

# Bound-state effects in $t\bar{t}b\bar{b}$ production at the LHC

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based on works in collaboration with  
K.Hagiwara(KEK) and Y.Sumino(Tohoku U)

GGI workshop, Firenze, 2011.09.28



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Outline :

1. Introduction : Top-Quark at hadron colliders
2. Bound-state effects on  $t\bar{t}$  production  
at hadron colliders
3. Differential cross-section / Event Generation
4. Summary



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# 1. Introduction

# 1. Introduction

## □ Properties of the top-quark

- **Mass :** (CDF and D0 combined, [arXiv:1107.5255](https://arxiv.org/abs/1107.5255))

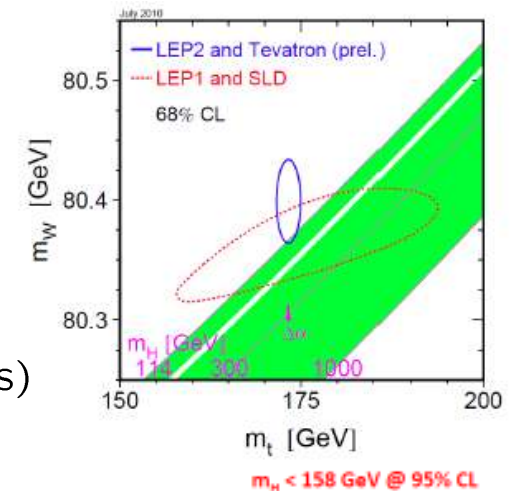
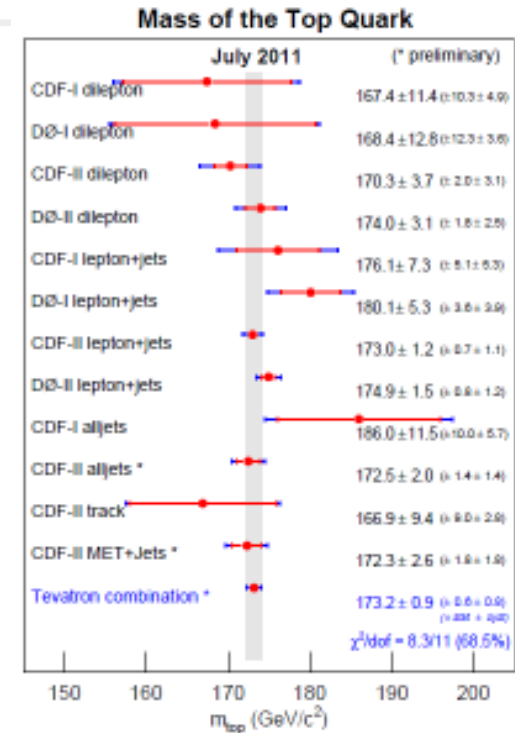
$$m_t = 173.2 \pm 0.6(\text{stat.}) \pm 0.8(\text{syst.}) [\text{GeV}]$$

- Heaviest fundamental particle ever found
- Important input for EW precision test
- $y_t \sim 1 \Rightarrow$  may be related to the BSM physics

- **Width :** SM prediction;  $\Gamma_t \simeq \frac{G_F m_t^3}{8\sqrt{2}\pi} |V_{tb}|^2 \sim 1.5 [\text{GeV}]$

$$\text{D0 measurement ('10)}; \Gamma_t = 1.99^{+0.69}_{-0.55} \text{ GeV}$$

- decay before hadronization,  $\Gamma_t \gg \Lambda_{\text{QCD}}$  ( $\tau \sim 10^{-25}\text{s}$ )  
thus spin information is preserved in decay products



# 1. Introduction

- Precise mass (and width) determination :

What mass of the top-quark is observed ?  
pole mass? a parameter in Monte-Carlo?

**Fact** : Pole mass is not well-defined.

Theoretical prediction using pole mass has bad pert. convergency

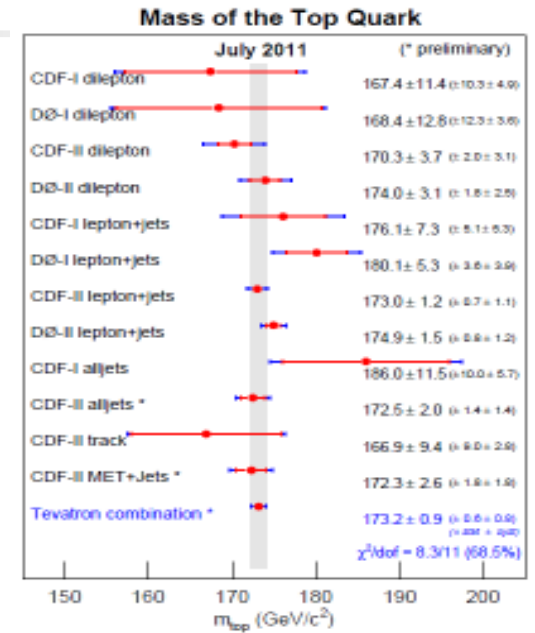
**Important** : definition of the mass in an IR-safe manner  
(short-distance mass)

Langenfeld, Moch, Uwer('09) determination of the  $\overline{MS}$  mass from the total cross-section at Tevatron.

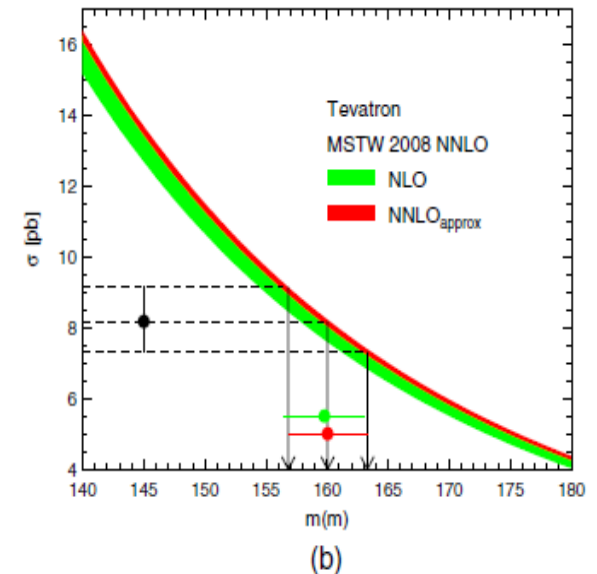
$$\bar{m}(\bar{m}) = 160.0^{+3.3}_{-3.2} \text{ [GeV]} \quad (\sim \text{NNLO})$$

→ pole mass  $m_t = 168.9 \pm 3.5 \text{ GeV}$

Another choice may be "threshold mass"



PHYSICAL REVIEW D **80**, 054009 (2009)



# 1. Introduction

- In  $e^+e^-$  colliders, "Threshold Scan" can be performed

precise determinations of the top-quark mass, width and strong coupling constant are possible

$$\delta m_t \sim 100 \text{ [MeV]} \quad \delta \Gamma_t / \Gamma_t \sim 20 \text{ [%]}$$

- Peak of the 1S resonance  $\Leftrightarrow$  Threshold mass (1S mass, PS mass, ...)

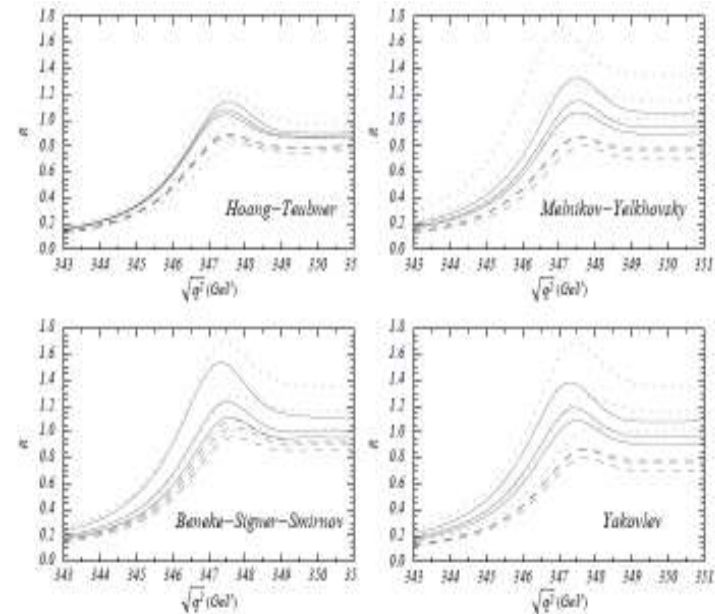
Relation between the threshold mass,  $\overline{\text{MS}}$  mass and pole mass is well-known to higher-orders

- However, at hadron colliders, (partonic) collision energy is not fixed, so one has to reconstruct top-pair invariant mass

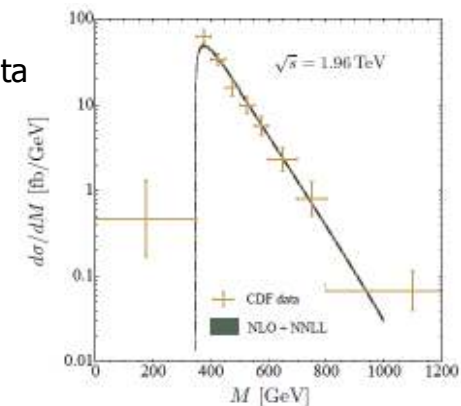
It would be also a challenging task to have good precision at the LHC

(p,v)NRQCD up to  $N^{\text{NLL}}/N^{\text{NLO}}$

hep-ph/0001286



CDF data

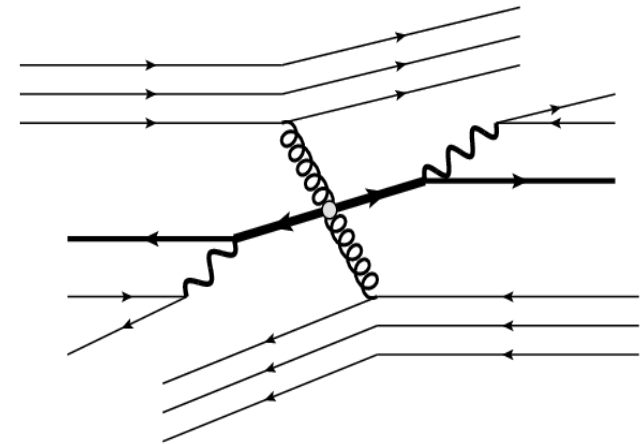


# 1. Introduction

- Top-quark production at Hadron Colliders

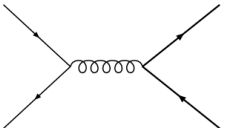
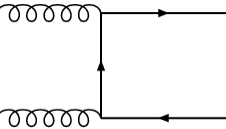
$$d\sigma_{t\bar{t}}(s) = \sum_i \int d\tau \frac{dL_i}{d\tau}(\tau) d\hat{\sigma}_i(\hat{s} = \tau s)$$

partonic luminosity      partonic cross-section

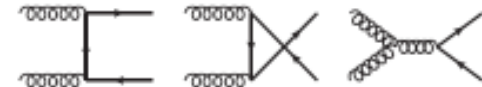


# 1. Introduction

- Partonic subprocess at Hadron Colliders

			Tevatron	LHC
$q\bar{q} \rightarrow t\bar{t}$		Color: <b>Octet</b> $ J =1$	85%	10%
			..	..
$gg \rightarrow t\bar{t}$		Color: <b>Singlet, Octet</b> $ J =0,1,2,...$	15%	90%

- Color decomposition in gg process is obtained from the color-matrix in the amplitude



$$\begin{aligned} \mathcal{M}_{gg \rightarrow t\bar{t}} &= (T^a T^b)_{ij} \mathcal{M}_1 + (T^b T^a)_{ij} \mathcal{M}_2 \\ &= \frac{1}{2} \{T^a, T^b\}_{ij} \mathcal{M}_S + \frac{1}{2} [T^a, T^b]_{ij} \mathcal{M}_A \end{aligned}$$

$$\frac{1}{2} \{T^a, T^b\} = \frac{1}{2N_c} \delta^{ab} \delta_{ij} + \frac{1}{2} d^{abc} T_{ij}^c$$

color-singlet

$$\left| \frac{1}{2N_c} \delta_{ab} \delta_{ij} \right|^2 / \left| \frac{1}{2} d^{abc} T_{ij}^c \right|^2 = \frac{2}{N_c^2 - 4}$$

~20-30% of ttbar are color-singlet in gg fusion





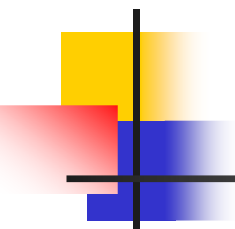
# 1. Introduction

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- Perturbative calculations for the **partonic cross-section** :

$$\hat{\sigma}_i(\hat{s}) = \hat{\sigma}_i^{(0)}(\hat{s}) + \frac{\alpha_s}{\pi} \hat{\sigma}_{ij}^{(1)}(\hat{s}; \mu) + \left(\frac{\alpha_s}{\pi}\right)^2 \hat{\sigma}_i^{(2)}(\hat{s}; \mu) + \dots$$

- Fixed-Order Calculation** :
- **NLO** : Dawson, Ellis, Nason('88), Beenakker et al.('90), Mangano, Nason, Ridolfi(92)  
(**analytic**) : Cazkon, Mitov('08)
  - Towards **Full NNLO** correction :  
Korner et al.('06), Dittmaier et al.('07), Cazkon et al.('07),,,
  - **ElectroWeak** correction : Bernreuther, Fuecker, Si('06),  
Kuhn, Scharf, Uwer('06), Moretti, Nolten, Ross('06)
- Resummation** :
- **Threshold Resummation**  
NLL : Bonciani et al.('98),,, ;  
NNLL : Moch, Uwer('08), Ahrens et al ('10), Beneke et al ('11)
  - **Coulomb Summation** : Catani, Mangano, Nason, Trentadue('96),  
Hagiwara, Sumino, HY('08), Kiyo et al('08), Sumino, HY('10)



# TOP 2011

**4th International Workshop on  
TOP QUARK PHYSICS**

**September 25-30, 2011**

**Sant Feliu de Guixols, Costa Brava, SPAIN**



## NLO predictions --- Theoretical accuracy



$m_t = 173.1 \text{ GeV}/c^2, m_t/2 < \mu < 2m_t$  MSTW2008 PDF(90%cl)

	Tevatron	LHC (7 TeV)
$\sigma_{\text{LO}}$ (pb)	$6.66^{+2.95+0.34}_{-1.87-0.27}$	$122^{+49+6}_{-32-7}$
$\sigma_{\text{NLO}}$	$6.72^{+0.36+0.37}_{-0.76-0.24}$	$159^{+20+8}_{-21-9}$

$$\text{NLO} \quad \frac{\Delta\sigma}{\sigma} \approx \begin{cases} +5\% \text{ (scale)} \quad +5\% \text{ (pdf) Tevatron} \\ -10\% \text{ (scale)} \quad -4\% \text{ (pdf) Tevatron} \\ +12\% \text{ (scale)} \quad \pm 5\% \text{ (pdf) LHC} \\ -13\% \text{ (scale)} \quad \pm 5\% \text{ (pdf) LHC} \end{cases}$$

## Top-quark pair production in NLO QCD



Experimental accuracy below 10 %

→ need to go beyond NLO accuracy

Possible corrections (percent level):

QCD NNLO

$$\sim \alpha_s^4$$

Bound state  
effects

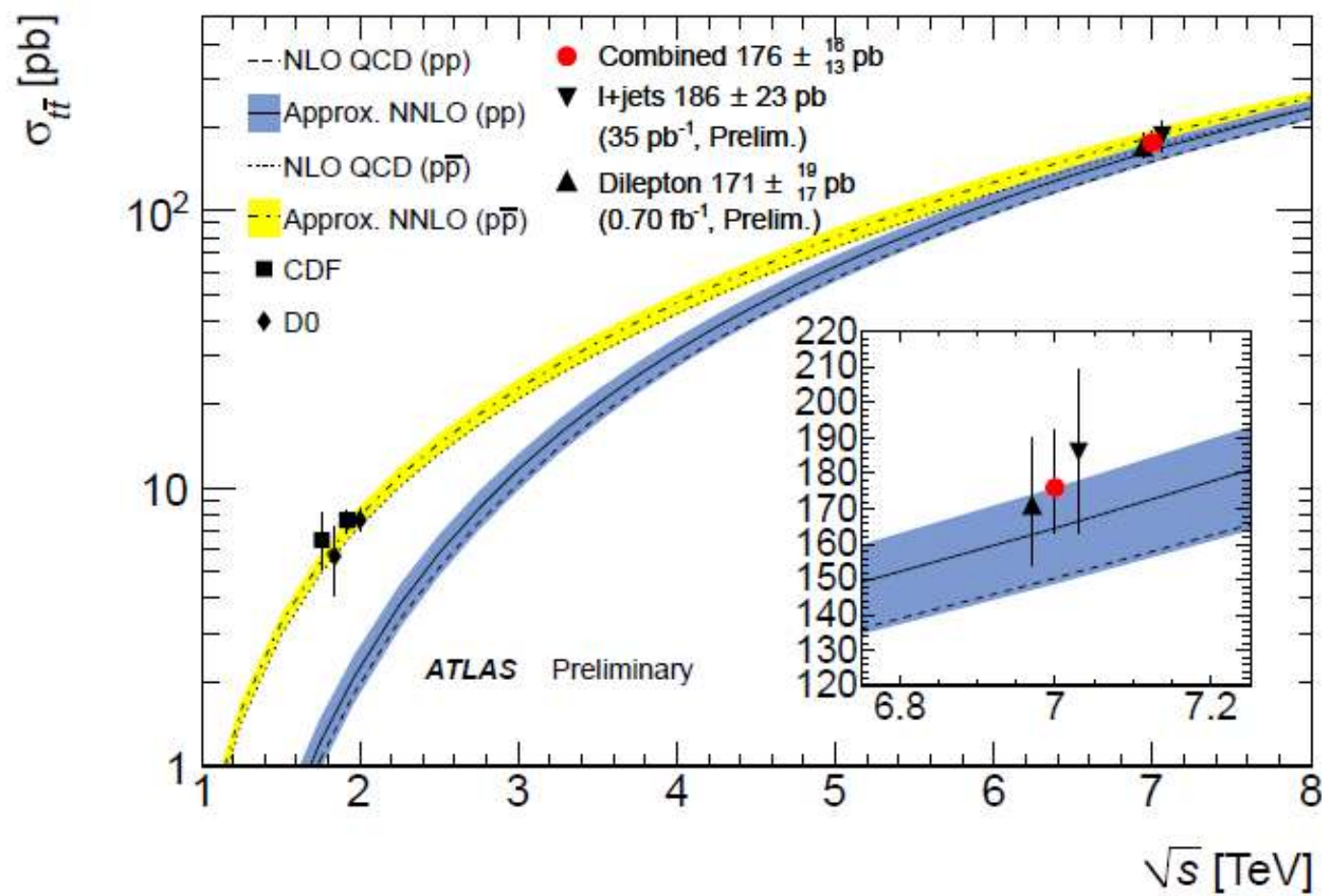
$$\sim (\alpha_s^3 \frac{1}{\beta}) \beta$$

Mixed weak-QCD  
corrections

$$\sim \alpha_s^2 \alpha$$

finite width

$$\Gamma_t / m_t$$





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## 2. Bound-state effects on $t\bar{t}$ production at Hadron Colliders

## 2. Bound-state effects at Hadron colliders

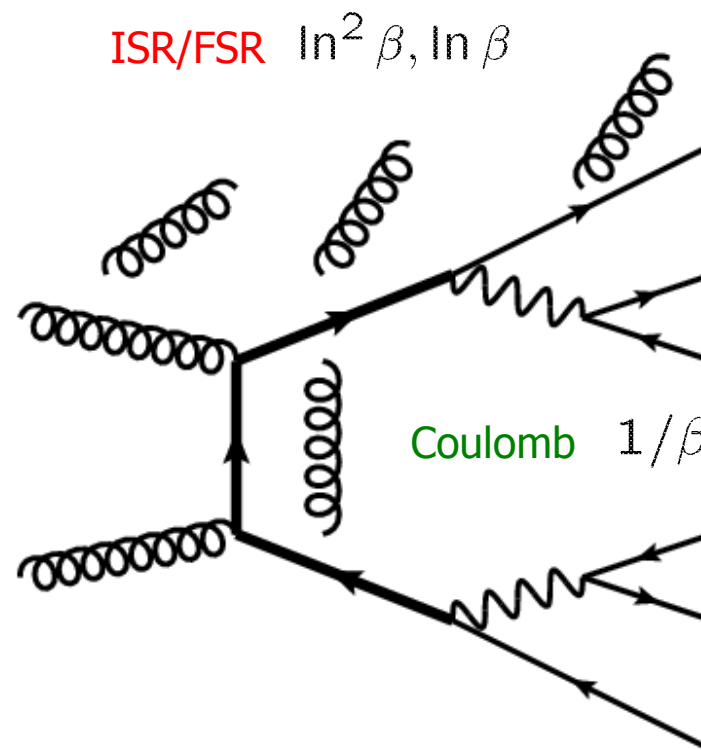
- QCD higher-order correction Nason, Dawson, Ellis

$$\hat{\sigma}_i^{(c), \text{NLO}} \sim \hat{\sigma}_i^{(c), \text{LO}} \frac{\alpha_s}{\pi} \left\{ A_i \ln^2(8\beta^2) + B_i^{(c)} \ln(8\beta^2) + C_i^{(c)} \frac{\pi^2}{\beta} + D_i^{(c)} + \mathcal{O}(\beta) \right\}$$

**Threshold logs** : emission of soft and/or collinear gluon in initial-state and final-state

**Coulomb singularity** :  
Coulomb gluon exchange between  $t$  and  $t$ -bar

**Hard correction** :  
process dependent virtual corr.



## 2. Bound-state effects at Hadron colliders

□ Factorization of each contribution :

Beneke,Falgari,Schwinn ('09)

⇒ Resummation of each correction separately

$$\begin{aligned}\hat{\sigma}_i^{\text{NLO}} &\sim \hat{\sigma}_i^{\text{LO}} \left[ 1 + \frac{\alpha_s}{\pi} \left( A_i \ln^2 \beta^2 + B_i^{(e)} \ln \beta^2 + C_i^{(e)} \frac{\pi^2}{\beta} + D_i^{(e)} + \mathcal{O}(\beta) \right) \right] \\ &\sim \hat{\sigma}_i^{\text{LO}} \left[ 1 + \frac{\alpha_s}{\pi} \left( A_i \ln^2 \beta^2 + B_i^{(e)} \ln \beta^2 \right) \right] \left[ 1 + \frac{\alpha_s}{\pi} C_i^{(e)} \frac{\pi^2}{\beta} \right] \left[ 1 + \frac{\alpha_s}{\pi} D_i^{(e)} \right] + \mathcal{O}(\beta)\end{aligned}$$

- **Threshold resummation**

Laenen,Smith,van Neerven('92),

Kidonakis Sterman('97),Bonciani et al.('98)

Moch,Uwer('08), Caccari et al('08),

Ahrens et al ('10), Beneke et al('11),,,

- **Coulomb summation,  
or Bound-state formation**

Fadin,Khoze,Sjostrand('90)

Catani,Mangano,Nason,Trentadue('96)

Hagiwara,Sumino,HY('08)

Kiyo,Kuhn,Moch,Steinhauser('09)

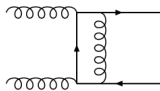
Sumino,HY('10)



## 2. Bound-state effects at Hadron colliders

### □ Coulomb corrections to all-orders

- Coulomb singularity



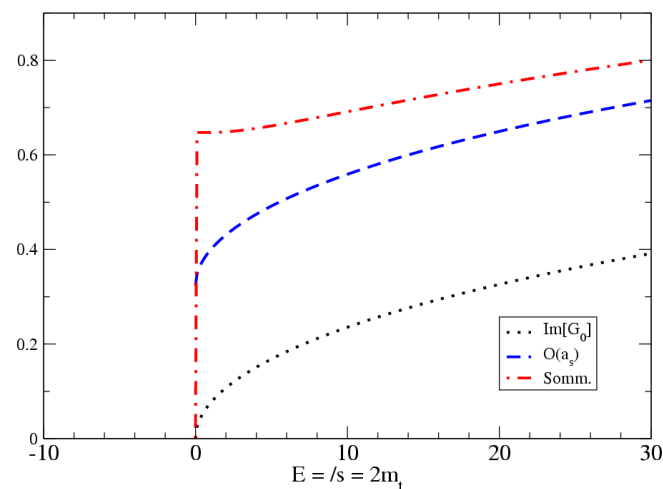
$$\propto C^{(c)} \frac{\alpha_s}{\beta} \quad \mathcal{O}(1) \text{ for } \beta \simeq \alpha_s$$

$$\text{color-factor} \begin{cases} \text{singlet} & C^{(1)} = -C_F \\ \text{octet} & C^{(8)} = C_A/2 - C_F \end{cases}$$

- Summation of ladder diagrams = Sommerfeld factor Sommerfeld, Sakharov (QED)

$$S(z) = \frac{z}{1 - \exp[z]}, \quad z = C^{(c)} \pi \alpha_s / \beta$$

$$= \text{tree} + \alpha_s / \beta \text{ (ladder)} + (\alpha_s / \beta)^2 \text{ (ladder)} + \dots$$



## 2. Bound-state effects at Hadron colliders

### □ Coulomb corrections to all-orders

- **Green's function formalism (NRQCD)** Fadin,Khoze('87),...

Schrodinger's Eq. 
$$\left[ (E + i\Gamma_t) - \left\{ -\frac{\nabla^2}{m_t} + V_{QCD}^{(c)}(r) \right\} \right] G^{(c)}(E, \vec{x}) = \delta^3(\vec{x})$$

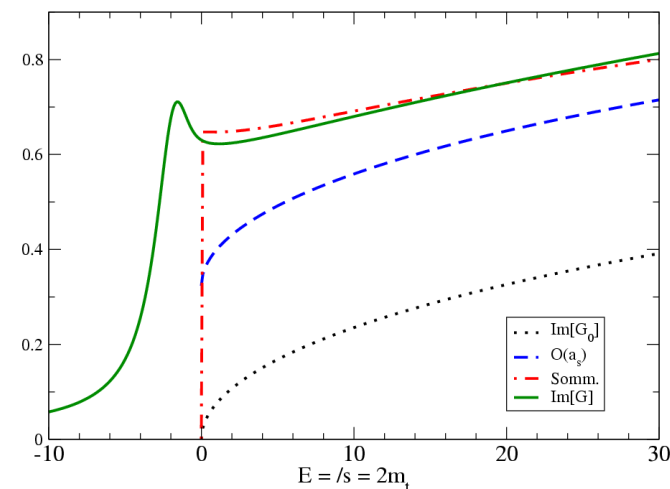


where  $E = m_{tt} - 2m_t$

**finite width effect** is incorporated by complex energy  
 $\Rightarrow$  Off-shellness of top-quarks

$$G(E, \vec{x}) = \sum_n \frac{\Psi_n(\vec{x}) \Psi_n^*(0)}{E - E_n + i\Gamma_n/2} + \text{continuum}$$

- Large width smears the multiple resonance structure, but only one broad peak can be seen as a remnant.



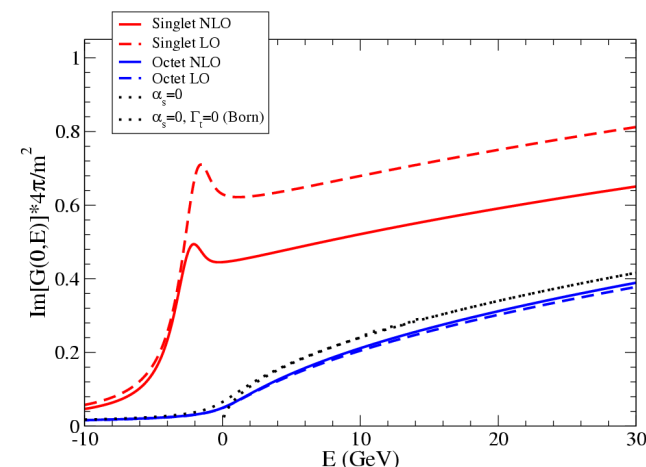
## 2. Bound-state effects at Hadron colliders

- Perturbative QCD potential (NLO), since an IR cut-off by  $r \lesssim \frac{1}{\Gamma_t}$

$$V_{\text{QCD}}^{(c)}(r) = C^{(c)} \frac{\alpha_s(\mu_B)}{r} \times \left[ 1 + \frac{\alpha_s}{\pi} v_1^{(c)}(r) + \dots \right]$$

$$\begin{cases} \text{singlet} & C^{(1)} = -4/3 \\ \text{octet} & C^{(8)} = 1/6 \end{cases}$$

Singlet is attractive,  
but octet is repulsive and small color-factor



- Toponium system :  $m_t \gg \mu_B > E_B \simeq \Gamma_t \gg \Lambda_{\text{QCD}}$

- Binding energy :  $E_B \simeq m_t \alpha_s^2 \simeq 2\text{GeV}$

If  $\Gamma_t > E_B$ , top-quark decays  
before bound-state formation

- Bohr radius :  $\mu_B \simeq m_t \alpha_s \simeq 20 - 30\text{GeV}$

(typical momentum of the Coulomb gluon)

## 2. Bound-state effects at Hadron colliders

- Cross-section is proportional to the Imaginary part of the Green function by Optical theorem.

$$\hat{\sigma}_{tt}^{(c)} \rightarrow \hat{\sigma}_{tt,\text{Born}}^{(c)} \cdot \text{Im}[G^{(c)}(E, \vec{x} = \vec{0})]$$

- Combining Initial-state/Final-state radiation effects,

$$\frac{d\sigma}{dm_{tt}}(s, m_{tt}^2) = \hat{\sigma}_i^{(c)}(m_{tt}^2) \cdot K_i^{(c)} \int_{\tau_0}^1 \frac{dz}{z} F_i^{(c)}(z) \frac{d\mathcal{L}_i}{d\tau}(\tau_0/z)$$

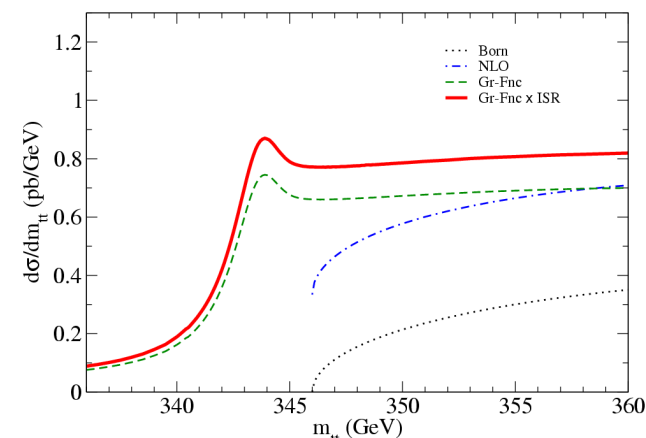
$$F_i^{(c)}(z) = \delta(1-z) + \frac{\alpha_s}{\pi} \left[ A_i \left\{ \left( \frac{\ln(1-z)}{1-z} \right)_+ - \left( \frac{1}{1-z} \right)_+ \ln \left( \frac{\mu_F}{2m_t} \right) \right\} + D_{tt}^{(c)} \left( \frac{1}{1-z} \right)_+ + k_i^{(c)} \delta(1-z) \right]$$

- Color-dependent **hard correction** :  $K_i^{(c)} = 1 + \frac{\alpha_s}{\pi} h_i^{(c)}$   
extracted from the Quarkonium cross-section at NLO

Petrelli,Cacciari,Greco,Maltoni,Mangano ('98)

+Non-decoupling term HSY('08), Czakon,Mitov('08)

gg->tt, color-singlet at the LHC



NLL : Kiyo et al. ('08)

$$h_{gg}^{(1)}\left(\frac{\mu_R}{m_t}\right) = C_A \left(1 + \frac{\pi^2}{12}\right) + C_F \left(-5 + \frac{\pi^2}{4}\right) + \beta_0 \ln\left(\frac{\mu_R}{2m_t}\right),$$

$$h_{gg}^{(8)}\left(\frac{\mu_R}{m_t}\right) = C_A \left(3 - \frac{\pi^2}{24}\right) + C_F \left(-5 + \frac{\pi^2}{4}\right) + \beta_0 \ln\left(\frac{\mu_R}{2m_t}\right),$$

$$h_{q\bar{q}}^{(8)}\left(\frac{\mu_R}{m_t}\right) = C_A \left(\frac{59}{9} - \frac{\pi^2}{4} + \frac{2 \ln 2}{3}\right) + C_F \left(-8 + \frac{\pi^2}{3}\right) - \frac{5}{9} n_q - \frac{8}{9} + \beta_0 \ln\left(\frac{\mu_R}{2m_t}\right).$$

## 2. Bound-state effects at Hadron colliders

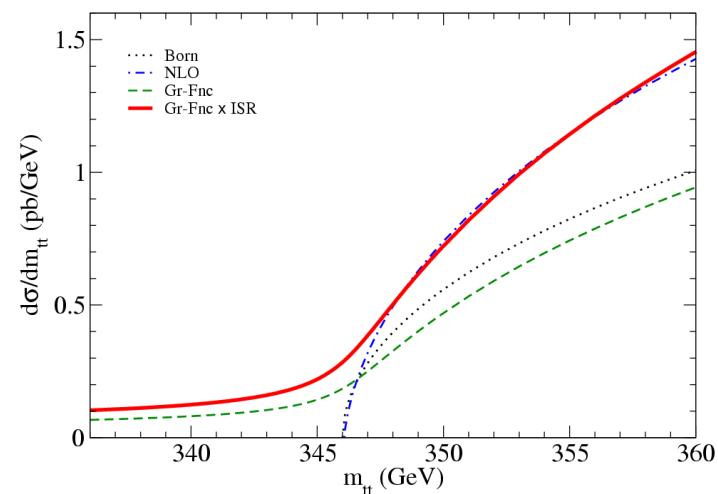
### □ $t\bar{t}$ invariant-mass distributions

Black : Born  
 Blue :  $O(\alpha_s)$  corr. (NLO)  
 Green : Gr-Fnc. without ISR  
 Red : Gr-Fnc. with ISR

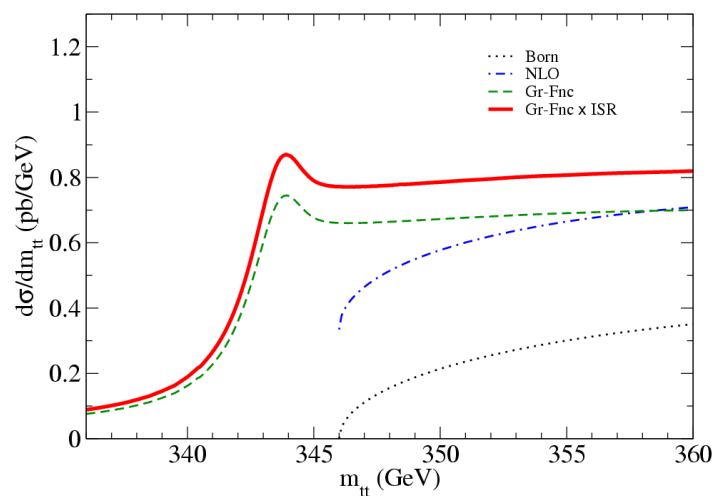
ISR : up to  $O(\alpha_s)$  (soft/collinear)

$m_t=173$  GeV,  $\Gamma_t=1.5$  GeV, CTEQ6M

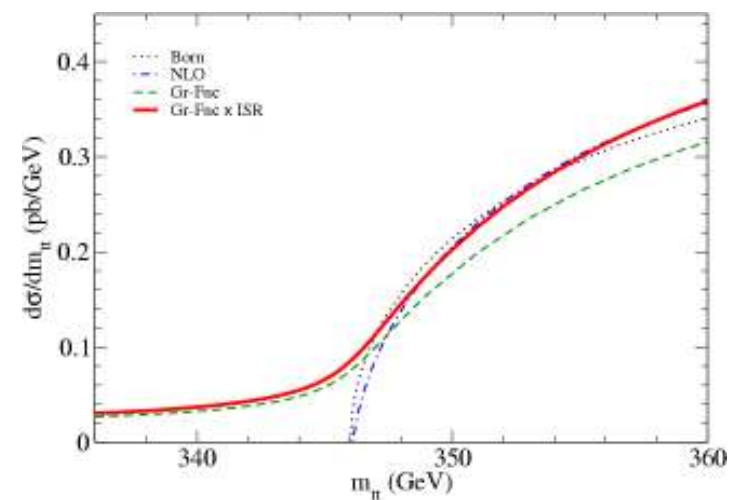
gg->tt, color-octet



gg->tt, color-singlet



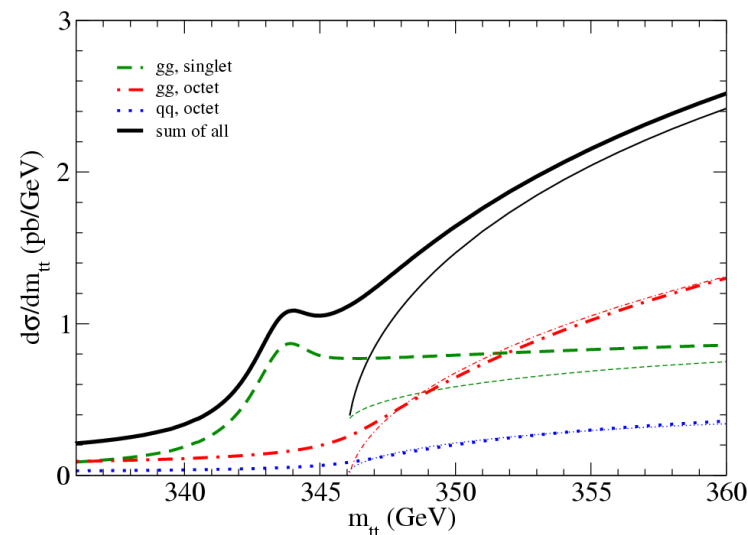
qq->tt, color-octet



## 2. Bound-state effects at Hadron colliders

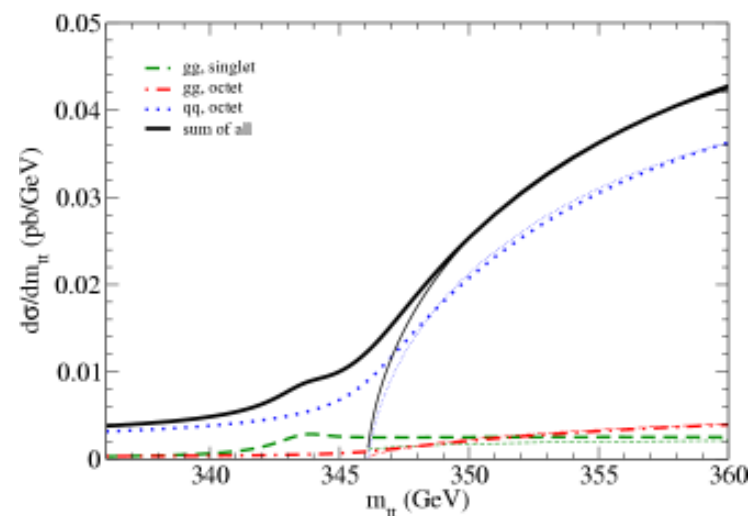
### ◆ In total at the LHC :

- BS effects **deform the invariant-mass distribution** near threshold
- Form **a broad resonance peak** below threshold (observable in principal)
- Enhance the total cross-section by **10 pb  $\sim O(1\%)$**

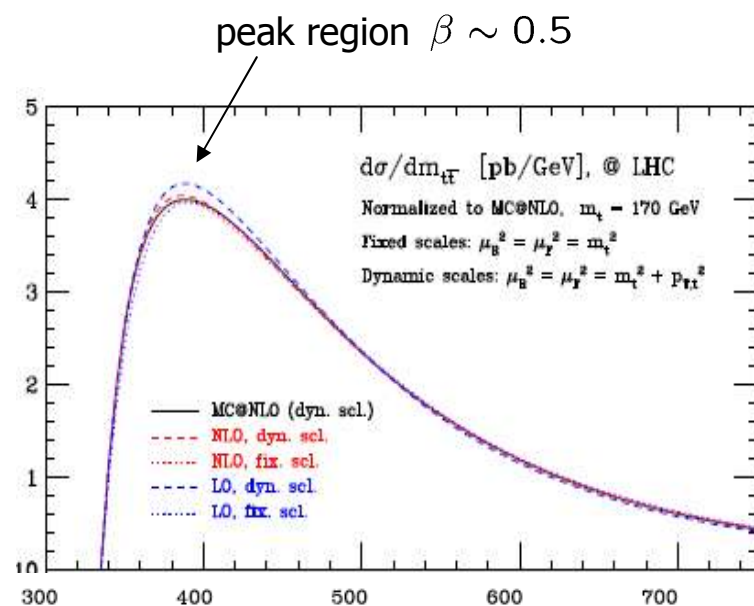


### ◆ On the other hand at the Tevatron :

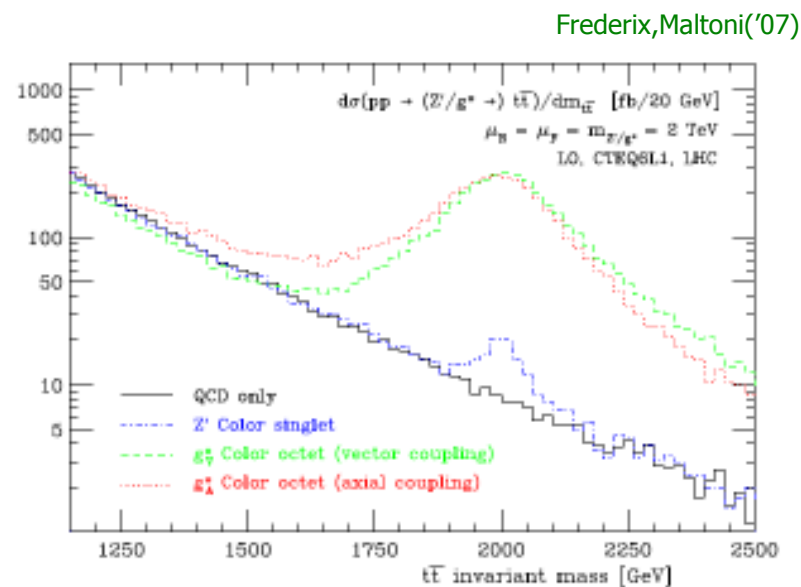
- Resonance can't be seen due to the color-octet dominance
- Correction is small due to the small color-factor.



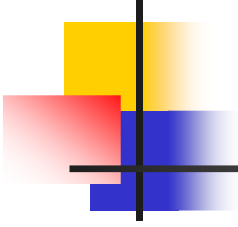
- ttbar Invariant-mass distribution :  $m_{t\bar{t}}^2 = (p_t + p_{\bar{t}})^2$



↑  
 threshold region  $\beta \leq 0.3$  ( $m_{t\bar{t}} < 360$  GeV)



High-energy region  $\beta \sim 1$   
 (Boosted top, **NP resonances**, FB asymmetries?,,,)



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### 3. Differential Cross-section / Event Generation with BS effects



### 3. Differential Cross-section

#### □ Coulomb correction in differential cross-sections

- Differential distributions are useful for the analysis with kinematical cuts
- Full differential distribution is required in Event-Generation by Monte-Carlo simulation
- Method to include BS effects is developed in  $e^+e^-$  collider study

Jezabek,Kuhn,Teubner('92)

Sumino,Fujii,Hagiwara,Murayama,Ng('93),,,

- Our Aim : take into account the "Leading-order" contribution in both region :

- **Threshold region** :  $(\alpha_s/\beta)^n$  but not  $\alpha_s^n, \beta^n$

Sumino,HY ('10)

- **High-energy region** :  $\beta^n$  but not  $\alpha_s^n$

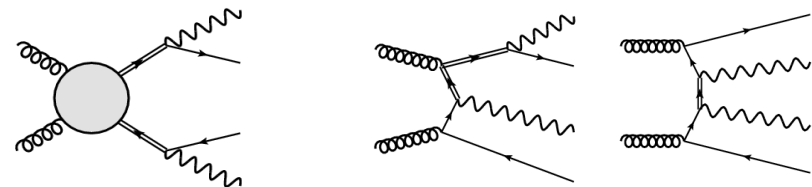
- plus some NLO effects (K-factor, width, QCD potential,,)      note,  $\Gamma_t/m_t \sim \alpha_W \sim \alpha_s^2$

### 3. Differential Cross-section

#### □ Coulomb correction in differential cross-sections

- $gg/qq$  to  $bWbW$  process to take into account the off-shellness of top-quarks
- Divide **resonant part** and **non-resonant part** of the amplitude (gauge dependent)

$$\mathcal{M}^{(c)}(I \rightarrow bWbW) = \mathcal{M}_{t\bar{t}}^{(c)} + \mathcal{M}_{nr}^{(c)}$$



- Double resonant part of the amplitude :

$$\mathcal{M}_{t\bar{t}} = \bar{D}_{t \rightarrow bW} \cdot \frac{i}{\not{p}_t - m_t + i\Gamma_t} \cdot P_{i \rightarrow t\bar{t}} \cdot \frac{i}{-\not{p}_{\bar{t}} - m_t + i\Gamma_t} \cdot D_{\bar{t} \rightarrow \bar{b}W}$$

### 3. Differential Cross-section

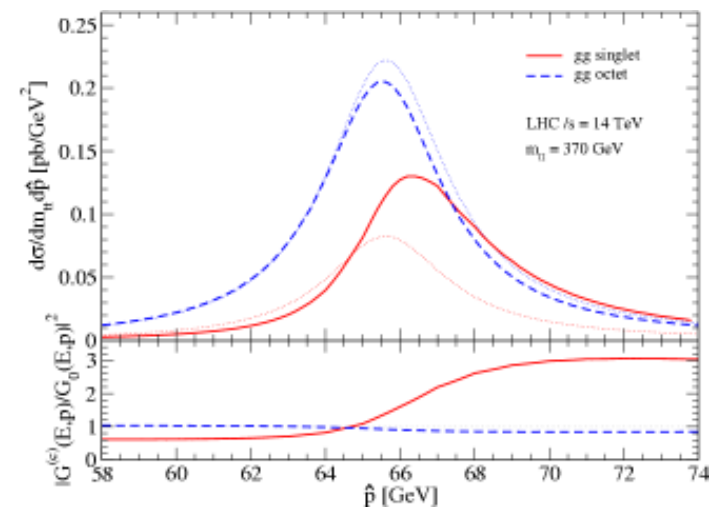
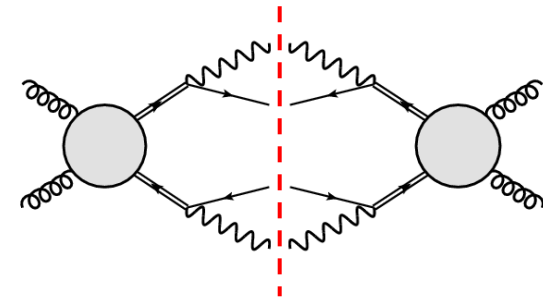
- A correction factor to the resonant part by the **Momentum-Space Green Functions**

$$\mathcal{M}_{t\bar{t}}^{(c)} \rightarrow \mathcal{M}_{t\bar{t}}^{(c)} \times \frac{\tilde{G}^{(c)}(E + i\Gamma_t, \vec{p})}{\tilde{G}_0(E + i\Gamma_t, \vec{p})}$$

(for S-wave only  $|\vec{p}|$  dependence)

- Top-quark momentum distribution (of the  $t\bar{t}$  cm frame) is affected by the Coulomb correction.

$$\frac{d\sigma}{d|\vec{p}|} \propto |\tilde{G}(E, \vec{p})|^2 \quad \begin{array}{l} \text{color-singlet : } \delta p > 0 \\ \text{color-octet : } \delta p < 0 \end{array}$$



### 3. Event Generation

“Toponium@LHC” Sumino, HY ('10)

□ Event Generator (LO + all-order Coulomb) :

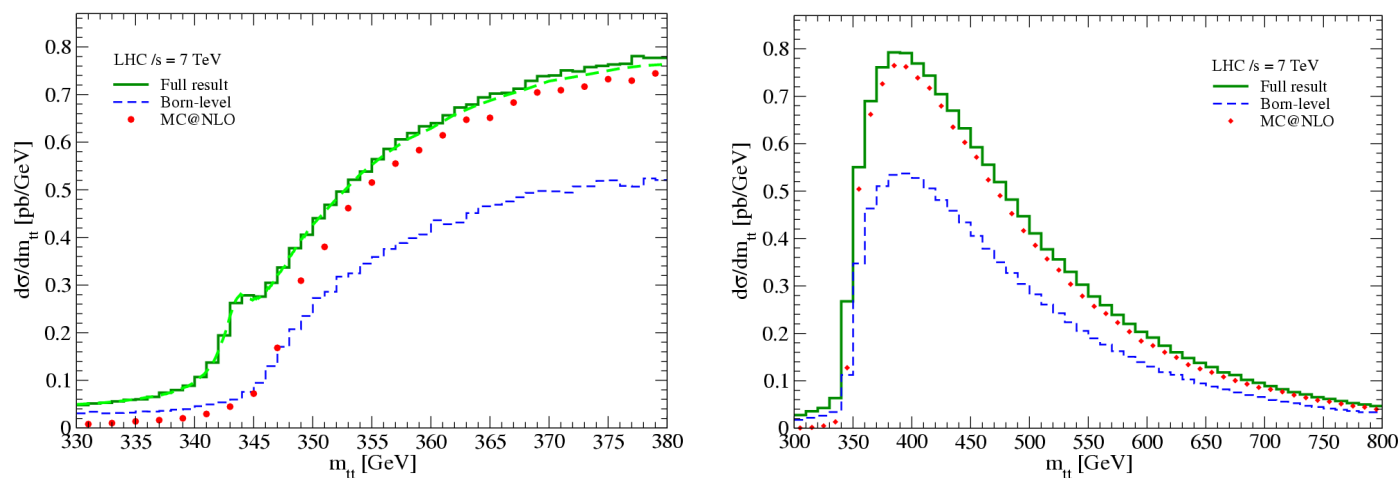
<http://madgraph.kek.jp/~yokoya/TopBS>

- Full  $gg/qq \rightarrow bWbW$  plus W-decays Matrix-Elements (6-bodys)
- Color-decomposition in  $gg \rightarrow bWbW$  process
- Bound-state correction to the double-resonant amplitudes
- Color-dependent K-factors to reproduce NLO  $m_{tt}$  dist. near threshold
- Difference from General-purpose Monte-Carlo's :
  - MadGraph/MadEvent, Sherpa,,, (PYTHIA, HERWIG,,,)
  - LO(Tree-level), ○ non-resonant effects, off-shell effect,,,
  - MCFM, MC@NLO,,,
  - NLO ( $\alpha_s/\beta$  term), × non-resonant effects, Breit-Wigner,,,

### 3. Event Generation

#### □ Some Examples (at partonic-level)

(1)  $t\bar{t}$  invariant-mass ( $m_{t\bar{t}}$ ) distribution :



- The only generator which describes the threshold enhancement and resonance
- Effectively, well reproduce MC@NLO results at large  $m_{t\bar{t}}$  by taking the scales as

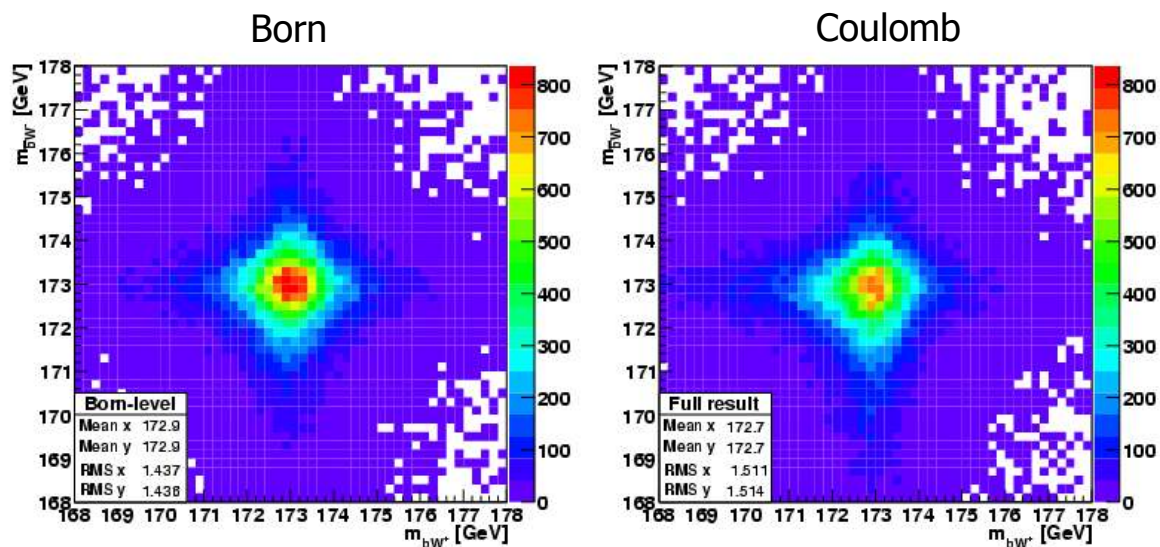
$$\mu = m_t \quad (\mu = \sqrt{m_t^2 + p_T^2} \quad \text{in MC@NLO})$$

### 3. Event Generation

#### □ Some Examples (at partonic-level)

(2) (bW)-(bW) double invariant-mass distribution of top-quarks ;  $m_{bW} = (p_b + p_W)^2$

limiting for the events with  $m_{tt} < 370$  GeV (10% of the total event)



$$\delta m \simeq -200 \text{ [MeV]}$$

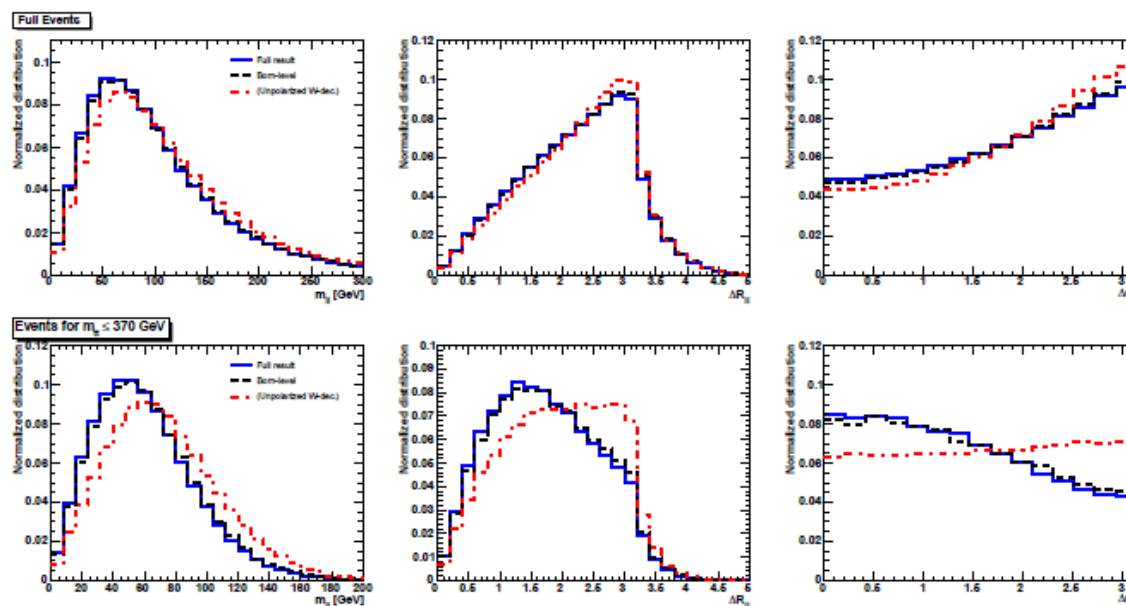
**correlated deviation** : one top-quark is still on-shell,  
but the other invariant-mass is reduced

## 3. Event Generation

### □ Some Examples (at partonic-level)

#### (3) lepton angular distributions (di-lepton case)

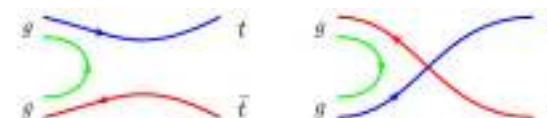
- Unfortunately, small bound-state effects in final lepton angular distributions
- bWbW MEs plus W-decay by parton-shower give wrong distribution (no W polarization)  
⇒ use 6-body MEs



### 3. Event Generation

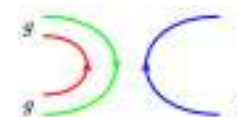
#### □ Color-flow assignment

in gluon-fusion process,



$$\begin{aligned}\mathcal{M}_{gg \rightarrow t\bar{t}} &= (T^a T^b)_{ij} \mathcal{M}_1 + (T^b T^a)_{ij} \mathcal{M}_2 \\ &= \frac{1}{2} \{T^a, T^b\}_{ij} \mathcal{M}_S + \frac{1}{2} [T^a, T^b]_{ij} \mathcal{M}_A\end{aligned}$$

symmetric part :  $\frac{1}{2} \{T^a, T^b\} = \frac{1}{2N_c} \delta^{ab} \delta_{ij} + \frac{1}{2} d^{abc} T_{ij}^c$



color-singlet

ratio of the amplitudes squared :  $\left| \frac{1}{2N_c} \delta_{ab} \delta_{ij} \right|^2 / \left| \frac{1}{2} d^{abc} T_{ij}^c \right|^2 = \frac{2}{N_c^2 - 4}$

this is zero in large-N limit, but not in QCD

our color-singlet events have correct color-flow assignment in the LHEF record





## 4. Summary

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Bound-state effects in the top-quark production at hadron colliders are studied

- At the LHC, **gluon-fusion process** dominates and substantial amount of  $t\bar{t}$  pair is **color-singlet**.
- The bound-state effects are calculated for the  $m_{t\bar{t}}$  distribution up to NLO (Green's func., gluon radiation, hard-correction)
- Large corrections in  $m_{t\bar{t}}$  dist. near threshold is predicted, and a broad resonance below the threshold may be seen
- Correction to the total cross-section  $\sim 1\%$



## 4. Summary

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- **Differential cross-sections** are also calculated including **BS effects**, non-resonant amp's as well as decays of Ws
  - incorporate momentum-space Green functions for color-singlet and octet
  - smooth interpolation to the high-energy region
  - non-resonant diagrams can be also taken into account
  - develop an Event-Generator to describe the threshold  $t\bar{t}$