Mass hierarchy and electroweak symmetry breaking

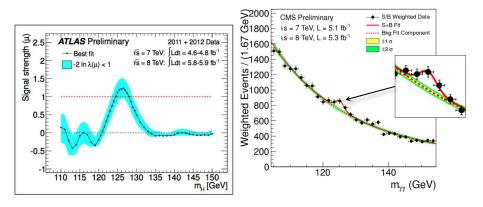
I. Antoniadis



Beyond the Standard Model after the first run of LHC GGI, Florence, 9-12 July 2013

- Low energy SUSY and 125 GeV Higgs
- Live with the hierarchy
- Extra *U*(1)'s
- Low scale strings and extra dimensions

Higgs boson discovery



 $m_H = 125.5 \pm 0.2 \,(\text{stat.}) \pm 0.5 \,(\text{syst.})$

 $m_H = 125.7 \pm 0.3 \pm 0.3$ GeV

Beyond the Standard Model of Particle Physics: driven by the mass hierarchy problem

Standard picture: low energy supersymmetry

Natural framework: Heterotic string (or high-scale M/F) theory

Advantages:

- natural elementary scalars
- gauge coupling unification
- LSP: natural dark matter candidate
- radiative EWSB

Problems:

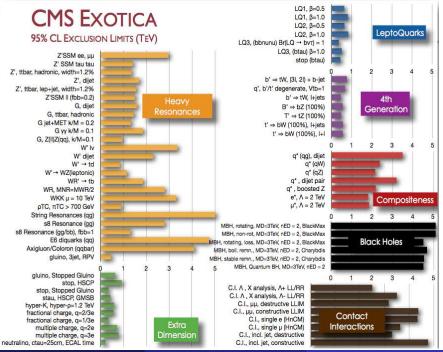
- too many parameters: soft breaking terms
- MSSM : already a % ‰ fine-tuning 'little' hierarchy problem

ATLAS SUSY Searches* - 95% CL Lower Limits (Status: March 26, 2013)

	MSUGRA/CMSSM : 0 lep + j's + E T miss	L+5.8 fb ⁻¹ , 8 TeV [ATLAS-CONF-2012-109]	t.so TeV g̃= g̃ mass	
	MSUGRA/CMSSM : 1 lep + j's + E - min	L=5.8 fb ⁻¹ , 8 TeV (ATLAS-CONF-2012-104)	1.24 TeV g = g mass	
60	Pheno model : 0 lep + j's + E T miss	L+5.8 fb ⁻¹ .8 TeV (ATLAS-CONF-2012-109)	1.18 TeV Q MASS (m(q) < 2 TeV, light 7	S ATLAS
16.	Pheno model : 0 lep + j's + E Trais	L#5.8 fb ⁻¹ , 8 TeV (ATLAS-CONF-2012-109)	1.38 TeV Q mass (m(g) < 2 TeV, lig	
2	Gluino med. $\tilde{\chi}^{\pm}(\tilde{g} \rightarrow q\bar{q}\tilde{\chi}^{\pm})$: 1 lep + j's + $E_{T miss}$	L=4.7 fb ⁻¹ , 7 TeV [1208.4688]	900 GeV \tilde{g} mass $(m(\tilde{\chi}^{0}) < 200 \text{ GeV}, m(\tilde{\chi}^{1}) =$	
993	GMSB (INI SP) : 2 len (OS) + i'e + E	L=4.7 fb ⁻¹ .7 TeV [1208.4688]	1.24 TeV Q MASS (tang < 15)	fore Lundan
Inclusive searches	GMSB (ÎNLSP) : 2 lep (OS) + j's + Ε GMSB (τ NLSP) : 1-2 τ + j's + Ε	L=20.7 fb ⁻¹ .8 TeV (1210.1314)	1.40 TeV ĝ mass (tanβ > 18)	
SIV	GGM (bino NLSP) : yy + E ^{+mss}	L=4.8 fb ⁻¹ , 7 TeV [1209.0753]	1.07 TeV Q MASS (m(q)) > 50 GeV)	ſ .
13	GGM (wino NLSP) : y + len + F	L#4.8 fb ⁻¹ , 7 TeV (ATLAS-CONF-2012-144)	619 GeV g mass	Ldt = (4.4 - 20.7) fb ⁻¹
ŝ	GGM (higgsing-bing NLSP) : $y + b + E$	L=4.8 fb ⁻¹ , 7 TeV [1211.1167]	900 GeV Q MASS (m(g) > 220 GeV)	J
	GGM (higgsino NLSP) : Z + jets + E Timiss	L+5.8 fb ⁻¹ .8 TeV IATLAS-CONF-2012-1521	690 GeV Q MASS (m(H) > 200 GeV)	s = 7, 8 TeV
	Gravitino LSP : 'monojet' + E Trias	L=10.5 fb ⁻¹ , 8 TeV [ATLAS-CONF-2012-147]	645 GeV F ^{1/2} scale (m(G) > 10 ⁻⁴ eV)	
~	$\tilde{g} \rightarrow b \tilde{b} \tilde{\chi}_{1}^{0}$: 0 lep + 3 b-j's + $E_{T,max}$	L=12.8 fb ⁻¹ .8 TeV (ATLAS-CONF-2012-145)	1.24 TeV \tilde{Q} mass $(m(\tilde{\chi}^2) < 200 \text{ GeV})$	
90. No	$\tilde{g} \rightarrow t \tilde{t} \tilde{\chi}$: 2 SS-lep + (0-3b-)j's + $E_{T miss}$	L=20.7 fb ⁻¹ , 8 TeV [ATLAS-CONF-2013-007]	900 GeV ĝ mass (any m(x))	8 TeV, all 2012 data
l gi	$\tilde{g} \rightarrow tt \tilde{\chi}_{1}^{0}$: 0 lep + multi-j's + $E_{T,miss}$	L=5.8 fb ⁻¹ .8 TeV (ATLAS-CONF-2012-103)	1.00 TeV Q Mass (m(\overline{z}^{0}) < 300 GeV)	
3rd gen. gluino mediated	g→tt ² _x 0 lep + 1 lef s + E _{T miss}	L=12.8 fb ⁻¹ , 8 TeV [ATLAS-CONF-2012-145]	1.15 TeV Q MASS (mQ ²) < 200 GeV)	8 TeV, partial 2012 data
	$bb, b \rightarrow b\chi^{-1} = 0 \text{ lep } + 2\text{ -b-jets } + E_{T \text{ mass}}$	L=12.8 fb ⁻¹ , 8 TeV [ATLAS-CONF-2012-165]	620 GeV b mass (m(x)) < 120 GeV)	7 TeV, all 2011 data
(0, ~	$\tilde{b}\tilde{b}, \tilde{b}_{1} \rightarrow t\tilde{\chi}_{1}^{1}$: 2 SS-lep + (0-3b-)]'s + $E_{T,miss}$	L=20.7 fb ⁻¹ , 8 TeV [ATLAS-CONF-2013-007] 430 GeV		r ice, air 2011 data
ior in	tt (light), $t \rightarrow b\tilde{\chi}_1^+$: 1/2 lep (+ b-jet) + E Triss) = 55 GeV)	
not	tt (medium), $t \rightarrow b\tilde{\chi}_{\pm}^{+}$: 12 lep (+ b-jet) + $E_{T,miss}$ tt (medium), $t \rightarrow b\tilde{\chi}_{\pm}^{+}$: 1 lep + b-jet + $E_{T,miss}$	L=20.7 fb ⁻¹ , 8 TeV [ATLAS-CONF-2013-037] 160 GeV [111033 [111]		
od s	tt (medium), $t \rightarrow b\chi_{z}^{*}$: 2 lep + $E_{T,max}$	L=13.0 fb ⁻¹ .8 TeV [ATLAS-CONF-2012-167] 160-440 Ge	\tilde{t} mass $(m(\chi^2) = 0 \text{ GeV}, m(\chi^2) = 10 \text{ GeV})$	
3rd gen. squarks direct production	$\tilde{t}t$ (heavy), $\tilde{t} \rightarrow t \tilde{\chi}_{1}^{0}$: 1 lep + b-jet + $E_{T miss}$		-610 GeV t mass (m(2 [*]) = 0)	
- di			20-660 GeV T mass $(m(\chi^2) = 0)$	
3rc	tt (natural GMSB) : $Z(\rightarrow II) + b_{jet} + E_{T,miss}$		Gev t mass (m(2)) > 150 GeV)	
	$\tilde{t}_1 \tilde{t}_2, \tilde{t}_2 \rightarrow \tilde{t}_1 + Z : Z(\rightarrow II) + 1 \text{ lep } + b \text{-jet} + E_{T,miss}$		GeV L, mass (m(L) = m(2 ⁴) + 180 GeV)	
	$[1,] \rightarrow [\chi]$: 2 lep + E _{1 max}		$m(\tilde{\chi}^2) = 0)$	
	$\chi_{1}^{+}\chi_{1}^{-}\chi_{2}^{+}\chi_{1}^{+}\chi$		$m(\chi_1) = 0$ $mass (m(\chi_1^2) \le 10 \text{ GeV}, m(\tilde{x}) = \frac{1}{2}(m(\chi_1^2) + m(\chi_1^2)))$	
EW direct	$\chi_{\chi_{\chi_{\chi_{\chi_{\chi_{\chi_{\chi_{\chi_{\chi_{\chi_{\chi_{\chi_{\chi$	L+20.7 fb ⁻¹ , 8 TeV [ATLAS-CONF-2013-028] 180-330 GeV Ž	mass $(m(\tilde{\chi}_{1}^{2}) < 10 \text{ GeV}, m(\tilde{\chi}_{1}^{2}) = \frac{1}{2}(m(\tilde{\chi}_{1}^{2}) + m(\tilde{\chi}_{1}^{2})))$	
비송	$\chi^{+}\chi^{-}\chi^{+}\to \tilde{V}(\tilde{V}): 2 \operatorname{lep} + E_{\tau,\operatorname{miss}}^{\tau,\operatorname{miss}}$ $\chi^{+}\chi^{+}\chi^{+}\to \tilde{V}(\tau\tilde{V}): 2\tau + E_{\tau,\operatorname{miss}}$ $\chi^{+}\chi^{0}_{-}\to \tilde{V}(\tilde{U}(\tilde{V}V), \tilde{V})[(\tilde{V}V): 3 \operatorname{lep} + E_{\tau,\operatorname{miss}}^{\tau,\operatorname{miss}}$		600 GeV $\tilde{\chi}_{i}^{\pm}$ mass $(m(\tilde{\chi}_{i}^{+}) = m(\tilde{\chi}_{i}^{0}), m(\tilde{\chi}_{i}^{0}) = 0, m(\tilde{\chi}_{i}^{0})$	a abarrat
	$\tilde{\chi}^{\dagger}\tilde{\chi}_{-}^{0} \rightarrow W^{(*)}\tilde{\chi}^{0}Z^{(*)}\tilde{\chi}^{0}$: 3 lep + $E_{T,miss}$		mass $(m(\chi^2) = m(\chi^2), m(\chi^2) = 0$, sleptons decoupled)	s above)
	Direct χ^{-} pair prod. (AMSB) : long-lived χ^{-}	L=4.7 fb ⁻¹ , 7 TeV [1210.2852] 220 GeV $\tilde{\chi}^2$ mass	$(m(\chi_1) - m(\chi_2), m(\chi_1) = 0, \text{ steptons decoupled})$ $(1 < \tau(\tilde{\chi}^2) < 10 \text{ ns})$	
pe s	Stable ζ, R-hadrons : low β, βy	L=4.7 fb ⁻¹ , 7 TeV [1210.2032] 220 GeV X ₁ [11G33	985 Gev g mass	
cle liv	GMSB, stable τ : low β	L=4.7 fb ⁻¹ , 7 TeV [1211.1597] 300 GeV $\tilde{\zeta}$ [1]		
Long-lived particles	GMSB, $\tilde{\chi}^0 \rightarrow \gamma \tilde{G}$: non-pointing photons	L=4.7 fb ⁻¹ .7 TeV [ATLAS-CONF-2013-016] 230 GeV $\tilde{\chi}_{1}^{0}$ mass		
9 d	$\tilde{\chi}^0 \rightarrow qq\mu (RPV): \mu + heavy displaced vertex$	L=4.4 fb ⁻¹ , 7 TeV [1210.7451]	700 GeV G MASS (1 mm < ct < 1 m, g decoupled)	
	$\chi_{,} \rightarrow qq\mu (RFV) : \mu + heavy displaced vertexLFV : pp \rightarrow \bar{v}, +X, \bar{v}, \rightarrow e+\mu resonance$	L=4.6 fb ⁻¹ , 7 TeV [1210.7451] L=4.6 fb ⁻¹ , 7 TeV [1212.1272]	1.61 TeV V, MASS (//m/sct //m/gdecoped)	1 -2020
	LFV : $pp \rightarrow \tilde{v}_{\tau} + X, \tilde{v}_{\tau} \rightarrow e^{\mu}$ resonance	L=4.6 fb ⁻¹ , 7 TeV [1212.1272]	1.10 TeV V, MASS (λ ₁₀ =0.10, λ ₁₀₀ =0	
RPV	Bilinear RPV CMSSM : 1 lep + 7 j's + E _{T miss}	L#4.6 fb . 7 feV [1212.1272] L#4.7 fb ⁻¹ , 7 TeV [ATLAS-CONF-2012-140]	1,2 TeV Q = Q MASS (x ₂₁₁ =0.10, x ₁₂₃₂ =0 1,2 TeV Q = Q MASS (ct ₁₂₀ < 1 mm)	(03)
	$\tilde{\chi}_{,,\tilde{\chi}_{,,\tilde{\chi}_{,,\tilde{\chi}_{,,\tilde{\chi}_{,,\tilde{\chi}_{,\tilde{\chi}},\tilde{\chi}_{,\tilde{\chi}},\tilde{\chi}},\tilde{\chi}},\tilde{\chi}}}}}}}}}}}}}}}}}$	L=20.7 fb ⁻¹ , 8 TeV [ATLAS-CONF-2012-140] L=20.7 fb ⁻¹ , 8 TeV [ATLAS-CONF-2013-036]	760 GeV $\widetilde{\chi}_{1}^{+}$ mass $(m(\widetilde{\chi}_{1}^{0}) > 300 \text{ GeV}, \lambda_{m} > 0)$	
	$\tilde{\chi}, \tilde{\chi},, \tilde{\chi}, \rightarrow \tau\tau v_e, e\tau v_\tau$: 3 lep + $1\tau + E_{T,miss}$		χ_1^+ mass $(m(\chi_1^+) > 300 \text{ GeV}, \chi_{121} > 0)$ χ_1^+ mass $(m(\chi_1^+) > 80 \text{ GeV}, \chi_{121} > 0)$	
	$\chi_1\chi_1,, \chi_1 \rightarrow crv_e, erv_1 = 5 \text{ lep } + 17 + 2_{T,miss}$ $\tilde{q} \rightarrow qqq : 3 let resonance pair$	L=4.6 fb ⁻¹ , 7 TeV [1210.4813]	666 GeV @ mass	
	$g \rightarrow qqq$: 3-jet resonance pair $\tilde{g}\rightarrow \tilde{t}t, \tilde{t}\rightarrow bs$: 2 SS-lep + (0-3b-)j's + E	L=4.6 fb ⁻¹ , 7 TeV [1210.4813] L=20.7 fb ⁻¹ , 8 TeV [ATLAS-CONF-2013-007]	666 GeV g mass 880 GeV g mass (any m(t))	
	g→it, t→bs : 2 33-lep + (0-30-)) s + E Scalar gluon : 2-jet resonance pair		ION MASS (incl. limit from 1110.2693)	
WIN	P interaction (D5, Dirac χ): 'monojet' + E	L=4.6 fb ⁻¹ , 7 TeV [1210.4526] 100-287 GeV SGIU L=10.5 fb ⁻¹ , 8 TeV [ATLAS-CONF-2012-147]	704 GeV M* scale (m. < 80 GeV, limit of < 687 Ge	10-20
				v tor Us)
		10-1		
		10 ⁻¹	1	10

 $^{*}Only$ a selection of the available mass limits on new states or phenomena shown. All limits quoted are observed minus 1 σ theoretical signal cross section uncertainty.

Mass scale [TeV]



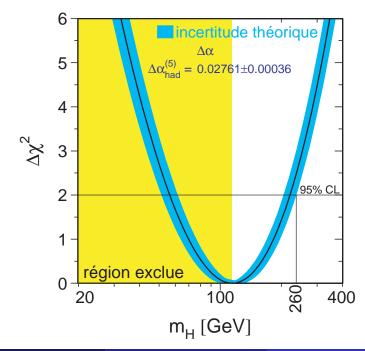
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Remarks on the value of the Higgs mass $\sim 125~\text{GeV}$

- consistent with expectation from precision tests of the SM
- favors perturbative physics quartic coupling $\lambda = m_H^2/v^2 \simeq 1/8$

Window to new physics

- compatible with supersymmetry
 - but appears fine-tuned in its minimal version [8]
 - early to draw a general conclusion before LHC13/14
 - e.g. an extra singlet or split families can alleviate the fine tuning [9]
- very important to measure its properties and couplings [13] any deviation of its couplings to top, bottom and EW gauge bosons implies new light states involved in the EWSB altering the fine-tuning



Fine-tuning in MSSM

Upper bound on the lightest scalar mass:

$$m_h^2 \lesssim m_Z^2 \cos^2 2\beta + \frac{3}{(4\pi)^2} \frac{m_t^4}{v^2} \left[\ln \frac{m_{\tilde{t}}^2}{m_t^2} + \frac{A_t^2}{m_{\tilde{t}}^2} \left(1 - \frac{A_t^2}{12m_{\tilde{t}}^2} \right) \right] \lesssim (130 \, GeV)^2$$

 $m_h \simeq 126 \,\, {
m GeV} \, \Rightarrow \, m_{ ilde{t}} \simeq 3 \,\, {
m TeV}$ or $A_t \simeq 3 m_{ ilde{t}} \simeq 1.5 \,\, {
m TeV}$

 \Rightarrow % to a few ‰ fine-tuning

minimum of the potential:
$$m_Z^2=2rac{m_1^1-m_2^2 an_\beta^2}{ an^2eta-1}\sim -2m_2^2+\cdots$$

 $\begin{array}{ll} \mathsf{RG evolution:} & m_2^2 = & m_2^2(M_{\mathrm{GUT}}) - \frac{3\lambda_t^2}{4\pi^2}m_{\tilde{t}}^2\ln\frac{M_{\mathrm{GUT}}}{m_{\tilde{t}}} + \cdots \ {}_{[25]} \\ & \sim & m_2^2(M_{\mathrm{GUT}}) - \mathcal{O}(1)m_{\tilde{t}}^2 + \cdots \end{array} \qquad {}_{[6]} \end{array}$

MSSM with dim-5 and 6 operators

I.A.-Dudas-Ghilencea-Tziveloglou '08, '09, '10

parametrize new physics above MSSM by higher-dim effective operators

relevant super potential operators of dimension-5:

$$\mathcal{L}^{(5)} = \frac{1}{M} \int d^2 \theta \left(\eta_1 + \eta_2 S \right) \left(H_1 H_2 \right)^2$$

 η_1 : generated for instance by a singlet

$$W = \lambda \sigma H_1 H_2 + M \sigma^2 \quad \rightarrow \quad W_{\text{eff}} = \frac{\lambda^2}{M} (H_1 H_2)^2$$

Strumia '99 ; Brignole-Casas-Espinosa-Navarro '03 Dine-Seiberg-Thomas '07

 η_1 : corresponding soft breaking term spurion $S \equiv m_S \theta^2$

Physical consequences of MSSM₅: Scalar potential

$$\begin{split} \mathcal{V} &= \ m_1^2 |h_1|^2 + m_2^2 |h_2|^2 + B \mu (h_1 h_2 + \mathrm{h.c.}) + \frac{g_2^2 + g_Y^2}{8} \left(|h_1|^2 - |h_2|^2 \right)^2 \\ &+ \left(|h_1|^2 + |h_2|^2 \right) \left(\eta_1 h_1 h_2 + \mathrm{h.c.} \right) + \frac{1}{2} \left[\eta_2 (h_1 h_2)^2 + \mathrm{h.c.} \right] + \mathcal{O} \left(\eta_i^2 \right) \end{split}$$

- $\eta_{1,2} \Rightarrow$ quartic terms along the D-flat direction $|h_1| = |h_2|$
- potential stability $\Rightarrow \eta_2 \ge 4|\eta_1|$

requiring η -corrections to be smaller than MSSM mass matrix elements \Rightarrow only η_2 can change the tree-level bound $m_h \leq m_Z$ but marginally Relaxing the condition on potential positivity: guaranteed by dim-6 ops

only one dim-6 along the D-flat direction induced by dim-5: $\propto \eta_1^2$

$$W = \eta_1(H_1H_2)^2 \longrightarrow V = \left|\frac{\partial W}{\partial H_i}\right|^2 \sim \eta_1^2 |H_1H_2|^2 \left(|H_1|^2 + |H_2|^2\right)$$

- tree-level mass can increase significantly
- bigger parameter space for LSP being dark matter

Bernal-Blum-Nir-Losada '09

MSSM Higss with dim-6 operators

dim-6 operators can have an independent scale from dim-5

Classification of all dim-6 contributing to the scalar potential (without SUSY) \Rightarrow

large tan β expansion: $\delta_6 m_h^2 = f v^2 + \cdots$ constant receiving contributions from several operators

$$f \sim f_0 imes \left(\mu^2/M^2, \ m_S^2/M^2, \ \mu m_S/M^2, \ v^2/M^2
ight)$$

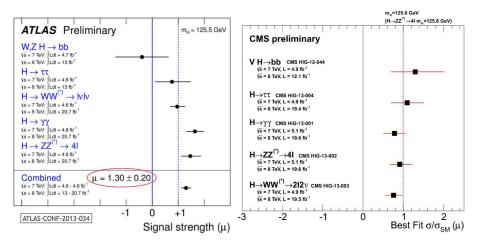
 $m_S=1$ TeV, M=10 TeV, $f_0\sim 1-2.5$ for each operator

 $\Rightarrow m_h \simeq 103 - 119 \text{ GeV}$

 \Rightarrow MSSM with dim-5 and dim-6 operators:

possible resolution of the MSSM fine-tuning problem [6]

Couplings of the new boson vs SM

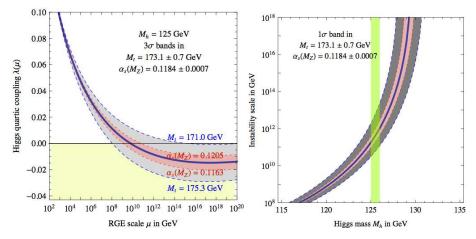


exclusion : spin 2 and pseudoscalar at $\gtrsim 95\%$ CL

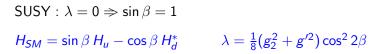
Agreement with Standard Model expectation at $\sim 2\,\sigma$

Can the SM be valid at high energies?

Degrassi-Di Vita-Elias Miró-Espinosa-Giudice-Isidori-Strumia '12

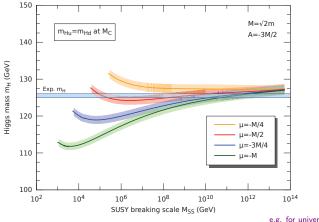


Instability of the SM Higgs potential \Rightarrow metastability of the EW vacuum



 $\lambda=0$ at a scale $\geq 10^{10}~{
m GeV} \Rightarrow m_{H}=126\pm 3~{
m GeV}$

Ibanez-Valenzuela '13



e.g. for universal $\sqrt{2}m=M=M_{SS},\;A=-3/2M$

If the weak scale is tuned \Rightarrow split supersymmetry is a possibility Arkani Hamed-Dimopoulos '04, Giudice-Romaninio '04

- natural splitting: gauginos, higgsinos carry R-symmetry, scalars do not
- main good properties of SUSY are maintained gauge coupling unification and dark matter candidate
- also no dangerous FCNC, CP violation, ...
- experimentally allowed Higgs mass \Rightarrow 'mini' split

 $m_S \sim {
m few}$ - thousands TeV

gauginos: a loop factor lighter than scalars ($\sim m_{3/2}$)

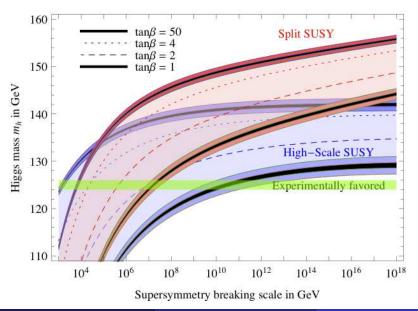
• natural string framework: intersecting (or magnetized) branes

IA-Dimopoulos '04

D-brane stacks are supersymmetric with massless gauginos intersections have chiral fermions with broken SUSY & massive scalars

Giudice-Strumia '11

Predicted range for the Higgs mass

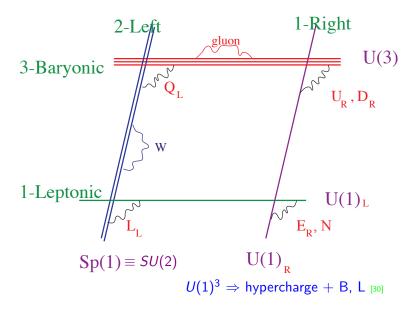


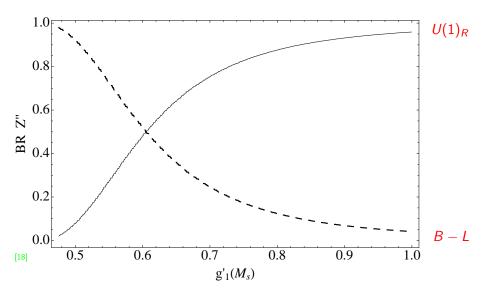
An extra U(1) can also cure the instability problem Anchordoqui-IA-Goldberg-Huang-Lüst-Taylor-Vicek '12

usually associated to known global symmetries of the SM: B, L, \ldots

- B anomalous and superheavy
- B L massless at the string scale (no associated 6d anomaly) but broken at TeV by a scalar VEV with the quantum numbers of N_R
- L-violation from higher-dim operators suppressed by the string scale
- U(3) unification, Y combination \Rightarrow 2 parameters: 1 coupling + $m_{Z''}$
- perturbativity $\Rightarrow 0.5 \lesssim g_{U(1)_R} \lesssim 1$ [20]
- interesting LHC phenomenology and cosmology [21]

Standard Model on D-branes : SM⁺⁺





- Rotation of U(1)'s from the string to low energy basis Z, Z', Z'':
 completely fixed in terms of the couplings
 - Decoupling of anomalous $Z' \simeq B$
 - Z'' linear combination of B L and $U(1)_R$
- Recent cosmological observations indicate extra relativistic component dark radiation parametrized by an effective ν -number close to 4 * \rightarrow use the 3 ν_R 's interacting with SM fermions via Z'' data: their decoupling during the quark-hadron transition

 \Rightarrow 3.5 $\lesssim M_{Z''} \lesssim$ 7 TeV (within LHC14 discovery potential) * before Planck results

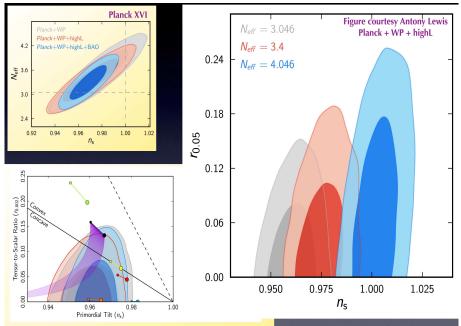


Fig. 1. Marginalized joint 68% and 95% CL regions for n_s and $r_{0.002}$ from *Planck* in combination with other data sets compared to the theoretical predictions of selected inflationary models.

Scalar potential:

 $V(H, H'') = \mu^{2} |H|^{2} + {\mu'}^{2} |H''|^{2} + \lambda_{1} |H|^{4} + \lambda_{2} |H''|^{4} + \lambda_{3} |H|^{2} |H''|^{2}$

5 parameters \Rightarrow v, m_h, v", m_{h"} + a scalar mixing angle α

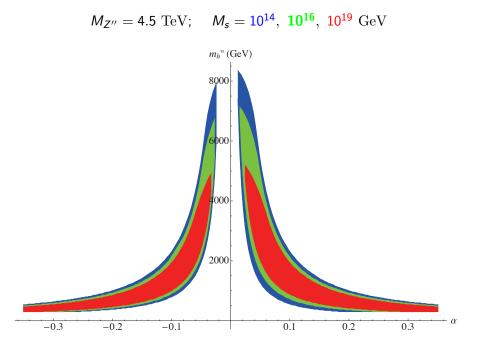
 \Rightarrow 3 free parameters : $m_{h''}, \alpha, v'' \leftrightarrow M_{Z''}$

Stability conditions: $\lambda_1 > 0$, $\lambda_2 > 0$, $\lambda_1 \lambda_2 > \frac{1}{4} \lambda_3^2$

RGE analysis up to $M_s \Rightarrow$ stability is possible in SM⁺⁺

for $0.02 \lesssim |\alpha| \lesssim 0.35$ and 500 GeV $\lesssim m_{h''} \lesssim 5$ TeV

AAGHLTV '12



Alternative answer: Low UV cutoff $\Lambda \sim \text{TeV}$

- low scale gravity \Rightarrow extra dimensions: large flat or warped
- low string scale \Rightarrow low scale gravity, ultra weak string coupling

 $M_s \sim 1 \text{ TeV} \Rightarrow \text{volume } R_{\perp}^n = 10^{32} l_s^n \ (R_{\perp} \sim .1 - 10^{-13} \text{ mm for } n = 2 - 6)$

- spectacular model independent predictions

- radical change of high energy physics at the TeV scale

Moreover no little hierarchy problem:

radiative electroweak symmetry breaking with no logs [8]

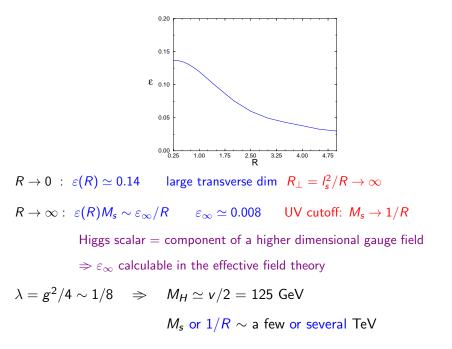
 $\Lambda \sim$ a few TeV and $m_{H}^{2} =$ a loop factor $imes \Lambda^{2}$

But unification has to be probably dropped

New Dark Matter candidates e.g. in the extra dims

Origin of EW symmetry breaking?

possible answer: radiative breaking I.A.-Benakli-Quiros '00 $V = \mu^2 H^{\dagger} H + \lambda (H^{\dagger} H)^2$ $\mu^2 = 0$ at tree but becomes < 0 at one loop non-susy vacuum simplest case: one scalar doublet from the same brane \Rightarrow tree-level V same as susy: $\lambda = \frac{1}{8}(g_2^2 + g'^2)$ D-terms $\mu^2 = -g^2 \varepsilon^2 M_s^2 \leftarrow \text{effective UV cutoff}$ $e^{2}(R) = \frac{R^{3}}{2\pi^{2}} \int_{0}^{\infty} dll^{3/2} \frac{\theta_{2}^{4}}{16l^{4}\eta^{12}} \left(il + \frac{1}{2}\right) \sum n^{2} e^{-2\pi n^{2}R^{2}l}$



Accelerator signatures: 4 different scales

- Gravitational radiation in the bulk \Rightarrow missing energy present LHC bounds: $M_* \gtrsim 3-5$ TeV
- Massive string vibrations \Rightarrow e.g. resonances in dijet distribution

 $M_j^2 = M_0^2 + M_s^2 j$; maximal spin: j + 1

higher spin excitations of quarks and gluons with strong interactions present LHC limits: $M_s\gtrsim 5~{
m TeV}$

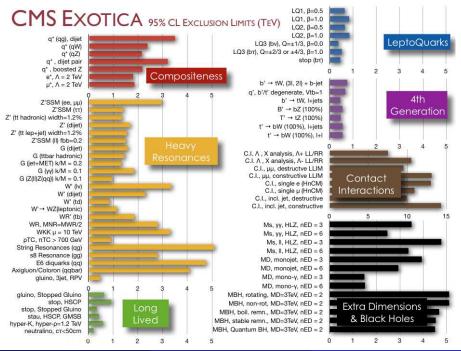
• Large TeV dimensions \Rightarrow KK resonances of SM gauge bosons I.A. '90

$$M_k^2 = M_0^2 + k^2/R^2$$
; $k = \pm 1, \pm 2, \dots$

experimental limits: $R^{-1} \gtrsim 0.5 - 4$ TeV (UED - localized fermions)

• extra U(1)'s and anomaly induced terms

masses suppressed by a loop factor from M_s [30]



Extra U(1)'s and anomaly induced terms

masses suppressed by a loop factor

usually associated to known global symmetries of the SM

(anomalous or not) such as (combinations of)

Baryon and Lepton number, or PQ symmetry

Two kinds of massive U(1)'s: I.A.-Kiritsis-Rizos '02

- 4d anomalous U(1)'s: $M_A \simeq g_A M_s$
- 4d non-anomalous U(1)'s: (but masses related to 6d anomalies)

 $M_{NA} \simeq g_A M_s V_2 \leftarrow (6d \rightarrow 4d)$ internal space $\Rightarrow M_{NA} \ge M_A$

or massless in the absence of such anomalies [19]

TeV string scale Anchordogui-IA-Goldberg-Huang-Lüst-Taylor '11

- B and L become massive due to anomalies Green-Schwarz terms
- the global symmetries remain in perturbation
 - Baryon number \Rightarrow proton stability
 - Lepton number \Rightarrow protect small neutrino masses

- Lepton number \Rightarrow process _ no Lepton number $\Rightarrow \frac{1}{M_s}LLHH \rightarrow$ Majorana mass: $\frac{\langle H \rangle^2}{M_s}LL$ $\swarrow \sim$ GeV

• $B, L \Rightarrow$ extra Z's

with possible leptophobic couplings leading to CDF-type Wij events $Z' \simeq B$ lighter than 4d anomaly free $Z'' \simeq B - L$

Conclusions

- Confirmation of the EWSB scalar at the LHC: important milestone of the LHC research program
- Precise measurement of its couplings is of primary importance
- Hint on the origin of mass hierarchy and of BSM physics
 - natural or unnatural SUSY?
 - low string scale in some realization?
 - something new and unexpected?

all options are still open

• LHC enters a new era with possible new discoveries