Integrability of Limit Shape Phenomena in Six Vertex Model

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- In the thermodynamic limit, the limiting average height function becomes deterministic and can be found by solving a certain boundary value problem.
- The six vertex model is quantum integrable in the sense that it admits commuting transfer matrices and can be solved by Bethe ansatz.
- What does the quantum integrability imply for the PDE governing the limiting height function?

Outline of Talk

- Quick Review of Six Vertex Model
- Thermodynamic Limit
- Integrability:
 - Transfer Matrices
 - Commuting Hamiltonians
- Examples
- Outlook

Configurations and Weights

• Let $S_T = [0, T] \times [0, 1]$, and let $S_T^{\epsilon} = \epsilon \mathbb{Z}^2$ be the scaled square lattice centered inside S_T .

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- Each vertex has a weight v(s).
- The Boltzmann weight of *s*:

$$w(s) = \prod_{\text{vertex } v} v(s)$$

Boundary Conditions

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The partition function and the normalized free energy are:

$$egin{aligned} Z^{\epsilon}_{\eta_1,\eta_2,\,\mathcal{T}} &= \sum_{\substack{s(0)=\eta_1\ s(1)=\eta_2}} w(s)\ f^{\epsilon}_{\eta_1,\eta_2,\,\mathcal{T}} &= \epsilon^2 \logig(Z_{\eta_1,\eta_2}ig) \end{aligned}$$

Height Function

- A height function is a function on faces satisfying a gradient constraint:
 - $0 \leq h(x,y) h(x + \epsilon, y) \leq 1$
 - $0 \leq h(x, y + \epsilon) h(x, y) \leq 1$

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3	3	3		
2	2	2		
2	2	1		
1	1	0		

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- The normalized height function $\bar{h} = \epsilon h$. The average height function $\langle \bar{h} \rangle$ is the ensemble average of the normalized height function.

Thermodynamic limit

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- In this case, there exist limiting free energy and limiting height function:

$$egin{aligned} &f_{\eta_1,\eta_2,\mathcal{T}} = \lim_{\epsilon o 0} f^\epsilon_{\eta_1,\eta_2,\mathcal{T}} \ &\langle h
angle = \lim_{\epsilon o 0} \langle h
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Variational Principle

• The limiting free energy and average height function can be computed by variational principle.

$$f_{\eta_1,\eta_2,T} = \max_{h \in \mathcal{H}} \int_0^1 \int_0^T \sigma_w(\partial_t h, \partial_y h) dt dy$$

where σ is called the surface tension function, and \mathcal{H} is the set of limiting height functions, $h: S_t \to \mathbb{R}$ satisfying: h(0,0) = 0, monotonicity, and Lipschitz continuity with constant 1.

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- Euler Lagrange equations:

$$\partial_{11}\sigma_{w} \partial_{t}^{2}h + 2 \partial_{12}\sigma_{w} \partial_{t}\partial_{y}h + \partial_{22}\sigma_{w} \partial_{y}^{2}h = 0$$

Integrability of the Six Vertex Mode Transfer Matrices

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- Define the transfer matrix $T_w: V \to V$ by its matrix elements:

$$\langle s_1 | T_w | s_2 \rangle = Z_{s_1, s_2, \epsilon}$$

(ie. the partition function for just one column).

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(ie. the partition function for just one column).Then:

$$Z_{\eta_1,\eta_2,t} = \langle \eta_1 | T_w^{\lfloor t/\epsilon \rfloor} | \eta_2 \rangle$$

Hamiltonian Formulation of Variational Principle

 Recast the variational problem in the Hamiltonian formulation by Legendre transform:

$$\mathcal{H}_w(\pi,t) = \max_s \pi s - \sigma_w(s,t)$$

The new variables are h and π , where π is conjugate to $\partial_t h$.

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$$f_{\eta_1,\eta_2,T} = \max_{\pi,h} S[\pi,h]$$
$$S[\pi,h] = \int_0^T \int_0^1 \pi \,\partial_t h - H_w(\pi,\partial_y h) \,dt \,dy$$

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These are equivalent to the Euler-Lagrange equations.

Commuting Transfer Matrices and Hamiltonians

• Recall
$$\Delta w = \frac{w_1^2 + w_2^2 - w_3^2}{2w_1w_2}$$

 Quantum Integrability: if w and w̃ satisfy Δw = Δw̃ then the transfer matrices commute:

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 Main result is semiclassical integrability: if Δw = Δw then the corresponding Hamiltonians Poisson commute:

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 Lemma 2: The Hessian of the surface tension σ_w of the six vertex model σ depends on w only via Δ(w).

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• Lemma: The hessian of σ is π^2 , independent of weights.

Hexagonal Dimer Model

• The six vertex model with weights

 $w_1 = 0$ $w_2 = a$ $w_3 = b$ $w_4 = c$ $w_5 = \sqrt{bc}$ $w_6 = \sqrt{bc}$

Corresponds to the dimer model on the hexagonal lattice with edge weights (a, b, c).



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• The Euler-Langrange equations for the limiting height function can be trasnformed to the Burger's equation, $\partial_t u + u \partial_y u = 0$, which admits many integrals of motion: $\int u(y)^n dy$.

• The surface tension function σ can be calculated in closed form, and the Hamiltonians can be shown directly to commute.

Free Fermion Point

 More generally, when Δw = 0, the six vertex model is equivalent to the dimer model on the graph:



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• This corresponds to gluing two regions together:

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• In the limit $\epsilon \rightarrow 0$, by large deviation principle:

$$f_{\eta_1,\eta_2,t,\widetilde{t}} = \max_{\eta} f_{\eta_1,\eta,t} + \widetilde{f}_{\eta,\eta_2,\widetilde{t}}$$

• The commutation of the transfer matrices implies:

$$\max_{\eta} f_{\eta_1,\eta,t} + \widetilde{f}_{\eta,\eta_2,\widetilde{t}} = \max_{\eta} \widetilde{f}_{\eta_1,\eta,\widetilde{t}} + f_{\eta,\eta_2,t}$$

for all t, \tilde{t} and boundary conditions η_1, η_2 .

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for all t, t̃ and boundary conditions η₁, η₂.
Recall that f is the Hamilton-Jacobi action.

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- Recall that f is the Hamilton-Jacobi action.
- Generally:

If the Hamilton-Jacobi actions of H and \widetilde{H} commute in the above sense, then does $\{H, \widetilde{H}\}$?

• Generally no, but under mild assumptions then yes,

Further Work Integrability

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Further Work Integrability

- The existence of commuting transfer matrices underlies the solvability of the six vertex model by Bethe Ansatz.
- In the infinite dimensional setting, the Liouville integrability (the existence of many commuting Hamiltonians) is not enough to have the complete solvability.
- The existence of commuting hamiltonians is first step towards showing the integrability of the limit shape PDE.

End!