

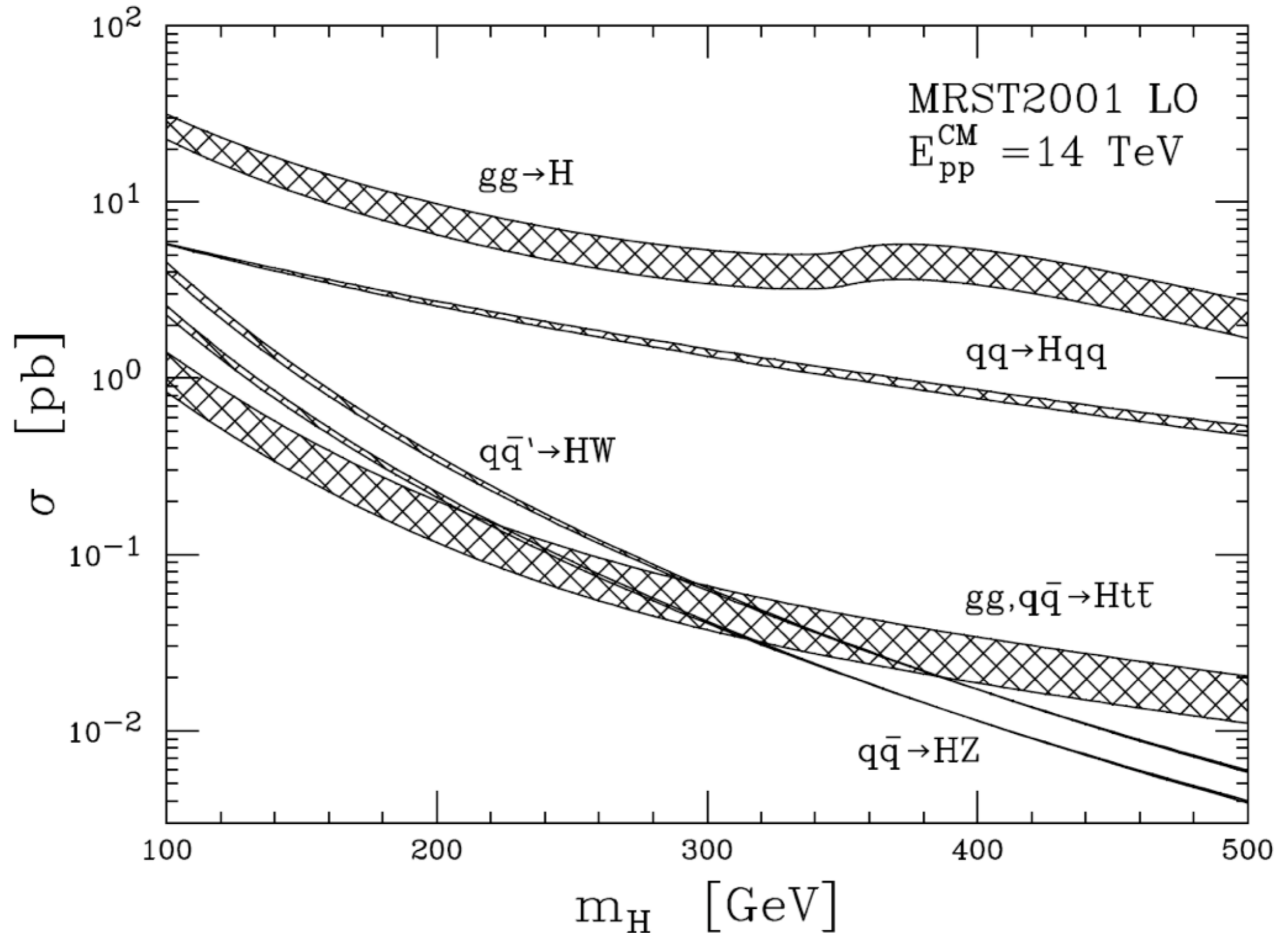
Higgs

in gluon fusion
and other NNLO
stories

by A. Lazopoulos
ETH ZURICH
15 september 2010

HP².3rd

Gluon fusion is the largest Higgs production mode at the LHC



Total cross section @ NLO
Dawson | Spira, Djouadi, Graudenz,
Zerwas

Total Cross section @ NNLO
Harlander, Kilgore | Anastasiou,
Melnikov | Ravindran, Smith, van
Neerven

Threshold resummation
Catani, de Florian, Grazzini, Nason |
Moch, Vogt | Laanen, Magnea | Kulesza,
Sterman | Idilbi, Xi, Ma, Juan | Ravindran
| Ahrens, Becher, Neubert

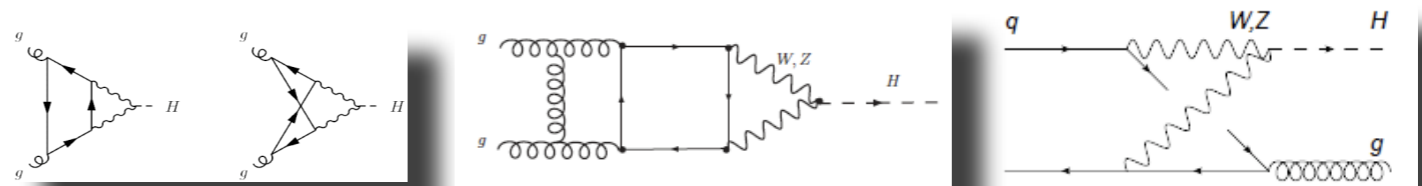
Pt resummation
Bozzi, Catani, de Florian, Grazzini

Two-loop Light fermions EW
Aglietti, Bonziani, Degrassi, Vicini

Two-loop EW
Actis, Passarino, Sturm, Ucciratti

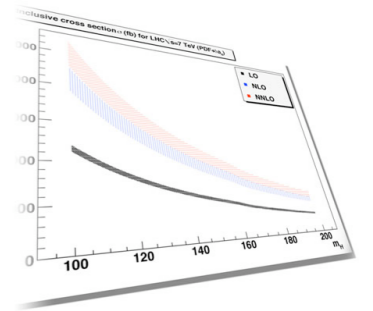
Three-loop mixed qcd and ewk
(Anastasiou, Boughezal, Petriello)

One-loop ewk, $P_t > 0$
Keung, Petriello




FeHiPro: next version

- FeHiPro is undergoing some major renovations. New version out soon.
- ZZ decay to leptons included
- Finite Higgs width effects included
- Flexible python interface for arbitrary cuts and histograms, mass scans, PDF and scale uncertainties etc.
- BSM physics with FeHiPro (see talk by [E.Furlan](#))

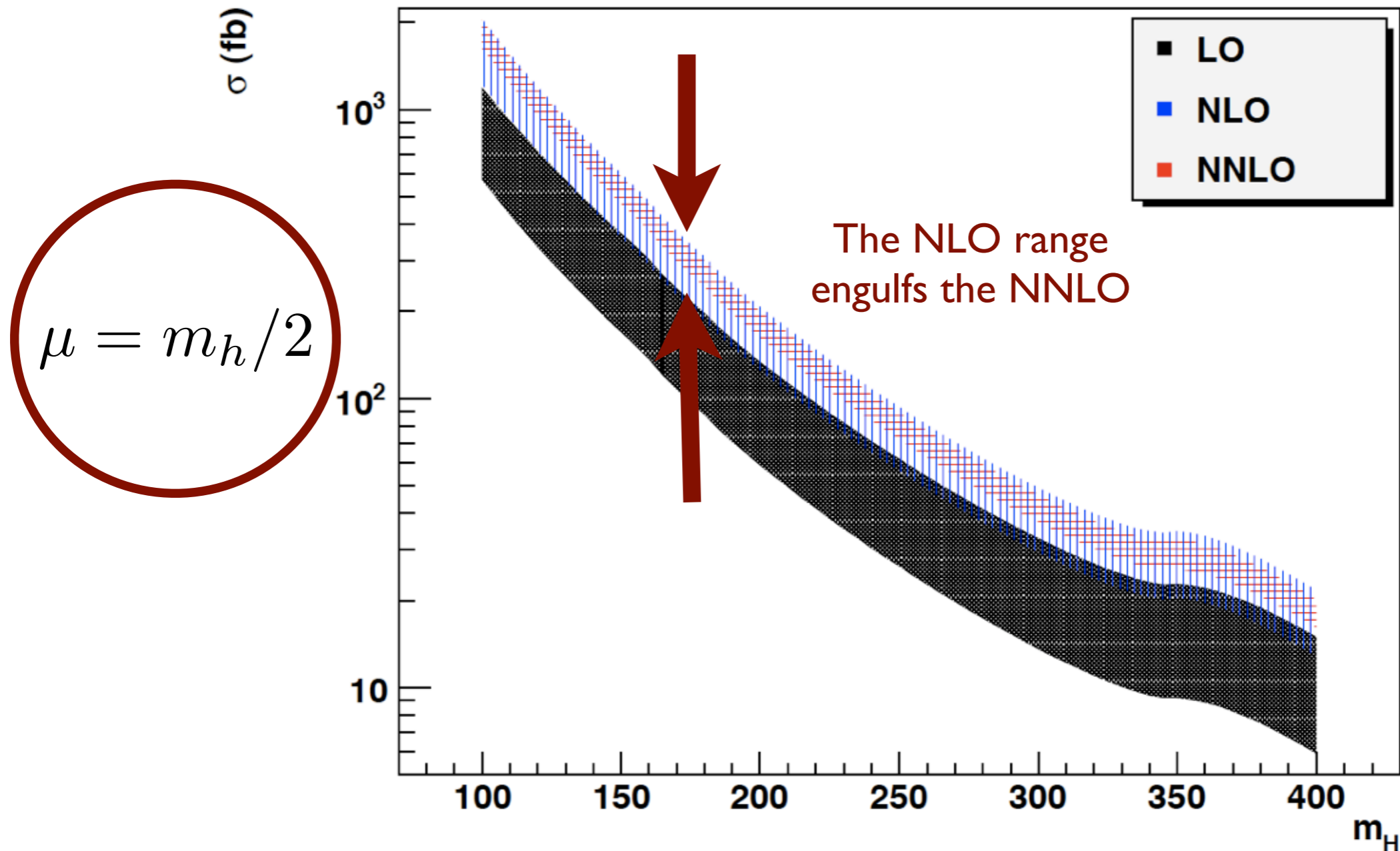


Total and differential cross section

- use NNLO computation
 - choose scales such that the physics is captured ($\mu_0 = m_h/2$) 
 - Include PDFs at NNLO
 - Exact quark mass effects at NLO
 - Include known e/w corrections
 - Supply PDF error (including a_s)
- Structure of logs
SCET studies
Perturbative convergence

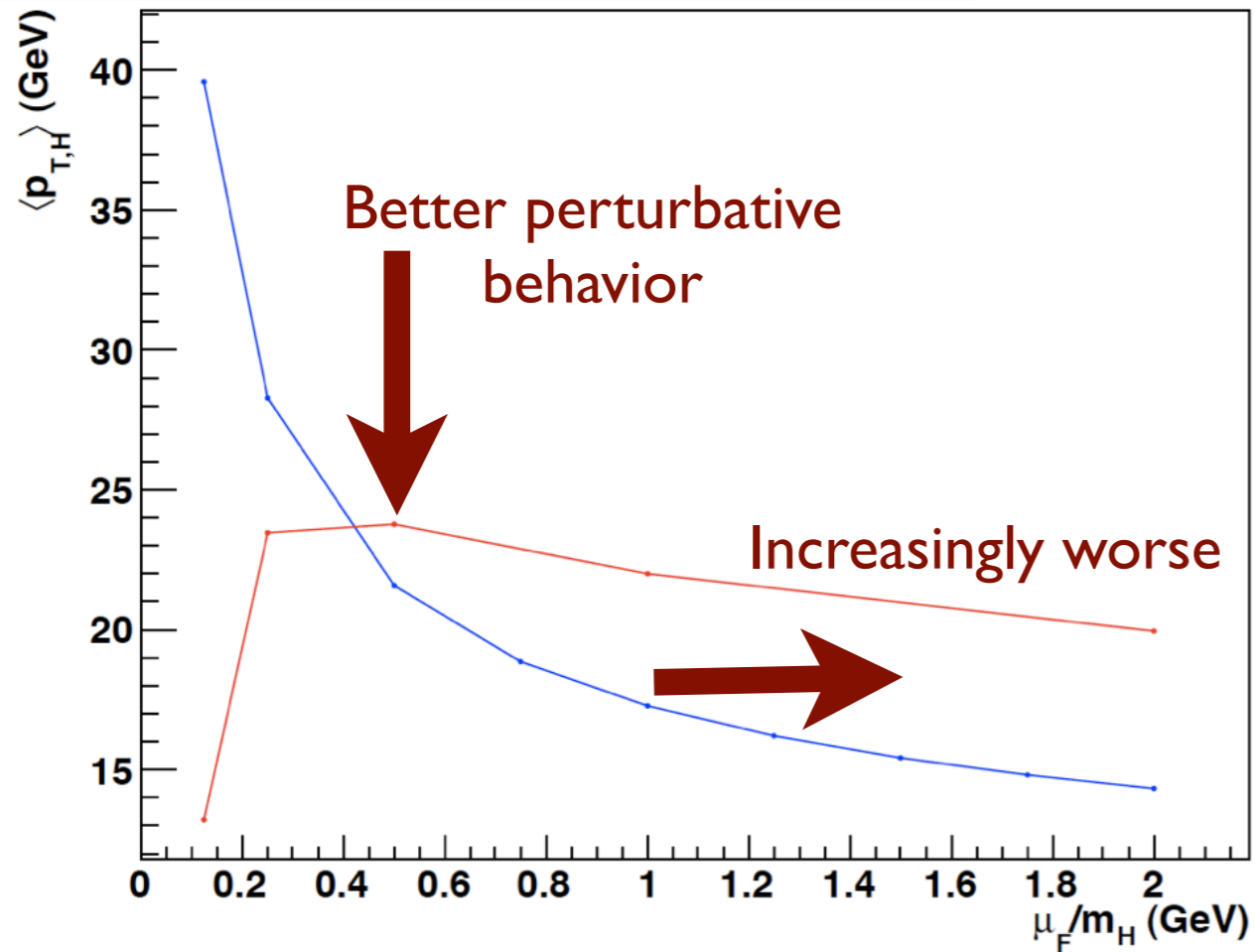
What scale ?

Total Inclusive cross section σ (fb) for Tevatron $\sqrt{s}=1.96$ TeV $\delta\mu$

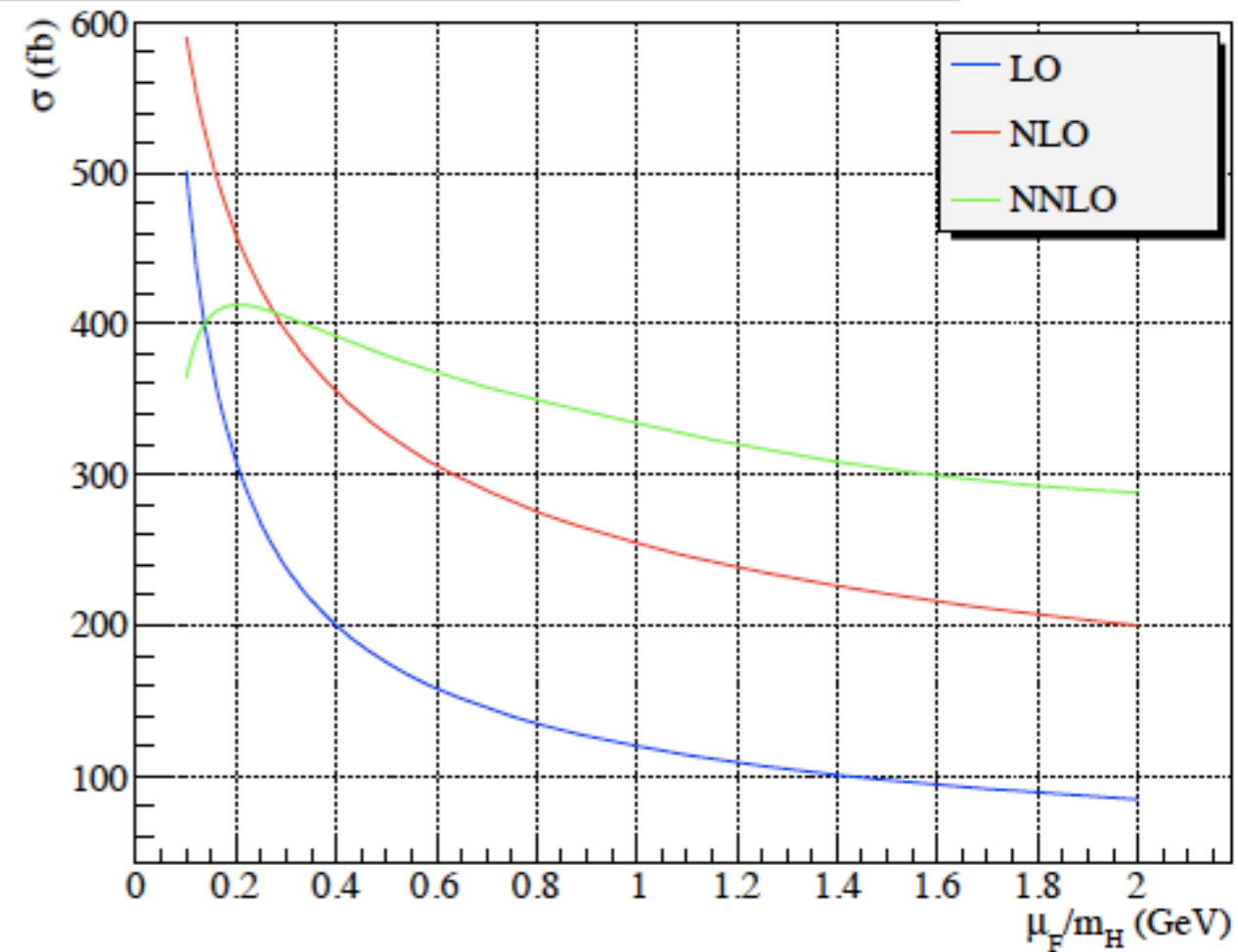


What scale?

$\langle p_{T,H} \rangle$ (GeV) as a function of μ_F/m_H at $m_H=165\text{GeV}$ for Tevatron ($\sqrt{s}=1.96\text{TeV}$)

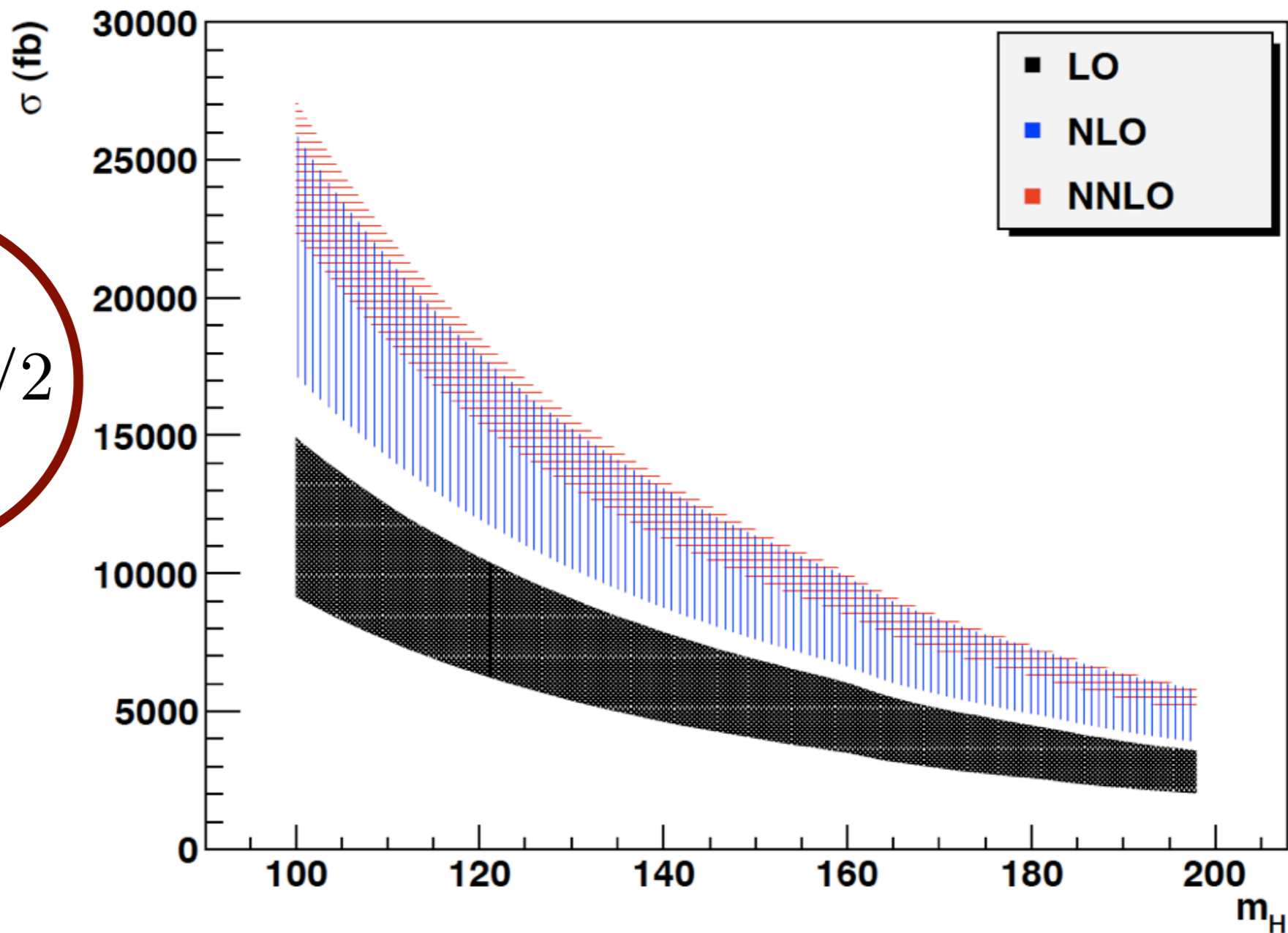


Total cross section as a function of μ_F/m_H at $m_H=165\text{GeV}$ for Tevatron ($\sqrt{s}=1.96\text{TeV}$)



What scale?

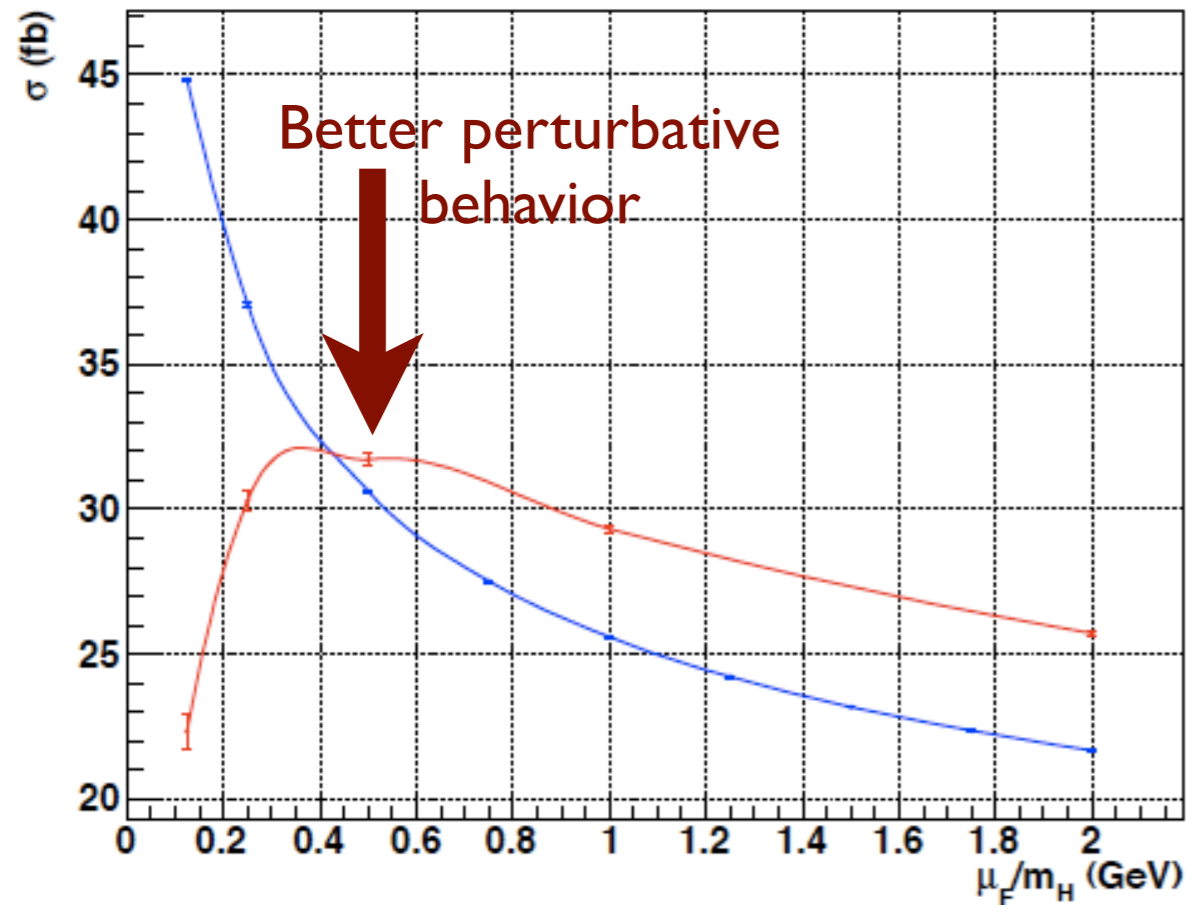
Total Inclusive cross section σ (fb) for LHC $\sqrt{s}=7$ TeV $\delta\mu$



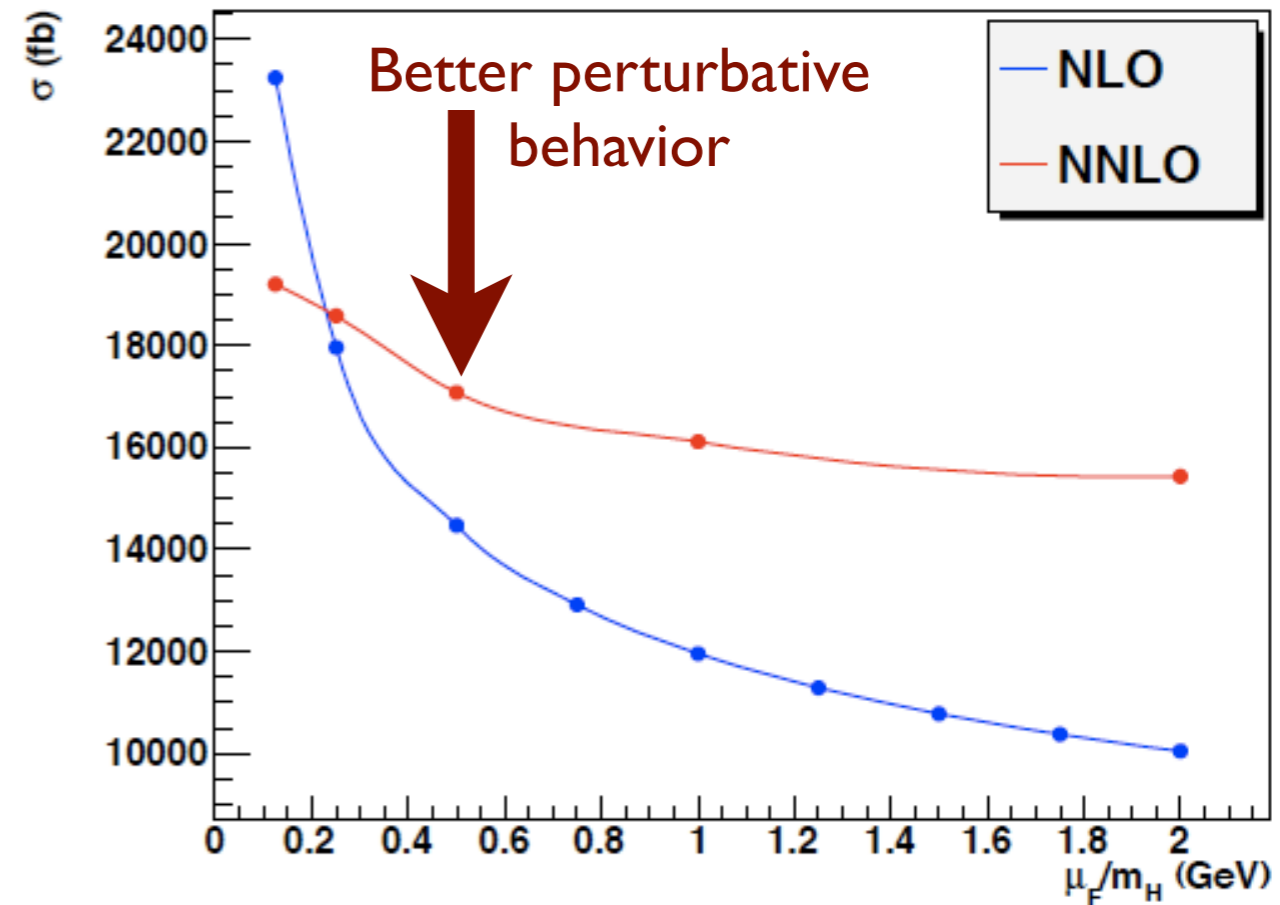
$$\mu = m_h/2$$

What scale?

$\langle p_{T,H} \rangle$ (GeV) as a function of μ_F/m_H at $m_H=120\text{GeV}$

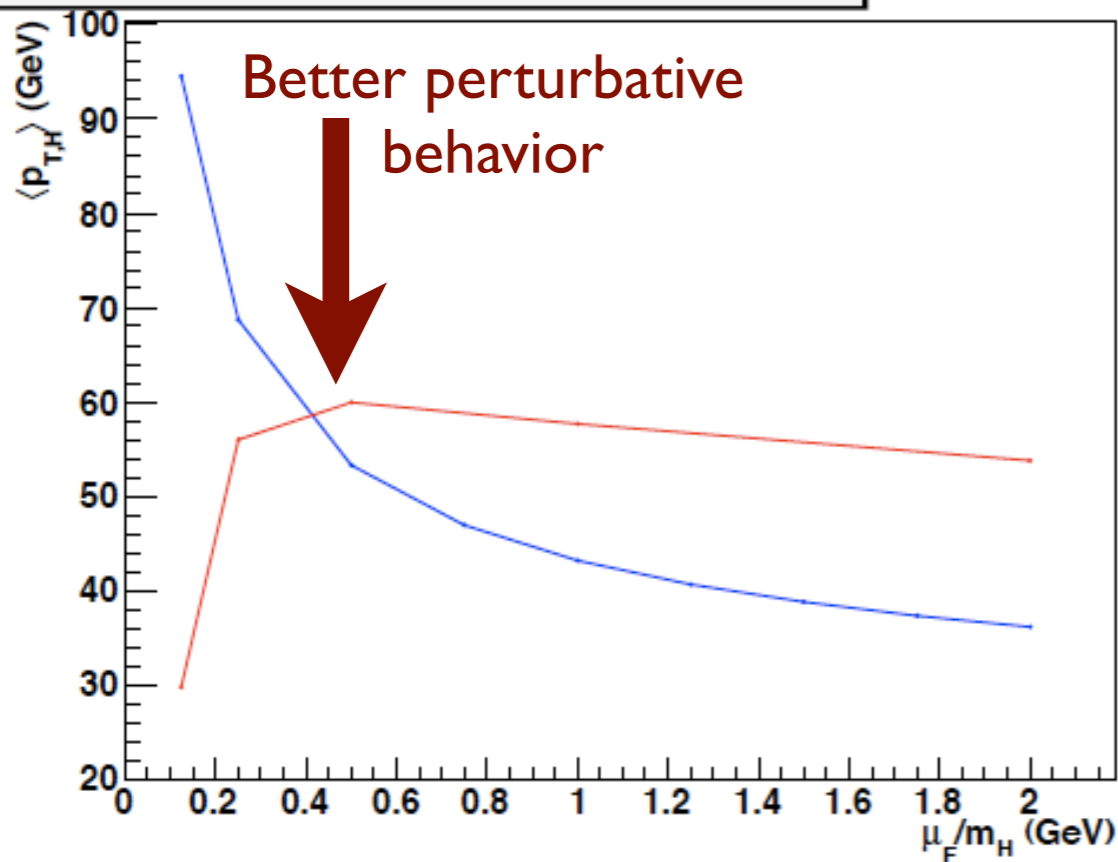


Total cross section as a function of μ_F/m_H at $m_H=120\text{GeV}$ for LHC $\sqrt{s}=7\text{TeV}$

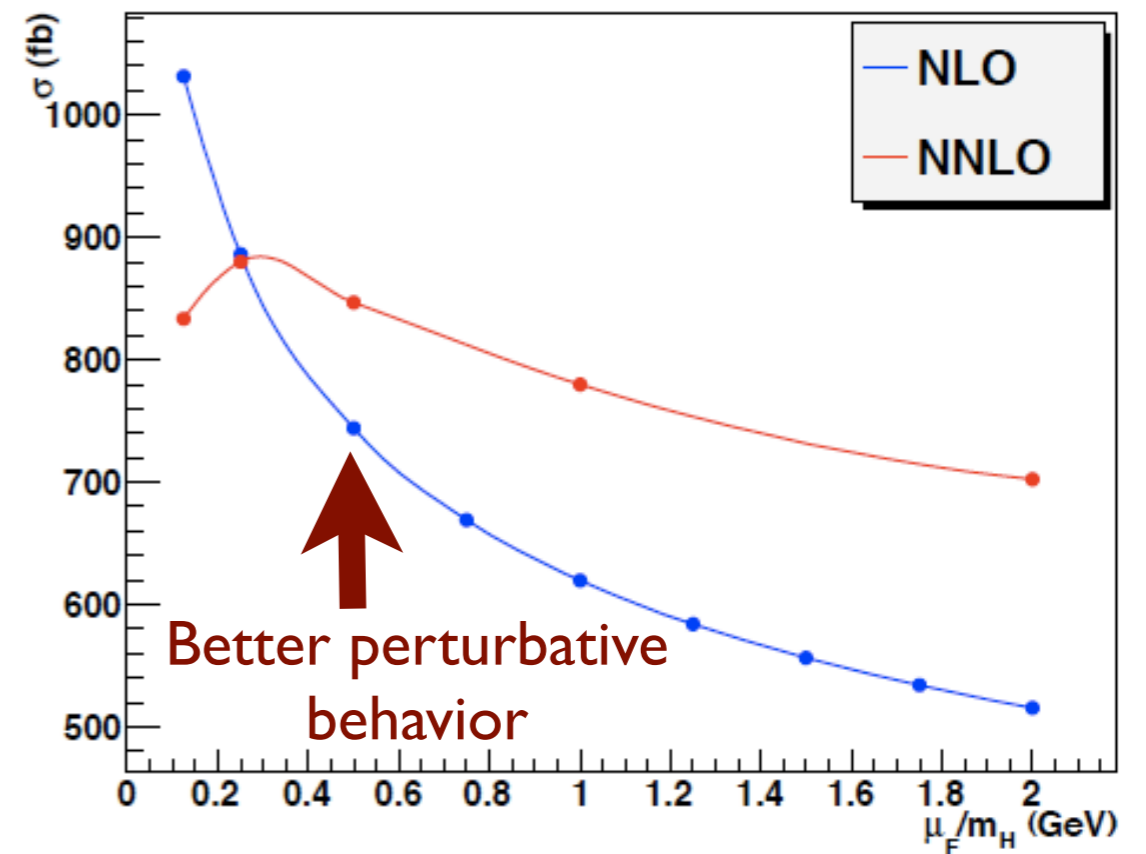


What scale?

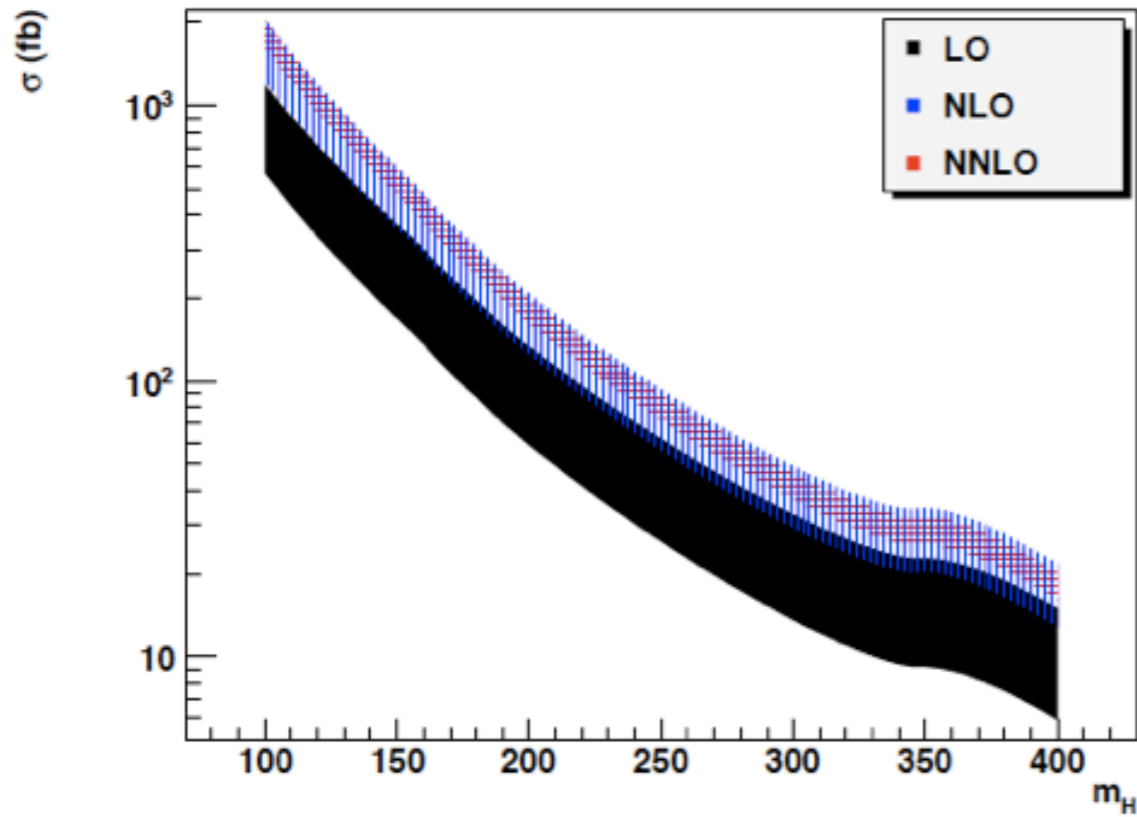
$\langle p_{T,H} \rangle$ (GeV) as a function of μ_F/m_H at $m_H=500\text{GeV}$ for LHC ($\sqrt{s}=7\text{TeV}$)



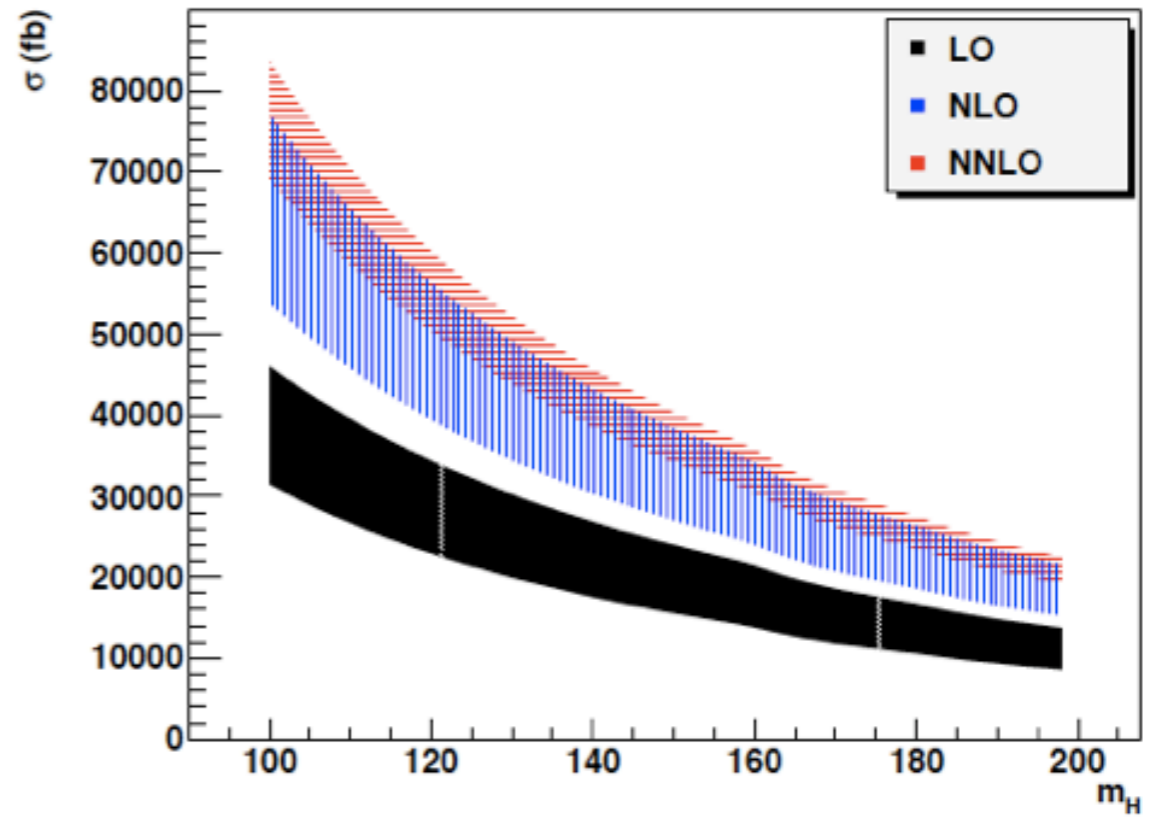
Total cross section as a function of μ_F/m_H at $m_H=500\text{GeV}$ for LHC ($\sqrt{s}=7\text{TeV}$)



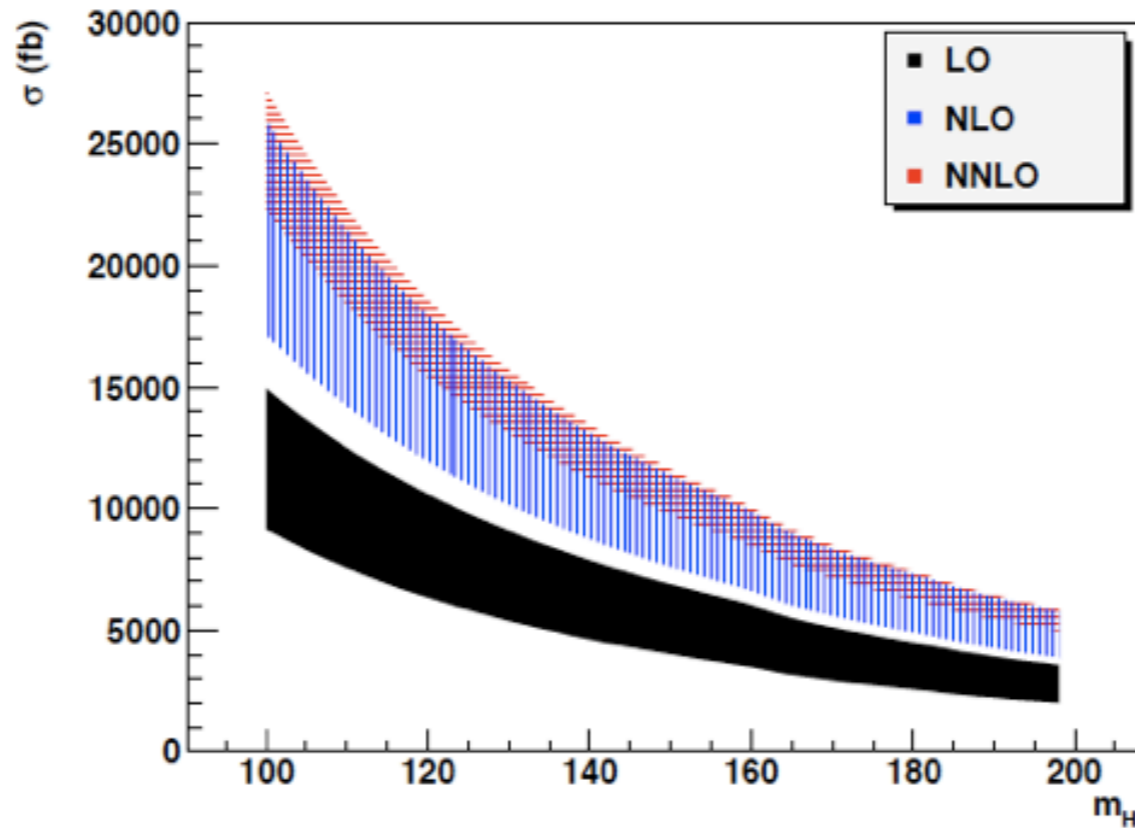
Total Inclusive cross section σ (fb) for Tevatron $\sqrt{s}=1.96$ TeV $\delta\mu$



Total Inclusive cross section σ (fb) for LHC $\sqrt{s}=14$ TeV $\delta\mu$



Total Inclusive cross section σ (fb) for LHC $\sqrt{s}=7$ TeV $\delta\mu$



Scale and PDF uncertainty correlations

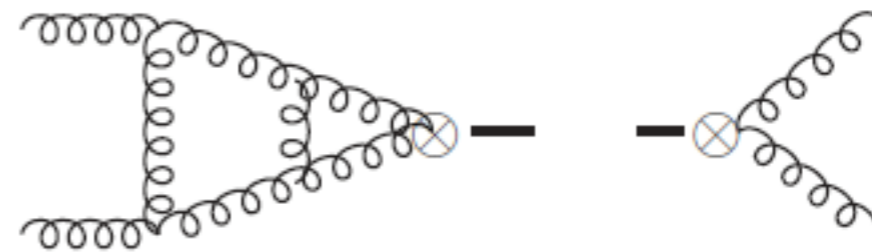
- Is the PDF error correlated with the choice of scale?
- Calculate PDF uncertainty with scales $[\mu = m_h/4, m_h]$
- Over a large mass range, for the Tevatron, the relative scale uncertainty $\frac{d\sigma(pdf, \mu)}{\sigma(pdf, \mu)}$ is constant to 0.3%

And now for something completely different: towards better NNLO calculations

In collaboration with C. Anastasiou and F. Herzog

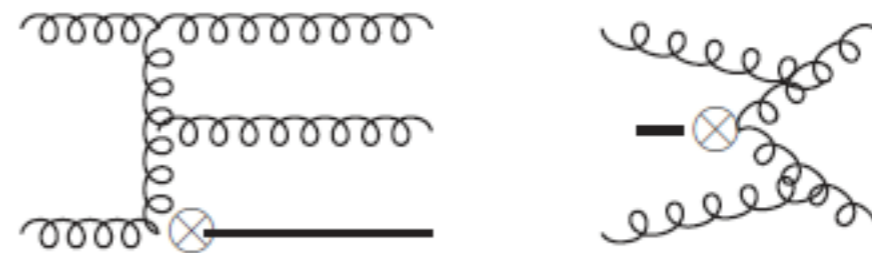
NNLO is different than NLO

Double virtual



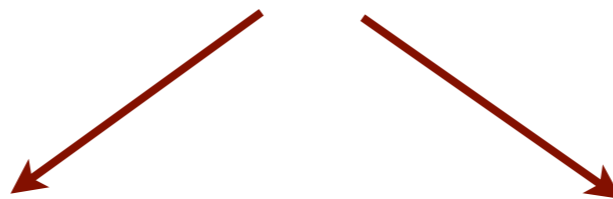
entangled singularities

Double real



The entangled singularities

$$I = \int_0^1 dx_1 dx_2 \frac{1}{(x_1 + cx_2)^{2-\epsilon}} \quad c > 0$$



$$x_1 > x_2$$

$$x_2 > x_1$$

$x_2 \rightarrow x_2 x_1$ (linear transformations)

$x_1 \rightarrow x_1 x_2$

$$I_1 = \int_0^1 \frac{dx_1 dx_2}{x_1^{1-\epsilon}} \frac{1}{(1 + cx_2)^{2-\epsilon}}$$

$$I_2 = \int_0^1 \frac{dx_1 dx_2}{x_2^{1-\epsilon}} \frac{1}{(x_1 + c)^{2-\epsilon}}$$

Disentangling them

$$x_2 = \frac{x_1 x'_2}{x_1 + (1 - x'_2)}$$



non-linear transformation
preserves the integration region

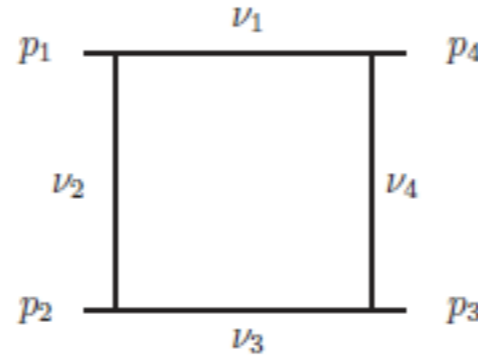
$$I_1 = \int_0^1 \frac{dx_1 dx_2}{\underline{x_1^{1-\epsilon}}} \frac{(1+x_1)(1-x_2+x_1)^\epsilon}{(1-x_2+x_1+cx_2)^{2-\epsilon}}$$

FINITE

Ready for plus-distribution expansion

$$x^{-1+\epsilon} = \frac{\delta(x)}{\epsilon} + \sum_{n=0}^{\infty} \frac{\epsilon^n}{n!} \left[\frac{\ln^n(x)}{x} \right]_+$$

The one loop box



$$I = \int_0^1 dx_1 \dots dx_4 \delta(1 - x_1 - \dots - x_4) f(x_1, \dots, x_4)$$

$$f(x_1, \dots, x_4) \equiv \frac{\Gamma(2 + \epsilon)}{[-sx_1x_3 - tx_2x_4 - i0]^{2+\epsilon}}$$

$$I = 2\Gamma(2 + \epsilon) \int_0^1 dy_1 dy_2 dy_3 (1 + y_1 + y_2 + y_3)^{2\epsilon} \\ \times \left\{ [-sy_1 - ty_2y_3]^{-2-\epsilon} + [-ty_1 - sy_2y_3]^{-2-\epsilon} \right\}$$

The one loop box

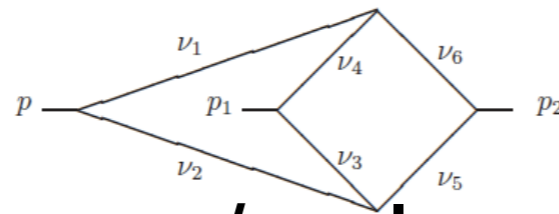
$$y_1 \rightarrow \frac{y_1 y_2 y_3}{1 - y_1 + y_2 y_3}$$

$$I = 2\Gamma(2 + \epsilon) \int_0^1 dy_1 dy_2 dy_3 \boxed{(y_2 y_3)^{-1-\epsilon}} (1 - y_1 + y_2 y_3)^{-\epsilon} [y_1 y_2 y_3 + (1 + y_2 + y_3)(1 - y_1 + y_2 y_3)]^{2\epsilon} \\ \times \left\{ [-s y_1 - t(1 - y_1) - t y_2 y_3]^{-2-\epsilon} + [-t y_1 - s(1 - y_1) - s y_2 y_3]^{-2-\epsilon} \right\} \quad (4.15)$$

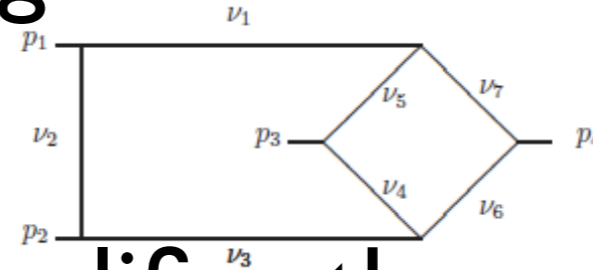
Singularities extracted!

The most complicated cases at 2-loops

- Massless non-planar integrals with all legs on-shell

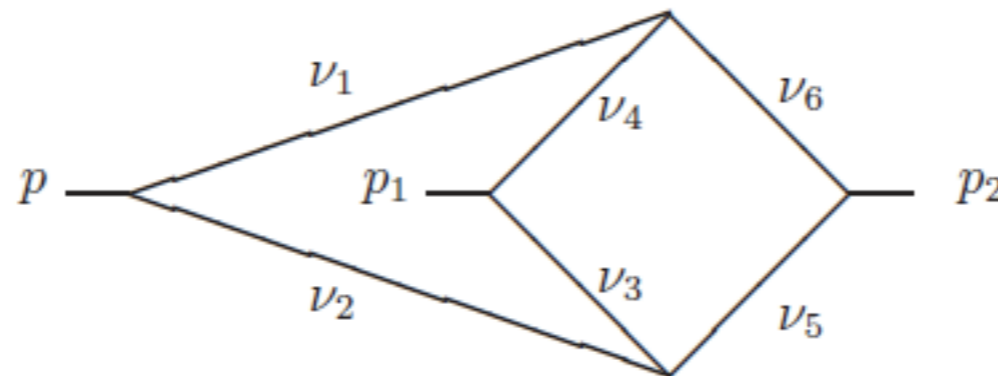


- Adding masses/scales simplifies the singularity structure.



- Threshold singularities that appear can be treated numerically with contour deformation.

The two loop Xtriangle



$$Xtri = 4^{2+2\epsilon} \int_0^1 dx_1 dx_2 dz dy dx \frac{zy^{1+\epsilon}(1-y)^{-1-\epsilon}(1-z)^{-1-\epsilon}}{[x(1-x) + yz(x-x_1)(x-x_2)]^{2+2\epsilon}}$$



Complicated singularity structure $x, y, z \rightarrow 0, 1$ etc.

Unavoidable split in 4 integrals to remove singularities at 1

The Xtriangle resolved

$$Xtri_{11} = 2^{6+9\epsilon} \int_0^1 dx_1 dx_2 dz dy dx \frac{zy^{1+\epsilon}(2-y)^{-1-\epsilon}(2-z)^{-1-\epsilon}}{[4x(2-x) + yz(2x_1-x)(2x_2-x)]^{2+2\epsilon}}$$

$$x \rightarrow \frac{xyzx_1x_2}{1-x+yzx_1x_2}$$



Trivial numerical evaluation
Number of integrals reduced by a factor of 16 with respect to sector decomposition

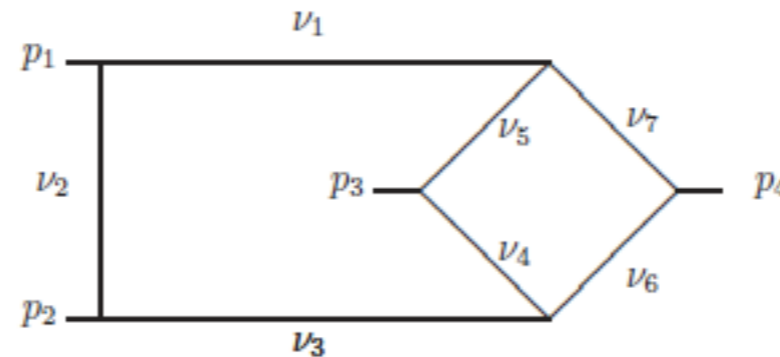
Identical procedure for the other integrals

$$Xtri_{12} = 2^{6+9\epsilon} \int_0^1 dx_1 dx_2 dz dy dx \frac{(2-z)y^{1+\epsilon}(2-y)^{-1-\epsilon}z^{-1-\epsilon}}{[4x(2-x) + y(2-z)(2x_1-x)(2x_2-x)]^{2+2\epsilon}} \quad (6.)$$

$$Xtri_{21} = 2^{6+9\epsilon} \int_0^1 dx_1 dx_2 dz dy dx \frac{z(2-y)^{1+\epsilon}y^{-1-\epsilon}(2-z)^{-1-\epsilon}}{[4x(2-x) + (2-y)z(2x_1-x)(2x_2-x)]^{2+2\epsilon}} \quad (6.)$$

$$Xtri_{22} = 2^{6+9\epsilon} \int_0^1 dx_1 dx_2 dz dy dx \frac{(2-z)(2-y)^{1+\epsilon}y^{-1-\epsilon}z^{-1-\epsilon}}{[4x(2-x) + (2-y)(2-z)(2x_1-x)(2x_2-x)]^{2+2\epsilon}}$$

The Xbox and alike



Same type of singularities

$$X_{box} = \Gamma(3 + 2\epsilon)(-s)^{-3-2\epsilon} \int \frac{dy dz dx_2 dx_4 [dx_{567}] y^{1+\epsilon} z(1-y)^{-1-\epsilon}(1-z)^{-1-\epsilon}}{[x_5 x_6 + zy(x_6 - x_2)(x_6 + x_7 - x_4) + zy(x_4 - x_2)x_7 t/s]^{3+2\epsilon}}$$

The number of integrals is reduced by a factor of 20 in comparison to sector decomposition.

Extending to double real

- We pursue an identical approach for the double real.
- The standard approach used in Fehip is sector decomposition and results in many sectors.
- The number of integrals is greatly reduced with our non-linear mappings.

Overview

Higgs in gluon fusion

- Developing the next version of FeHiPro for NNLO predictions of Higgs production at the fully differential level.
- Perturbative evidence that the correct central scale should be chosen at values lower than m_h
- Studies of average pT also suggests low central scale
- Studies of PDF and scale uncertainties show little correlation

Towards new NNLO

- New approach to disentangling singularities for NNLO double virtual and double real amplitudes.
- It avoids the proliferation of integrals.
- Very promising first results in the most difficult 2-loop cases.
- Similar treatment for the double real.
- Great optimism for better differential NNLO calculations.

What scale?

Tevatron

68% CL

$m_H (GeV)$	$\sigma (fb)$	$\delta_{PDF+a_s}^+ (\%)$	$\delta_{PDF+a_s}^- (\%)$	$\delta\mu^+ (\%)$	$\delta\mu^- (\%)$
100	1851.39	5.40	4.75	9.55	12.22
105	1605.82	5.47	4.84	9.37	12.16
150	536.96	6.28	5.70	8.37	11.90
155	482.56	6.38	5.80	8.30	11.88
160	431.60	6.49	5.90	8.21	11.87
165	379.16	6.60	6.01	8.18	11.88
170	338.92	6.71	6.11	8.17	11.90
175	305.40	6.81	6.21	8.13	11.91
180	275.91	6.93	6.31	8.11	11.92