
$e^+e^- \rightarrow 3 \text{ jets at NNLO}$

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RADCOR 2007

$e^+e^- \rightarrow 3$ jets and event shapes

Classical QCD observable

- testing ground for QCD: perturbation theory, power corrections and logarithmic resummation
- precision measurement of strong coupling constant α_s
- current error on α_s from jet observables dominated by theoretical uncertainty:
S. Bethke, 2006

$$\alpha_s(M_Z) = 0.121 \pm 0.001(\text{experiment}) \pm 0.005(\text{theory})$$

- theoretical uncertainty largely from missing higher orders
- current status: NLO plus NLL resummation

Theoretical description

- easier than at hadron colliders, since coloured partons only in final state:
no initial state emission, no parton distributions
- new calculational methods first developed for e^+e^- , then extended to hadronic processes

$e^+e^- \rightarrow 3$ jets and event shapes

Event shape variables

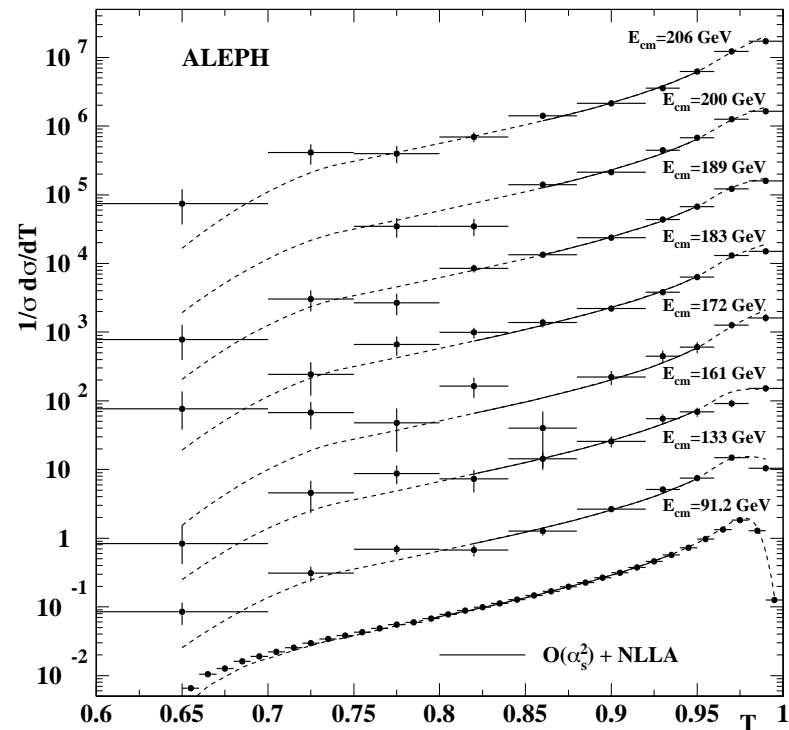
assign a number x to a set of final state momenta: $\{p\}_i \rightarrow x$

e.g. Thrust in e^+e^-

$$T = \max_{\vec{n}} \frac{\sum_{i=1}^n |\vec{p}_i \cdot \vec{n}|}{\sum_{i=1}^n |\vec{p}_i|}$$

limiting values:

- back-to-back (two-jet) limit: $T = 1$
- spherical limit: $T = 1/2$



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Standard Set of LEP

- Thrust (E. Farhi)

$$T = \max_{\vec{n}} \left(\sum_{i=1}^n |\vec{p}_i \cdot \vec{n}| \right) / \left(\sum_{i=1}^n |\vec{p}_i| \right)$$

- Heavy jet mass (L. Clavelli, D. Wyler)

$$M_i^2/s = \frac{1}{E_{\text{vis}}^2} \left(\sum_{k \in H_i} |\vec{p}_k| \right)^2$$

- C -parameter: eigenvalues of the tensor (G. Parisi)

$$\Theta^{\alpha\beta} = \frac{1}{\sum_k |\vec{p}_k|} \frac{\sum_k p_k^\alpha p_k^\beta}{\sum_k |\vec{p}_k|}$$

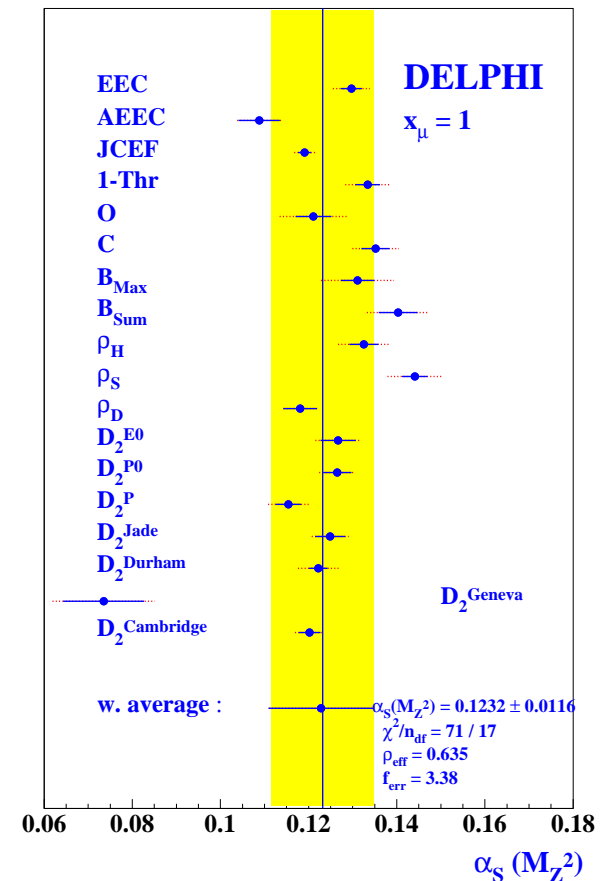
- Jet broadenings (S. Catani, G. Turnock, B. Webber)

$$B_i = \left(\sum_{k \in H_i} |\vec{p}_k \times \vec{n}_T| \right) / \left(2 \sum_k |\vec{p}_k| \right)$$

$$B_W = \max(B_1, B_2) \quad B_T = B_1 + B_2$$

- $3j \rightarrow 2j$ transition parameter in Durham algorithm y_{23}^D

S.Catani, Y.L.Dokshitzer, M.Olsson, G.Turnock, B.Webber



$e^+e^- \rightarrow 3$ jets and event shapes

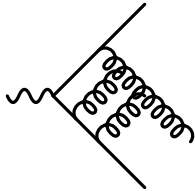
Current status: NLO and NLL

- NLO calculations of event shapes and $3j$
R.K. Ellis, D.A. Ross, A.E. Terrano; Z. Kunszt
J. Vermaseren, K.F. Gaemers, S.J. Oldham; L. Clavelli, D. Wyler
K. Fabricius, I. Schmitt, G. Kramer, G. Schierholz
- NLO parton level event generators for $3j$
EVENT: Z. Kunszt, P. Nason
EERAD: W. Giele, E.W.N. Glover
EVENT2: S. Catani, M. Seymour
- NLO parton level event generators for $4j$
MenloParc: L.D. Dixon, A. Signer
EERAD2: J. Campbell, M. Cullen, E.W.N. Glover
Debrecen: Z. Nagy, Z. Trocsanyi
Mercurito: D. Kosower, S. Weinzierl
- NLL resummation
S. Catani, L. Trentadue, G. Turnock, B. Webber
- Power corrections
G. Korchemsky, G. Sterman; Y. Dokshitzer, B.R. Webber

Ingredients to NNLO $e^+e^- \rightarrow 3\text{-jet}$

Two-loop matrix elements

$|\mathcal{M}|_{2\text{-loop},3\text{ partons}}^2$



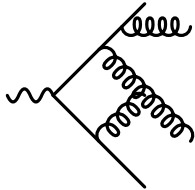
explicit infrared poles from loop integrals

L. Garland, N. Glover, A. Koukoutsakis, E. Remiddi, TG
(RADCOR 00/02);

S. Moch, P. Uwer, S. Weinzierl

One-loop matrix elements

$|\mathcal{M}|_{1\text{-loop},4\text{ partons}}^2$



explicit infrared poles from loop integral and

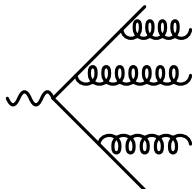
implicit infrared poles due to single unresolved radiation

Z. Bern, L. Dixon, D. Kosower, S. Weinzierl;

J. Campbell, D.J. Miller, E.W.N. Glover

Tree level matrix elements

$|\mathcal{M}|_{\text{tree},5\text{ partons}}^2$



implicit infrared poles due to double unresolved radiation

K. Hagiwara, D. Zeppenfeld;

F.A. Berends, W.T. Giele, H. Kuijf;

N. Falck, D. Graudenz, G. Kramer

Infrared Poles cancel in the sum

NNLO Infrared Subtraction

Structure of NNLO m -jet cross section:

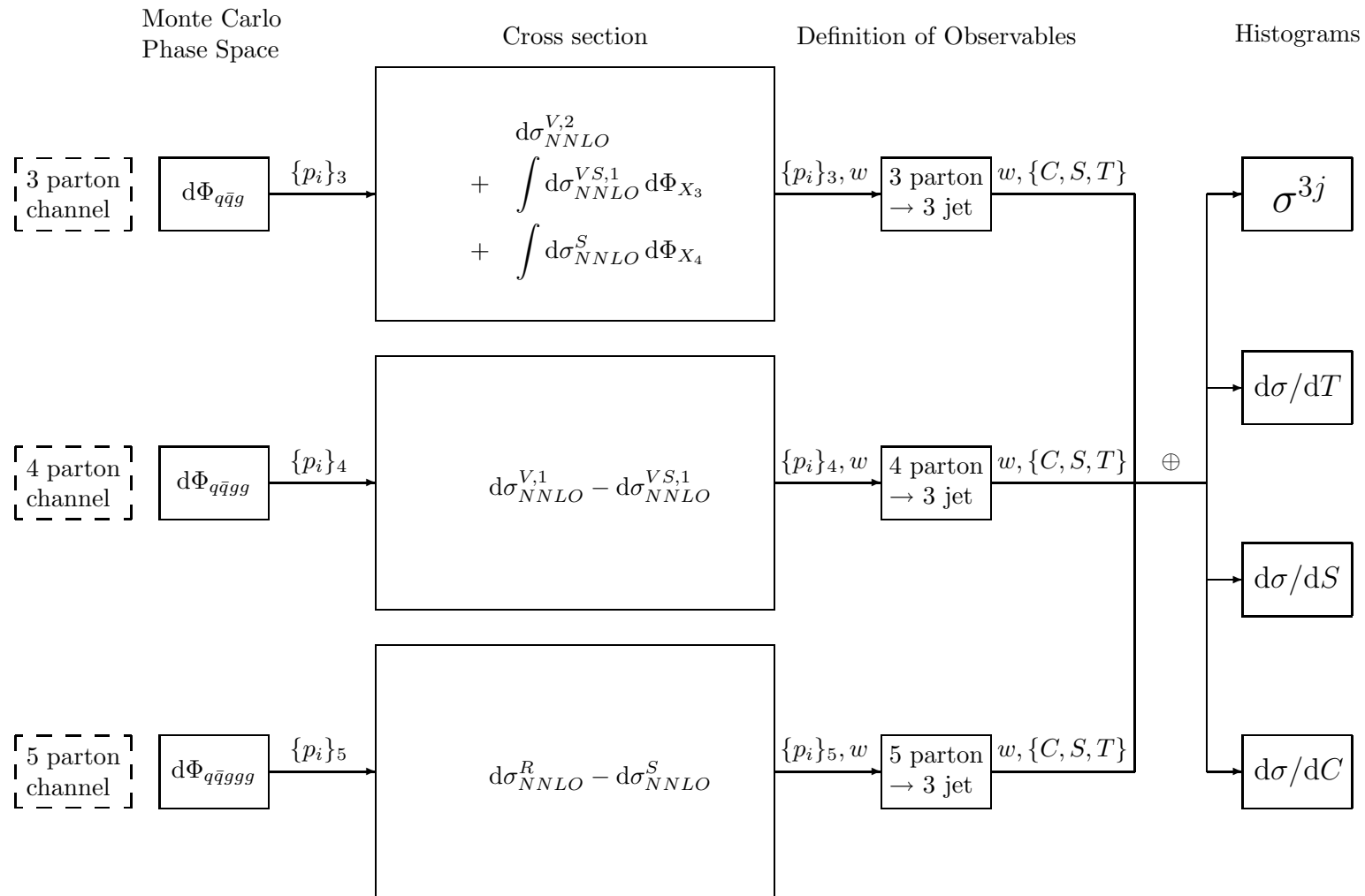
$$\begin{aligned} d\sigma_{NNLO} = & \int_{d\Phi_{m+2}} \left(d\sigma_{NNLO}^R - d\sigma_{NNLO}^S \right) \\ & + \int_{d\Phi_{m+1}} \left(d\sigma_{NNLO}^{V,1} - d\sigma_{NNLO}^{VS,1} \right) \\ & + \int_{d\Phi_m} d\sigma_{NNLO}^{V,2} + \int_{d\Phi_{m+2}} d\sigma_{NNLO}^S + \int_{d\Phi_{m+1}} d\sigma_{NNLO}^{VS,1} , \end{aligned}$$

- $d\sigma_{NNLO}^S$: real radiation subtraction term for $d\sigma_{NNLO}^R$
- $d\sigma_{NNLO}^{VS,1}$: one-loop virtual subtraction term for $d\sigma_{NNLO}^{V,1}$
- $d\sigma_{NNLO}^{V,2}$: two-loop virtual corrections

Each line above is finite numerically and free of infrared ϵ -poles \longrightarrow numerical programme

Numerical Implementation

Structure of $e^+e^- \rightarrow 3$ jets program:



Numerical Implementation

Antenna subtraction

NLO: M. Cullen, J. Campbell, E.W.N. Glover; D. Kosower; A. Daleo, D. Maitre, TG

NNLO: A. Gehrmann-De Ridder, E.W.N. Glover, TG (RADCOR 05)

- construct subtraction terms from physical $1 \rightarrow 3$ and $1 \rightarrow 4$ matrix elements
- each antenna function interpolates between all limits associated to one or two unresolved partons
- integrated subtraction terms cancel infrared pole structure of two-loop matrix element

S. Catani; G. Sterman, M.E. Yeomans-Tejeda

Checks

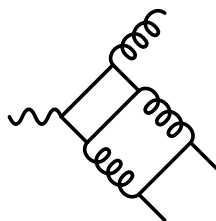
- cancellation of infrared poles in 3-parton and 4-parton channel
- convergence of subtraction terms towards matrix elements along phase space trajectories
- distributions in raw phase space variables
- independence on phase space cut y_0

Colour structure at NNLO

Decomposition into leading and subleading colour terms

$$d\sigma_{NNLO} = (N^2 - 1) \left[N^2 A_{NNLO} + B_{NNLO} + \frac{1}{N^2} C_{NNLO} + NN_F D_{NNLO} \right. \\ \left. + \frac{N_F}{N} E_{NNLO} + N_F^2 F_{NNLO} + N_{F,\gamma} \left(\frac{4}{N} - N \right) G_{NNLO} \right]$$

- last term: closed quark loop coupling to vector boson



$$N_{F,\gamma} = \frac{\left(\sum_q e_q \right)^2}{\sum_q e_q^2}$$

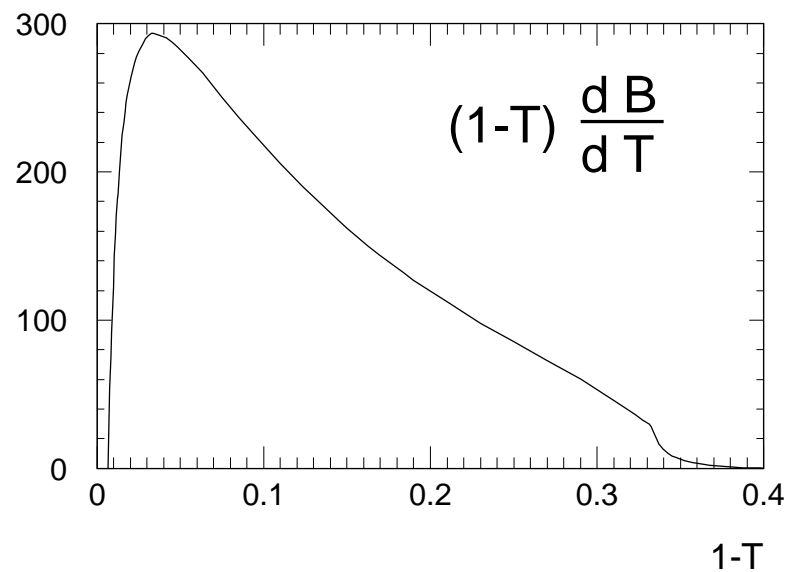
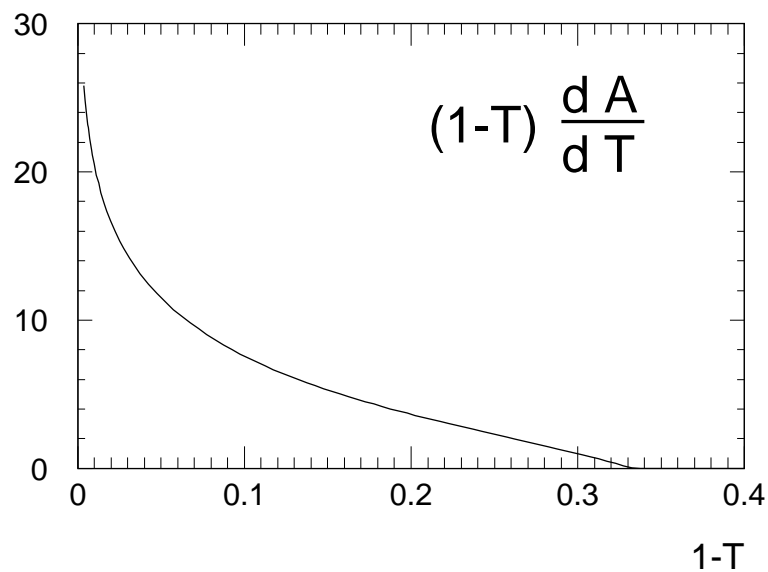
- was found to be $\mathcal{O}(10^{-4})$ in NLO $4j$ final states
L.D. Dixon, A. Signer
- will be neglected here

Event shapes at NNLO

NNLO expression for Thrust

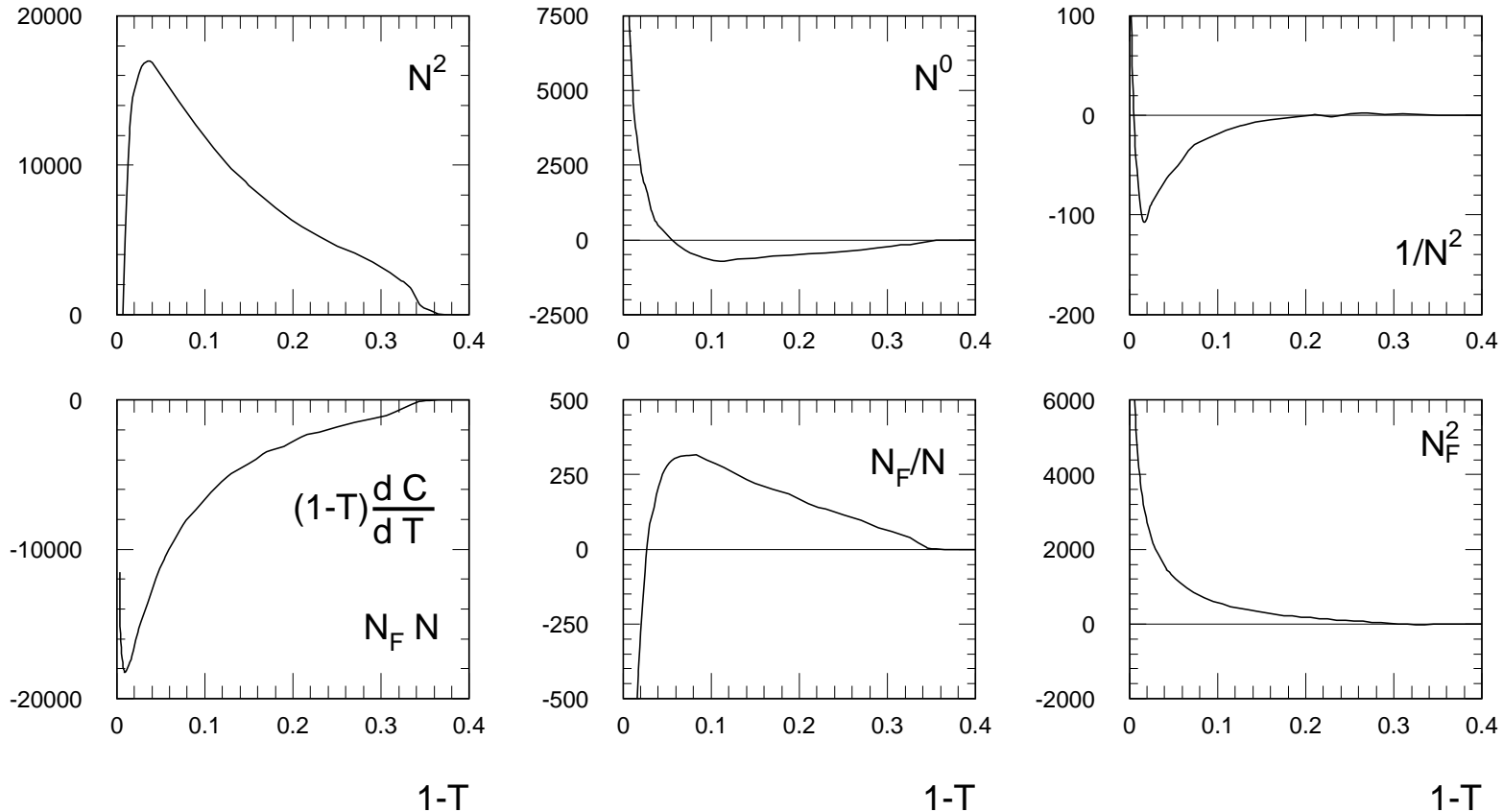
$$(1-T) \frac{1}{\sigma_{\text{had}}} \frac{d\sigma}{dT} = \left(\frac{\alpha_s}{2\pi}\right) A(T) + \left(\frac{\alpha_s}{2\pi}\right)^2 (B(T) - 2A(T)) \\ + \left(\frac{\alpha_s}{2\pi}\right)^3 (C(T) - 2B(T) - 1.64A(T))$$

with LO contribution $A(T)$, NLO contribution $B(T)$



Event shapes at NNLO

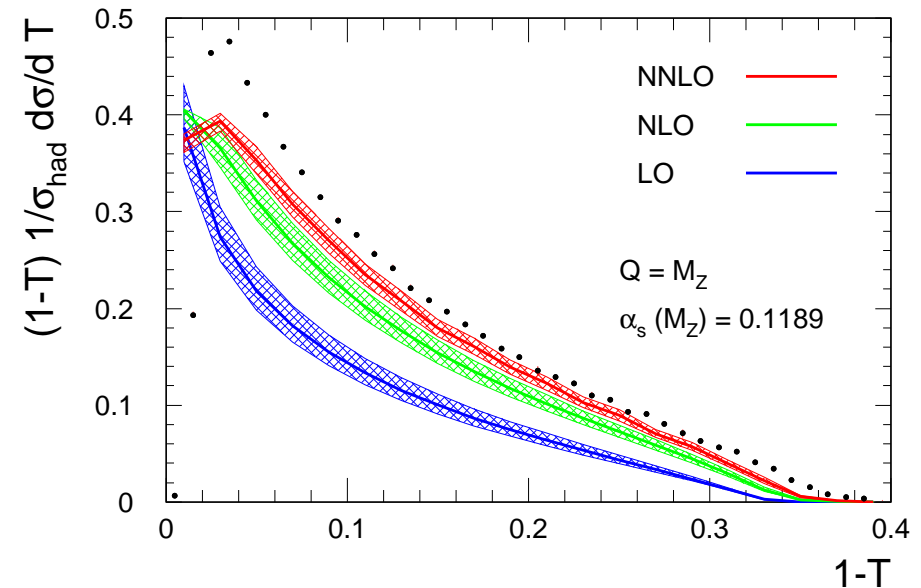
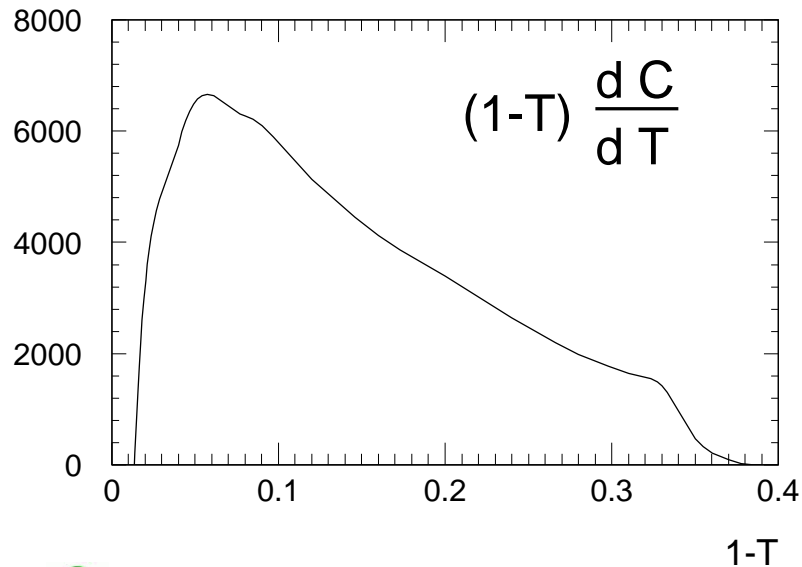
Individual colour structures



- dominated by leading colour N^2 and $N_F N$
- sizable contributions from N^0 , N_F/N and N_F^2
- negligible contribution from $1/N^2$

Results

NNLO thrust distribution



- NNLO corrections sizable
- theory error reduced by 30–40 %
- large $1 - T$: need hadronization corrections
- small $1 - T$: two-jet region, need matching onto NLL resummation

Work in progress: G. Luisoni, TG

● mean value $\langle 1 - T \rangle$: $\mathcal{A} = 2.101$ $\mathcal{B} = 44.98$ $\mathcal{C} = 1095 \pm 130$

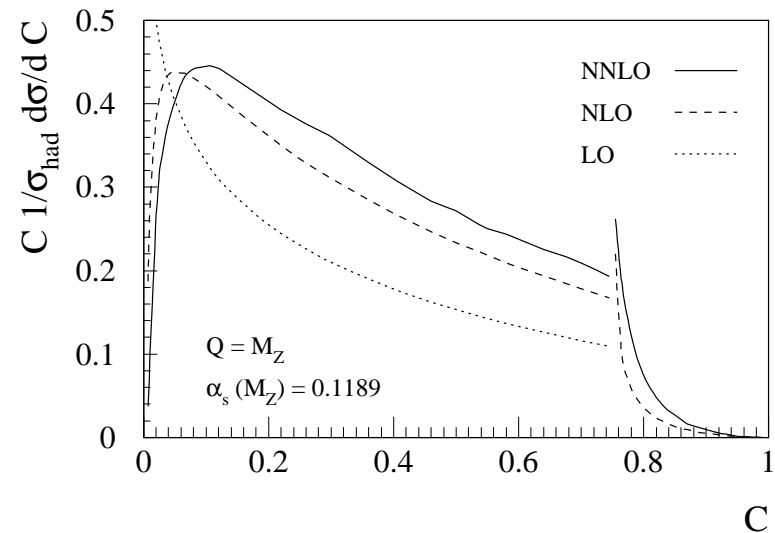
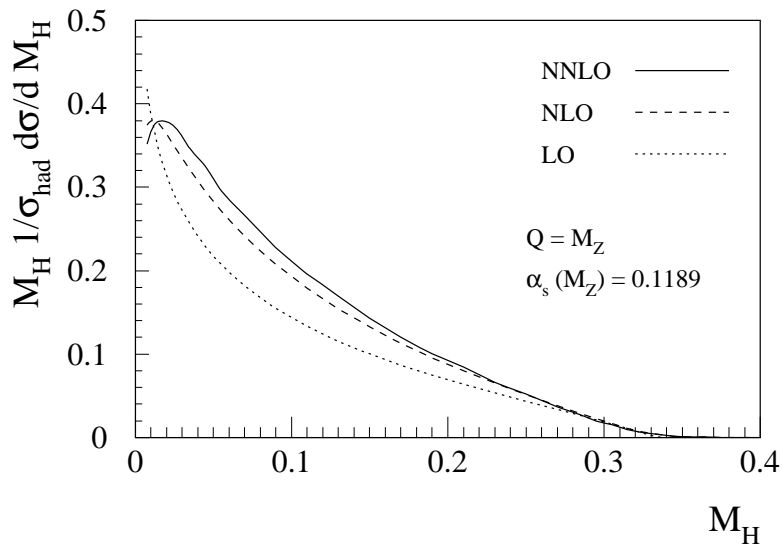
$$\langle 1 - T \rangle(\alpha_s = 0.1189) = 0.0398 + 0.0146 + 0.0068$$

Results

NNLO heavy jet mass and C -parameter

heavy jet mass M_H^2/s

C -parameter



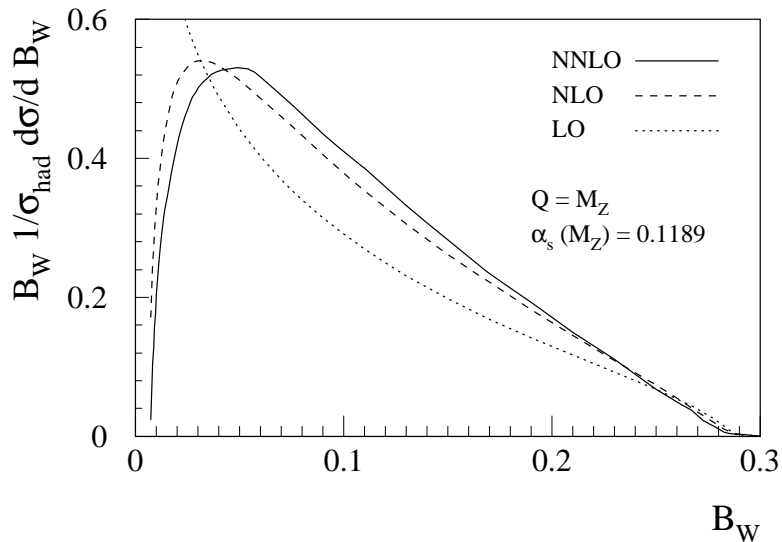
- heavy jet mass (closely related to thrust) has very small NNLO corrections
- NNLO corrections for C large
- again require matching onto NLL resummation and hadronization corrections
- Sudakov shoulder in $C = 0.75$ also requires resummation

S. Catani, B. Webber

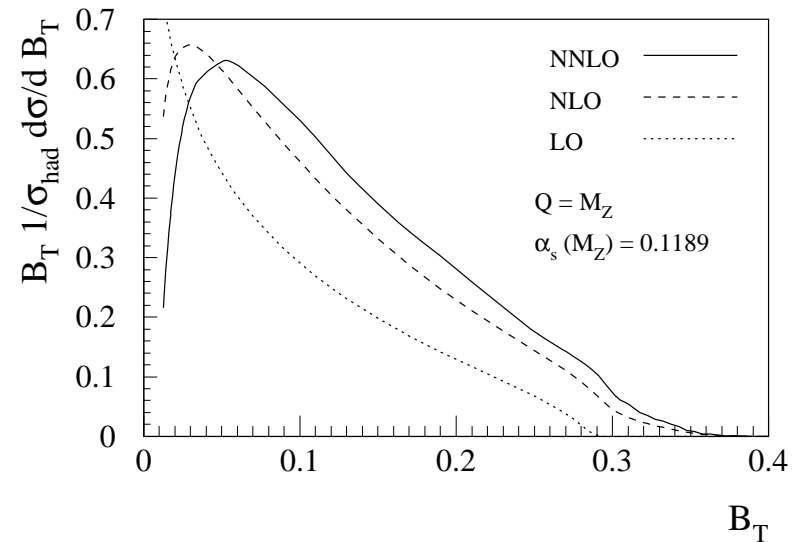
Results

NNLO jet broadenings

wide jet boadening B_W



total jet boadening B_T



- relative magnitude of NNLO corrections smaller than for thrust
- NNLO corrections for B_W smaller than for B_T
- again require matching onto NLL resummation and hadronization corrections

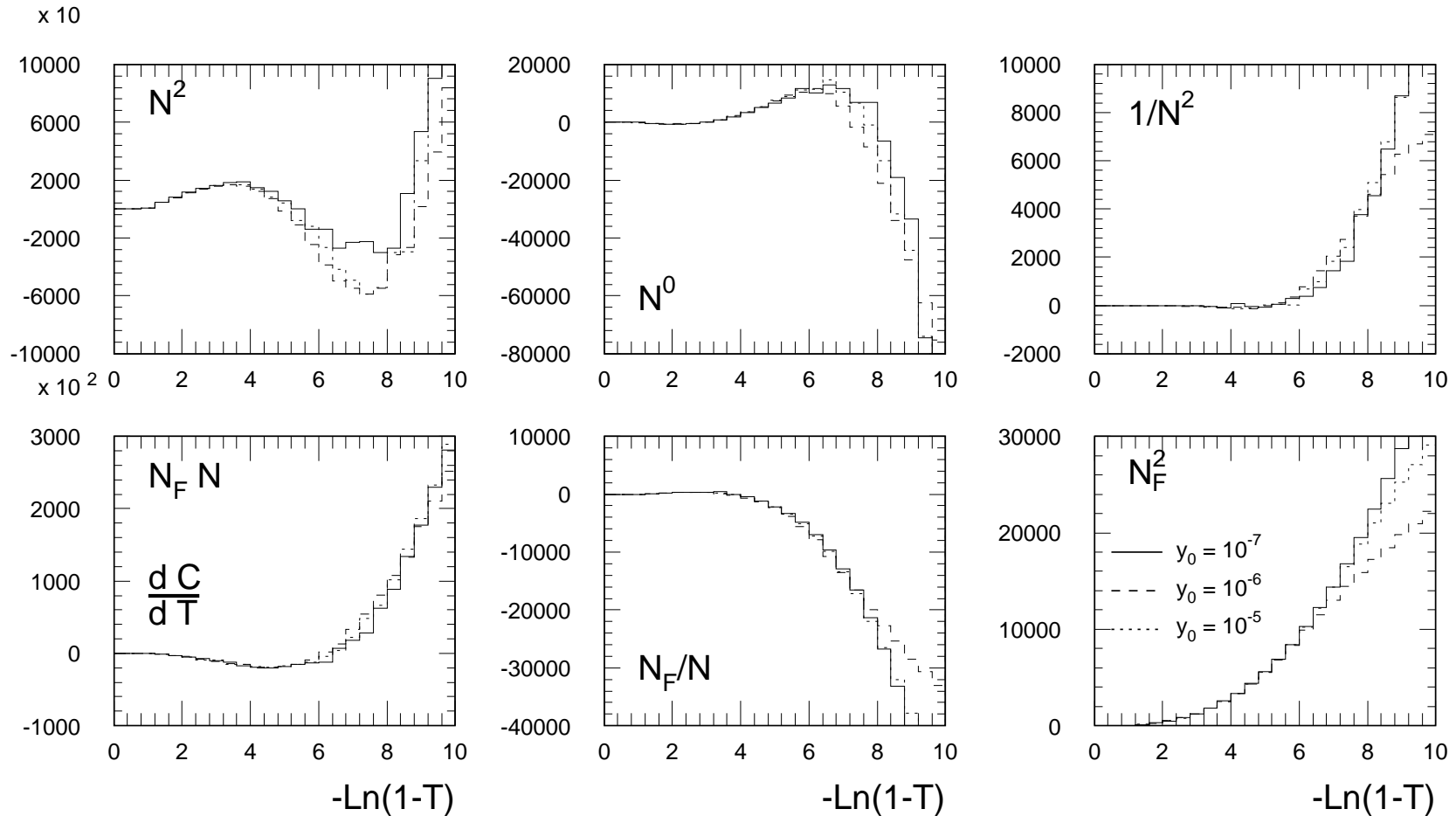
Summary and Outlook

- completed calculation of NNLO corrections to event shapes related to $e^+e^- \rightarrow 3j$
- constructed parton-level event generator, based on antenna subtraction method
- size of NNLO corrections not uniform:
small: $B_W, M_H/s$, moderate: B_T , substantial: C, T
- still running: y_{23}^D, R_{3j}^D
- comparison with data just started
- next steps: matching onto resummation, hadronization corrections

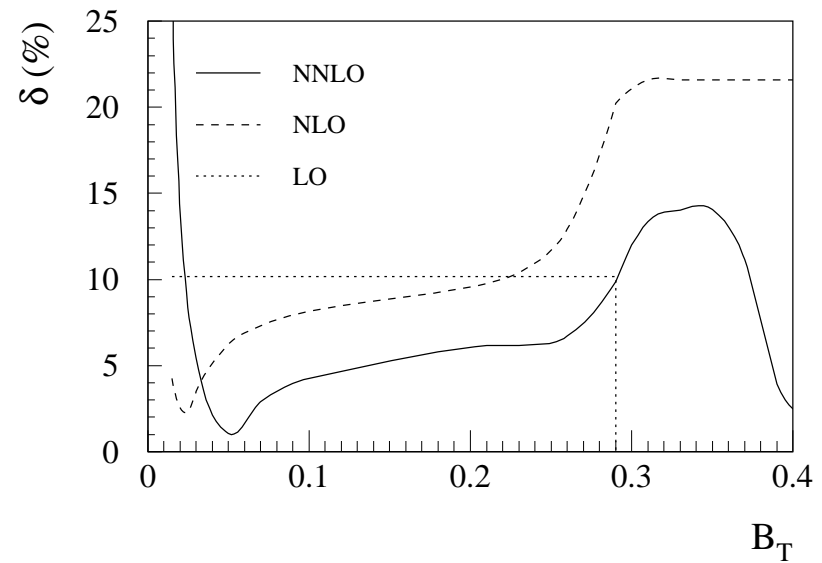
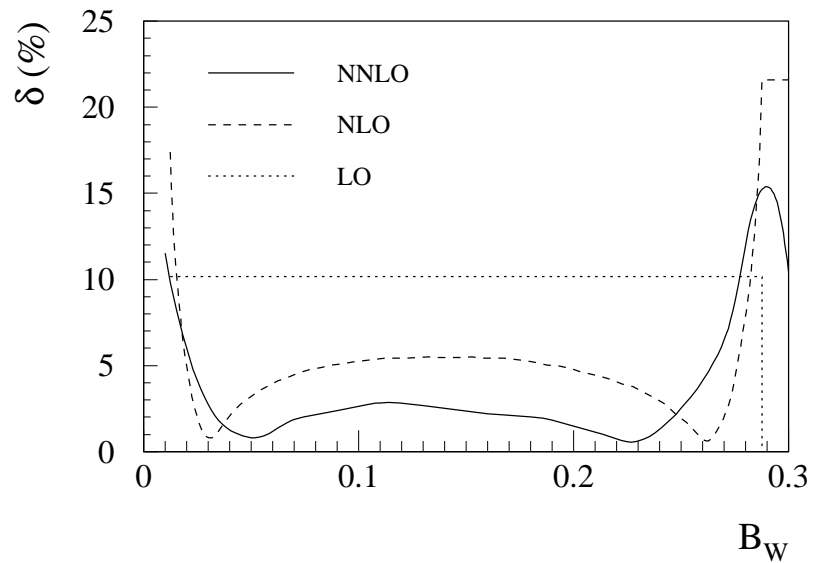
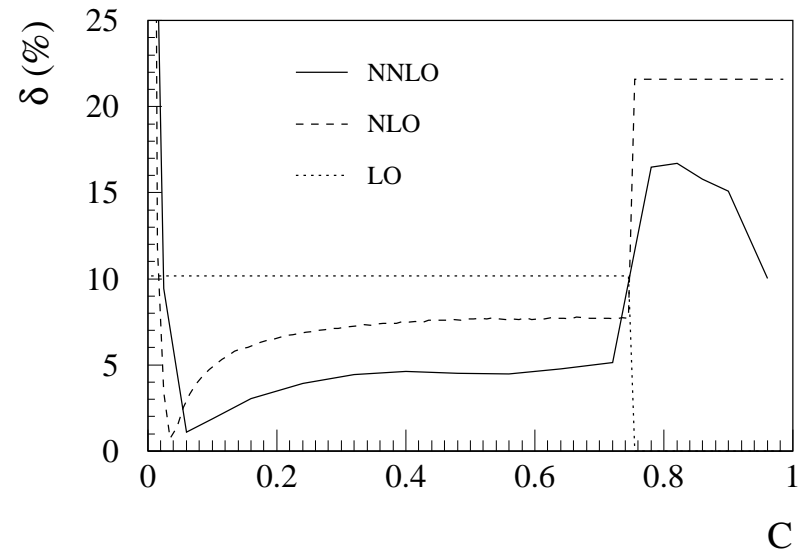
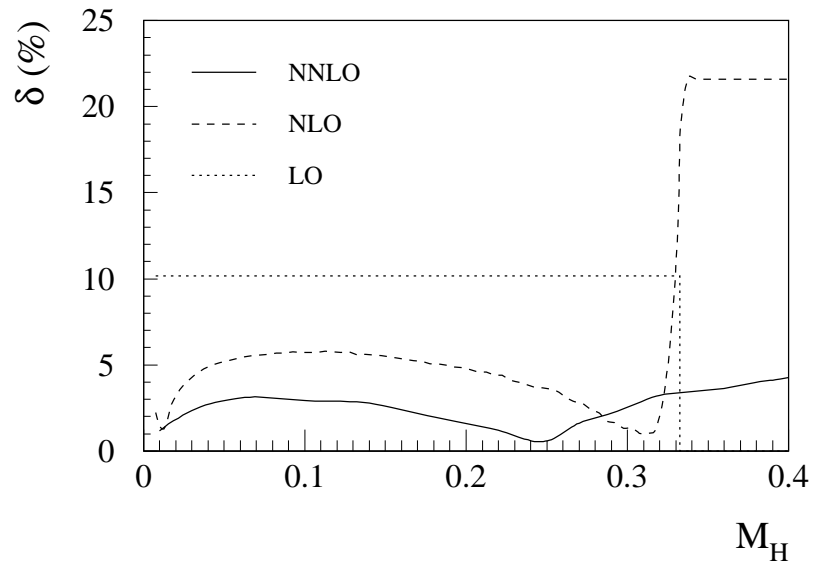
Backup

Thrust at NNLO

Checks

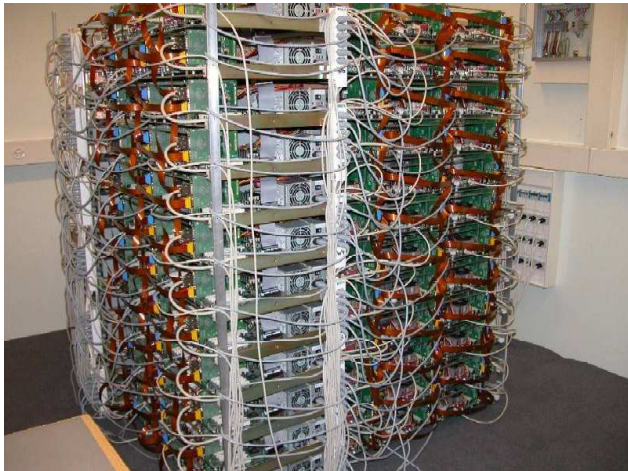


Errors at NNLO



Numerical computation

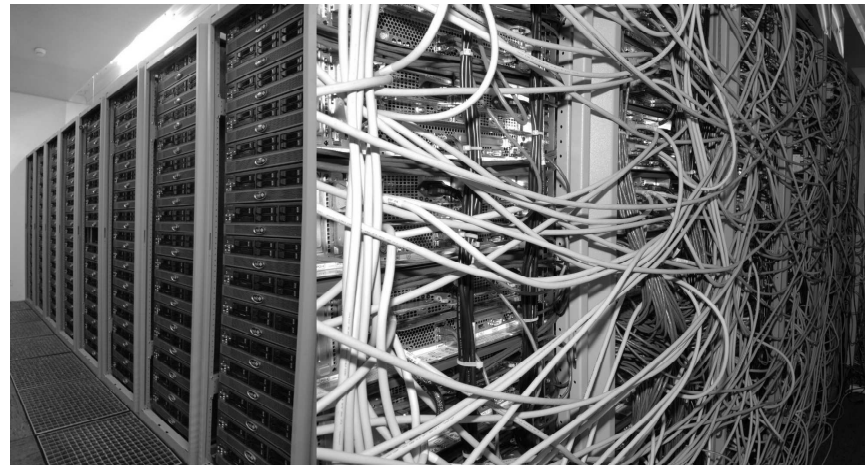
zBox1 and zBox2 supercomputers



zBox1

- 288 processors,
2.2 GHz AMD Athlon
- 0.57 TFlops
- built in-house from
off-the-shelf components

J. Stadel, B. Moore



zBox2

- 500 processors,
2.6 GHz Opteron 852
- 5.2 TFlops
- built by Sun
microsystems

used mostly by our computational astrophysics group