#### Mass hierarchy and electroweak symmetry breaking

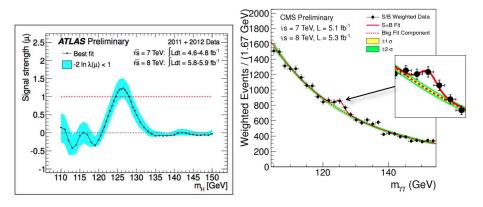
#### I. Antoniadis



Beyond the Standard Model after the first run of LHC GGI, Florence, 9-12 July 2013

- Low energy SUSY and 125 GeV Higgs
- Live with the hierarchy
- Extra *U*(1)'s
- Low scale strings and extra dimensions

# Higgs boson discovery



 $m_H = 125.5 \pm 0.2 \,(\text{stat.}) \pm 0.5 \,(\text{syst.})$ 

 $m_H = 125.7 \pm 0.3 \pm 0.3$  GeV

# Beyond the Standard Model of Particle Physics: driven by the mass hierarchy problem

Standard picture: low energy supersymmetry

Natural framework: Heterotic string (or high-scale M/F) theory

Advantages:

- natural elementary scalars
- gauge coupling unification
- LSP: natural dark matter candidate
- radiative EWSB

Problems:

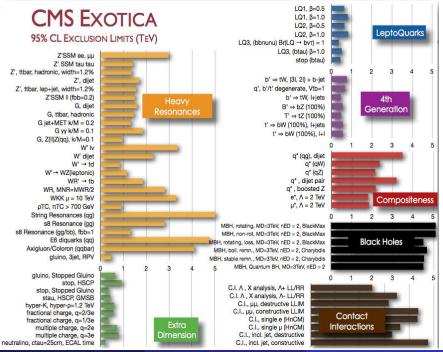
- too many parameters: soft breaking terms
- MSSM : already a % ‰ fine-tuning 'little' hierarchy problem

#### ATLAS SUSY Searches\* - 95% CL Lower Limits (Status: March 26, 2013)

	MSUGRA/CMSSM : 0 lep + j's + E T miss	L+5.8 fb <sup>-1</sup> , 8 TeV [ATLAS-CONF-2012-109]	t.so TeV g̃= g̃ mass	
	MSUGRA/CMSSM : 1 lep + j's + E - min	L=5.8 fb <sup>-1</sup> , 8 TeV (ATLAS-CONF-2012-104)	1.24 TeV g = g mass	
60	Pheno model : 0 lep + j's + E T miss	L+5.8 fb <sup>-1</sup> .8 TeV (ATLAS-CONF-2012-109)	1.18 TeV Q MASS (m(q) < 2 TeV, light 7	S ATLAS
16.	Pheno model : 0 lep + j's + E Trais	L#5.8 fb <sup>-1</sup> , 8 TeV (ATLAS-CONF-2012-109)	1.38 TeV Q mass (m(g) < 2 TeV, lig	
2	Gluino med. $\tilde{\chi}^{\pm}(\tilde{g} \rightarrow q\bar{q}\tilde{\chi}^{\pm})$ : 1 lep + j's + $E_{T miss}$	L=4.7 fb <sup>-1</sup> , 7 TeV [1208.4688]	900 GeV $\tilde{g}$ mass $(m(\tilde{\chi}^{0}) < 200 \text{ GeV}, m(\tilde{\chi}^{1}) =$	
993	GMSB (INI SP) : 2 len (OS) + i'e + E	L=4.7 fb <sup>-1</sup> .7 TeV [1208.4688]	1.24 TeV Q MASS (tang < 15)	fore Lundan
Inclusive searches	GMSB (ÎNLSP) : 2 lep (OS) + j's + Ε GMSB (τ NLSP) : 1-2 τ + j's + Ε	L=20.7 fb <sup>-1</sup> .8 TeV (1210.1314)	1.40 TeV ĝ mass (tanβ > 18)	
SIV	GGM (bino NLSP) : yy + E <sup>+mss</sup>	L=4.8 fb <sup>-1</sup> , 7 TeV [1209.0753]	1.07 TeV Q MASS (m(q)) > 50 GeV)	ſ .
13	GGM (wino NLSP) : y + len + F	L#4.8 fb <sup>-1</sup> , 7 TeV (ATLAS-CONF-2012-144)	619 GeV g mass	Ldt = (4.4 - 20.7) fb <sup>-1</sup>
ŝ	GGM (higgsing-bing NLSP) : $y + b + E$	L=4.8 fb <sup>-1</sup> , 7 TeV [1211.1167]	900 GeV Q MASS (m(g) > 220 GeV)	J
	GGM (higgsino NLSP) : Z + jets + E Timiss	L+5.8 fb <sup>-1</sup> .8 TeV IATLAS-CONF-2012-1521	690 GeV Q MASS (m(H) > 200 GeV)	s = 7, 8 TeV
	Gravitino LSP : 'monojet' + E Trias	L=10.5 fb <sup>-1</sup> , 8 TeV [ATLAS-CONF-2012-147]	645 GeV F <sup>1/2</sup> scale (m(G) > 10 <sup>-4</sup> eV)	
~	$\tilde{g} \rightarrow b \tilde{b} \tilde{\chi}_{1}^{0}$ : 0 lep + 3 b-j's + $E_{T,max}$	L=12.8 fb <sup>-1</sup> .8 TeV (ATLAS-CONF-2012-145)	1.24 TeV $\tilde{Q}$ mass $(m(\tilde{\chi}^2) < 200 \text{ GeV})$	
90. No	$\tilde{g} \rightarrow t \tilde{t} \tilde{\chi}$ : 2 SS-lep + (0-3b-)j's + $E_{T miss}$	L=20.7 fb <sup>-1</sup> , 8 TeV [ATLAS-CONF-2013-007]	900 GeV ĝ mass (any m(x))	8 TeV, all 2012 data
l gi	$\tilde{g} \rightarrow tt \tilde{\chi}_{1}^{0}$ : 0 lep + multi-j's + $E_{T,miss}$	L=5.8 fb <sup>-1</sup> .8 TeV (ATLAS-CONF-2012-103)	1.00 TeV Q Mass (m(\overline{z}^{0}) < 300 GeV)	
3rd gen. gluino mediated	g→tt <sup>2</sup> <sub>x</sub> 0 lep + 1 lef s + E <sub>T miss</sub>	L=12.8 fb <sup>-1</sup> , 8 TeV [ATLAS-CONF-2012-145]	1.15 TeV Q MASS (mQ <sup>2</sup> ) < 200 GeV)	8 TeV, partial 2012 data
	$bb, b \rightarrow b\chi^{-1} = 0 \text{ lep } + 2\text{ -b-jets } + E_{T \text{ mass}}$	L=12.8 fb <sup>-1</sup> , 8 TeV [ATLAS-CONF-2012-165]	620 GeV b mass (m(x)) < 120 GeV)	7 TeV, all 2011 data
(0, ~	$\tilde{b}\tilde{b}, \tilde{b}_{1} \rightarrow t\tilde{\chi}_{1}^{1}$ : 2 SS-lep + (0-3b-)]'s + $E_{T,miss}$	L=20.7 fb <sup>-1</sup> , 8 TeV [ATLAS-CONF-2013-007] 430 GeV		r ice, air 2011 data
ior in	tt (light), $t \rightarrow b\tilde{\chi}_1^+$ : 1/2 lep (+ b-jet) + E Triss		) = 55 GeV)	
not	tt (medium), $t \rightarrow b\tilde{\chi}_{\pm}^{+}$ : 12 lep (+ b-jet) + $E_{T,miss}$ tt (medium), $t \rightarrow b\tilde{\chi}_{\pm}^{+}$ : 1 lep + b-jet + $E_{T,miss}$	L=20.7 fb <sup>-1</sup> , 8 TeV [ATLAS-CONF-2013-037] 160 GeV [111033 [111]		
od s	tt (medium), $t \rightarrow b\chi_{z}^{*}$ : 2 lep + $E_{T,max}$	L=13.0 fb <sup>-1</sup> .8 TeV [ATLAS-CONF-2012-167] 160-440 Ge	$\tilde{t}$ mass $(m(\chi^2) = 0 \text{ GeV}, m(\chi^2) = 10 \text{ GeV})$	
3rd gen. squarks direct production	$\tilde{t}t$ (heavy), $\tilde{t} \rightarrow t \tilde{\chi}_{1}^{0}$ : 1 lep + b-jet + $E_{T miss}$		-610 GeV t mass (m(2 <sup>*</sup> ) = 0)	
- di			20-660 GeV T mass $(m(\chi^2) = 0)$	
3rc	tt (natural GMSB) : $Z(\rightarrow II) + b_{jet} + E_{T,miss}$		Gev t mass (m(2)) > 150 GeV)	
	$\tilde{t}_1 \tilde{t}_2, \tilde{t}_2 \rightarrow \tilde{t}_1 + Z : Z(\rightarrow II) + 1 \text{ lep } + b \text{-jet} + E_{T,miss}$		GeV L, mass (m(L) = m(2 <sup>4</sup> ) + 180 GeV)	
	$[1, ] \rightarrow [\chi]$ : 2 lep + E <sub>1 max</sub>		$m(\tilde{\chi}^2) = 0)$	
	$\chi_{1}^{+}\chi_{1}^{-}\chi_{2}^{+}\chi_{1}^{+}\chi$		$m(\chi_1) = 0$ $mass (m(\chi_1^2) \le 10 \text{ GeV}, m(\tilde{x}) = \frac{1}{2}(m(\chi_1^2) + m(\chi_1^2)))$	
EW direct	$\chi_{\chi_{\chi_{\chi_{\chi_{\chi_{\chi_{\chi_{\chi_{\chi_{\chi_{\chi_{\chi_{\chi$	L+20.7 fb <sup>-1</sup> , 8 TeV [ATLAS-CONF-2013-028] 180-330 GeV Ž	mass $(m(\tilde{\chi}_{1}^{2}) < 10 \text{ GeV}, m(\tilde{\chi}_{1}^{2}) = \frac{1}{2}(m(\tilde{\chi}_{1}^{2}) + m(\tilde{\chi}_{1}^{2})))$	
비송	$\chi^{+}\chi^{-}\chi^{+}\to \tilde{V}(\tilde{V}): 2 \operatorname{lep} + E_{\tau,\operatorname{miss}}^{\tau,\operatorname{miss}}$ $\chi^{+}\chi^{+}\chi^{+}\to \tilde{V}(\tau\tilde{V}): 2\tau + E_{\tau,\operatorname{miss}}$ $\chi^{+}\chi^{0}_{-}\to \tilde{V}(\tilde{U}(\tilde{V}V), \tilde{V})[(\tilde{V}V): 3 \operatorname{lep} + E_{\tau,\operatorname{miss}}^{\tau,\operatorname{miss}}$		600 GeV $\tilde{\chi}_{i}^{\pm}$ mass $(m(\tilde{\chi}_{i}^{+}) = m(\tilde{\chi}_{i}^{0}), m(\tilde{\chi}_{i}^{0}) = 0, m(\tilde{\chi}_{i}^{0})$	a abarrat
	$\tilde{\chi}^{\dagger}\tilde{\chi}_{-}^{0} \rightarrow W^{(*)}\tilde{\chi}^{0}Z^{(*)}\tilde{\chi}^{0}$ : 3 lep + $E_{T,miss}$		mass $(m(\chi^2) = m(\chi^2), m(\chi^2) = 0$ , sleptons decoupled)	s above)
	Direct $\chi^{-}$ pair prod. (AMSB) : long-lived $\chi^{-}$	L=4.7 fb <sup>-1</sup> , 7 TeV [1210.2852] 220 GeV $\tilde{\chi}^2$ mass	$(m(\chi_1) - m(\chi_2), m(\chi_1) = 0, \text{ steptons decoupled})$ $(1 < \tau(\tilde{\chi}^2) < 10 \text{ ns})$	
pe s	Stable ζ, R-hadrons : low β, βy	L=4.7 fb <sup>-1</sup> , 7 TeV [1210.2032] 220 GeV X <sub>1</sub> [11G33	985 Gev g mass	
cle liv	GMSB, stable τ : low β	L=4.7 fb <sup>-1</sup> , 7 TeV [1211.1597] 300 GeV $\tilde{\zeta}$ [1]		
Long-lived particles	GMSB, $\tilde{\chi}^0 \rightarrow \gamma \tilde{G}$ : non-pointing photons	L=4.7 fb <sup>-1</sup> .7 TeV [ATLAS-CONF-2013-016] 230 GeV $\tilde{\chi}_{1}^{0}$ mass		
9 d	$\tilde{\chi}^0 \rightarrow qq\mu (RPV): \mu + heavy displaced vertex$	L=4.4 fb <sup>-1</sup> , 7 TeV [1210.7451]	700 GeV G MASS (1 mm < ct < 1 m, g decoupled)	
	$\chi_{,} \rightarrow qq\mu (RFV) : \mu + heavy displaced vertexLFV : pp \rightarrow \bar{v}, +X, \bar{v}, \rightarrow e+\mu resonance$	L=4.6 fb <sup>-1</sup> , 7 TeV [1210.7451] L=4.6 fb <sup>-1</sup> , 7 TeV [1212.1272]	1.61 TeV V, MASS (//m/sct //m/gdecoped)	1 -2020
	LFV : $pp \rightarrow \tilde{v}_{\tau} + X, \tilde{v}_{\tau} \rightarrow e^{\mu}$ resonance	L=4.6 fb <sup>-1</sup> , 7 TeV [1212.1272]	1.10 TeV V, MASS (λ <sub>10</sub> =0.10, λ <sub>100</sub> =0	
RPV	Bilinear RPV CMSSM : 1 lep + 7 j's + E <sub>T miss</sub>	L#4.6 fb . 7 feV [1212.1272] L#4.7 fb <sup>-1</sup> , 7 TeV [ATLAS-CONF-2012-140]	1,2 TeV Q = Q MASS (x <sub>211</sub> =0.10, x <sub>1232</sub> =0 1,2 TeV Q = Q MASS (ct <sub>120</sub> < 1 mm)	(03)
	$\tilde{\chi}_{,,\tilde{\chi}_{,,\tilde{\chi}_{,,\tilde{\chi}_{,,\tilde{\chi}_{,,\tilde{\chi}_{,\tilde{\chi}},\tilde{\chi}_{,\tilde{\chi}},\tilde{\chi}},\tilde{\chi}},\tilde{\chi}}}}}}}}}}}}}}}}}$	L=20.7 fb <sup>-1</sup> , 8 TeV [ATLAS-CONF-2012-140] L=20.7 fb <sup>-1</sup> , 8 TeV [ATLAS-CONF-2013-036]	760 GeV $\widetilde{\chi}_{1}^{+}$ mass $(m(\widetilde{\chi}_{1}^{0}) > 300 \text{ GeV}, \lambda_{m} > 0)$	
	$\tilde{\chi}, \tilde{\chi},, \tilde{\chi}, \rightarrow \tau\tau v_e, e\tau v_\tau$ : 3 lep + $1\tau + E_{T,miss}$		$\chi_1^+$ mass $(m(\chi_1^+) > 300 \text{ GeV}, \chi_{121} > 0)$ $\chi_1^+$ mass $(m(\chi_1^+) > 80 \text{ GeV}, \chi_{121} > 0)$	
	$\chi_1\chi_1,, \chi_1 \rightarrow crv_e, erv_1 = 5 \text{ lep } + 17 + 2_{T,miss}$ $\tilde{q} \rightarrow qqq : 3  let resonance pair$	L=4.6 fb <sup>-1</sup> , 7 TeV [1210.4813]	666 GeV @ mass	
	$g \rightarrow qqq$ : 3-jet resonance pair $\tilde{g}\rightarrow \tilde{t}t, \tilde{t}\rightarrow bs$ : 2 SS-lep + (0-3b-)j's + E	L=4.6 fb <sup>-1</sup> , 7 TeV [1210.4813] L=20.7 fb <sup>-1</sup> , 8 TeV [ATLAS-CONF-2013-007]	666 GeV g mass 880 GeV g mass (any m(t))	
	g→it, t→bs : 2 33-lep + (0-30-)) s + E Scalar gluon : 2-jet resonance pair		ION MASS (incl. limit from 1110.2693)	
WIN	P interaction (D5, Dirac $\chi$ ): 'monojet' + E	L=4.6 fb <sup>-1</sup> , 7 TeV [1210.4526] 100-287 GeV SGIU L=10.5 fb <sup>-1</sup> , 8 TeV [ATLAS-CONF-2012-147]	704 GeV M* scale (m. < 80 GeV, limit of < 687 Ge	10-20
				v tor Us)
		10-1		
		10 <sup>-1</sup>	1	10

 $^{*}Only$  a selection of the available mass limits on new states or phenomena shown. All limits quoted are observed minus 1 $\sigma$  theoretical signal cross section uncertainty.

Mass scale [TeV]



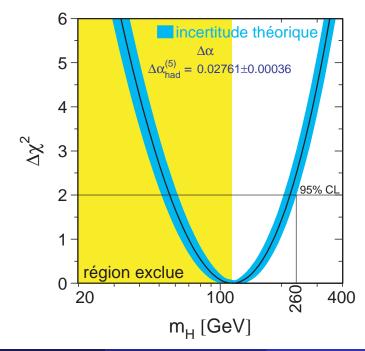
I. Antoniadis (CERN)

#### Remarks on the value of the Higgs mass $\sim 125~\text{GeV}$

- consistent with expectation from precision tests of the SM
- favors perturbative physics quartic coupling  $\lambda = m_H^2/v^2 \simeq 1/8$

#### Window to new physics

- compatible with supersymmetry
  - but appears fine-tuned in its minimal version [8]
  - early to draw a general conclusion before LHC13/14
  - e.g. an extra singlet or split families can alleviate the fine tuning [9]
- very important to measure its properties and couplings [13] any deviation of its couplings to top, bottom and EW gauge bosons implies new light states involved in the EWSB altering the fine-tuning



### Fine-tuning in MSSM

Upper bound on the lightest scalar mass:

$$m_h^2 \lesssim m_Z^2 \cos^2 2\beta + \frac{3}{(4\pi)^2} \frac{m_t^4}{v^2} \left[ \ln \frac{m_{\tilde{t}}^2}{m_t^2} + \frac{A_t^2}{m_{\tilde{t}}^2} \left( 1 - \frac{A_t^2}{12m_{\tilde{t}}^2} \right) \right] \lesssim (130 \, GeV)^2$$

 $m_h \simeq 126 \,\, {
m GeV} \, \Rightarrow \, m_{ ilde{t}} \simeq 3 \,\, {
m TeV}$  or  $A_t \simeq 3 m_{ ilde{t}} \simeq 1.5 \,\, {
m TeV}$ 

 $\Rightarrow$  % to a few ‰ fine-tuning

minimum of the potential: 
$$m_Z^2=2rac{m_1^1-m_2^2 an_\beta^2}{ an^2eta-1}\sim -2m_2^2+\cdots$$

 $\begin{array}{ll} \mathsf{RG evolution:} & m_2^2 = & m_2^2(M_{\mathrm{GUT}}) - \frac{3\lambda_t^2}{4\pi^2}m_{\tilde{t}}^2\ln\frac{M_{\mathrm{GUT}}}{m_{\tilde{t}}} + \cdots \ {}_{[25]} \\ & \sim & m_2^2(M_{\mathrm{GUT}}) - \mathcal{O}(1)m_{\tilde{t}}^2 + \cdots \end{array} \qquad {}_{[6]} \end{array}$ 

#### MSSM with dim-5 and 6 operators

I.A.-Dudas-Ghilencea-Tziveloglou '08, '09, '10

parametrize new physics above MSSM by higher-dim effective operators

relevant super potential operators of dimension-5:

$$\mathcal{L}^{(5)} = \frac{1}{M} \int d^2 \theta \left( \eta_1 + \eta_2 S \right) \left( H_1 H_2 \right)^2$$

 $\eta_1$ : generated for instance by a singlet

$$W = \lambda \sigma H_1 H_2 + M \sigma^2 \quad \rightarrow \quad W_{\text{eff}} = \frac{\lambda^2}{M} (H_1 H_2)^2$$

Strumia '99 ; Brignole-Casas-Espinosa-Navarro '03 Dine-Seiberg-Thomas '07

 $\eta_1$ : corresponding soft breaking term spurion  $S \equiv m_S \theta^2$ 

# Physical consequences of MSSM<sub>5</sub>: Scalar potential

$$\begin{split} \mathcal{V} &= \ m_1^2 |h_1|^2 + m_2^2 |h_2|^2 + B \mu (h_1 h_2 + \mathrm{h.c.}) + \frac{g_2^2 + g_Y^2}{8} \left( |h_1|^2 - |h_2|^2 \right)^2 \\ &+ \left( |h_1|^2 + |h_2|^2 \right) \left( \eta_1 h_1 h_2 + \mathrm{h.c.} \right) + \frac{1}{2} \left[ \eta_2 (h_1 h_2)^2 + \mathrm{h.c.} \right] + \mathcal{O} \left( \eta_i^2 \right) \end{split}$$

- $\eta_{1,2} \Rightarrow$  quartic terms along the D-flat direction  $|h_1| = |h_2|$
- potential stability  $\Rightarrow \eta_2 \ge 4|\eta_1|$

requiring  $\eta$ -corrections to be smaller than MSSM mass matrix elements  $\Rightarrow$ only  $\eta_2$  can change the tree-level bound  $m_h \leq m_Z$  but marginally Relaxing the condition on potential positivity: guaranteed by dim-6 ops

only one dim-6 along the D-flat direction induced by dim-5:  $\propto \eta_1^2$ 

$$W = \eta_1(H_1H_2)^2 \longrightarrow V = \left|\frac{\partial W}{\partial H_i}\right|^2 \sim \eta_1^2 |H_1H_2|^2 \left(|H_1|^2 + |H_2|^2\right)$$

- tree-level mass can increase significantly
- bigger parameter space for LSP being dark matter

Bernal-Blum-Nir-Losada '09

# MSSM Higss with dim-6 operators

#### dim-6 operators can have an independent scale from dim-5

Classification of all dim-6 contributing to the scalar potential (without SUSY)  $\Rightarrow$ 

large tan  $\beta$  expansion:  $\delta_6 m_h^2 = f v^2 + \cdots$ constant receiving contributions from several operators

$$f \sim f_0 imes \left( \mu^2/M^2, \ m_S^2/M^2, \ \mu m_S/M^2, \ v^2/M^2 
ight)$$

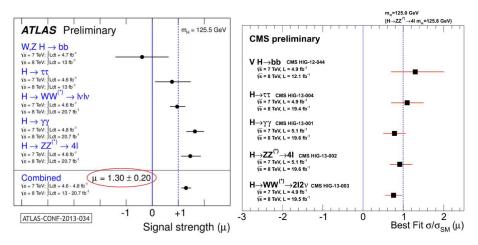
 $m_S=1$  TeV, M=10 TeV,  $f_0\sim 1-2.5$  for each operator

 $\Rightarrow m_h \simeq 103 - 119 \text{ GeV}$ 

 $\Rightarrow$  MSSM with dim-5 and dim-6 operators:

possible resolution of the MSSM fine-tuning problem [6]

### Couplings of the new boson vs SM

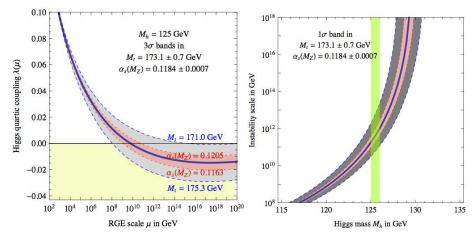


exclusion : spin 2 and pseudoscalar at  $\gtrsim 95\%$  CL

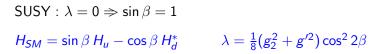
Agreement with Standard Model expectation at  $\sim 2\,\sigma$ 

### Can the SM be valid at high energies?

Degrassi-Di Vita-Elias Miró-Espinosa-Giudice-Isidori-Strumia '12

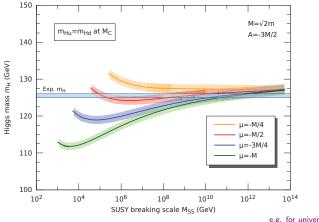


Instability of the SM Higgs potential  $\Rightarrow$  metastability of the EW vacuum



 $\lambda=0$  at a scale  $\geq 10^{10}~{
m GeV} \Rightarrow m_{H}=126\pm 3~{
m GeV}$ 

Ibanez-Valenzuela '13



e.g. for universal  $\sqrt{2}m=M=M_{SS},\;A=-3/2M$ 

If the weak scale is tuned  $\Rightarrow$  split supersymmetry is a possibility Arkani Hamed-Dimopoulos '04, Giudice-Romaninio '04

- natural splitting: gauginos, higgsinos carry R-symmetry, scalars do not
- main good properties of SUSY are maintained gauge coupling unification and dark matter candidate
- also no dangerous FCNC, CP violation, ...
- experimentally allowed Higgs mass  $\Rightarrow$  'mini' split

 $m_S \sim {
m few}$  - thousands TeV

gauginos: a loop factor lighter than scalars ( $\sim m_{3/2}$ )

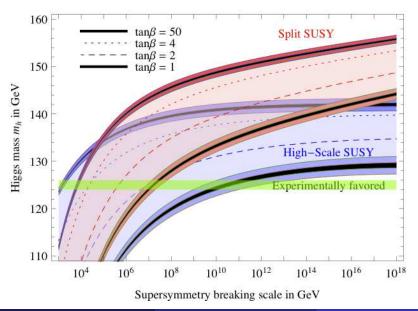
• natural string framework: intersecting (or magnetized) branes

IA-Dimopoulos '04

D-brane stacks are supersymmetric with massless gauginos intersections have chiral fermions with broken SUSY & massive scalars

#### Giudice-Strumia '11

#### Predicted range for the Higgs mass

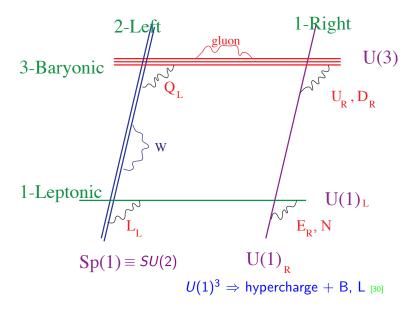


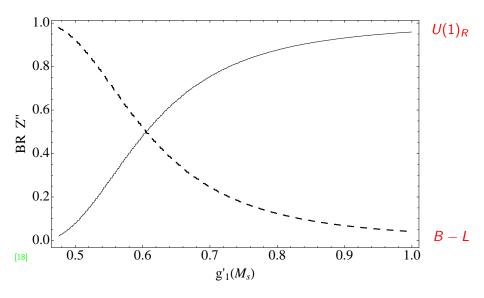
# An extra U(1) can also cure the instability problem Anchordoqui-IA-Goldberg-Huang-Lüst-Taylor-Vicek '12

usually associated to known global symmetries of the SM:  $B, L, \ldots$ 

- B anomalous and superheavy
- B L massless at the string scale (no associated 6d anomaly) but broken at TeV by a scalar VEV with the quantum numbers of  $N_R$
- L-violation from higher-dim operators suppressed by the string scale
- U(3) unification, Y combination  $\Rightarrow$  2 parameters: 1 coupling +  $m_{Z''}$
- perturbativity  $\Rightarrow 0.5 \lesssim g_{U(1)_R} \lesssim 1$  [20]
- interesting LHC phenomenology and cosmology [21]

#### Standard Model on D-branes : SM<sup>++</sup>





- Rotation of U(1)'s from the string to low energy basis Z, Z', Z'':
   completely fixed in terms of the couplings
  - Decoupling of anomalous  $Z' \simeq B$
  - Z'' linear combination of B L and  $U(1)_R$
- Recent cosmological observations indicate extra relativistic component dark radiation parametrized by an effective  $\nu$ -number close to 4 \*  $\rightarrow$  use the 3  $\nu_R$ 's interacting with SM fermions via Z'' data: their decoupling during the quark-hadron transition

 $\Rightarrow$  3.5  $\lesssim M_{Z''} \lesssim$  7 TeV (within LHC14 discovery potential) \* before Planck results

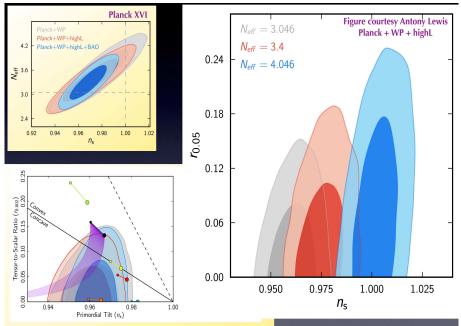


Fig. 1. Marginalized joint 68% and 95% CL regions for  $n_s$  and  $r_{0.002}$  from *Planck* in combination with other data sets compared to the theoretical predictions of selected inflationary models.

Scalar potential:

 $V(H, H'') = \mu^{2} |H|^{2} + {\mu'}^{2} |H''|^{2} + \lambda_{1} |H|^{4} + \lambda_{2} |H''|^{4} + \lambda_{3} |H|^{2} |H''|^{2}$ 

5 parameters  $\Rightarrow$  v, m<sub>h</sub>, v", m<sub>h"</sub> + a scalar mixing angle  $\alpha$ 

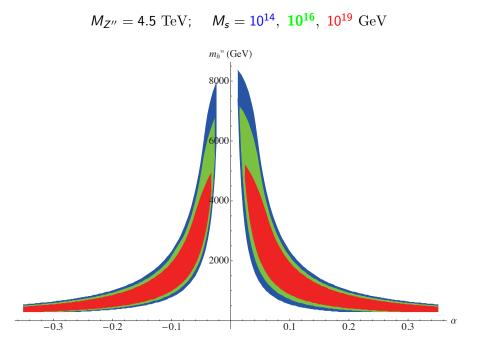
 $\Rightarrow$  3 free parameters :  $m_{h''}, \alpha, v'' \leftrightarrow M_{Z''}$ 

Stability conditions:  $\lambda_1 > 0$ ,  $\lambda_2 > 0$ ,  $\lambda_1 \lambda_2 > \frac{1}{4} \lambda_3^2$ 

RGE analysis up to  $M_s \Rightarrow$  stability is possible in SM<sup>++</sup>

for  $0.02 \lesssim |\alpha| \lesssim 0.35$  and 500 GeV  $\lesssim m_{h''} \lesssim 5$  TeV

AAGHLTV '12



#### Alternative answer: Low UV cutoff $\Lambda \sim \text{TeV}$

- low scale gravity  $\Rightarrow$  extra dimensions: large flat or warped
- low string scale  $\Rightarrow$  low scale gravity, ultra weak string coupling

 $M_s \sim 1 \text{ TeV} \Rightarrow \text{volume } R_{\perp}^n = 10^{32} l_s^n \ (R_{\perp} \sim .1 - 10^{-13} \text{ mm for } n = 2 - 6)$ 

- spectacular model independent predictions

- radical change of high energy physics at the TeV scale

Moreover no little hierarchy problem:

radiative electroweak symmetry breaking with no logs [8]

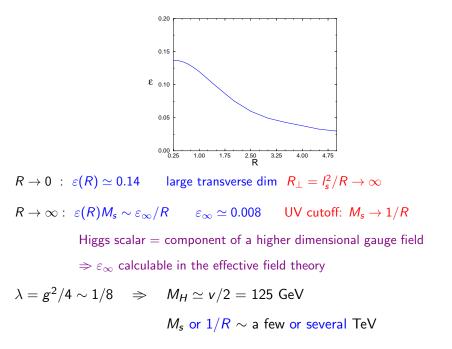
 $\Lambda \sim$  a few TeV and  $m_{H}^{2} =$  a loop factor  $imes \Lambda^{2}$ 

But unification has to be probably dropped

New Dark Matter candidates e.g. in the extra dims

# Origin of EW symmetry breaking?

possible answer: radiative breaking I.A.-Benakli-Quiros '00  $V = \mu^2 H^{\dagger} H + \lambda (H^{\dagger} H)^2$  $\mu^2 = 0$  at tree but becomes < 0 at one loop non-susy vacuum simplest case: one scalar doublet from the same brane  $\Rightarrow$  tree-level V same as susy:  $\lambda = \frac{1}{8}(g_2^2 + g'^2)$ D-terms  $\mu^2 = -g^2 \varepsilon^2 M_s^2 \leftarrow \text{effective UV cutoff}$  $e^{2}(R) = \frac{R^{3}}{2\pi^{2}} \int_{0}^{\infty} dll^{3/2} \frac{\theta_{2}^{4}}{16l^{4}\eta^{12}} \left(il + \frac{1}{2}\right) \sum n^{2} e^{-2\pi n^{2}R^{2}l}$ 



#### Accelerator signatures: 4 different scales

- Gravitational radiation in the bulk  $\Rightarrow$  missing energy present LHC bounds:  $M_* \gtrsim 3-5$  TeV
- Massive string vibrations  $\Rightarrow$  e.g. resonances in dijet distribution

 $M_j^2 = M_0^2 + M_s^2 j$ ; maximal spin: j + 1

higher spin excitations of quarks and gluons with strong interactions present LHC limits:  $M_s\gtrsim 5~{
m TeV}$ 

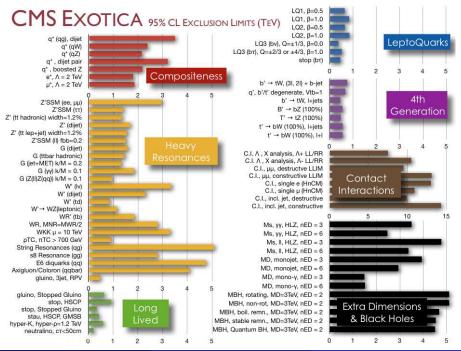
• Large TeV dimensions  $\Rightarrow$  KK resonances of SM gauge bosons I.A. '90

$$M_k^2 = M_0^2 + k^2/R^2$$
;  $k = \pm 1, \pm 2, \dots$ 

experimental limits:  $R^{-1} \gtrsim 0.5 - 4$  TeV (UED - localized fermions)

• extra U(1)'s and anomaly induced terms

masses suppressed by a loop factor from  $M_s$  [30]



# Extra U(1)'s and anomaly induced terms

#### masses suppressed by a loop factor

usually associated to known global symmetries of the SM

(anomalous or not) such as (combinations of)

Baryon and Lepton number, or PQ symmetry

Two kinds of massive U(1)'s: I.A.-Kiritsis-Rizos '02

- 4d anomalous U(1)'s:  $M_A \simeq g_A M_s$
- 4d non-anomalous U(1)'s: (but masses related to 6d anomalies)

 $M_{NA} \simeq g_A M_s V_2 \leftarrow (6d \rightarrow 4d)$  internal space  $\Rightarrow M_{NA} \ge M_A$ 

or massless in the absence of such anomalies [19]

#### TeV string scale Anchordogui-IA-Goldberg-Huang-Lüst-Taylor '11

- B and L become massive due to anomalies Green-Schwarz terms
- the global symmetries remain in perturbation
  - Baryon number  $\Rightarrow$  proton stability
  - Lepton number  $\Rightarrow$  protect small neutrino masses

- Lepton number  $\Rightarrow$  process \_ no Lepton number  $\Rightarrow \frac{1}{M_s}LLHH \rightarrow$  Majorana mass:  $\frac{\langle H \rangle^2}{M_s}LL$  $\swarrow \sim$  GeV

•  $B, L \Rightarrow$  extra Z's

with possible leptophobic couplings leading to CDF-type Wij events  $Z' \simeq B$  lighter than 4d anomaly free  $Z'' \simeq B - L$ 

# Conclusions

- Confirmation of the EWSB scalar at the LHC: important milestone of the LHC research program
- Precise measurement of its couplings is of primary importance
- Hint on the origin of mass hierarchy and of BSM physics
  - natural or unnatural SUSY?
  - low string scale in some realization?
  - something new and unexpected?

all options are still open

• LHC enters a new era with possible new discoveries