

Quantum information processing with cold Fermi gases in the fast pairing regime

Razvan Teodorescu

Center for Nonlinear Studies,
Los Alamos National Laboratory

September 19, 2008



Schrödinger's ket

Quantum Information ...

Quantum information theory ...

Quantum information theory ...



$$|\text{State}\rangle = c_1|\text{Alive}\rangle + c_2|\text{Not}\rangle, \quad c_{1,2} \in \mathbb{C}, \quad |c_1|^2 + |c_2|^2 = 1.$$

Topological quantum computation: why?

Topological quantum computation: why?

- The “other” qubits (spin, flux, charge) - two-level systems:

$$H = \Delta\sigma_x + [\epsilon + \eta(t)]\sigma_z, \quad \langle \eta(t)\eta(t') \rangle = C(t, t'),$$

Topological quantum computation: why?

- The “other” qubits (spin, flux, charge) - two-level systems:

$$H = \Delta\sigma_x + [\epsilon + \eta(t)]\sigma_z, \quad \langle \eta(t)\eta(t') \rangle = C(t, t'),$$

- Decoherence and damping: for $C(t, t') \sim \delta(t, t')$

Topological quantum computation: why?

- The “other” qubits (spin, flux, charge) - two-level systems:

$$H = \Delta\sigma_x + [\epsilon + \eta(t)]\sigma_z, \quad \langle \eta(t)\eta(t') \rangle = C(t, t'),$$

- Decoherence and damping: for $C(t, t') \sim \delta(t, t')$
- Effects of time-correlated noise ?

$$C(t, t') = V e^{-|t-t'|/\tau}$$

The trouble with “usual” quantum computation

The trouble with “usual” quantum computation

- Locality

$$H = H_0 + H_1$$

The trouble with “usual” quantum computation

- Locality

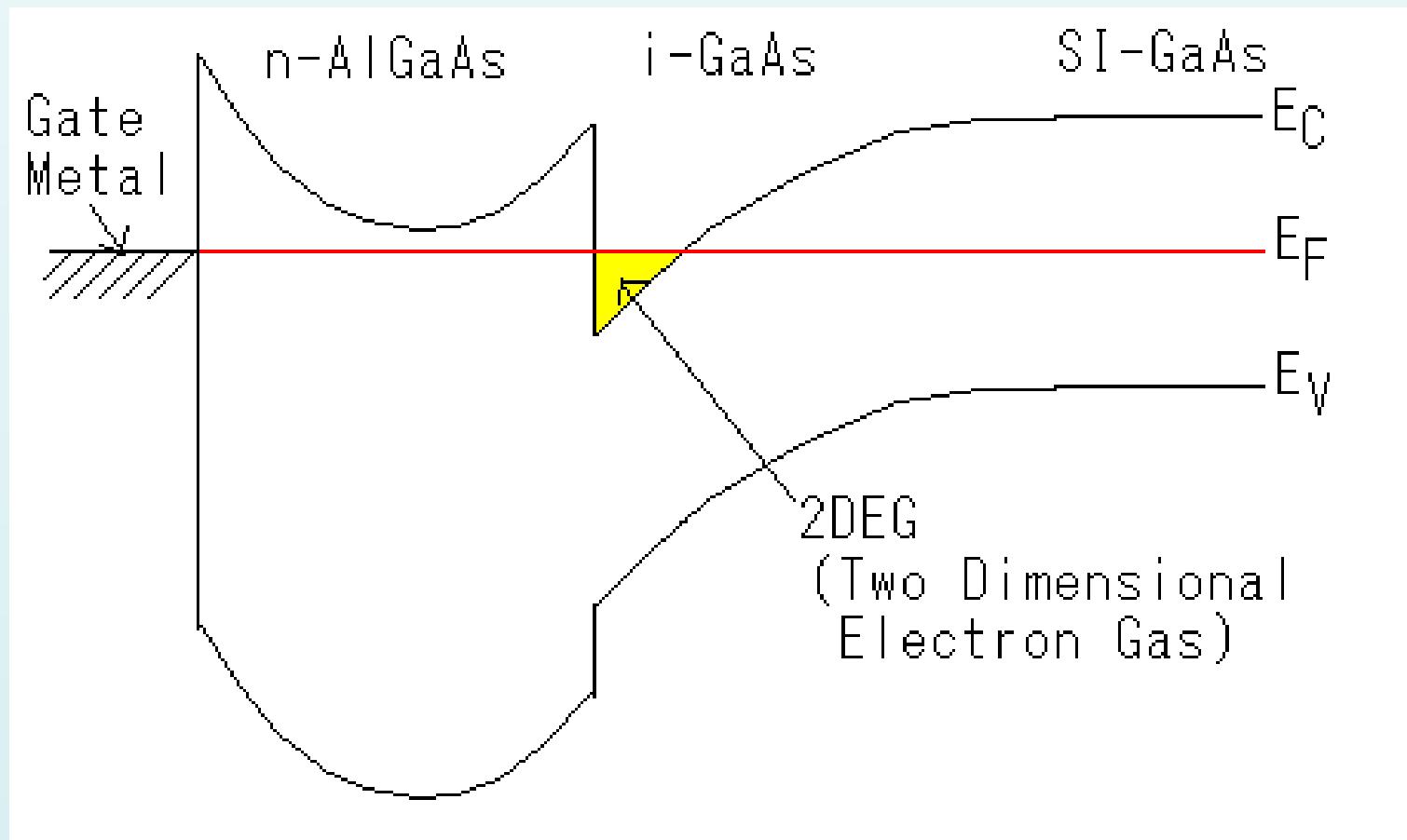
$$H = H_0 + H_1$$

- Decay channels

$$E_{0,1} = \pm \sqrt{\epsilon^2 + |\Delta|^2}, \quad E_2, E_3, \dots$$

Non-abelian anyons, non-local excitations from FQHE

Non-abelian anyons, non-local excitations from FQHE



Fractional Quantum Hall Effect: a refresher

- Berry phase, Aharonov-Bohm: $\Psi(\vec{r}) \sim \exp \left[\frac{ie}{\hbar c} \int_{\vec{r}_0}^{\vec{r}} \vec{A} \cdot d\vec{\ell} \right]$

Fractional Quantum Hall Effect: a refresher

- Berry phase, Aharonov-Bohm: $\Psi(\vec{r}) \sim \exp \left[\frac{ie}{\hbar c} \int_{\vec{r}_0}^{\vec{r}} \vec{A} \cdot d\vec{\ell} \right]$

$$\oint_{\Gamma} \vec{A} \cdot d\vec{\ell} = \iint_{\Omega} \vec{\nabla} \times \vec{A} \cdot dS = N \frac{hc}{e} \rightarrow \Phi = N\Phi_0$$

Fractional Quantum Hall Effect: a refresher

- Berry phase, Aharonov-Bohm: $\Psi(\vec{r}) \sim \exp \left[\frac{ie}{\hbar c} \int_{\vec{r}_0}^{\vec{r}} \vec{A} \cdot d\vec{\ell} \right]$

$$\oint_{\Gamma} \vec{A} \cdot d\vec{\ell} = \iint_{\Omega} \vec{\nabla} \times \vec{A} \cdot dS = N \frac{hc}{e} \rightarrow \Phi = N\Phi_0$$

- Quantization (Streda's formula): $\sigma = \frac{\delta A}{\delta J} \sim \frac{\delta N_\phi}{\delta N_e} = \nu$
- The $\nu = 5/2$ state: Laughlin's wavefunction:

$$\Psi_{GS}(z_1, \dots, z_{2n}) = \text{Pf}[(z_i - z_j)^{-1}] \prod_{k < l} (z_k - z_l)^2 \prod_j e^{-|z_j|^2/4},$$

$$\text{Pf}(A) = \sqrt{\text{Det}(A)}$$

Two-point functions in the $\nu = 5/2$ state

$$\Psi_{2qh}(z_1, z_2; z_3, \dots, z_{2n}) = \text{Pf} \left[\frac{z_1 - z_2}{z_1 z_2 (z_i - z_j)} \right] \prod_m z_m \prod_{k < l} (z_k - z_l)^2 \prod_j e^{-|z_j|^2/4},$$

$$\Psi(z_2, z_1) = \Psi(z_1, z_2) e^{\frac{2i\pi}{\kappa}}, \quad \kappa \in \mathbb{Z}.$$

- $\kappa = 1$: bosons
- $\kappa = 2$: fermions
- $\kappa = 4$: Pfaffians (for $\nu = 5/2$)

Kitaev's proposal: the good news ...

- Non-local (noise-protected)

Kitaev's proposal: the good news ...

- Non-local (noise-protected)
- Non-decaying (zero modes, topological)

Kitaev's proposal: the good news ...

- Non-local (noise-protected)
- Non-decaying (zero modes, topological)
- Non-abelian (non-trivial information)

... and the bad

- Difficult experiment: attempts for $\nu = 5/2$

... and the bad

- Difficult experiment: attempts for $\nu = 5/2$
- Worse yet - not universal: representations of $SU(2)_2$ not dense in the Hilbert state-space

... and the bad

- Difficult experiment: attempts for $\nu = 5/2$
- Worse yet - not universal: representations of $SU(2)_2$ not dense in the Hilbert state-space
- Ideas?

... and the bad

- Difficult experiment: attempts for $\nu = 5/2$
- Worse yet - not universal: representations of $SU(2)_2$ not dense in the Hilbert state-space
- Ideas?
- 2D quantum field \rightarrow CFT \rightarrow Virasoro algebra (conformal invariance)

... and the bad

- Difficult experiment: attempts for $\nu = 5/2$
- Worse yet - not universal: representations of $SU(2)_2$ not dense in the Hilbert state-space
- Ideas?
- 2D quantum field \rightarrow CFT \rightarrow Virasoro algebra (conformal invariance)

$$[L_n, L_m] = (n - m)L_{n+m} + \frac{c}{12}n(n^2 - 1)\delta_{n+m,0}$$

- Bosons: $c = 1$;

... and the bad

- Difficult experiment: attempts for $\nu = 5/2$
- Worse yet - not universal: representations of $SU(2)_2$ not dense in the Hilbert state-space
- Ideas?
- 2D quantum field \rightarrow CFT \rightarrow Virasoro algebra (conformal invariance)

$$[L_n, L_m] = (n - m)L_{n+m} + \frac{c}{12}n(n^2 - 1)\delta_{n+m,0}$$

- Bosons: $c = 1$; fermions: $c = 1/2$

... and the bad

- Difficult experiment: attempts for $\nu = 5/2$
- Worse yet - not universal: representations of $SU(2)_2$ not dense in the Hilbert state-space
- Ideas?
- 2D quantum field \rightarrow CFT \rightarrow Virasoro algebra (conformal invariance)

$$[L_n, L_m] = (n - m)L_{n+m} + \frac{c}{12}n(n^2 - 1)\delta_{n+m,0}$$

- Bosons: $c = 1$; fermions: $c = 1/2$; non-abelian anyons?

Foundations

Quantum Information ...

Kac-Moody-Virasoro algebras (loop groups)

Kac-Moody-Virasoro algebras (loop groups)

- Loop current algebra (local anomaly):

$$J^i = \frac{\delta W}{\delta A_i} = (\bar{\psi} \gamma^i \psi) = \sum_{-\infty}^{\infty} J_n^i z^{-n-1},$$

$$[J_n^i, J_m^k] = i f^{ik\ell} J_{n+m}^\ell$$

Kac-Moody-Virasoro algebras (loop groups)

- Loop current algebra (local anomaly):

$$J^i = \frac{\delta W}{\delta A_i} = (\bar{\psi} \gamma^i \psi) = \sum_{-\infty}^{\infty} J_n^i z^{-n-1},$$

$$[J_n^i, J_m^k] = i f^{ik\ell} J_{n+m}^\ell + \frac{k}{2} \delta^{ij} \delta_{n+m,0}$$

- Wess-Zumino-Novikov-Witten theory (maps from space-time to $SU(2)$)

Kac-Moody-Virasoro algebras (loop groups)

- Loop current algebra (local anomaly):

$$J^i = \frac{\delta W}{\delta A_i} = (\bar{\psi} \gamma^i \psi) = \sum_{-\infty}^{\infty} J_n^i z^{-n-1},$$

$$[J_n^i, J_m^k] = i f^{ik\ell} J_{n+m}^\ell + \frac{k}{2} \delta^{ij} \delta_{n+m,0}$$

- Wess-Zumino-Novikov-Witten theory (maps from space-time to $SU(2)$)

$$c = \frac{3k}{k+2},$$

Kac-Moody-Virasoro algebras (loop groups)

- Loop current algebra (local anomaly):

$$J^i = \frac{\delta W}{\delta A_i} = (\bar{\psi} \gamma^i \psi) = \sum_{-\infty}^{\infty} J_n^i z^{-n-1},$$

$$[J_n^i, J_m^k] = i f^{ik\ell} J_{n+m}^\ell + \frac{k}{2} \delta^{ij} \delta_{n+m,0}$$

- Wess-Zumino-Novikov-Witten theory (maps from space-time to $SU(2)$)

$$c = \frac{3k}{k+2}, \quad k=1 \rightarrow c=1; \quad k=2 \rightarrow c=\frac{3}{2}=1+\frac{1}{2}.$$

Foundations

Quantum Information ...

WZNW and Knizhnik-Zamolodchikov equations

WZNW and Knizhnik-Zamolodchikov equations

- WZNW: maps $g(z)$ from \mathbb{C} to $SU(2)$ (generically Lie group G):

$$S = \textcolor{red}{k} \cdot \frac{1}{4\pi} \left[\text{Tr} \int d^2\xi \frac{1}{2} \partial^\mu g^{-1} \partial_\mu g + \epsilon^{\mu\nu} \int_0^1 d\tau \int d^2\xi g^{-1} \partial_\tau g g^{-1} \partial_\mu g g^{-1} \partial_\nu g \right]$$

- Equations for field g (K-Z equations):

WZNW and Knizhnik-Zamolodchikov equations

- WZNW: maps $g(z)$ from \mathbb{C} to $SU(2)$ (generically Lie group G):

$$S = \textcolor{red}{k} \cdot \frac{1}{4\pi} \left[\text{Tr} \int d^2\xi \frac{1}{2} \partial^\mu g^{-1} \partial_\mu g + \epsilon^{\mu\nu} \int_0^1 d\tau \int d^2\xi g^{-1} \partial_\tau g g^{-1} \partial_\mu g g^{-1} \partial_\nu g \right]$$

- Equations for field g (K-Z equations):

$$\left[(k+2) \frac{\partial}{\partial z_i} + \sum_{j \neq i}^n \frac{\tau_i^a \tau_j^a}{z_i - z_j} \right] \langle g(z_1) g(z_2) \dots g(z_n) \rangle = 0.$$

The solution in search of a model

Wanted:

- Non-local (multi-particle)

The solution in search of a model

Wanted:

- Non-local (multi-particle)
- Non-abelian anyons

The solution in search of a model

Wanted:

- Non-local (multi-particle)
- Non-abelian anyons
- Degenerate (zero modes)

The solution in search of a model

Wanted:

- Non-local (multi-particle)
- Non-abelian anyons
- Degenerate (zero modes)
- Dense representations

Is there such a physical system within reach ?

The pairing model

Richardson (1964), Gaudin (1972)

$$\hat{H} = \sum_{\mathbf{p}, \sigma} \epsilon_{\mathbf{p}} \hat{c}_{\mathbf{p}, \sigma}^\dagger \hat{c}_{\mathbf{p}, \sigma} - g \sum_{\mathbf{p}, \mathbf{k}} \hat{c}_{\mathbf{p} \uparrow}^\dagger \hat{c}_{-\mathbf{p} \downarrow}^\dagger \hat{c}_{-\mathbf{k} \downarrow} \hat{c}_{\mathbf{k} \uparrow}$$

The pairing model

Richardson (1964), Gaudin (1972)

$$\hat{H} = \sum_{\mathbf{p}, \sigma} \epsilon_{\mathbf{p}} \hat{c}_{\mathbf{p}, \sigma}^\dagger \hat{c}_{\mathbf{p}, \sigma} - g \sum_{\mathbf{p}, \mathbf{k}} \hat{c}_{\mathbf{p} \uparrow}^\dagger \hat{c}_{-\mathbf{p} \downarrow}^\dagger \hat{c}_{-\mathbf{k} \downarrow} \hat{c}_{\mathbf{k} \uparrow}$$

$$\mathbf{c} \rightarrow \mathbf{t}, \quad [t_i^3, t_j^\pm] = \pm \delta_{ij} t_j^\pm, \quad [t_i^+, t_j^-] = 2\delta_{ij} t_j^3,$$

The pairing model

Richardson (1964), Gaudin (1972)

$$\hat{H} = \sum_{\mathbf{p}, \sigma} \epsilon_{\mathbf{p}} \hat{c}_{\mathbf{p}, \sigma}^\dagger \hat{c}_{\mathbf{p}, \sigma} - g \sum_{\mathbf{p}, \mathbf{k}} \hat{c}_{\mathbf{p}}^\dagger \hat{c}_{-\mathbf{p}}^\dagger \hat{c}_{-\mathbf{k}} \hat{c}_{\mathbf{k}}^\dagger$$

$$\mathbf{c} \rightarrow \mathbf{t}, \quad [t_i^3, t_j^\pm] = \pm \delta_{ij} t_j^\pm, \quad [t_i^+, t_j^-] = 2\delta_{ij} t_j^3,$$

$$H_P = \sum_{l \in \Lambda} 2\epsilon_l t_l^3 - g \sum_{l, l'} t_l^+ t_{l'}^- = \sum_{l \in \Lambda} 2\epsilon_l t_l^3 - g \mathbf{t}^+ \cdot \mathbf{t}^-$$

Solution of the Richardson model

$$H_P = 2 \sum_{l \in \Lambda} \epsilon_l R_l + \text{const}, \quad R_l = t_l^3 - \frac{g}{2} \sum_{l' \neq l} \frac{\mathbf{t}_l \cdot \mathbf{t}_{l'}}{\epsilon_l - \epsilon_{l'}}$$

Solution of the Richardson model

$$H_P = 2 \sum_{l \in \Lambda} \epsilon_l R_l + \text{const}, \quad R_l = t_l^3 - \frac{g}{2} \sum_{l' \neq l} \frac{\mathbf{t}_l \cdot \mathbf{t}_{l'}}{\epsilon_l - \epsilon_{l'}}$$

$$[R_i, R_k] = 0$$

Solution of the Richardson model

$$H_P = 2 \sum_{l \in \Lambda} \epsilon_l R_l + \text{const}, \quad R_l = t_l^3 - \frac{g}{2} \sum_{l' \neq l} \frac{\mathbf{t}_l \cdot \mathbf{t}_{l'}}{\epsilon_l - \epsilon_{l'}}$$

$[R_i, R_k] = 0$

$$|\Psi\rangle = \prod_{k=1}^N b_k^\dagger |0\rangle, , \quad b_k^\dagger = \sum_l \frac{t_l^\dagger}{2\epsilon_l - e_k}, \frac{1}{g} = \sum_{p \neq k} \frac{2}{e_k - e_p} + \sum_l \frac{1}{2\epsilon_l - e_k}$$

Solution of the Richardson model

$$H_P = 2 \sum_{l \in \Lambda} \epsilon_l R_l + \text{const}, \quad R_l = t_l^3 - \frac{g}{2} \sum_{l' \neq l} \frac{\mathbf{t}_l \cdot \mathbf{t}_{l'}}{\epsilon_l - \epsilon_{l'}}$$

$[R_i, R_k] = 0$

$$|\Psi\rangle = \prod_{k=1}^N b_k^\dagger |0\rangle, , \quad b_k^\dagger = \sum_l \frac{t_l^\dagger}{2\epsilon_l - e_k}, \quad \frac{1}{g} = \sum_{p \neq k} \frac{2}{e_k - e_p} + \sum_l \frac{1}{2\epsilon_l - e_k}$$

$$R_k \Psi = 0, \quad H_P \Psi = 0.$$

Richardson-Gaudin as singular limit of WZNW

- String of results (1991 – 1999):

Richardson-Gaudin as singular limit of WZNW

- String of results (1991 – 1999):
- Reformulate Richardson-Gaudin on a torus, group G

Richardson-Gaudin as singular limit of WZNW

- String of results (1991 – 1999):
- Reformulate Richardson-Gaudin on a torus, group G
- Torus \rightarrow Cylinder \rightarrow Plane = Elliptic \rightarrow Trigonometric \rightarrow Rational

Richardson-Gaudin as singular limit of WZNW

- String of results (1991 – 1999):
- Reformulate Richardson-Gaudin on a torus, group G
- Torus \rightarrow Cylinder \rightarrow Plane = Elliptic \rightarrow Trigonometric \rightarrow Rational
- $k + 2$ = torus modular parameter:

Richardson-Gaudin as singular limit of WZNW

- String of results (1991 – 1999):
- Reformulate Richardson-Gaudin on a torus, group G
- Torus \rightarrow Cylinder \rightarrow Plane = Elliptic \rightarrow Trigonometric \rightarrow Rational
- $k + 2 =$ torus modular parameter:

$$(k+2) \frac{\partial \Psi}{\partial \epsilon_i} + \tilde{R}_i \Psi = 0, \quad \tilde{R}_i \sim \sum_{i \neq j} \tau_i^a \tau_j^a \zeta(\epsilon_i - \epsilon_j)$$

$$c = \frac{3k}{k+2} \rightarrow -\infty$$

The limit $c \rightarrow -\infty$: how to approach it?

- Problem: $e^{i\pi/(k+2)}$ - anything

The limit $c \rightarrow -\infty$: how to approach it?

- Problem: $e^{i\pi/(k+2)}$ - anything
- Special choice in elliptic formulation gives $\frac{c}{12} \in \mathbb{Z}$: $c = -12m$

The limit $c \rightarrow -\infty$: how to approach it?

- Problem: $e^{i\pi/(k+2)}$ - anything
- Special choice in elliptic formulation gives $\frac{c}{12} \in \mathbb{Z}$: $c = -12m$

$$k = \frac{-8m}{4m + 1}, \quad e^{i\pi/(k+2)} = e^{i\pi/2}$$

The limit $c \rightarrow -\infty$: how to approach it?

- Problem: $e^{i\pi/(k+2)}$ - anything
- Special choice in elliptic formulation gives $\frac{c}{12} \in \mathbb{Z}$: $c = -12m$

$$k = \frac{-8m}{4m + 1}, \quad e^{i\pi/(k+2)} = e^{i\pi/2}$$

- Same commutation relations like $k = 2$, but different level of $SU(2)_k$
- Fractional level representations: Gaberdiel et. al.

Summary of properties

We have:

- Non-local (collective modes)

Summary of properties

We have:

- Non-local (collective modes)
- Non-abelian anyons with $k = 2$ statistics

Summary of properties

We have:

- Non-local (collective modes)
- Non-abelian anyons with $k = 2$ statistics
- Degenerate (zero modes)

Summary of properties

We have:

- Non-local (collective modes)
- Non-abelian anyons with $k = 2$ statistics
- Degenerate (zero modes)
- Dense representations $k \neq 1, 2, 4$

Semiclassical approximation

$$H_{MF} = \sum_{l \in \Lambda} \epsilon_l S_l^3 - \frac{g}{4} |J^-|^2, \quad \mathbf{J} = \sum_{l \in \Lambda} \mathbf{S}_l$$

Commutators give Poisson brackets, nonlinear Bloch equations

$$\{S_i^\alpha, S_j^\beta\} = 2\epsilon^{\alpha\beta\gamma} S_i^\gamma \delta_{ij}, \quad \boxed{\dot{\vec{S}}_i = 2(-\vec{\Delta} + \epsilon_i \hat{z}) \times \vec{S}_i}$$

Classical gap parameter, constants of motion:

$$\vec{\Delta} = \frac{1}{2}(gJ_x, gJ_y, 0), \quad r_i = \frac{1}{2} \left[S_i^z - \frac{g}{2} \sum_{j \neq i} \frac{\vec{S}_i \cdot \vec{S}_j}{\epsilon_i - \epsilon_j} \right]$$

Elliptic solutions

Levitov-Barankov-Spivak Ansatz:

$$\epsilon_k = -\epsilon_{-k}, \quad S_k^{y,z} = -S_{-k}^{y,z}, \quad S_k^x = S_{-k}^x$$

$$S_k^x = A_k \Omega, \quad S_k^y = B_k \dot{\Omega}, \quad S_k^z = C_k + D_k \Omega^2, \quad \Omega = |\Delta|$$

$$|\vec{S}_k| = 1 \Rightarrow (\dot{\Omega})^2 = (\Omega^2 - \Omega_-^2)(\Omega_+^2 - \Omega^2).$$

Constants Ω_{\pm} fixed by initial conditions and

$$1 = g \sum_k A_k$$

Multi-frequency (multi-cut) hyperelliptic solutions

Dubrovin equations:

$$i\frac{\dot{\Omega}}{\Omega} = \sum_{k=1}^{n-1} u_k, \quad \dot{u}_i = \frac{2iy(u_i)}{\prod_{j \neq i} (u_i - u_j)},$$

Spectral curve:

$$y^2(\lambda) = \prod_{i=1}^n (\lambda - \epsilon_i)^2 \det \left[\sigma_3 + \frac{g}{2} \sum_{i=0}^n \frac{\vec{S}_i \cdot \vec{\sigma}}{\lambda - \epsilon_i} \right] = \prod_{i=0}^{2n} (\lambda - E_i)$$

Open questions

- Imposing additional (symmetry) constraints?

Open questions

- Imposing additional (symmetry) constraints?
- Practical gate design?

Open questions

- Imposing additional (symmetry) constraints?
- Practical gate design?
- Using excited states?

Open questions

- Imposing additional (symmetry) constraints?
- Practical gate design?
- Using excited states?