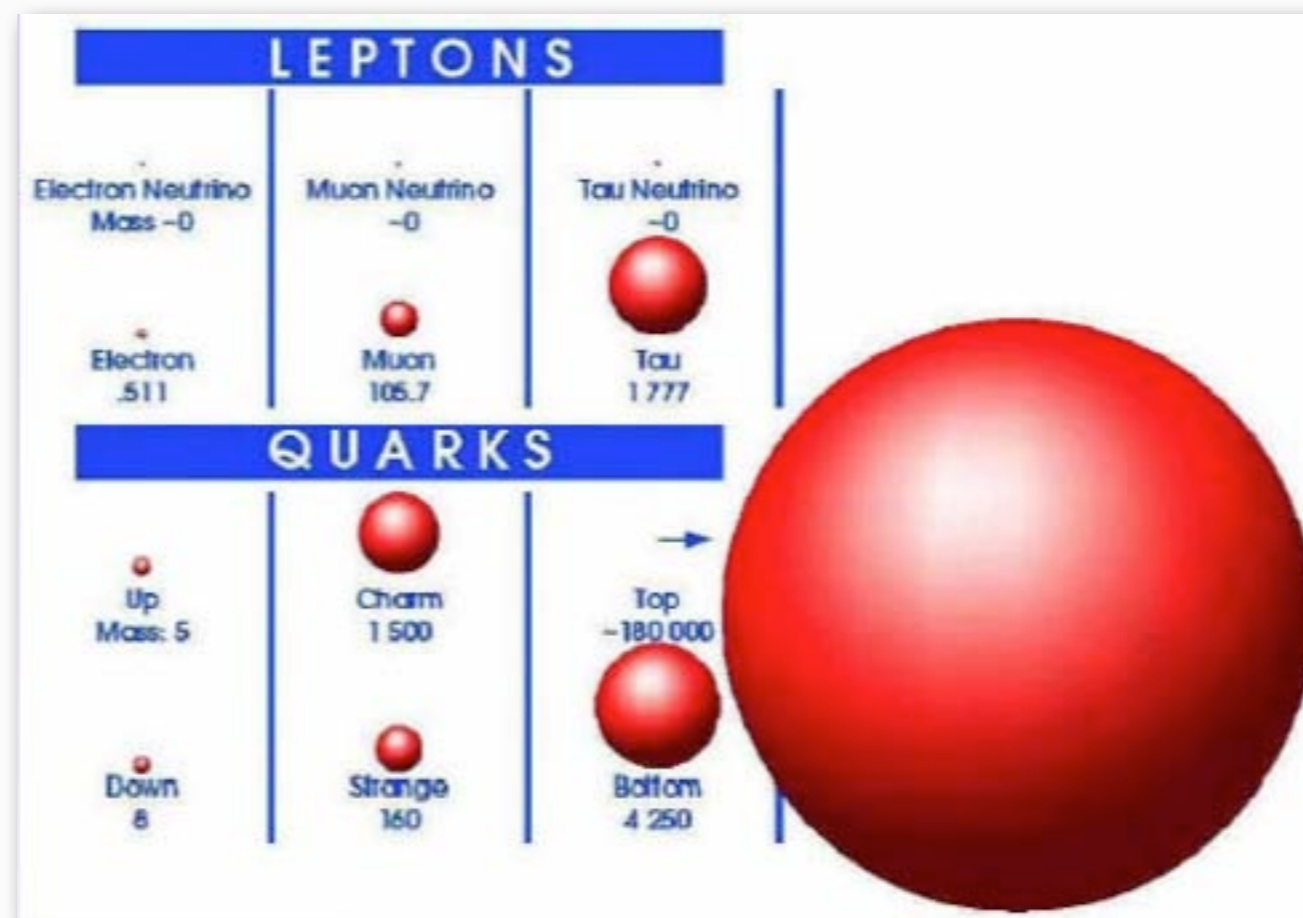
Experimental Z counting precision of 1 - 2 % appears feasible, even after 1. year



- Extremely **important background** for many searches
  - in particular for SUSY searches in the “jets+lepton+E<sub>Tmiss</sub>” channel
- Remember : **Jet scale uncertainty** extremely important (xsec as function of jet p<sub>T</sub>), also here
  - can expect some 30 % uncertainty from that. Probably less in case of rate measurements.
- Should also have a **more “inclusive” look** at it : Measuring the **Z p<sub>T</sub>** can be done with a relative precision at the **per-cent level (leptons (!) again)**, will be invaluable for checking predictions and tuning MCs

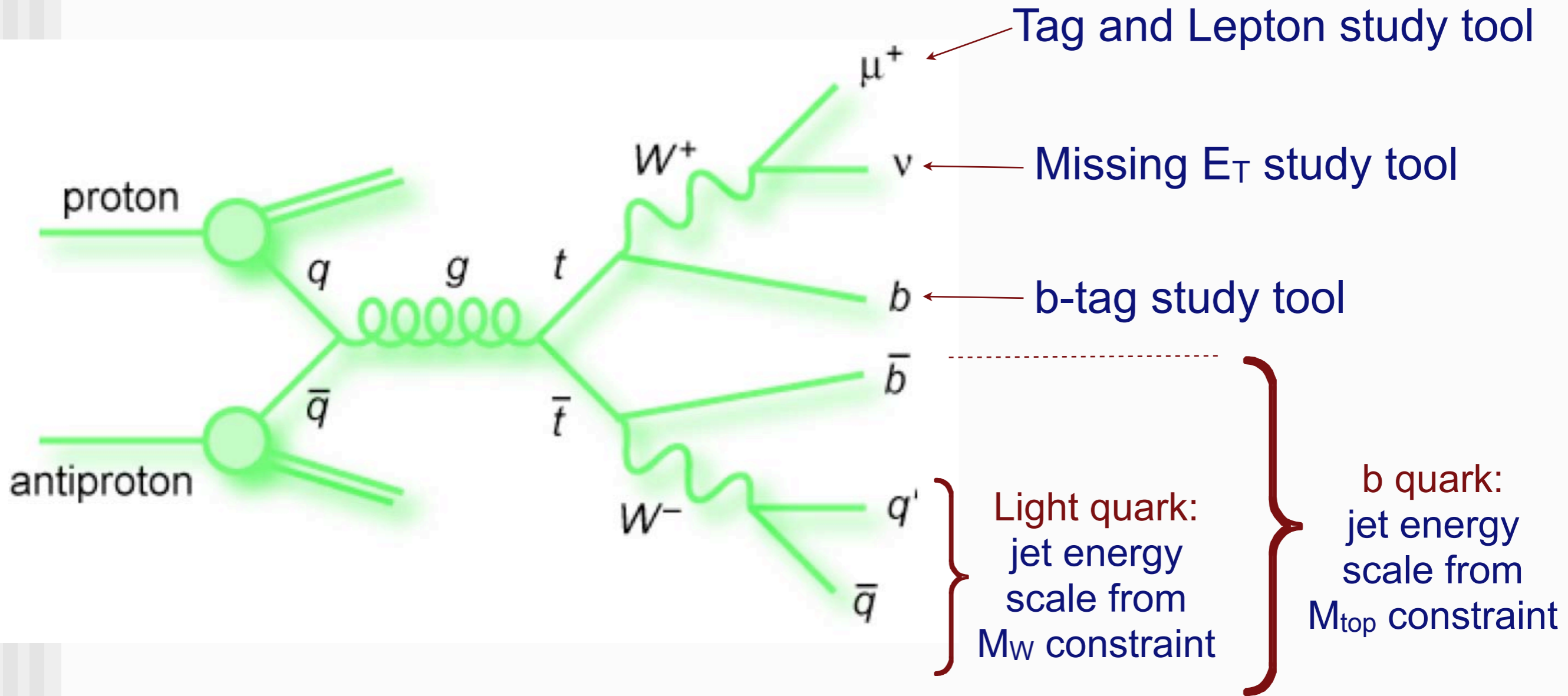


# Top production





# Top as a "Tool"



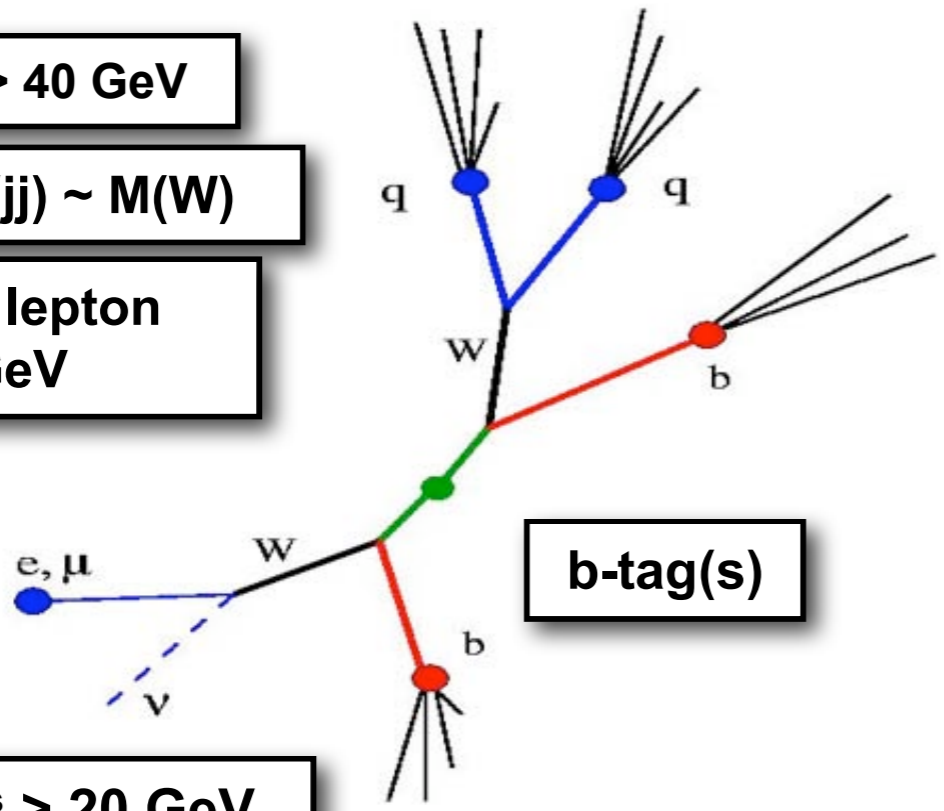


# Top identification

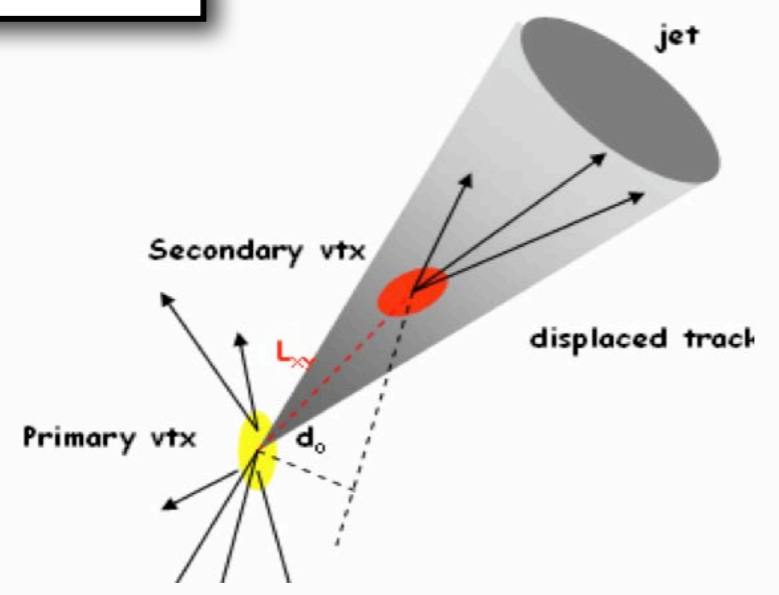
4 jets  $p_T > 40$  GeV

2 jets  $M(jj) \sim M(W)$

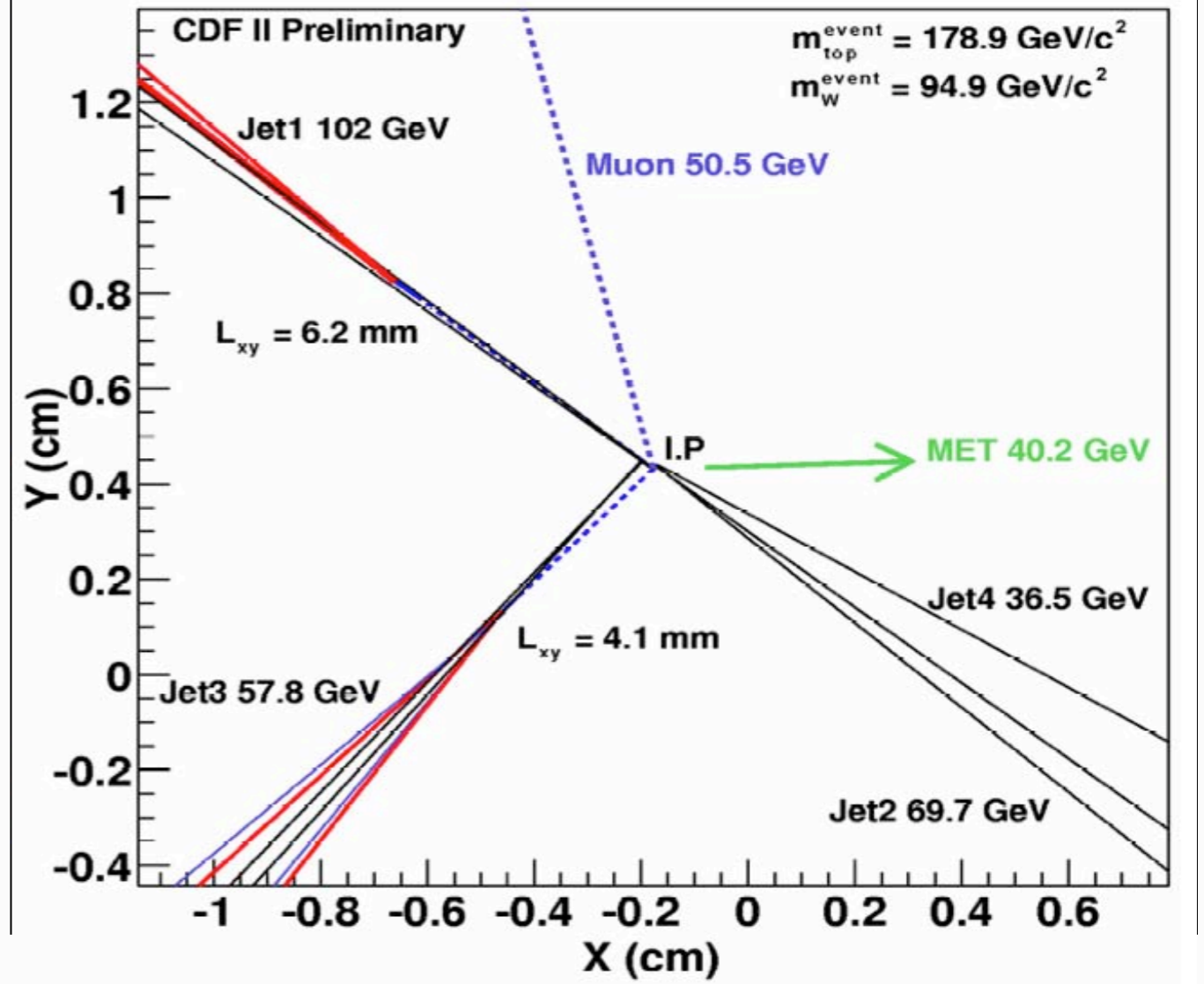
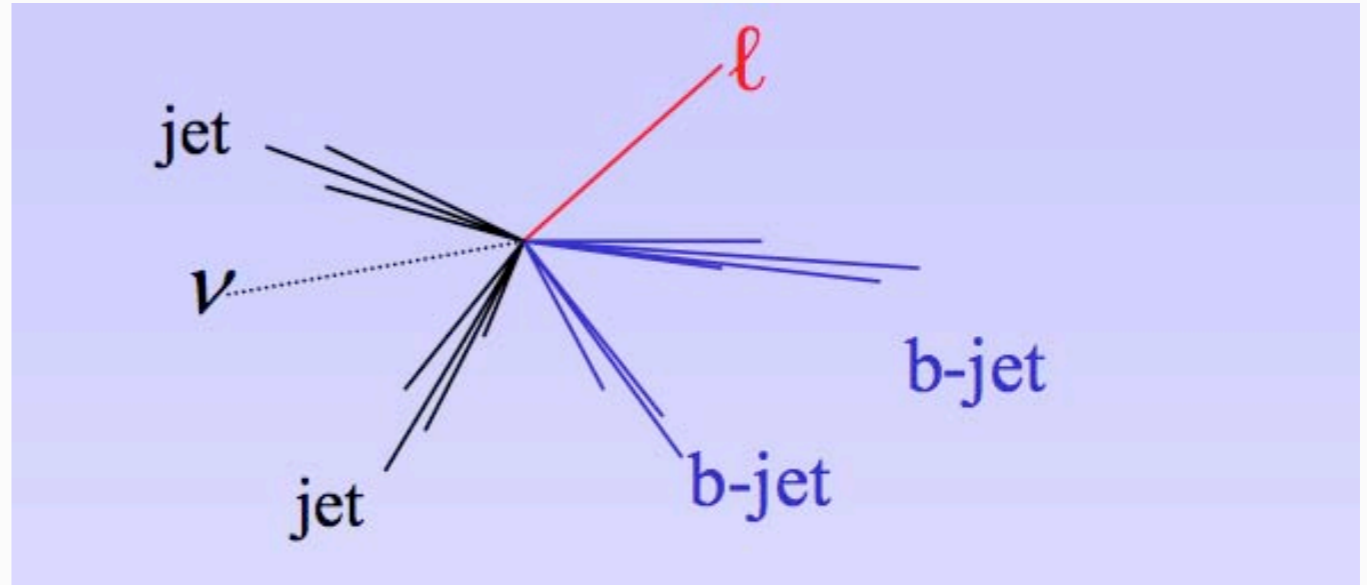
Isolated lepton  $p_T > 20$  GeV



$E_{T,miss} > 20$  GeV



b-tagging important:  
Need excellent Silicon Vertex  
and Pixel Detectors







# Top Production (example : semi-leptonic case)

See the top immediately with simple selection :  
Missing  $E_T$ , 1 lepton,  $\geq 4$  jets ,  
**even without b-tag (!)**,  
cut on hadronic W mass

**Example** (ATLAS study):

- Observe it with  $30 \text{ pb}^{-1}$
- $\sigma(tt)$  to 20 % with  $100 \text{ pb}^{-1}$
- $M(t)$  to 7-10 GeV

**Once b-tagging is understood:**

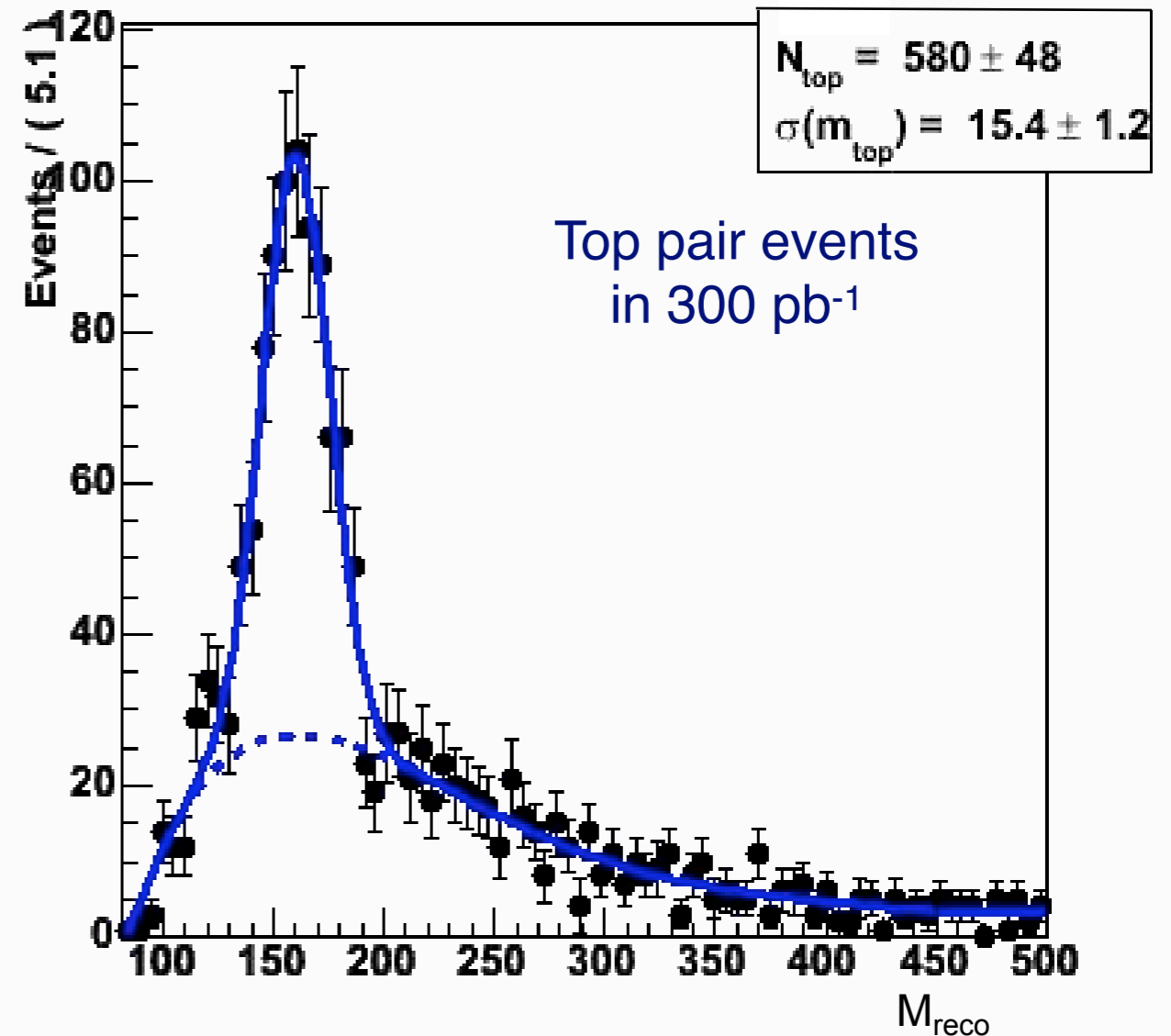
Very high **S/B** achievable  $\sim 27$  !

**Backgrounds :**

W+4j, Wbb+2j(3j) (minor here)

relevant also for single-top

Atlas FullSim Preliminary



Study the **top quark properties**

mass, charge, spin, couplings, production and decay,  
 $\Delta M_{top} \sim 1 \text{ GeV}$  ?

important background for searches

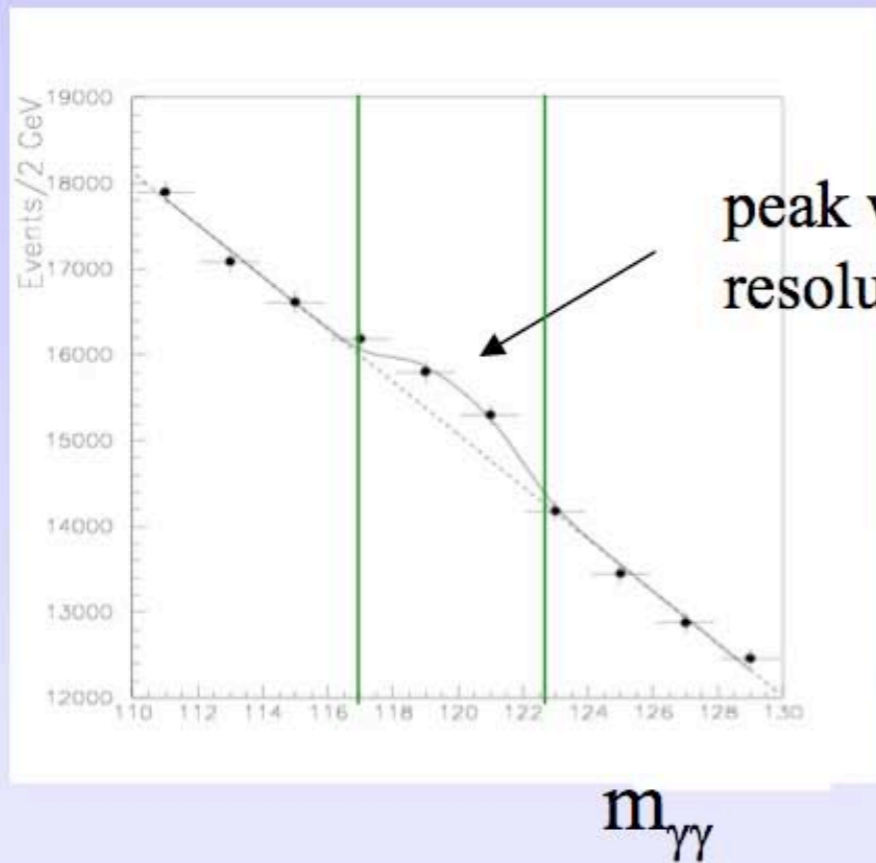
**Jet energy scale** from  $W \rightarrow \text{jet jet}$ ,

commission b-tagging



# General remarks on searches for new physics

Suppose a new narrow particle  $X \rightarrow \gamma\gamma$  is produced:



Signal significance:

$$S = \frac{N_S}{\sqrt{N_B}}$$

$N_S$  = number of signal events

$N_B$  = number of background events

} in peak region

$\sqrt{N_B} \equiv$  error on number of background events, for large numbers  
otherwise: use Poisson statistics

$S > 5$  : signal is larger than 5 times error on background.  
Gaussian probability that background fluctuates up by more than  $5\sigma$  :  $10^{-7} \rightarrow$  **discovery**



# Signal Significance : Issues

## Detector resolution (eg. mass resolution $\sigma_m$ )

If  $\sigma_m$  increases by e.g. a factor of two, then need to enlarge peak region by a factor of two to keep the same number of signal events

→  $N_B$  increases by  $\sim 2$   
(assuming background flat)

⇒  $S = N_S / \sqrt{N_B}$  decreases by  $\sqrt{2}$

⇒  $S \sim 1 / \sqrt{\sigma_m}$

“A detector with better resolution has larger probability to find a signal”

Note: only valid if  $\Gamma_H \ll \sigma_m$ . If Higgs is broad detector resolution is not relevant.

$m_H = 100 \text{ GeV} \rightarrow \Gamma_H \sim 0.001 \text{ GeV}$

$m_H = 200 \text{ GeV} \rightarrow \Gamma_H \sim 1 \text{ GeV}$

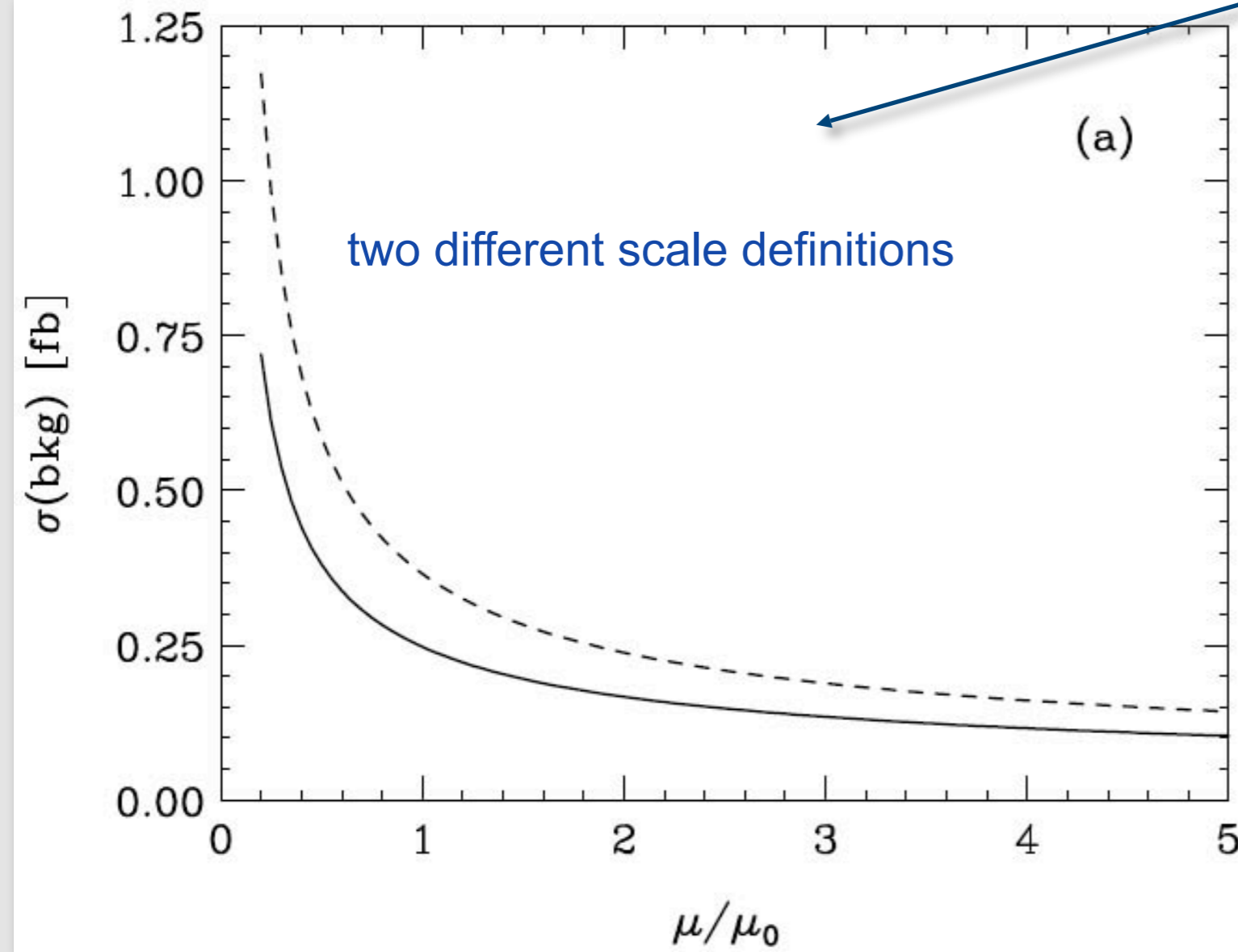
$m_H = 600 \text{ GeV} \rightarrow \Gamma_H \sim 100 \text{ GeV} \quad \Gamma_H \sim m_H^3$

## Luminosity

$N_S \sim L$   
 $N_B \sim L$  }

⇒  $S \sim \sqrt{L}$

Backgrounds to  $H \rightarrow WW \rightarrow e\nu e\nu$  :  $WW$ ,  $t\bar{t}$  for gluon fusion,  $t\bar{t}j$  for  $qqH$



Idea of extrapolation:

Kauer et al.

$$\sigma_{bkg} \approx \underbrace{\left( \frac{\sigma_{bkg, LO}}{\sigma_{ref, LO}} \right)}_{\text{low theoret. uncertainty}} \cdot \underbrace{\sigma_{ref}}_{\text{low experim. uncertainty}}$$

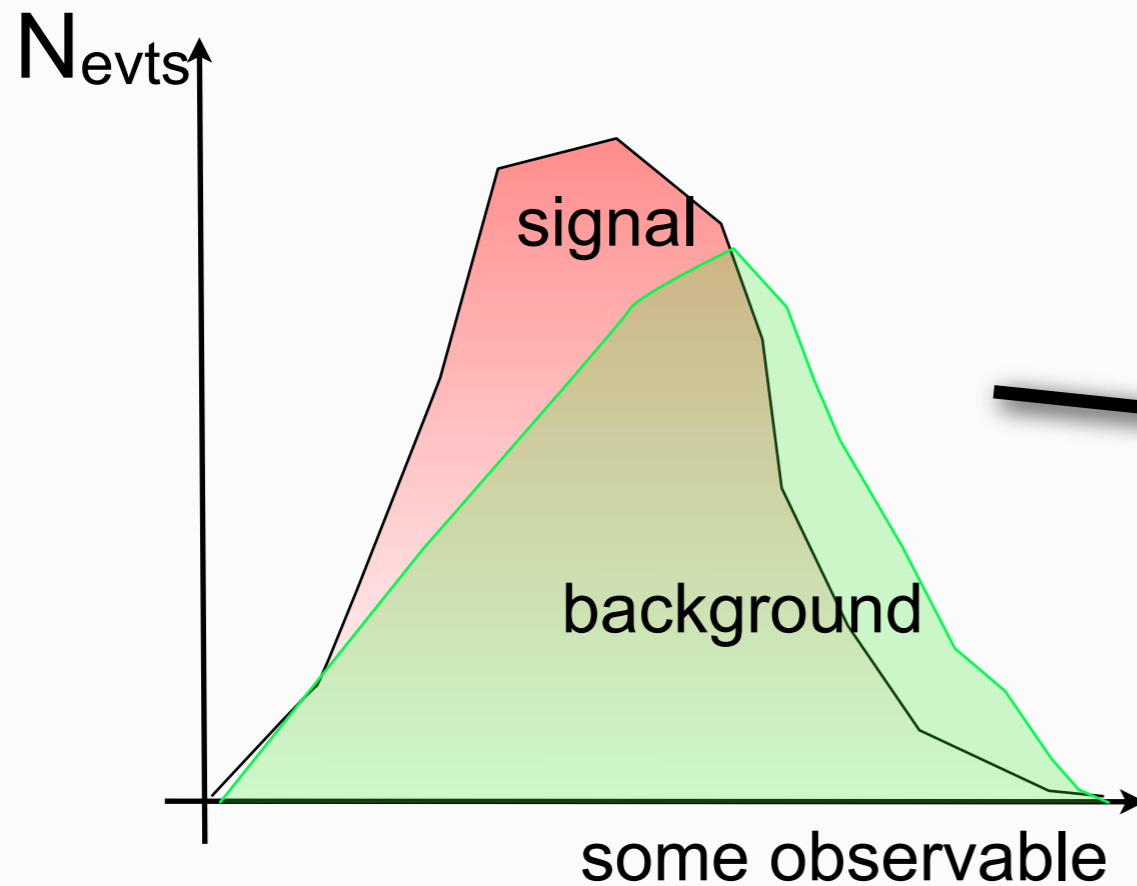
$\sigma_{bkg}$  : background with cuts optimized for finding signal

$\sigma_{ref}$  : background with cuts to enrich background (eg. revert the cuts above)

40-50% scale uncertainty at LO

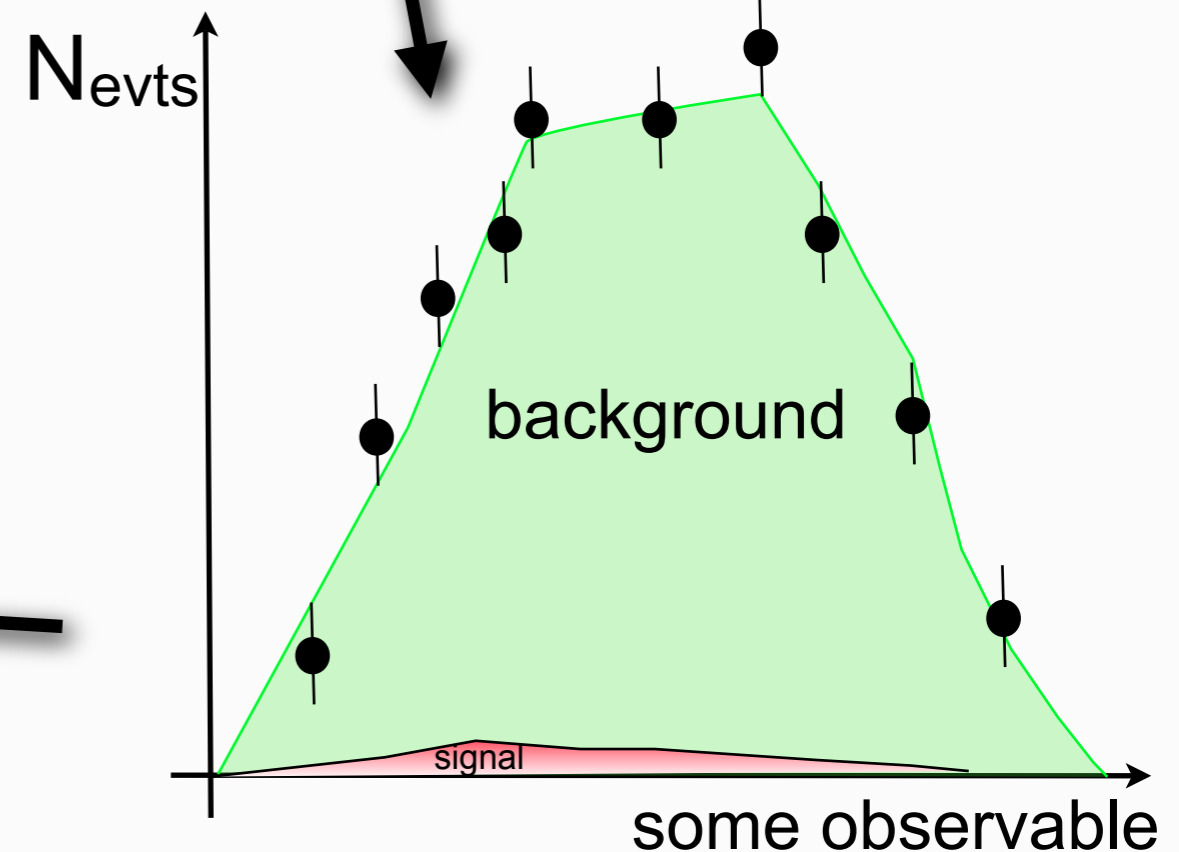


# Background extrapolation



invert cuts :  
from signal enhancement to  
background enhancement

use data to  
normalize background



theory :  
use theory to compute  
change in background  
when inverting cuts

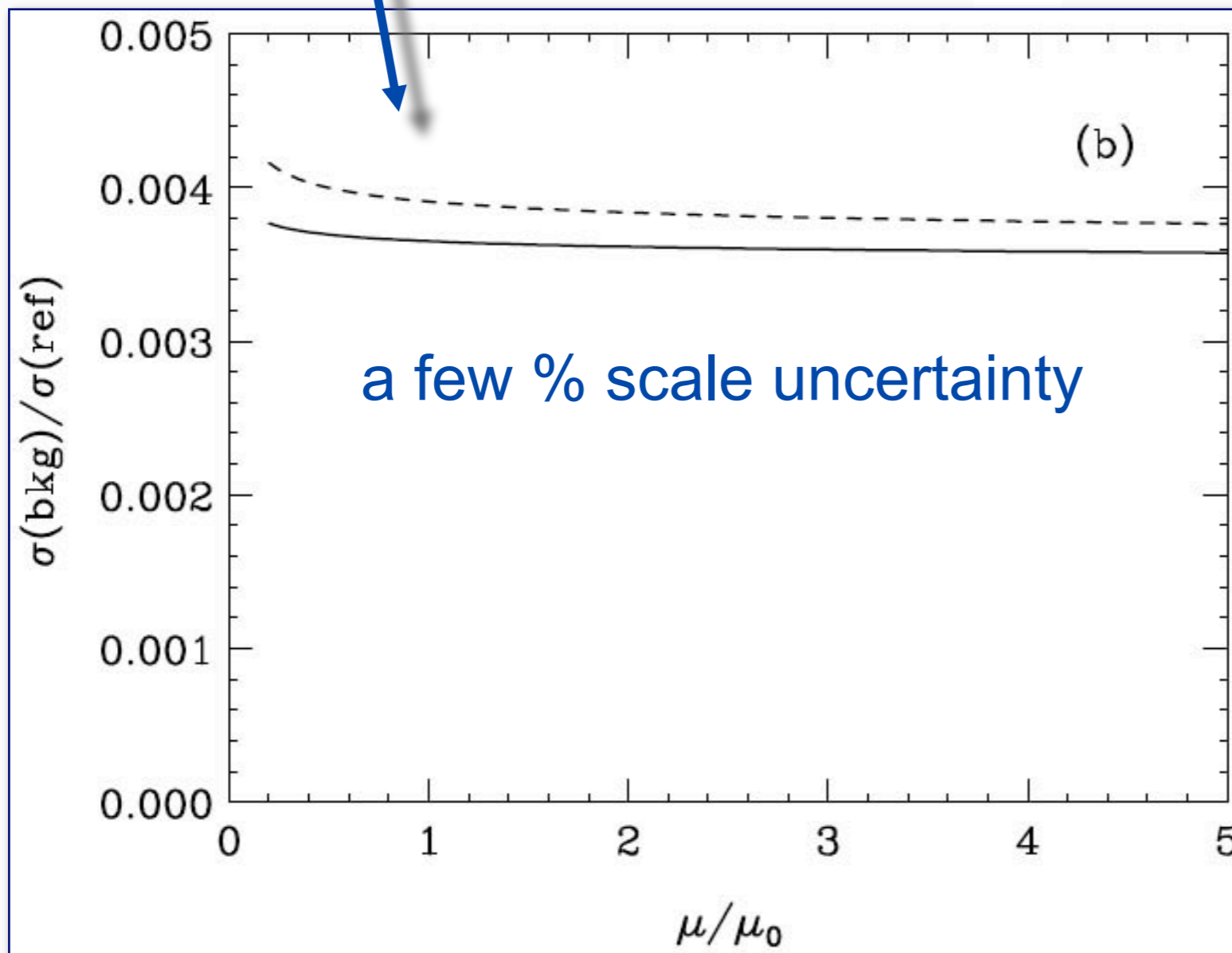
$$\sigma_{bkg} \approx \underbrace{\left( \frac{\sigma_{bkg, LO}}{\sigma_{ref, LO}} \right)}_{\text{low theoret. uncertainty}} \cdot \underbrace{\sigma_{ref}}_{\text{low experim. uncertainty}}$$



## Reference selection :

- like signal, but
- no central jet veto
- no lepton pair cuts
- require b-tag

This enriches the data sample in pure bckg (SM processes), with **high statistics**.





- Always try to be as independent from the Monte Carlo as possible!
  - eg. find a “Standard Model candle” for calibration
  - Obtain **backgrounds from the data** whenever possible
    - Easy if we have mass peak (from sidebands)
    - More difficult in case of excess in high-energy tails, in particular in relation to MET or high- $E_T$  jets
  - Study carefully the validity of a Monte Carlo, and what it is exactly based on
    - eg. LO 2-to-2 process + parton shower, or 2-to-n + parton shower, or NLO+parton shower, or ...
- Worry in particular about systematic errors in your search analysis when  **$S/B \ll 1$  !!**
  - be careful with calculation of significance



- as significance (number of 'sigmas') one usually sees the definition (  $\sigma_{\text{stat(background)}} = \sqrt{n_b}$  for large enough statistics )

$$S = n_\sigma = \frac{n_s}{\sqrt{n_b}}$$

- Adding a relative **systematic uncertainty  $f$** ,  $\sigma_{\text{syst}(n_b)} = f n_b$ , in quadrature to the statistical uncertainty, this becomes:

$$\tilde{n}_\sigma = \frac{n_s}{\sqrt{n_b + f^2 n_b^2}}$$

- this can be rewritten as

$$\tilde{n}_\sigma = n_\sigma \cdot \left[ 1 + \left( \frac{f \cdot n_\sigma}{n_s / n_b} \right)^2 \right]^{-\frac{1}{2}}$$

- limiting cases:

$$n_s / n_b \ll f \cdot n_\sigma \quad \Rightarrow \quad \tilde{n}_\sigma \approx \frac{n_s / n_b}{f} \quad \text{dominated by systematics}$$

$$n_s / n_b \gg f \cdot n_\sigma \quad \Rightarrow \quad \tilde{n}_\sigma \approx n_\sigma \quad \text{dominated by statistics}$$

- a concrete example (10% background uncertainty)

$n_s$	$n_b$	$n_s / n_b$	$n_\sigma$	$\tilde{n}_\sigma$
50	100	0.5	5	3.5
500	10000	0.05	5	0.5

- in the second case, more luminosity will not improve the significance!  
(unless more data help to better understand the background)

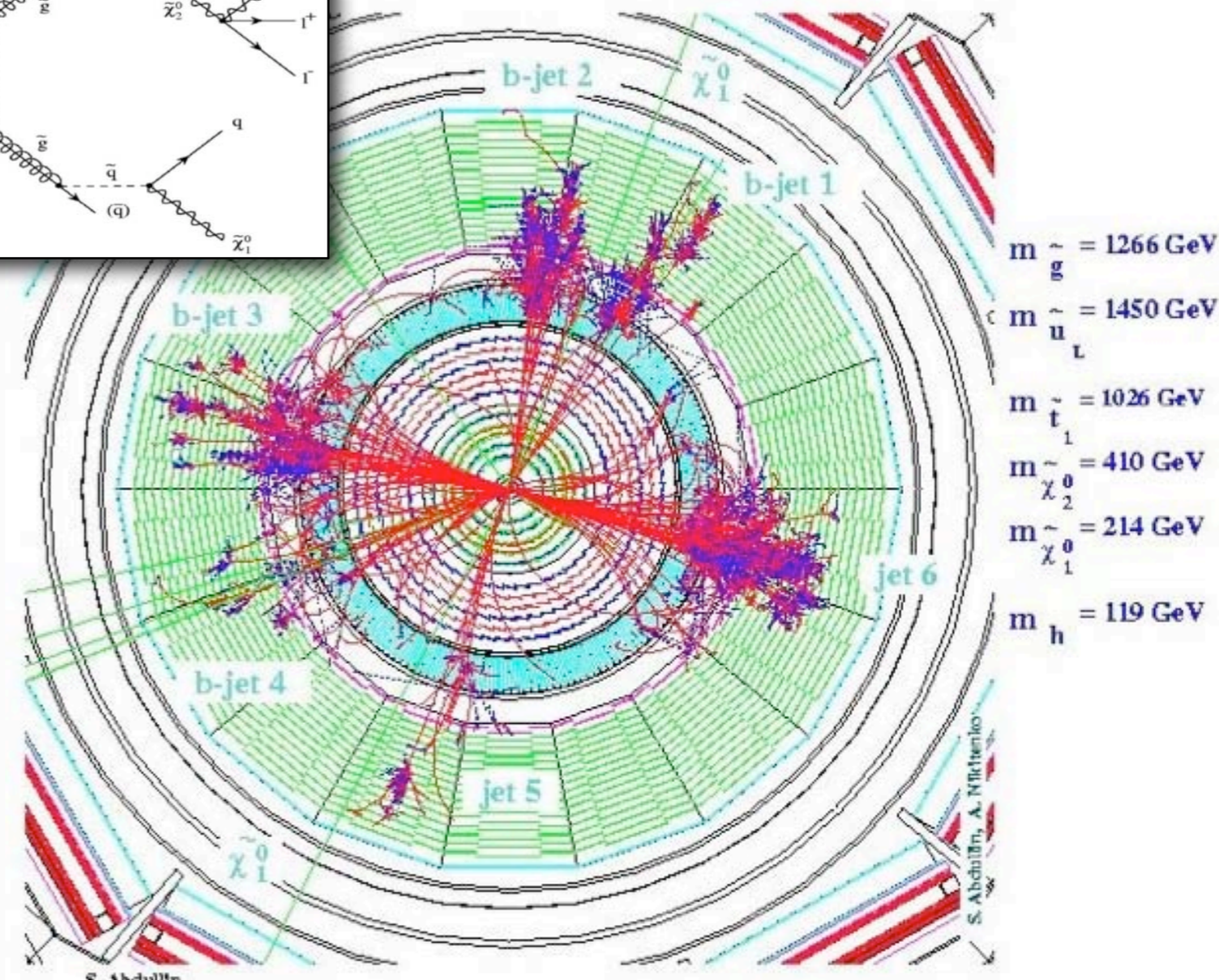
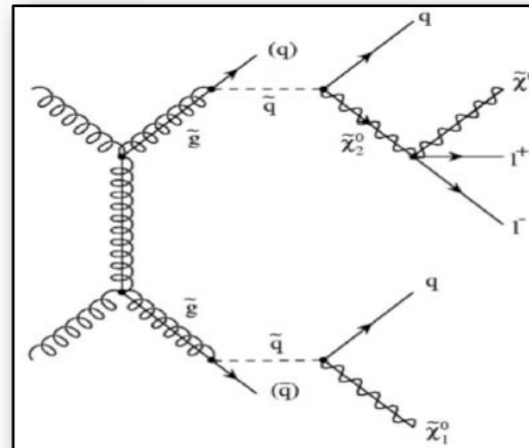
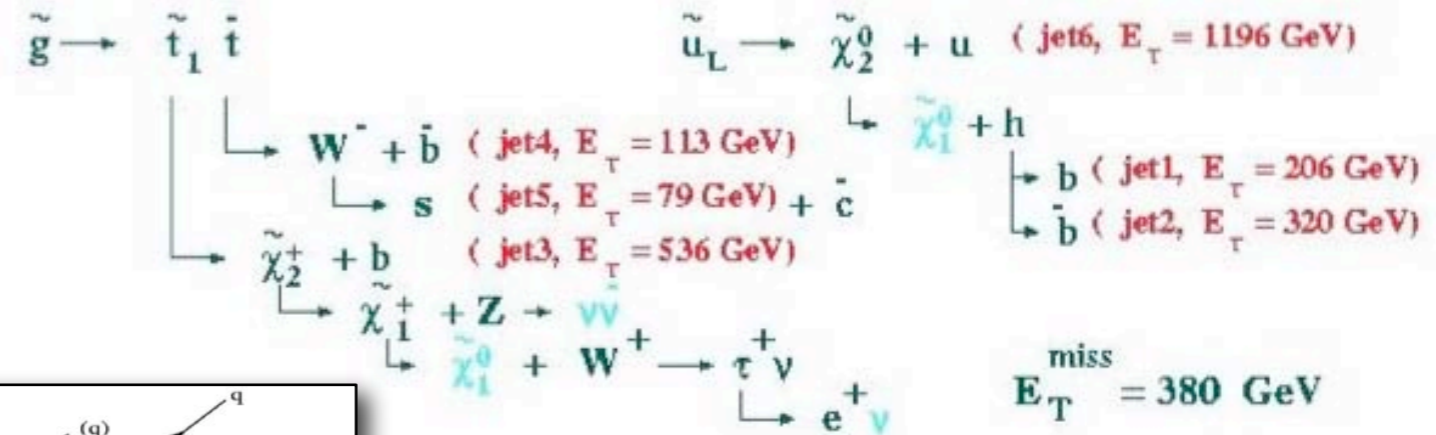


# SUSY and Missing $E_T$



# SUSY signatures

- Many hard jets
- Large missing energy
  - 2 LSPs
  - Many neutrinos
- Many leptons
- ie. : **Spectacular**



$M_{sp}(\text{GeV})$	$\sigma \text{ (pb)}$	$\text{Evs/yr}$
500	100	$10^6 - 10^7$
1000	1	$10^4 - 10^5$
2000	0.01	$10^2 - 10^3$

- for low masses the LHC becomes a real **SUSY factory**

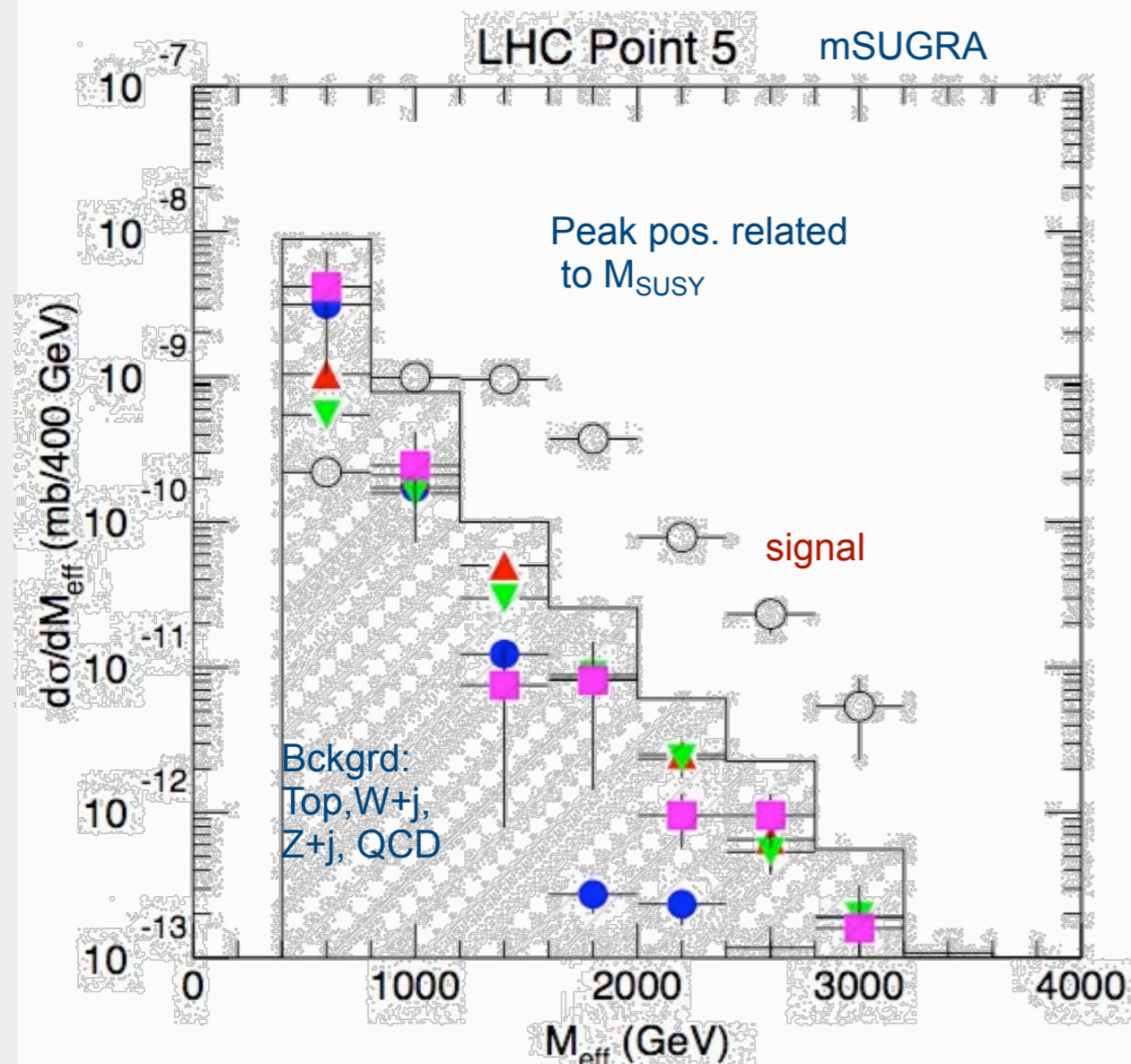


# Seems easy, but....

## Relies on

good reconstruction and understanding of

- multi-jet backgrounds
- Missing transverse Energy



$$M_{eff} = E_T^{miss} + P_T^1 + P_T^2 + P_T^3 + P_T^4$$

## Typical selection

- $N_{jet} > 4$
- $E_T > 100, 50, 50, 50$  GeV
- $E_{T,miss} > 100$  GeV

## Warning

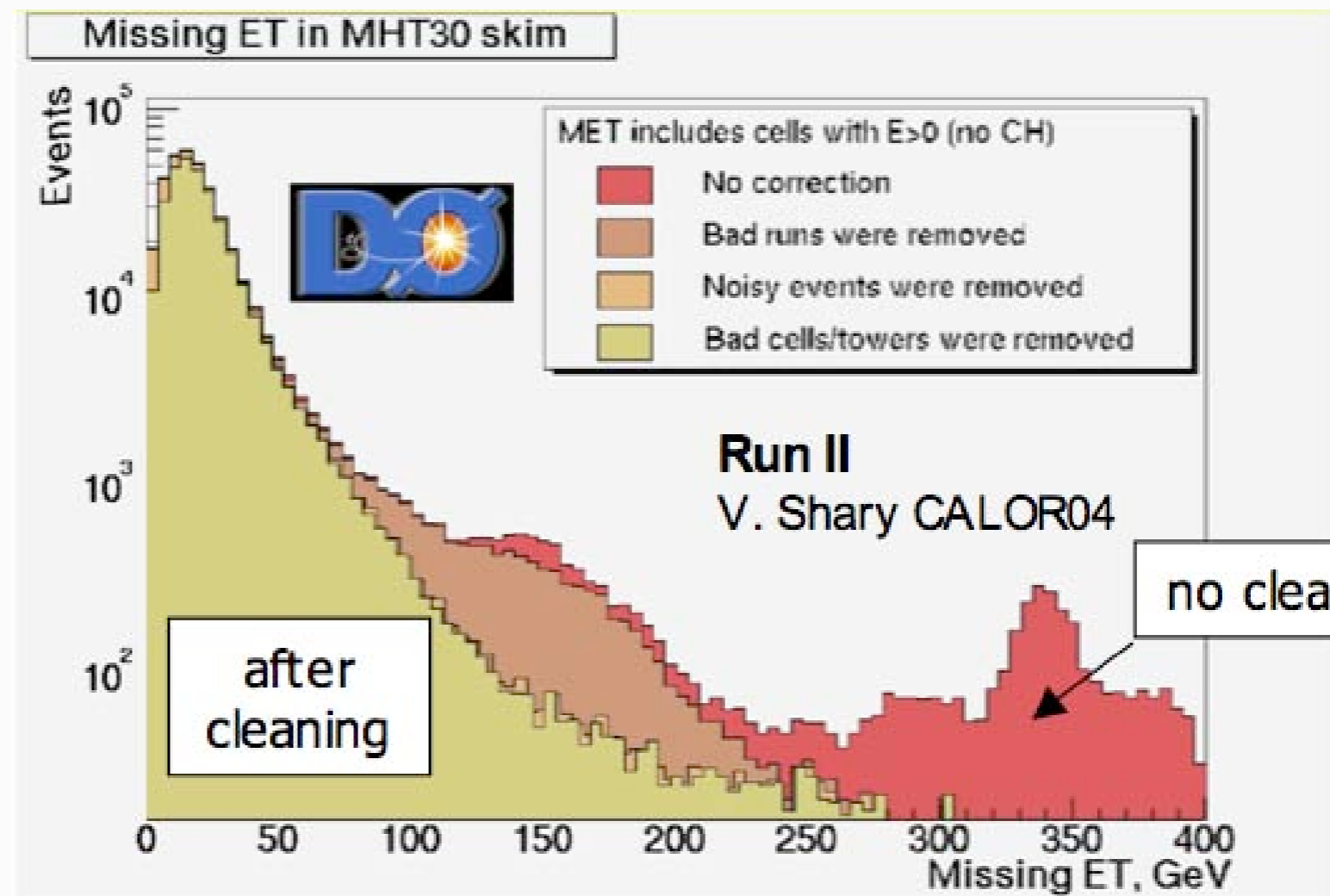
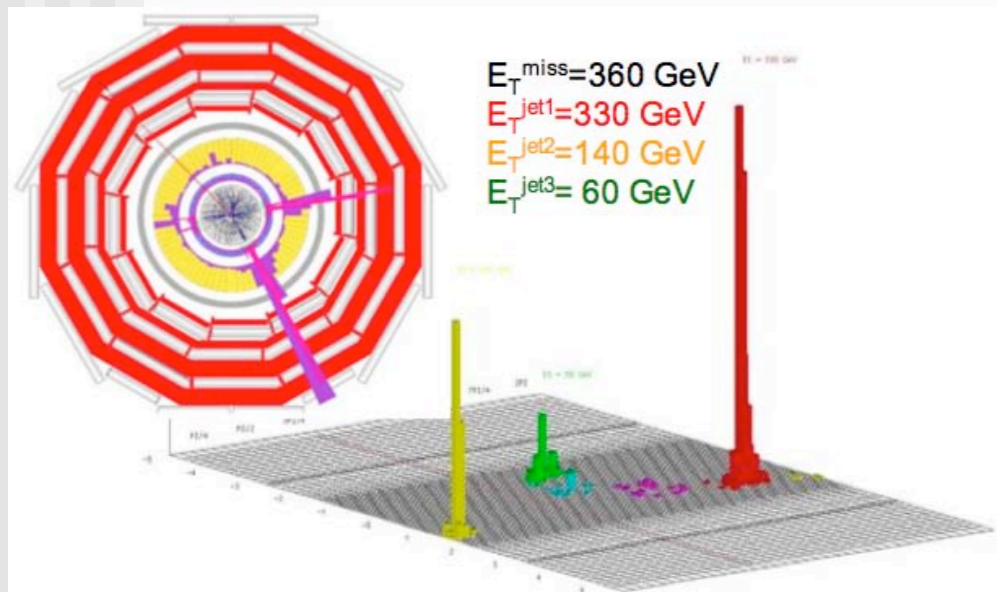
- For description of multi-jet backgrounds a simple Parton-shower MC is not good enough
- Have to combine with matrix elements, eg. ALPGEN



# Seems easy, but....

## Relies on

- good reconstruction and understanding of
  - multi-jet backgrounds
  - Missing transverse Energy



$E_{T,miss}$  is a very tricky variable to measure. Mis-measurements can easily fake "signals".

Have to fight backgrounds, noise, use control samples, eg. Z+jets



# Other discoveries





# The easy case: a $Z'$ (or similar)

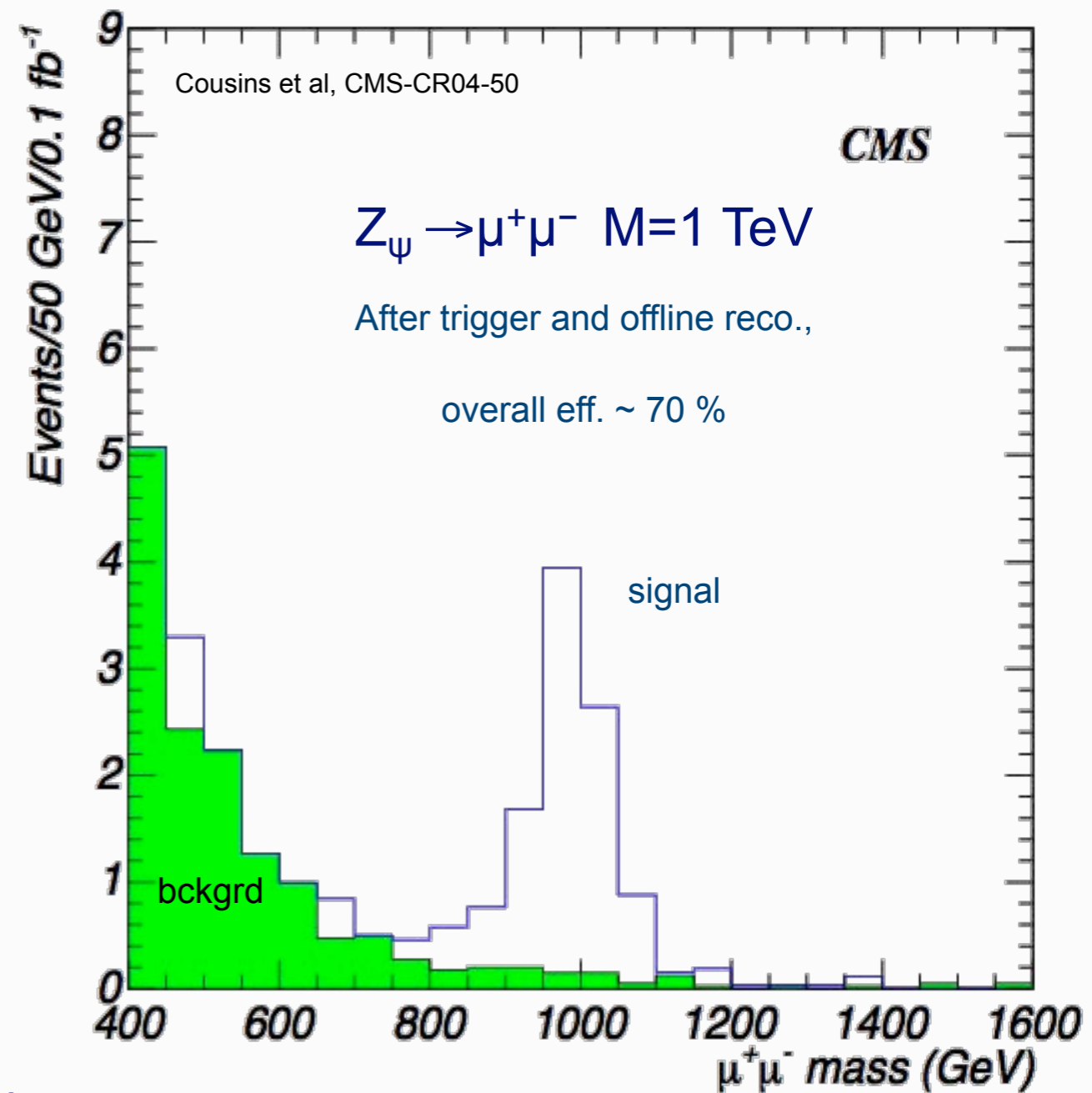
- $Z'$  : generic for new heavy gauge bosons
  - GUT, dynamical EWSB, little Higgs, ...
  - Clear signature
  - low background, mainly Drell-Yan

## One of main issues

- early control of **lepton reconstruction**, eg.
- alignment effects** reduce sensitivity by  $\sim 50\%$  in the early days ( $< 100 \text{ pb}^{-1}$ )

## Similar ATLAS study for $Z' \rightarrow e^+e^-$

- In SSM, SM-like couplings
- $\sim 1.5 \text{ fb}^{-1}$  needed for discovery up to 2 TeV
- $Z \rightarrow \ell\ell + \text{jet}$  and DY needed to get E-calibration and understand lepton eff.

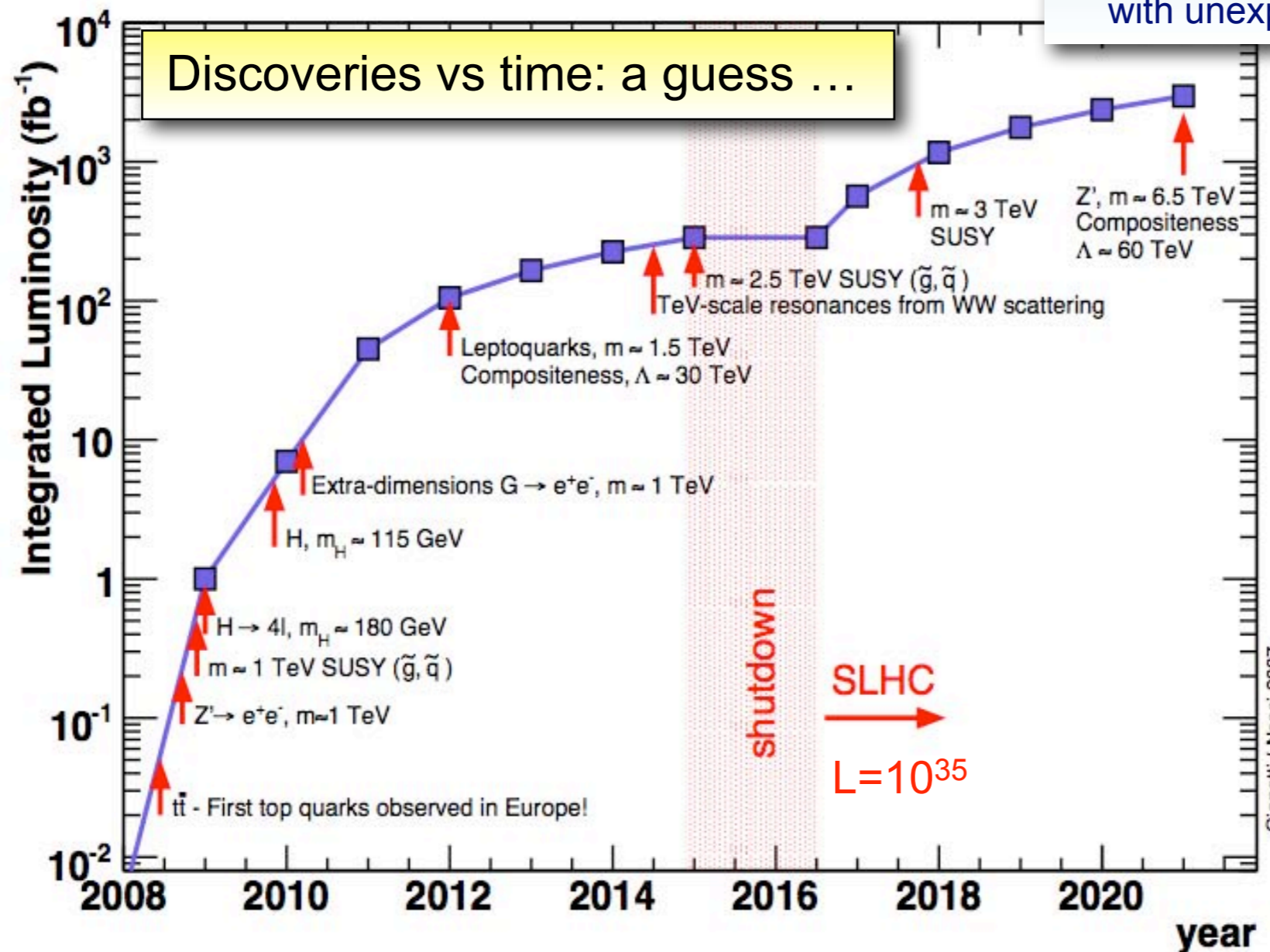




# What LHC could discover (besides Higgs and SUSY)

Excited quarks  $q^* \rightarrow \gamma q$ : up to  $m \approx 6$  TeV  
 Leptoquarks: up to  $m \approx 1.5$  TeV  
 Monopoles  $pp \rightarrow \gamma \gamma pp$ : up to  $m \approx 20$  TeV  
 Compositeness: up to  $\Lambda \approx 40$  TeV  
 $Z' \rightarrow ll, jj$ : up to  $m \approx 5$  TeV  
 $W' \rightarrow l \nu$ : up to  $m \approx 6$  TeV  
 etc.... etc....

Large number of scenarios studied  
 Main conclusions:  
 $\Rightarrow$  LHC direct discovery reach up to  $m \sim 5-6$  TeV  
 $\Rightarrow$  demonstrated detector sensitivity to many signatures  $\rightarrow$  robustness, ability to cope with unexpected scenarios





# Summary of Part 3

*“The only place where success comes before work is the dictionary”*



- **SM physics at the LHC:**  
we will have to re-discover the SM before going to other discoveries
- **Test the SM at an unprecedented energy scale**
  - lots of highly exciting and interesting physics
    - Jets, Ws and Zs, tops, ...
- **These are also important tools to**
  - understand, study, calibrate and improve the detector performance
  - constrain physics input (pdfs, underlying event)
  - necessary input for all other measurements
- **We are getting ready now to be able to perform all these measurements and run these tools as early as possible, once the data start flowing in....**

- Prospects for discoveries are very good
  
- At the LHC we have
  - large cross sections
  - spectacular signals
    - for many new signals
  
- But
  - have to understand the detector first
  - as well as the SM backgrounds
  - Thus : be careful not to claim discovery too early
  - but also not too late ... ;-)
  
- In any case, extremely exiting years are ahead of us!

