TOWARDS DIFFERENTIAL TOP PAIR PRODUCTION AT NNLO

Jet

Jet

Gabriel Abelof Northwestern University

Based on work done in collaboration with: O. Dekkers, A. Gehrmann-De Ridder, P. Meierhöfer, S. Pozzorini

HP2.5 - September 4, 2014 - GGI, Firenze

Top Pair Production At The LHC

- Interesting signal. Rich phenomenology. Important in new physics searches, PDF fits, ...
- Top quark pairs are copiously produced at the LHC

 $\sigma_{t\bar{t}+X}(\sqrt{s} = 7 \text{ TeV}) \sim 170 \text{ pb}$ $\sigma_{t\bar{t}+X}(\sqrt{s} = 8 \text{ TeV}) \sim 250 \text{ pb}$ $\sigma_{t\bar{t}+X}(\sqrt{s} = 14 \text{ TeV}) \sim 950 \text{ pb}$

- Abundant statistics. Expected experimental error ~5%
- Need theoretical predictions with similar accuracy
 - Requires computations through higher orders in perturbation theory

EUROPEAN ORGANIZATION FOR NUCLEAR RESEARCH (CERN) CERN-PH-EP/2013-234 2014/02/17 CMS-TOP-12-007 Measurement of the tt production cross section in the dilepton channel in pp collisions at $\sqrt{s} = 8$ TeV arXiv:1312.7582v2 [hep-ex] 14 Feb 2014 The CMS Collaboration* Abstract The top-antitop quark (tt) production cross section is measured in proton-proton collisions at $\sqrt{s} = 8$ TeV with the CMS experiment at the LHC, using a data sample corresponding to an integrated luminosity of 5.3 fb⁻¹. The measurement is performed by analysing events with a pair of electrons or muons, or one electron and one muon, and at least two jets, one of which is identified as originating from hadronisation of a bottom quark. The measured cross section is 239 ± 2 (stat.) ± 11 (syst.) ± 6 (lum.) pb for an assumed top-quark mass of 172.5 GeV, in agreement with the prediction of the standard model. Published in the Journal of High Energy Physics as doi:10.1007/JHEP02(2014)024. 239 ± 2 (stat.) ± 11 (syst.) ± 6 (lum.) © 2014 CERN for the benefit of the CMS Collaboration. CC-BY-3.0 license *See Appendix A for the list of collaboration members

Top Pair Production At The LHC: State Of The Art

•NLO QCD corrections: Ellis, Dawson, Nason; Beenakker, Kuijf, van Neerven, Smith '89

•NLO EW corrections: Beenakker, Bernreuther, Denner, Fuecker, Hollik, Kao, Kollar, Kühn, Ladinsky, Mertig, Moretti, Nolten, Ross, Sack, Scharf, Si, Uwer, Wackenroth, Yuan

• Threshold resummation and Coulomb corrections: Ahrens, Banfi, Berger, Bonciani, Catani, Contopanagos, Czakon, Ferroglia, Frixione, Kidonakis, Kiyo, Kühn, Laenen, Mangano, Mitov, Moch, Nason, Neubert, Pecjak Ridolfi, Steinhauser, Sterman, Uwer, Vogt, Yang

Yield a theoretical uncertainty of ~10%

To match theory and experimental accuracies at the LHC, cross sections for top pair production must be calculated through NNLO in pQCD

G. Abelof

Top Pair Production At The LHC: State Of The Art

- Calculation of the total NNLO cross section completed [Czakon, Fiedler, Mitov '13]
 - Combined with NNLL resummation
 - Theoretical and experimental uncertainties of similar sizes (percent level)



G. Abelof



Differential Top Pair Production

• Goal: fully differential event generator for $t\bar{t}$ production at NNLO

• This talk:

• Status of our NNLO calculation for the $q\bar{q}$ channel

$$d\hat{\sigma}_{q\bar{q},NNLO} = C_A C_F \left[N_c^2 A + N_c B + C + \frac{D}{N_c} + \frac{E}{N_c^2} + N_l \left(N_c F_l + \frac{G_l}{N_c} \right) + N_h \left(N_c F_h + \frac{G_h}{N_c} \right) + N_l^2 H_l + N_l N_h H_{lh} + N_h^2 H_h \right]$$

Preliminary differential distributions. N_l contributions only

G. Abelof

Ingredients For Top Pair Production At NNLO

• LO and NLO fully differential cross sections

- Known [Ellis, Dawson, Nason '89; Beenakker, Kuijf, van Neerven, Smith '89]
- Re-derived using NLO antenna subtraction [GA, Gehrmann-De Ridder '11]
- Two-loop $2 \rightarrow 2$ matrix elements for $q\bar{q} \rightarrow t\bar{t}$
 - Use analytic results [Bonciani, Ferroglia, Gehrmann, Maître, v. Manteuffel, Studerus]
- One-loop 2→3 matrix elements for qq̄ → tt̄g
 Obtained numerically with OpenLoops [Cascioli, Meierhöfer, Pozzorini]
 ✓ Color structure handled algebraically
 ✓ Quadruple precision evaluation in soft limit (P. Meierhöfer's Talk)
- Tree-level 2 → 4 matrix elements for $q\bar{q} \rightarrow t\bar{t}gg \quad q\bar{q} \rightarrow t\bar{t}q'\bar{q}' \quad q\bar{q} \rightarrow t\bar{t}q\bar{q}$







G. Abelof

Ingredients For Top Pair Production At NNLO

$$\mathrm{d}\hat{\sigma}_{NNLO} = \int_{\mathrm{d}\Phi_4} \mathrm{d}\hat{\sigma}_{NNLO}^{RR} + \int_{\mathrm{d}\Phi_3} \left(\mathrm{d}\hat{\sigma}_{NNLO}^{RV} + \mathrm{d}\hat{\sigma}_{NNLO}^{MF,1} \right) + \int_{\mathrm{d}\Phi_2} \left(\mathrm{d}\hat{\sigma}_{NNLO}^{VV} + \mathrm{d}\hat{\sigma}_{NNLO}^{MF,2} \right)$$

•
$$\int_{d\Phi_4} d\hat{\sigma}_{NNLO}^{RR}$$
, $\int_{d\Phi_3} d\hat{\sigma}_{NNLO}^{RV}$ \longrightarrow implicit IR poles from PS integration over single and double unresolved regions

Need a procedure to isolate and cancel all IR singularities, and assemble all parts in a parton-level event generator

G. Abelof

Antenna Subtraction At NNLO

$$\begin{split} \mathrm{d}\hat{\sigma}_{NNLO} &= \int_{\mathrm{d}\Phi_4} \left(\mathrm{d}\hat{\sigma}_{NNLO}^{RR} - \mathrm{d}\hat{\sigma}_{NNLO}^S \right) \\ &+ \int_{\mathrm{d}\Phi_3} \left(\mathrm{d}\hat{\sigma}_{NNLO}^{RV} - \mathrm{d}\hat{\sigma}_{NNLO}^{VS} + \mathrm{d}\hat{\sigma}_{NNLO}^{MF,1} + \int_1 \mathrm{d}\hat{\sigma}_{NNLO}^{S,1} \right) \\ &+ \int_{\mathrm{d}\Phi_2} \left(\mathrm{d}\hat{\sigma}_{NNLO}^{VV} + \mathrm{d}\hat{\sigma}_{NNLO}^{MF,2} + \int_1 \mathrm{d}\hat{\sigma}_{NNLO}^{VS} + \int_2 \mathrm{d}\hat{\sigma}_{NNLO}^{S,2} \right) \end{split}$$

- Introduce double real and real-virtual subtraction terms $d\hat{\sigma}_{NNLO}^S$, $d\hat{\sigma}_{NNLO}^{VS}$ and add them back in integrated form
- The integrated double real subtraction term is split as

$$\int_{\mathrm{d}\Phi_4} \mathrm{d}\hat{\sigma}_{NNLO}^S = \int_{\mathrm{d}\Phi_3} \int_1 \mathrm{d}\hat{\sigma}_{NNLO}^{S,1} + \int_{\mathrm{d}\Phi_2} \int_2 \mathrm{d}\hat{\sigma}_{NNLO}^{S,2}$$

• Each PS integrand is free of explicit poles, well behaved in singular regions, and can be integrated numerically in D=4

G. Abelof

Antenna Subtraction At NNLO

Antenna subtraction terms constructed with

• Antenna functions

G. Abelof

- Smoothly interpolate many unresolved limits
- Constructed from physical matrix elements

$$\begin{split} X_3^0(i,j,k) &= S_{ijk,IK} \frac{|\mathcal{M}_3^0(i,j,k)|^2}{|\mathcal{M}_2^0(I,K)|^2} \\ X_3^1(i,j,k) &= S_{ijk,IK} \frac{|\mathcal{M}_3^1(i,j,k)|^2}{|\mathcal{M}_2^0(I,K)|^2} - X_3^0(i,j,k) \frac{|\mathcal{M}_2^1(I,K)|^2}{|\mathcal{M}_2^0(I,K)|^2} \\ X_4^0(i,j,k,l) &= S_{ijkl,IL} \frac{|\mathcal{M}_4^0(i,j,k,l)|^2}{|\mathcal{M}_2^0(I,L)|^2} \end{split}$$

• 3 \rightarrow 2 and 4 \rightarrow 2 on-shell momentum mappings, (different for FF, IF, II configurations)

 $\{p_i, p_j, p_k\} \to \{p_I, p_K\} \qquad \{p_i, p_j, p_k, p_l\} \to \{p_I, p_L\}$

Conserve momentum in reduce matrix elements

Collapse to appropriate kinematics in each unresolved limit

Antenna Subtraction At NNLO

Integrated form of subtraction terms obtained with

• Phase space factorisations (different for FF, IF, II configurations). E.g. (FF)

$$\mathrm{d}\Phi_m(\ldots,p_i,p_j,p_k,\ldots) = \mathrm{d}\Phi_{m-1}(\ldots,p_I,p_K,\ldots) \times \mathrm{d}\Phi_{X_{ijk}}(p_i,p_j,p_k)$$

$$\mathrm{d}\Phi_m(\ldots,p_i,p_j,p_k,p_l,\ldots) = \mathrm{d}\Phi_{m-2}(\ldots,p_I,p_L,\ldots) \times \mathrm{d}\Phi_{X_{ijkl}}(p_i,p_j,p_k,p_l)$$

• Integrated antennae. Inclusive phase space integrals

$$\mathcal{X}_{ijk}^{0,1}(\epsilon, s_{IK}) = \left(8\pi^2 \,(4\pi)^\epsilon \,e^{\epsilon\gamma_E}\right) \int \mathrm{d}\Phi_{X_{ijk}} X_3^{0,1}(i,j,k)$$
$$\mathcal{X}_{ijkl}^0(\epsilon, s_{IL}) = \left(8\pi^2 \,(4\pi)^\epsilon \,e^{\epsilon\gamma_E}\right)^2 \int \mathrm{d}\Phi_{X_{ijkl}} X_4^0(i,j,k,l)$$

- Phase space integrals reduced to masters
- Master integrals evaluated directly when possible. With differential equations in external kinematic invariants otherwise

G. Abelof

NNLO Antenna Subtraction With Massive Quarks

All tools available for processes with massless partons

- Phase space mappings [Kosower '03; Daleo, Gehrmann, Maître '07]
- Antenna functions X_3^0 , X_4^0 , X_3^1 [Gehrmann-De Ridder, Gehrmann, Glover '04, '05]
- Integrated antennae:
 - ▶ FF X_3^0 , X_4^0 , X_3^1 [Gehrmann-De Ridder, Gehrmann, Glover '05]
 - ▶ IF, II X_3^0 [Daleo, Gehrmann, Maître '07]
 - ▶ IF X_4^0 , X_3^1 [Daleo, Gehrmann-De Ridder, Gehrmann, Luisoni '10]
 - II X_3^1 [Gehrmann, Monni '11]
 - ► II X⁰₄ [Boughezal, Gehrmann-De Ridder, Ritzmann '11; Gehrmann-De Ridder, Gehrmann, Ritzmann '12]

Challenge: extend NNLO antenna subtraction method to treat massive quarks.

- Re-derive phase space mappings and factorizations [G.A., Gehrmann-De Ridder '11]
- Compute and integrate NLO and NNLO massive antennae. So far incomplete

G. Abelof

Integrated Massive Antennae

- All massive three-parton tree-level antennae (X_3^0) known
 - FF [Gehrmann-De Ridder, Ritzmann '09]
 - ► IF [GA, Gehrmann-De Ridder '11]

• Some four-parton antennae (X_4^0) known

▶ FF: A⁰_{QggQ̄} B⁰_{Qqq̄Q̄} [Bernreuther, Bogner, Dekkers '11, '13]
▶ IF: B⁰_{q,Qq'q̄'} E⁰_{g,Qq̄} E⁰_{g,Qq̄} [GA, Dekkers, Gehrmann-De Ridder '12]

• In progress (almost there) $\mathcal{A}^{0}_{q,Qgg}$ $\mathcal{A}^{1}_{q,Qgg}$



G. Abelof

Integrated Massive Antennae

• Many new master integrals involved



G. Abelof

Double Real Contributions

• Subtraction terms for partonic processes

- ▶ $q\bar{q} \rightarrow t\bar{t}q'\bar{q}' \quad q\bar{q} \rightarrow t\bar{t}q\bar{q}$ [GA, Gehrmann-De Ridder '11]
- $q\bar{q} \rightarrow t\bar{t}gg$ (leading-color only) [GA, Gehrmann-De Ridder, Meierhöfer, Pozzorini '14]

• Check of convergence

- Generate events near every singular region
- Control proximity to singularities with a control variable x (specific to each limit)
- For each event, compute

$$\delta_{RR} = \left| \frac{\mathrm{d}\hat{\sigma}_{NNLO}^{RR}}{\mathrm{d}\hat{\sigma}_{NNLO}^{S}} - 1 \right|$$

• Convergence of $d\hat{\sigma}_{NNLO}^S$ to $d\hat{\sigma}_{NNLO}^{RR}$ observed in cumulative histograms in δ_{RR}

• Good convergence observed in all single and double unresolved limits

G. Abelof

Double Real Contributions



G. Abelof

Real Virtual Contributions

- Partonic process $q\bar{q} \rightarrow t\bar{t}g$ at one-loop
- One-loop amplitudes computed
 - Numerically with OpenLoops for leading-color contributions
 - Analytically for Nl, Nh pieces
- Subtraction terms constructed and implemented
 - Leading-color: [GA, Gehrmann-De Ridder, Meierhöfer, Pozzorini '14]
 - Fermionic contributions: [GA, Gehrmann-De Ridder (in preparation)]
- Pointwise cancellation of explicit IR poles checked analytically

$$\mathcal{P}oles\left(\mathrm{d}\hat{\sigma}_{NNLO}^{RV} - \mathrm{d}\hat{\sigma}_{NNLO}^{VS} + \mathrm{d}\hat{\sigma}_{NNLO}^{MF,1} + \int_{1} \mathrm{d}\hat{\sigma}_{NNLO}^{S,1}\right) = 0$$

G. Abelof

Real Virtual Contributions

- Check of convergence. Analogous to double real check
 - Good convergence observed in fermionic pieces in soft and collinear limits
 - Good convergence observed in collinear limit of leading color piece
 - ▶ Convergence in soft limit of leading-color piece only achieved evaluating d
 ^{RV}_{NNLO} in quadruple precision

$$\delta_{RV} = \left| \frac{\mathcal{F}inite(\mathrm{d}\hat{\sigma}_{NNLO}^{RV})}{\mathcal{F}inite(\mathrm{d}\hat{\sigma}_{NNLO}^{T})} - 1 \right|$$



Real Virtual Contributions $y_{cut} = p_T^g / \sqrt{\hat{s}}$

Only "bad points" are (re)evaluated by OpenLoops in quadruple precision σ — σ Γ — σ Γ = σ Γ

Stability check: Evaluate $R = \left(\sigma_{NNLO}^{RV} - \sigma_{NNLO}^{T}\right)/\sigma_{LO}$ as a function of $y_{cut} = p_T^g/\sqrt{\hat{s}}$

- Integration is stable
- R has a plateau for $y_{cut} < y_{cut}^{max} \sim 10^{-3}$
- Strong check of our subtraction terms
- We can run with $y_{cut} \sim 10^{-4}$. Only ~0.01% points require quadruple precision.
- Efficient evaluation in double precision for the vast majority of points



 $y_{cut} < y_{cut}^{max} = 10^{-3}$

Double Virtual Contributions

The ultimate check:

$$\mathcal{P}oles\left(\mathrm{d}\hat{\sigma}_{NNLO}^{VV} + \mathrm{d}\hat{\sigma}_{NNLO}^{MF,2} + \int_{1}\mathrm{d}\hat{\sigma}_{NNLO}^{VS} + \int_{2}\mathrm{d}\hat{\sigma}_{NNLO}^{S,2}\right) = 0$$

• Pole cancellation verified analytically in N1 piece

```
Pole1LC =
Simplify[
Simplify[Coefficient[TwoLoopPolesLC, ep, -1] Delta[1 - x1] Delta[1 - x2] + Coefficient[OneTimesOneLoopPolesLC, ep, -1] Delta[1 - x1] Delta[1 - x2] -
Coefficient[SubTermLC, ep, -1]] //. RepsLogsLC]
0
```

- Non-trivial check on new integrated massive antennae
- Proves applicability of NNLO antenna subtraction to reactions with massive fermions
- Pole cancellation in remaining color factors in progress

G. Abelof

Preliminary Results

- Preliminary results for $pp(q\bar{q}) \rightarrow t\bar{t} + X$ (N_l only)
 - $\sqrt{s} = 7 \,\mathrm{TeV}$
 - $m_{top} = 173.5 \, \text{GeV}$
 - $\mu = m_{top}$
 - PDF sets: MSTW2008LO, MSTW2008NLO, MSTW2008NNLO









Preliminary Results



G. Abelof

Summary And Outlook

- Fully differential NNLO calculation for $t\bar{t}$ production in the $q\bar{q}$ channel within reach (leading-color + fermionic contributions)
- Double real contributions: subtraction terms implemented and tested
- Real-virtual contributions:
 - Subtraction terms implemented and tested
 - Precise and stable one-loop amplitudes from OpenLoops in leading-color part
- Double virtual contributions:
 - Two-loop amplitudes available (for leading color and fermionic pieces)
 - Analytic pole cancelation in N1 part
- \bullet Event generator implemented and working (with N_l part for the moment)

<u>Outlook</u>

- Complete leading-color double virtual contributions in $q\bar{q}$ channel
- Phenomenology in $q\bar{q}$ channel. A_{FB}, main goal
- Include gg and qg channels

G. Abelof