

ON A CANONICAL QUANTIZATION OF PURE 3D ADS GRAVITY

Work in progress with Jihun Kim

OUTLINE

- CLASSICAL ADS GRAVITY AS CHERN-SIMONS $SL(2, \mathbb{R}) \times SL(2, \mathbb{R})$
- CLASSICAL PHASE SPACE OF $SL(2, \mathbb{R})$ CONTAINS TEICHMULLER SPACE
- HOLOMORPHIC QUANTIZATION: WHICH NORM?
- A CONNECTION WITH CFTs

- GLOBAL Diffeomorphisms and the Modular Group Action on Wave Functions
- Normalizability of the Wave Function
- An Improved Connection with CFTs: Holography as Superselection Projection
- The Case of $c < 1$
- Tentative Conclusions

CLASSICAL GRAVITY AS CHERN-SIMONS

$$A = e/l - \omega, \quad \tilde{A} = e/l + \omega$$

$$S_E = S(A) - S(\tilde{A})$$

$$S = \frac{\kappa}{4\pi} \int \text{tr} \left(AdA + \frac{2}{3} AAA \right), \quad \kappa = \frac{l}{4G}$$

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IN CANONICAL QUANTIZATION 3D SPACE IS

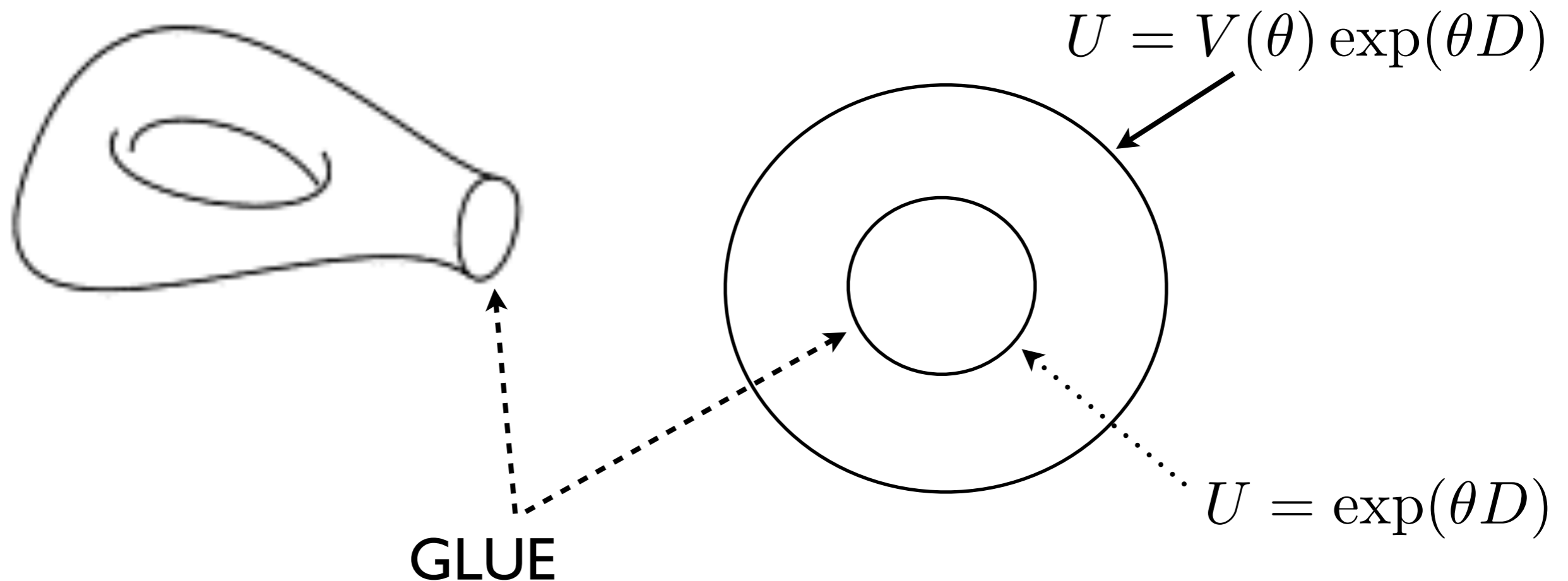
$$M = \Sigma \times R$$

CONSTRAINT EQUATION (GAUSS LAW)

$$F|_{\Sigma} = 0 \rightarrow A = dUU^{-1} \quad (\text{locally})$$

THE SPACE OF FLAT CONNECTIONS MODULO GAUGE TRANSFORMATIONS IS A DIRECT PRODUCT OF TWO SPACES: (EQUIVALENCE CLASSES OF) BOUNDARY GAUGE TRANSFORMATIONS TIMES A FINITE DIMENSIONAL SPACE

GEOMETRICALLY:



FINITE DIMENSIONAL SPACE WITH SEVERAL
CONNECTED COMPONENTS.

WHEN ALL HOLONOMIES ARE HYPERBOLIC AND
MANIFOLD HAS ONE BOUNDARY COMPONENT
THIS SPACE IS

$$T_{\Sigma} \times \widehat{SL}(2, R) / S_1$$

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RESTRICT

$$A|_{\partial\Sigma} = \begin{pmatrix} 0 & L(t + \phi) \\ 1 & 0 \end{pmatrix}$$

MODULI SPACE IS

$$T_{\Sigma} \times \text{Diff}(S_1) / S_1$$

THIS SPACE ADMITS A KÄHLER STRUCTURE AND A
KÄHLER FORM: THE WEYL-PETERSSON FORM

THE WEIL-PETERSSON FORM HAS A KÄHLER
POTENTIAL WHICH ALLOWS TO QUANTIZE THE
TEICHMÜLLER SPACE IN HOLOMORPHIC
QUANTIZATION.

THE RESULT IS JUST A (SUM OF PRODUCT)
HILBERT SPACES

$$\mathcal{H} = \sum_L V_L \otimes H_T,$$

H_T = holomorphic functions on T , $\text{tr exp } D = 2 \cosh L$

V_L = Irrep of Virasoro

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HOW DO STATE VECTORS

$$\psi \in H_T$$

TRANSFORM UNDER THE MODULAR GROUP?

NORM:

$$\langle \psi_I | \psi_J \rangle = \int_T \bar{\psi}_I \exp K \psi_J$$

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SINCE

$$K \rightarrow K + F + \bar{F}, \quad F = \text{holomorphic}$$

THEN

$$\psi_I \rightarrow U_I^J e^{-F} \psi_J, \quad U_I^L U_L^J = \delta_I^J$$

THE KÄHLER POTENTIAL FOR THE WEYL
PETERSSON FORM OF THE TEICHMÜLLER SPACE
FOR PUNCTURED SURFACE IS KNOWN: ZOGRAF
AND TAKHTAJAN PROVED THAT IT EQUALS THE
(REGULARIZED) LIOUVILLE ACTION COMPUTED
ON SHELL:

$$K_0 = S_L|_{on\ shell} + H(t, \bar{t})$$

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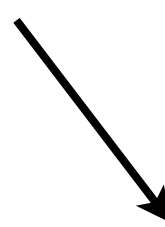
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FOR THE PUNCTURED SPHERE THIS IS THE POLYAKOV
CONJECTURE (TAKHTAJAN, ZOGRAF, MENOTTI,..)

THE SYMPLECTIC FORM FOR A RIEMANN SURFACE WITH A BOUNDARY OF LENGTH L IS (CO-HOMOLOGOUS TO) THE SUM OF THE SYMPLECTIC FORM FOR A SURFACE WITH A PUNCTURE PLUS THE FIRST CHERN CLASS OF THE TANGENT BUNDLE AT THE PUNCTURE (MIRZAKHANI, USING DUISTERMAAT-HECKMAN)

$$\omega_L \approx \omega_0 + \frac{1}{2}L^2 c_1(T)$$

ITS KÄHLER POTENTIAL IS THEN (SEE ALSO DIJKGRAAF, VERLINDE, VERLINDE, 1990)


$$K_L = K_0 + \frac{1}{2}L^2 \log g_{z\bar{z}}|_{z=w}$$

THE MEASURE THAT WE GET FROM THE KÄHLER
NORM THAT WE FIND IS

$$e^K = e^{-S_L^{-6k} + 2h\phi + H(t, \bar{t})} \Big|_{\phi} |F|^{-2}$$

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NOTICE:
TIMELIKE
LIOUVILLE



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$$h \equiv \frac{1}{2} k L^2$$

SQUARE OF HOLOMORPHIC
FUNCTION

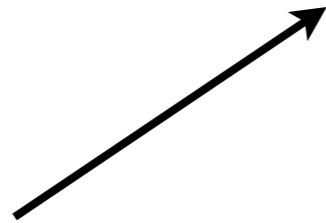
$$g_{z\bar{z}} \equiv e^{2\phi}$$

THIS NORM CAN BE PROMOTED TO A NORM VALID FOR
ANY VALUE OF $c=3l/2G$

$$e^{-S_L^{26-c} + 2h\phi + H(t, \bar{t})} \Big|_{\phi} |F|^{-2} \rightarrow \langle \int_{\Sigma} e^{2\alpha\phi} \rangle_{26-c}^L Z^{bc} (\det \mathfrak{S}\Omega_{ij})^{c/2} Z_c^S$$

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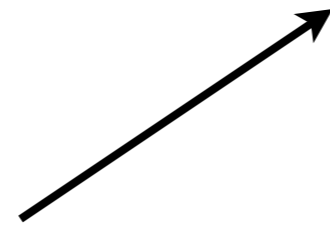
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ONE-POINT VEV IN (TIMELIKE) LIOUVILLE
WITH CENTRAL CHARGE 26-c

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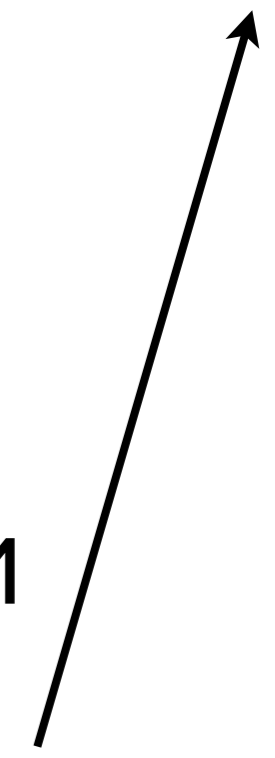
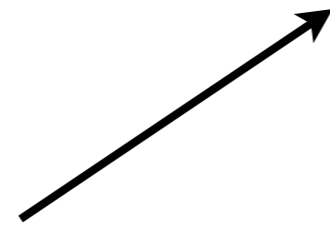


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PARTITION FUNCTION OF bc SYSTEM

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ONE-POINT VEV IN (TIMELIKE) LIOUVILLE
 WITH CENTRAL CHARGE 26-c

PARTITION FUNCTION OF bc SYSTEM

PARTITION FUNCTION OF c SCALARS

$$\alpha(\alpha - Q) = -h(L) + 1, \quad 26 - c = 1 - 6Q^2$$

CONFORMAL WEIGHT OF IRREP V_L

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THIS PRODUCT FACTORIZES INTO

$$e^{H(t, \bar{t})} |F|^{-2}$$

(QUILLEN)

REDEFINE THE WAVE FUNCTION

$$\chi_I = \psi_I / F,$$

UNDER MODULAR TRANSFORMATIONS:

$$dz^{kL^2/2} \chi_I \rightarrow dz^{kL^2/2} U_I^J \chi_J$$

SCALAR PRODUCT

$$\langle \chi_I | \chi_J \rangle = \int_T \langle \int_\Sigma e^{2\alpha\phi} \rangle \frac{L}{26-c} Z^{bc} e^{H(t, \bar{t})} \bar{\chi}_I \chi_J$$

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OBTAINED BY H. VERLINDE IN '89 BY QUANTIZING FIRST CS
SL(2,R), THEN SOLVING THE CONSTRAINTS

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THE WAVE FUNCTIONS χ_I TRANSFORM AS ONE-POINT
CONFORMAL BLOCKS OF LIOUVILLE THEORY (VERLINDE,
TESCHNER)

QG IN 3D ADS IS DEFINED BY $SL(2,R) \times SL(2,R)$
SO WE NEED TO COMBINE TOGETHER THE
MODULI SPACES OF BOTH CS

$$\Psi_{QG} = \chi_I(t) \chi_J(\bar{t}')?$$

NO: THE WAVE FUNCTION OF QG MUST BE INVARIANT ALSO
UNDER GLOBAL DIFFEOMORPHISMS

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UNDER GLOBAL DIFFS

$$\phi : \Sigma \rightarrow \Sigma$$

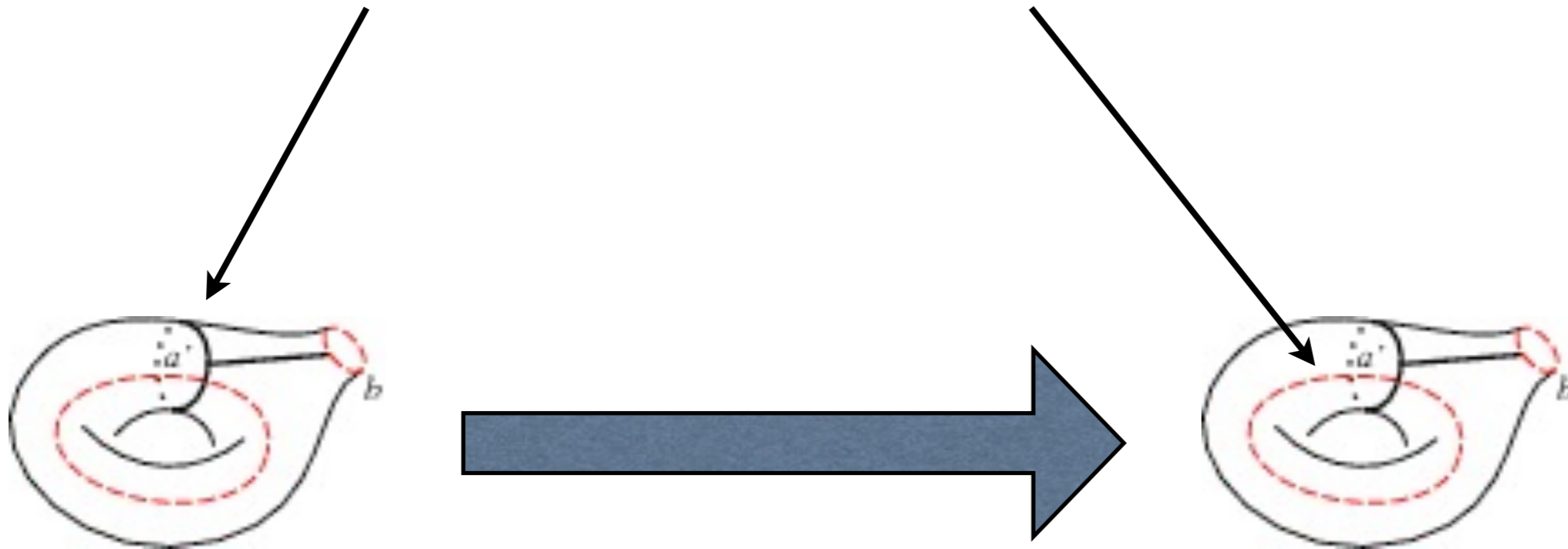
$$\text{tr } P e^{\oint_{\gamma} A} \in T \rightarrow \text{tr } P e^{\oint_{\phi^{-1}(\gamma)} \phi_* A}$$

$$\text{tr } P e^{\oint_{\gamma} \tilde{A}} \in \bar{T}' \rightarrow \text{tr } P e^{\oint_{\phi^{-1}(\gamma)} \phi_* \tilde{A}}$$

DIAGONAL ACTION ON

$$T \times \bar{T}'$$

THE BLACK CYCLE MAPS INTO THE RED CYCLE UNDER A GLOBAL DIFFEOMORPHISM



$$\text{tr } Pe^{\oint_{\gamma} A} \in T \rightarrow \text{tr } Pe^{\oint_{\phi^{-1}(\gamma)} \phi_* A}$$

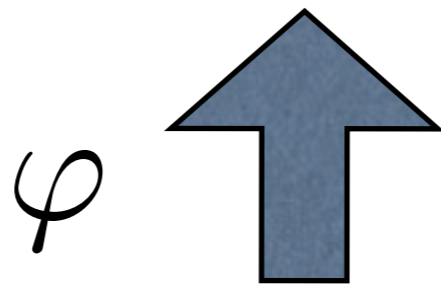
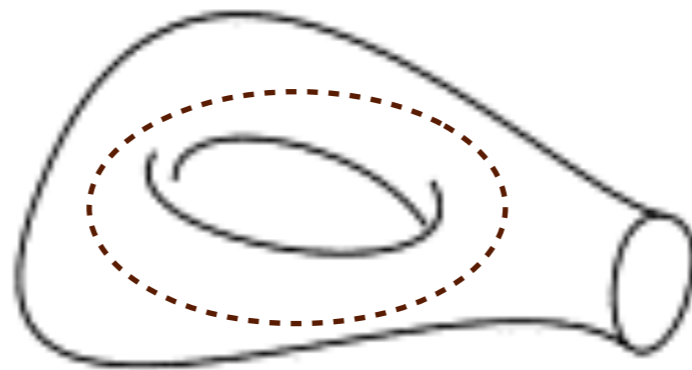
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$$T \times \bar{T}' \rightarrow T \times \bar{T}' / M$$

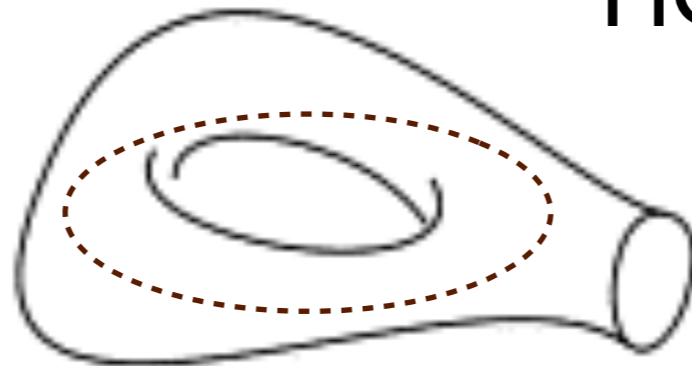
IN SECOND ORDER (METRIC) FORMULATION, THE
TEICHMULLER SPACES LABEL INITIAL AND FINAL
2D METRICS.

HOMOLOGY CANNOT CHANGE UNDER A 3D
DIFFEOMORPHISM THAT EXTENDS TO

$$\Sigma \times R$$



A 3D DIFF CANNOT
CHANGE THE
HOLONOMY BASIS



SO THE WAVE FUNCTION IS

$$\Psi = |v\rangle \sum_{IJ} \chi_I^h(t) \bar{\chi}_J^{\bar{h}}(\bar{t}') N^{IJ}, \quad |v\rangle \in V_h \otimes V_{\bar{h}}$$

$dz^h d\bar{z}^{\bar{h}} \Psi$ INVARIANT UNDER DIAGONAL MAPPING CLASS GROUP

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ALL ONE-POINT FUNCTIONS OF CFT BELONG TO THIS SPACE
I.E. HILBERT SPACE IS TARGET SPACE OF CFT'S ONE POINT
FUNCTIONS (SEE WITTEN '07)

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ACTUALLY ANALYTIC CONTINUATION TO
INDEPENDENT LEFT AND RIGHT MODULI
(SEE WITTEN '07, SEGAL)

NEW RESTRICTIONS COME FROM NORMALIZABILITY

$$\Psi = |v\rangle \sum_{IJ} \chi_I^h(t) \bar{\chi}_J^{\bar{h}}(\bar{t}') N^{IJ}, \quad |v\rangle \in V_h \otimes V_{\bar{h}}$$

$$\left\langle \sum_{IJ} \chi_I \bar{\chi}_J N^{IJ} \middle| \sum_{KL} \chi_K \bar{\chi}_L N^{KL} \right\rangle < \infty$$

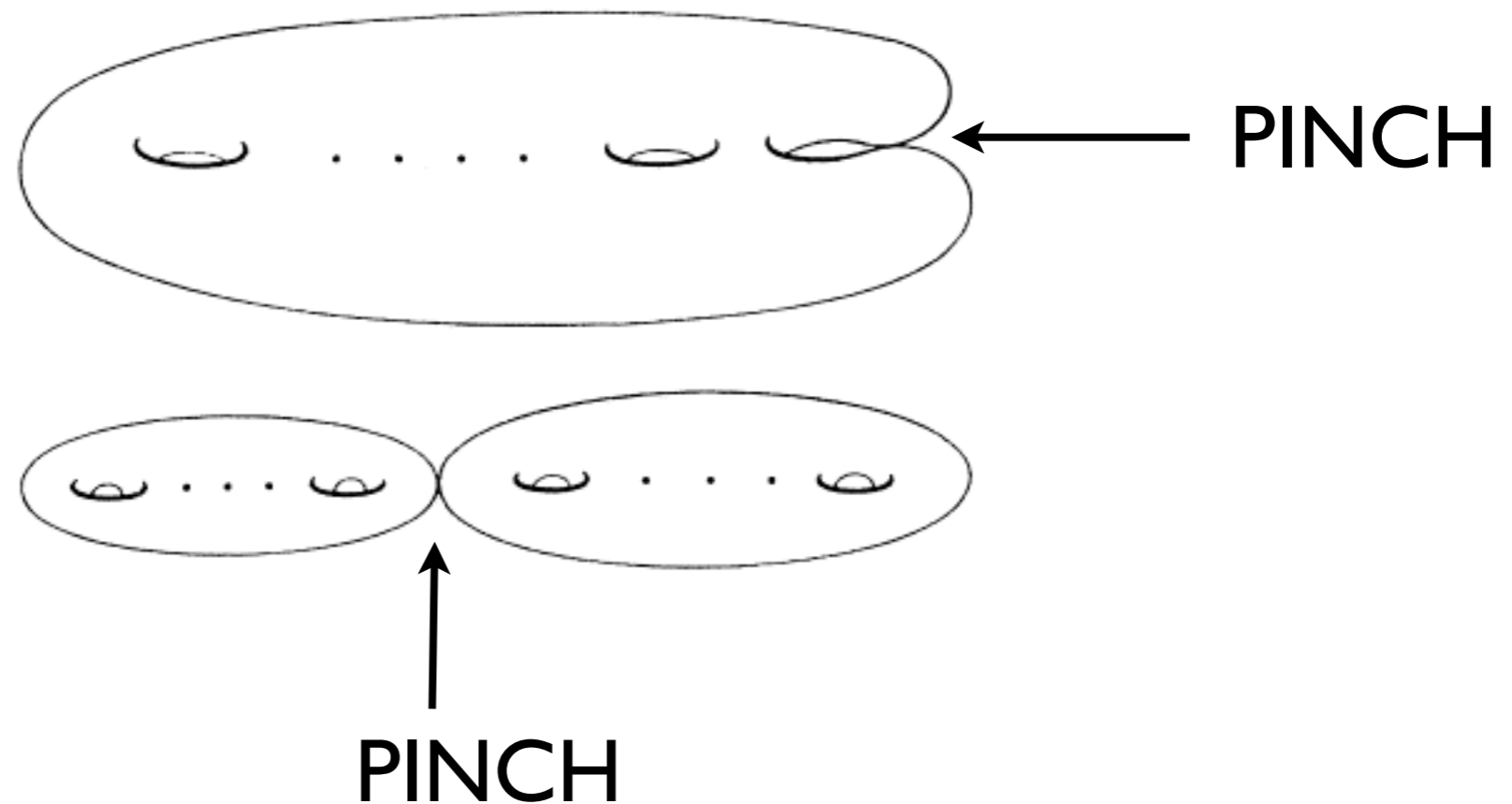
DANGEROUS REGIONS: DEGENERATING RIEMANN SURFACES

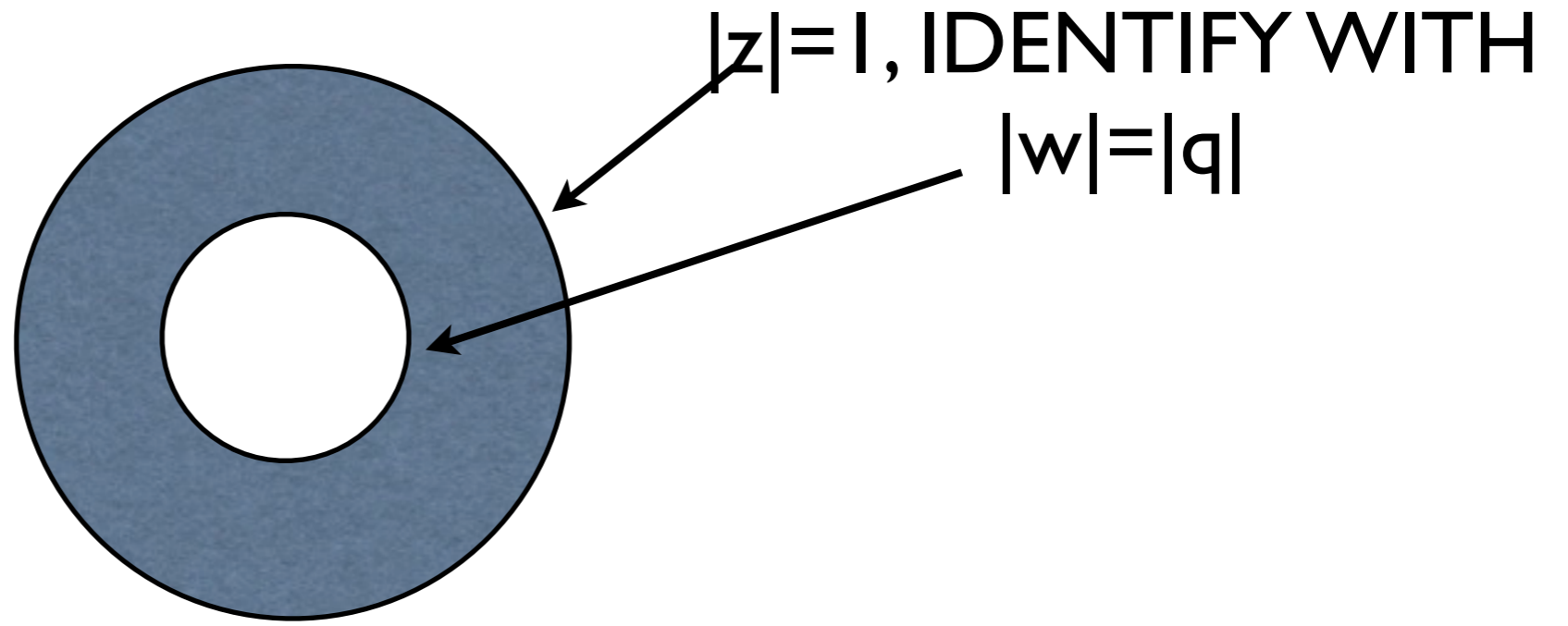
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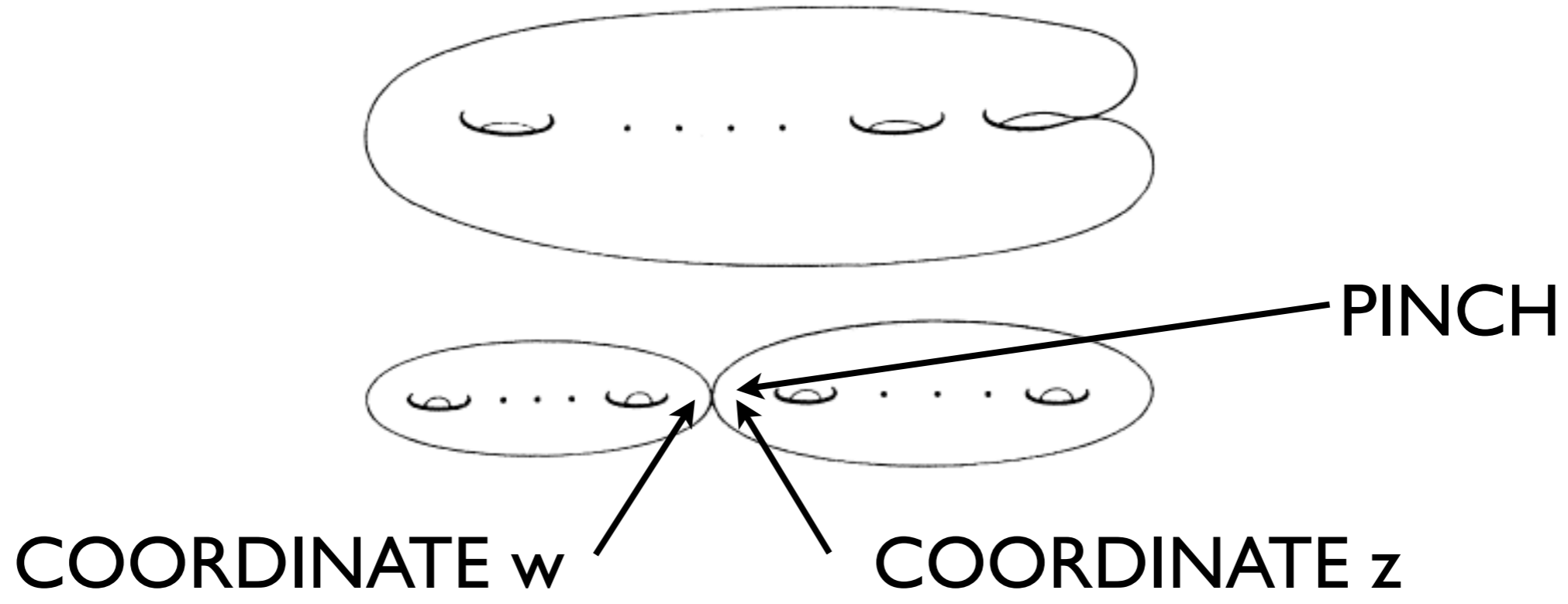
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DANGEROUS REGIONS: DEGENERATING RIEMANN SURFACES





$z = qw,$ $|q| < |z| < 1,$ $|q| < |w| < 1,$ $q = e^{2\pi i\tau}$



NEAR A NODE THE KÄHLER FORM AND CERTAIN MODULAR TRANSFORMATIONS ARE KNOWN AND SIMPLE

$$(\tau, \tau') \rightarrow (\tau + 1, \tau' + 1) \text{ (Dehn Twist)}$$

$$\omega_{WP} \propto \frac{d\tau \wedge d\bar{\tau}}{(\Im\tau)^3} \quad \longrightarrow \quad K \propto \frac{1}{\Im\tau}$$

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INVARIANCE UNDER THE DIAGONAL DEHN TWIST:

$$\Psi = |v\rangle \times \langle V^{h, \bar{h}} \rangle, \quad \langle V^{h, \bar{h}} \rangle = \sum_n \int d\mu_n(\Delta) q^{\Delta - c/24} \bar{q}'^{\Delta + n - c/24}$$

NORMALIZABILITY OF THE WAVE FUNCTION IMPOSES NEW CONSTRAINTS

DEFINE:

$$\tau = i\rho + (\theta + \zeta)/2\pi, \quad \tau' = i\rho' + (\theta - \zeta)/2\pi$$

NORM:

$$\|\langle V^{h, \bar{h}} \rangle\|^2 \approx \sum_n \int^\infty d\rho \int^\infty d\rho' \int_{-\infty}^\infty d\zeta \int d\mu_n(\Delta) d\mu_n(\Delta') e^{-2\pi\rho(\Delta + \Delta' - \frac{c-1}{12})} e^{-2\pi\rho'(\Delta + \Delta' + 2n - \frac{c-1}{12})} e^{i\zeta(\Delta - \Delta')}$$

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INTEGRATION OVER REAL LINE IMPLIES $\Delta = \Delta'$

NORM PROPORTIONAL TO DIRAC DELTA: CONTINUOUS SPECTRUM

$$||\langle V^{h, \bar{h}} \rangle||^2 \approx \sum_n \int d\mu_n(\Delta) d\mu_n(\Delta') F(\Delta, n) \delta(\Delta - \Delta')$$

NORM PROPORTIONAL TO DIRAC DELTA:
CONTINUOUS SPECTRUM

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SO HILBERT SPACE OF PURE 3D ADS GRAVITY
IS TARGET SPACE FOR CFTs WITH CONTINUOUS
SPECTRUM AND STATES WITH CONFORMAL WEIGHT

$$\Delta > (c - 1)/24$$

SPACE OF SUCH CFTs IS NONEMPTY
ONE SOLUTION:
LIOUVILLE CFT

DIDN'T WE KNOW IT FROM THE FACT THAT
WAVE FUNCTIONS OF $SL(2, \mathbb{R})$ TRANSFORM AS
CONFORMAL BLOCKS OF LIOUVILLE?

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NOT SAY ANYTHING ABOUT ITS RADIUS OF
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NOT SAY ANYTHING ABOUT ITS RADIUS OF
COMPACTIFICATION

IN PARTICULAR THEY ARE THE SAME FOR THE
UNCOMPACTIFIED BOSON (CONTINUOUS SPECTRUM)
AND FOR THE COMPACT BOSON (DISCRETE SPECTRUM)

HILBERT SPACE IS TOO LARGE

$$H_{CQG} = V_h \otimes V_{\bar{h}} \otimes \sum_g H_g$$

INFINITE MULTIPLICITY FOR ANY VIRASORO IRREP.

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INFINITE MULTIPLICITY FOR ANY VIRASORO IRREP.

PROPOSAL: CFTs AS SUPERSELECTION SECTORS

CLOSED UNDER OPE

$$O_\alpha : V_\beta \rightarrow V_\gamma$$

$$O_\alpha v_\beta = \sum_{v_\gamma} C_{\alpha v_\beta v_\gamma} v_\gamma$$

COEFFICIENTS OF OPE (FUSION RULES)

COEFFICIENTS OF FUSION RULES FIXED UP TO AN OVERALL CONSTANT

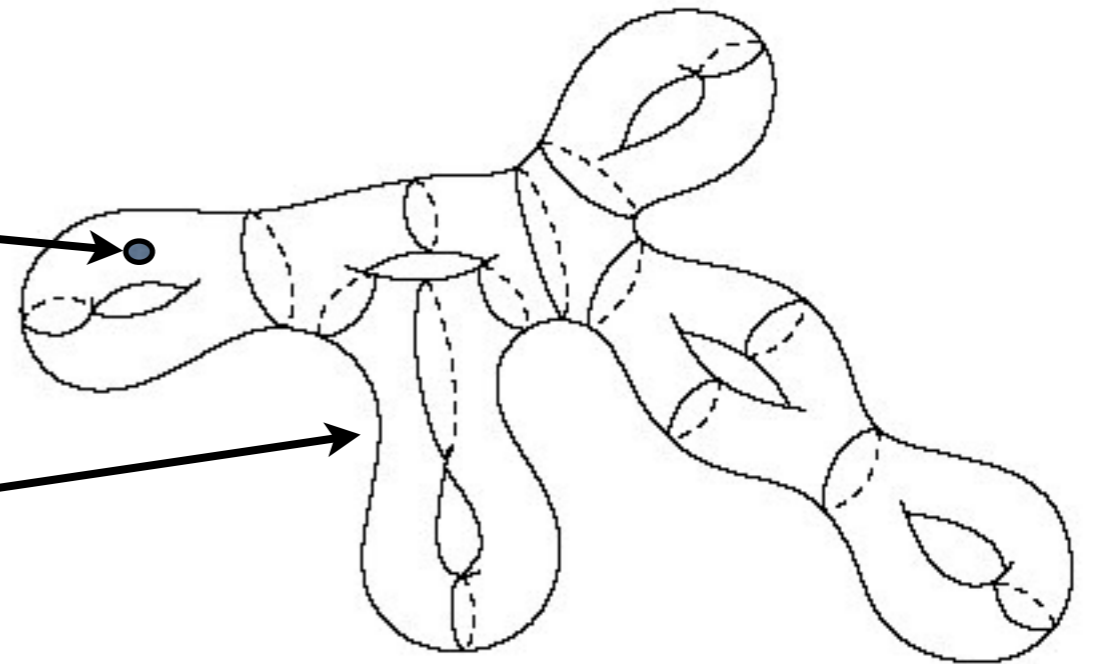
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INSERT VERTEX HERE

PANT $\approx C_{\alpha v_{\beta} v_{\gamma}}$

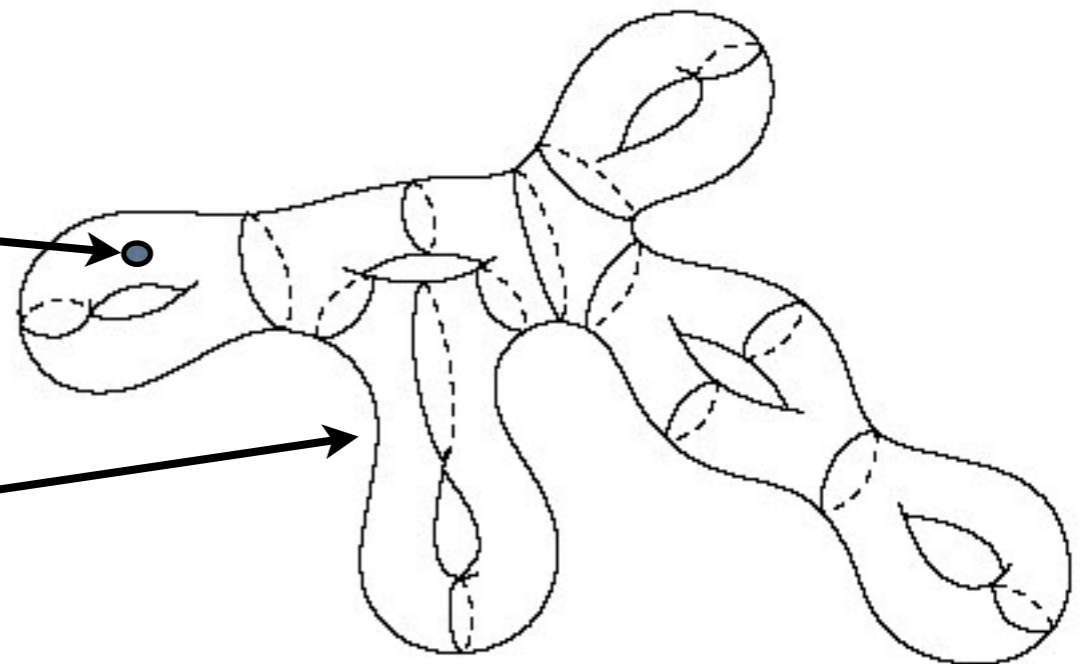


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SO

$$\langle V_{h,\bar{h}} \rangle_{\Sigma_g}$$

SCALES AS

$$\lambda^{2g-1} \langle V_{h,\bar{h}} \rangle_{\Sigma_g}$$

NEW RULE: ASSOCIATE STATE IN CFT TO SPECIAL VECTOR IN HILBERT SPACE OF CQG BY:

$$w \in H_{h, \bar{h}}^{CFT} \rightarrow \Psi = v \times (\langle V_{h\bar{h}} \rangle_{\Sigma_1}, \langle V_{h\bar{h}} \rangle_{\Sigma_2}, \langle V_{h\bar{h}} \rangle_{\Sigma_3}, \dots), \quad v \in V_h \otimes V_{\bar{h}}$$

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BAD FOR PURE GRAVITY!**

ALSO CONTINUOUS SPECTRUM=INFINITE ENTROPY

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ANY DISCRETIZATION GIVES EFFECTIVE CENTRAL CHARGE

$$c_{eff} = c - 24\Delta_{min} = c - 24 \frac{c-1}{24} = 1$$

MULTIPLICITY OF STATES NOT ENOUGH TO GIVE
BEKENSTEIN-HAWKING ENTROPY

LASCIATE OGNI SPERANZA?

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$$T/\mathcal{N} = \text{finite cover of } \mathcal{M}$$

$(T/\mathcal{N} \times \bar{T}'/\mathcal{N})/M$ HAS FINITE VOLUME = DISCRETE
SPECTRUM

FOR $c < 1$ MINIMAL MODELS WE CAN GET A
REASONABLE THEORY, WITH DISCRETE SPECTRUM
AND A DUAL TO THE ADS VACUUM IF WE IMPOSE
A NEW NON-GEOMETRICAL GAUGE SYMMETRY
THAT ACTS ON THE MODULI SPACE AS THE
NORMAL, FINITE INDEX KERNEL OF THE
MODULAR GROUP

(SEE CASTRO ET AL.)

TENTATIVE CONCLUSIONS

- A CANONICAL QUANTIZATION OF PURE GRAVITY IN 3D ADS YIELDS A HILBERT SPACE THAT CAN BE INTERPRETED AS THE TARGET SPACE FOR LIOUVILLE-LIKE CFTs
- CFTs APPEAR IN THE HILBERT SPACE AS SUPERSELECTION SECTORS CLOSED UNDER FUSION RULES
- FOR $c > 1$ THERE ARE SEVERAL PROBLEMS WITH THIS PICTURE SUCH AS ABSENCE OF THE DUAL TO ADS SPACE AND CONTINUOUS SPECTRUM
- FOR $c < 1$ WE FOUND THAT IMPOSING A NEW, NON GEOMETRIC GAUGE SYMMETRY GIVES A REASONABLE PICTURE (DISCRETE SPECTRUM, VACUUM)

MANY OPEN QUESTIONS

- DOES ALL THIS MAKE SENSE?
- CAN ONE EXTEND THE MINIMAL MODEL PICTURE TO HIGH SPIN 3D ADS THEORIES, OR $SL(N,R) \times SL(N,R)$ CS? (ANY RCFT ADMITS A NONTRIVIAL KERNEL OF THE MODULAR GROUP)
- IS $c > 1$ DOOMED? MAYBE WE CAN MAKE SENSE OF LIOUVILLE-LIKE CFTs: LIOUVILLE THEORY POPS UP TOO MANY TIMES IN 3D PURE GRAVITY TO BE JUST BACKGROUND NOISE!