BEIWEEN JEIS IN QUU



ANDREA BANFI ETH ZURLCH

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OUTLINE

Azimuthal decorrelations in QCD hard processes

Dijets in the back-to-back region

Phenomenology of azimuthal decorrelations in DIS

Joining soft-collinear and BFKL resummation?

- Azimuthal decorrelations between jets are directly sensitive to QCD radiation
- Tuning of MC event generators (PYTHIA DW)



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The distribution in Δφ can gives an incredible amount of information on QCD dynamics

Back-to-back region: multiple soft-collinear emissions



 $\Delta \phi \simeq \pi - \frac{|k_t \times p_{t1}|}{p_{t1}^2}$

Small-x regime: large phase space for multiple hard gluons



 $\Delta\phi\simeq\frac{2}{5}\pi\lesssim1$



Neither NLO nor BFKL account for large $\ln(\pi - \Delta \phi)$ terms which dominate the distribution at $\Delta \phi \sim \pi$

IR SENSITIVE JET OBSERVABLES

- The physics underlying azimuthal decorrelations at $\Delta \phi \sim \pi$ is present in many jet cross sections
 [AB Dasgupta '03]
 - Θ dijet rates with symmetric E_t cuts
 - Θ jet transverse energy imbalance $\Delta E_t = |E_{t1} E_{t2}|$
 - event shape distributions with jets as inputs
 - Jet cross sections developing an infrared sensitivity



Second Representation State and $p_{\text{jet},2} = p_1$ and $p_{\text{jet},2} = p_2$



$$\Delta \phi \simeq \pi - \frac{k_t}{p_{t1}} |\sin \phi|$$

^{Qq}Multiple gluon emissions collinear to the beam give



SOFT-COLLINEAR EFFECTS (II)

♀ Radiation inside one jet $\Leftrightarrow p_{jet,1} = p_1$ but $p_{jet,2} \neq p_2$



- If one uses a vectorial recombination scheme $\vec{p}_{t,jet} = \vec{p}_{t2} + \vec{k}_t$ the two jets remain back-to-back, i.e. $\Delta \phi = \pi$
- Azimuthal decorrelation is insensitive to radiation inside the jets, i.e. it is a non-global observable
- What is so special about non-global observables?

construction for approximation for 200 radiation



This picture breaks down for non-global observables
[Dasgupta Salam '01]



New soft large-angle contributions arise when harder gluons emit a softer gluon in the observed region

Monday, October 3, 2011

NON-GLOBAL LOGS IN JETS

- All observables involving jets are in principle affected by non-global logarithms
 [AB Dasgupta '03]
- Non-global logarithms depend on the jet algorithm and on the recombination scheme (maximal for anti-kt and reduced for clustering algorithms)
 [Appleby Seymour '02]



For clustering algorithms there are further clustering logarithms that spoil independent emission as well!
 [AB Dasgupta '05]

WHO'S AFRAID OF NG LOGS?



WHO'S AFRAID OF NG LOGS?

Non-global logs are soft large-angle non-abelian contributions resummed by solving a non-linear equation

[Dasgupta Salam '01] [AB Marchesini Smye '02]

 Dynamics of NG logs is understood both in the context of interjet energy flow and of jet cross sections (e.g. dependence on jet algorithm and jet radius)

> [Dasgupta Salam '02] [AB Dasgupta Khelifa-Kerfa Marzani '10]

- Clustering logarithms can be computed numerically and analytically estimated as a power series in the jet radius [AB Dasgupta '05]
 [Delenda Appleby AB Dasgupta '06]
- NG and clustering logs have a small impact on transverse momentum resummations
 [AB Dasgupta '03]

GLOBAL IS BETTER

- Non-global jet cross sections have various drawbacks
 - Resummation of NG logs can only be performed in the large N_c limit
 [Dasgupta Salam '01]

[AB Marchesini Smye '02]

 Predictions for non-global observables depend heavily on the details of the observation region (jet algorithm, recombination scheme, jet radius, etc.)
 [Dasgupta Salam '02]

[AB Dasgupta Khelifa-Kerfa Marzani '10]

Global observables are not sensitive to the details of the interjet region, except from universal soft anomalous dimensions, a.k.a. the fifth form factor

[Kidonakis Oderda Sterman '98] [AB Salam Zanderighi '05] [Dokshitzer Marchesini '06]

Is it possible to devise global dijet cross sections?

PT RECOMBINATION SCHEME

 $\Delta \phi$

 $-p_{1}$

Sensitivity to_pradiction inside the jets with the p_t weighted recombination scheme³ → global observable!

[AB Dasgupta '08]



$$\phi_{\text{jet},2} = \frac{p_{t2}\phi_2 + k_t\phi}{p_{t2} + k_t} \quad \Rightarrow \quad \Delta\phi = \pi - \frac{k_t}{p_{t1}}|\sin\phi - (\pi - \phi)|$$

Sovel kind of one-dimensional momentum cancellation

$$|\pi - \Delta \phi|_{\text{out}} \simeq \frac{1}{p_{t,\text{jet}}} \left| \sum_{i \in \text{jets}} k_{ti} \sin^3 \phi_i \right|$$

SOFT GLUON RESUMMATION

Resummation of soft-collinear gluons is performed in b
 space with the same phylosophy as for φ* in DY

$$\frac{1}{\sigma} \frac{d\sigma}{d\Delta} = \frac{2}{\pi} \int_0^\infty db \cos(b\Delta) \times \underbrace{\frac{f_1(\mu_F/\bar{b})}{f_1(\mu_F)} \frac{f_2(\mu_F/\bar{b})}{f_2(\mu_F)} e^{-R_{\rm in}(\bar{b})}}_{\rm initial \ \text{state} \ (DY-type)} \times \underbrace{\frac{e^{-R_{\rm out}(\bar{b})}}{f_{\rm initial \ \text{state}}} \times \underbrace{\frac{\mathrm{Tr}[He^{-\Gamma^{\dagger}t(\bar{b})}Me^{-\Gamma t(\bar{b})}]}{\mathrm{Tr}[HM]}}_{\rm soft \ \mathrm{large-angle}}$$

Solution Sudakov exponent $R_{out}(\bar{b})$ accounting for soft-collinear radiation inside the outgoing jets

$$R_{\text{out}}(\bar{b}) = \sum_{i} C_{i} \int_{0}^{Q} \frac{dk_{t}}{k_{t}} \frac{2\alpha_{s}(k_{t})}{\pi} \int_{0}^{\ln\frac{E_{i}}{k_{t}}} d\eta \times$$
$$\times \Theta \left[\frac{2}{3} \left(\frac{E_{i}}{p_{t}} \right)^{2} \frac{k_{t}}{p_{t}} e^{-2\eta} |\sin\phi|^{3} - \frac{1}{\bar{b}} \right] \simeq \frac{C_{\text{out}}}{3} \frac{\alpha_{s}}{\pi} \ln^{2} \bar{b}$$





Generation Both LO and NLO predictions diverge at Δφ ~ π
 Resummation approaches a plateau for Δφ ~ π, as expected for one-dimensional cancellations





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resummation coefficient constant

 $\Sigma(L) \sim (1 + C_1(L)) \Sigma_{\rm res}(L)$ $C_1(L) = \Sigma_1(L) - H_{12}L^2 - H_{11}L$



• At the largest value of x_B the coefficient constant becomes larger than one \Rightarrow resummation of small-x effects in C_1 ?

COMPLEMENTARY APPROACHES

- Summary of all-order resummation of azimuthal decorrelations at small and large $\Delta \phi$
 - Soft-collinear resummation
 - See Plateau (random walk)
 around $\Delta \phi = \pi$
 - Hard emissions only at fixed order
 - Described by shower MC at Tevatron (large x), not at HERA (small x)
 - Analytical calculations at NLL+NLO

Small-x resummation

- Unintegrated parton
 distributions at $\Delta \phi = \pi$
- Multiple hard emissions
 via BFKL or CCFM
- Described at HERA by CASCADE, no results for hadron colliders yet
- No analytical control over MC accuracy

CONCLUSIONS

- Azimuthal dijet decorrelations are rich in information on QCD dynamics at soft and hard scales
- In the quasi back-to-back region it is possible to resum multiple soft-collinear enhancements ⇒ first transverse momentum resummation with final-state coloured particles
- Sensitivity to intrinsic transverse momentum of the proton and to multiple hard gluon emissions (BFKL)

Description of LHC data requires a joint effort from different QCD communities