#### Towards Realistic Stringy Models of Particle Physics & Cosmology

as viewed by

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#### What is String Phenomenology?

## Particle Physics & Cosmology

- Deep connection, e.g., inflation, dark matter, neutrinos...
- Both study the universe in the extreme conditions.





## The Standard Model(s)





## Hierarchy problem SUSY?

Flatness, horizon, anisotropy Inflation? Dark Energy?

....











#### Are we ready for String Phenomenology?

#### The beginning of the unexpected ...



#### String Theory as a model of hadrons



String theory began as a phenomenological model.

Massless spin 2 particle: graviton!

#### Lessons

- Ideas driven by phenomenological questions.
- Need explicit models (c.f. QFT versus the Standard Model).
- Fixing problems that plague the theory often leads to new and far-reaching ideas:
  Extra spin-2 particle praviton Tachyon
  SUSY
- Works better than expected.

#### Meet the Quintuplets



#### The Heterotic Supremacy



 Type IIA/IIB: Difficult to implement non-Abelian gauge groups and chiral fermions. In fact, a no-go theorem for constructing the Standard Model. [Dixon, Kaplunovsky, Vafa]



- Heterotic E8xE8: naturally contains GUTs (E6, SO(10), SU(5),...) and hidden sectors.
- Type I and Heterotic SO(32): two other siblings that are largely ignored ...

#### String Phenomenology Begins



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- Low energy physics (spectrum, couplings,...) determined by topology & geometry of M.
- Building realistic heterotic string models: a huge industry beginning in the mid 80s.

## The Score Card

- E6, SO(10), SU(5) GUTs & MSSM-like vacua.
- Rank  $\leq 22$ .
- Constraints on gauge groups & matter reps.
- Gauge unification.
- Exotic matter: Schellekens' theorem.

Internal consistencies + phenomenological constraints



a very tight system!

However, two nagging problems ...



## Moduli Problem

#### Loss of predictivity

• Different  $< \phi_i >$  give inequivalent physics (e.g., couplings, particle masses, ...)

Phenomenological problems

- Existence of light scalars:
  - Equivalence principle violations.
  - Time varying  $\, lpha$  .
  - Energy in  $\phi_i$  can ruin cosmology.

## SUSY Breaking

- Assumptions:
  - Non-perturbative effects (e.g., gaugino and/or matter condensate) break SUSY.
  - The same NP effects also lift all moduli.
  - SUST scale ~ TeV (hierarchy problem).

#### • But ...

SUSY breaking effects on SM and moduli lifting potential not computed in a *controlled* stringy way.

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"When you come to a fork in the road, take it." Yogi Berra

#### Return of the Lost Family



### The Post-1995 Picture



#### Worth taking a fresh look at these long-standing problems.

#### All (new) roads lead to branes





## **Open Strings**

....

• Pioneering work (before 1995)



Bianchi, Pradisi, Sagnotti, ... Polchinski

• Recent review articles

Formalism:

Angelantonj, Sagnotti



**Model Building:** 

Blumenhagen, Cvetic, Langacker, Shiu

## Flux Compactification

- Just like particle couples to gauge field via  $\int_{\text{worldline}} A$
- Dp-brane couples to p+1 index gauge fields:  $\int_{\text{worldvolume}} A_{p+1}$
- Thus p+2-form field strengths:

$$F_{p+2} = dA_{p+1}$$



## Moduli Stabilization

• The energy cost of a given flux depends on detailed geometry of M:

$$V_{n_1,n_2,\ldots,n_k}(\phi_i)$$

where

$$n_j = \int_{\Sigma_j} F$$
,  $j = 1, \dots, k$ .

• Lift moduli  $\phi_i$  !

## Type IIB Flux Vacua

• Superpotential induced by  $G_3 = F_3 - \tau H_3$ 

$$W = \int_{\mathcal{M}} G \wedge \Omega$$

Gukov, Vafa, Witten

 Stabilizes the dilaton and complex structure moduli (shape) of M.
Dasgupta, Rajesh, Seth

Dasgupta, Rajesh, Sethi Greene, Schalm, Shiu Taylor, Vafa Giddings, Kachru, Polchinski

 Additional mechanism stabilizes the Kahler moduli (size).
Kachru, Kallosh, Linde, Trivedi

...









# Can the Standard Model fit into this picture?

## Chiral D-brane Models

Two known ways to obtain chiral fermions:

• Branes at singularities







## The Recipe



- Pick your  ${\cal M}$  , and the associated sLAG  $\Pi_a$
- Chiral spectrum:



- Tadpole cancellation (Gauss's law):
  - $\sum N_a \left( \Pi_a + \Pi'_a \right) 4 \Pi_O = 0$
- K-theory constraints

## K-theory Constraints

- D-brane charges are classified by K-theory. Minasian & Moore Witten
- Discrete charges invisible in SUGRA, forbid certain non-BPS branes to decay. Sen
- Uncanceled K-theory charges can manifest as Witten anomalies on D-brane probes.

Uranga

- Implications to the statistics of string vacua. Blumenhagen et al
- Direct construction of such discrete charged branes.
  Schellekens et al discrete charged Maiden, Shiu, Stefanski

#### Toward Realistic D-brane Models

For a review, see, e.g., Blumenhagen, Cvetic, Langacker, Shiu, hep-th/0502005.

- Many toroidal orbifold/orientifold models.
- MSSM flux vacua.

Marchesano, Shiu

- D-branes in general Calabi-Yau (less is known about supersymmetric  $\Pi_a$ ).
- Gepner orientifolds

Angelantonj, Bianchi, Pradisi, Sagnotti, Stanev Dijkstra, Huiszoon, Schellekens Blumenhagen, Weigard

#### How about Cosmology?

#### Inflation as a probe of stringy physics



- Almost scale invariant, Gaussian primordial spectrum predicted by inflation is in good agreement with data.
- A tantalizing upper bound on the energy density during inflation:

 $V \sim M_{GUT}^4 \sim (10^{16} \text{GeV})^4$  i.e.,  $H \sim 10^{14} \text{GeV}$ 







### Brane Inflation

Stringy signatures, e.g., gravitational waves ...



## Brane Inflation



Are the branes moving slowly enough?



Is reheating efficient?

Can the cosmic strings be stable?



#### Warping by Fluxes



• Fluxes back-react on the metric:





**DBI** inflation

Silverstein and Tong

$$S = -\int d^4x \sqrt{-g} \left( f(\phi)^{-1} \sqrt{1 - f(\phi)\dot{\phi}^2} - V(\phi) - f(\phi)^{-1} \right)$$

Casual speed limit:  $\dot{\phi}^2 \leq f(\phi)^{-1}$  warp factor

## Warped Throats



- Cosmic strings spatially separated from SM branes: not susceptible to breakage.
- Reheating via tunneling is efficient due to KK versus graviton wavefunctions. Barneby, Burgess, Cline Kofman and Yi Chialva, Shiu, Underwood Frey, Mazumdar, Myers

### Non-Gaussianities

Large 3-point correlations that are potentially observable. Moreover, distinctive shape. [Figures from Chen, Huang, Kachru, Shiu]

 $-54 < f_{NL} < 114 \; (WMAP3) \quad f_{NL} \sim 5 \; (PLANCK)$ 



#### Have we gone too far?

#### The Landscape





## Sightseeing in the Landscape

- These are *candidate* vacua (very few known examples where *all* moduli are stabilized.)
- The open string landscape (relevant to phenomenology!) is less understood.
- Realistic models are rare (QFT vs the Standard Model).



## The Wave Function?



Hartle, Hawking, Vilenkin, Linde, ...

...

In the context of string landscape

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Sarangi, Tye Kane, Perry, Zytkow Ooguri, Vafa, Verlinde

## Summary

- String phenomenology ~ 20+ year old baby --not fully accomplished but no longer naive.
- Too early for string phenomenology? Part of the SM was developed before gauge theories were shown to be renormalizable.
- Spin-off results (e.g., Calabi-Yau, G2, mirror symmetry, duality, topology change, ...).

## Summary

• Fountain of new ideas/scenarios for particle physics and cosmology:

SUSY: high/low, split, ...

Extra dimensions: large/small, warped/unwarped, universal/brane world.

Technicolor: AdS/CFT

Brane universe: brane inflation, DBI inflation, ...

#### 2005 +



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http://science/plants

Looking back to the down of time Unsecond vers l'aube du temps







## Thank you