

Aspects of the δN formalism

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 - Cf. spectral tilt: $n - 1 = 2\eta - 6\epsilon$ (Liddle/DHL 1992)
- Trispectrum, even higher correlators, could be as important as the bispectrum
- Need to specify box size L (infrared cutoff)
 - But parameters run with L

The correlators

Spectrum \mathcal{P} , bispectrum $^\dagger f_{\text{NL}}$, trispectrum $^{\dagger\dagger} \tau_{\text{NL}}$:

$$\langle \zeta_{\mathbf{k}} \zeta_{\mathbf{k}'} \rangle = (2\pi)^3 \delta(\mathbf{k} + \mathbf{k}') K_1 \mathcal{P}$$

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Spectrum \mathcal{P} , bispectrum[†] f_{NL} , trispectrum^{††} τ_{NL} :

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where the kinematic factors depend on the wave-vectors:

$$\begin{aligned}K_1 &\equiv 2\pi^2 / k^3 \\ K_2 &\equiv K_1(k) K_1(k') + 5\text{perms} \\ K_3 &\equiv K_2 K_1(|\mathbf{k} + \mathbf{k}''|) + 23\text{perms}\end{aligned}$$

[†] Komatsu/Spergel 2000; Maldacena 2003

^{††} Boubekur/DHL 2005

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$$N = \int_{t_1}^t \frac{d \ln a(\mathbf{x}, t)}{dt} dt$$

Salopek & Bond 1990; DHL, Malik & Sasaki 2005

(non-perturbative refs.)

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- Then

$$N(\mathbf{x}, t) = N(\phi_i(\mathbf{x}), \rho(t))$$

the expansion of a family of unperturbed universes

DHL, Malik & Sasaki 2005 (non-perturbative)

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- Slow-roll, GR $\Rightarrow \mathcal{P}_{\delta\phi} = (H/2\pi)^2$ and $\partial N/\partial\phi = V/V'$
 - First term of ζ dominates

$$\mathcal{P}(k) = \frac{1}{2\epsilon_*} \left(\frac{H_*}{2\pi} \right)^2$$

$$n - 1 = 2\eta_* - 6\epsilon_*$$

Non-gaussianity in the standard scenario

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Comparison with Maldacena:

- He uses comoving slicing, computes $\mathcal{R} \rightarrow \zeta$ directly
- The δN approach instead uses two stages
 - Vacuum fluctuation converted to classical $\delta\phi$ (flat slicing)
 - Then δN gives ζ in terms of $\delta\phi$

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- or at a reheating by other mechanisms
(Dvali/Gruzinov/Zaldarriaga 2004, Kofman 2004, Bauer/Graesser/Salem 2005)

Linear approximation; the spectrum

Linear in $\delta\phi_i$, NOT first-order cosmological perturbation theory

$$\zeta(\mathbf{x}, t) = \sum N_i(t) \delta\phi_i(\mathbf{x})$$
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$$\mathcal{P}_\zeta(k, t) = \left(\frac{H_*}{2\pi}\right)^2 \left[\frac{1}{2\epsilon_*} + \sum N_{\sigma_i}^2(t) \right]$$

$$r \equiv \mathcal{P}_{\text{tensor}} / \mathcal{P}_\zeta \leq 16\epsilon_* \quad (\epsilon \equiv -\dot{H}/H^2)$$

Starobinsky 1985, Sasaki & Stewart 1996

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- If one $\sigma_i \equiv \sigma$ dominates, $n - 1 = 2\eta_{\sigma\sigma} - 2\epsilon_*$.

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- For slow-roll, correlator small (Seery & Lidsey 2005) leading to

$$\frac{3}{5} f_{\text{NL}} = \frac{r}{32} f, \quad 1 < f(k_1, k_2, k_3) < \frac{11}{6}$$

making f_{NL} unmeasurable (DHL & Zaballa 2005, Vernizzi & Wands 2006).

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- SL (06) same result for trispectrum
- For k - and ghost inflation, f_{NL} probably measurable.

Quadratic approximation

$$\zeta(\mathbf{x}, t) = \sum N_i \delta\phi_i(\mathbf{x}) + \frac{1}{2} \sum N_{ij} \delta\phi_i \delta\phi_j$$

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$$\mathcal{P}_\zeta = \left(\frac{H_*}{2\pi}\right)^2 \sum N_i^2 + \ln(kL) \left(\frac{H_*}{2\pi}\right)^4 \text{Tr } N^2$$
$$\frac{3}{5} f_{\text{NL}} = \frac{\sum N_i N_j N_{ij}}{2(\sum N_i^2)^2} + \ln(kL) \mathcal{P}_\zeta \frac{\text{Tr } N^3}{(\sum N_i^2)^3}$$
$$\tau_{\text{NL}} = 2 \frac{N_i N_{ij} N_{jk} N_k}{(\sum N_i^2)^3} + \ln(kL) \mathcal{P}_\zeta \frac{\text{Tr } N^4}{(\sum N_i^2)^4}$$

Infrared running

$$\zeta = \delta\phi + b\delta\sigma + (\delta\sigma)^2 \quad \text{with } \overline{\delta\sigma} = 0$$

$$\mathcal{P}_\zeta = \mathcal{P}_{\delta\phi} + b^2\mathcal{P}_{\delta\sigma} + \mathcal{P}_{(\delta\sigma)^2}$$

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Assume slow-roll and minimal box size

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If middle term negligible, non-gaussianity depends on the box size with f_{NL} and $\tau_{\text{NL}} \propto \ln(kL)$.

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- Simplest case: $\frac{3}{5} f_{\text{NL}} = -\frac{3}{4}$

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(ii) Two-component modular inflation Kadota & Stewart 03

$N \propto 1/\sigma$ gives negligible non-gaussianity DHL & Rodriguez 05

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4. Is Starobinsky's stochastic formalism an approximation to a non-perturbative version of Weinberg's analysis?