

Systèmes de Référence Temps-Espace

# Fundamental Physics Tests using Rubidium and Cesium Fountains

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# Introduction

- LLI test using a Cryogenic Sapphire Oscillator
- LLI test using a Cs fountain
- LPI test: Stability of fundamental constants

# Prospects



- Experimental tests of fundamental physical laws
  - □ Einstein Equivalence Principle
    - Focus on LLI and LPI
- Contribute to constraining unification theories
   String theories, loop gravity,...
- LLI experiments analyzed within the SME framework
  - A general Lorentz violating extension of the Standard Model
  - □ Large number of parameters
  - Better insight of which part of the standard model is tested by a given experiment
  - Photon sector Maxwell equations with modified coefficients, 19 parameters
  - □ Matter sector: 44 parameters per particle  $(p^+,e^-,n,...)$



# TIME AND FREQUENCY METROLOGY APPLICATIONS

TAI calibration, more than 15 over the past 4 years



Secondary representation of the SI second (2004)
 □ Rb(hfs)
 □ ν<sub>Rb</sub>(2002) = 6 834 682 610.904 324(4)(7) Hz (1.3×10<sup>-15</sup>) (CCTF: 3×10<sup>-15</sup>)

## Support to the development of PHARAO/ACES

- □ Test of µW synthesizer IM, Ramsey cavity FM,...
- PHARAO EM is now operated as a clock, poster at this conference

# LLI test using a Cryogenic Sapphire Oscillator

P. Wolf, S. Bize, A. Clairon, A. Luiten, G. Santarelli, M. Tobar, Phys. Rev. Lett. 90, 060402 (2003) Gen. Rel. Grav. 36, 2351 (2004) Phys. Rev. D70 051902(R) (2004)

# SME ANALYSIS OF A MICROWAVE RESONATOR

- The mode frequency is perturbed by a term involving 7 relevant SME coefficients
- Earth motion induces modulations of the SME term (SME coefficients are tensor components attached to a supposedly preferred frame)

$$\frac{\Delta f}{f}(t) = \sum_{\omega_i} \left( C_{\omega_i} \cos(\omega_i t) + S_{\omega_i} \sin(\omega_i t) \right)$$

with  $\omega_1 = \omega_1 2\omega_1, \omega \pm \Omega, 2\omega \pm \Omega$ .

(sidereal, semi-sidereal pulsations with orbital sidebands

Detected wrt H-maser





# SUMMARY OF DATA ANALYSIS AND RESULTS

- 222 days, spanning from Sept. 2002 to Jan. 2004
- Analysis accounted for
  - 2 different methods
  - Non-white noise
  - Contamination by diurnal modulation
  - Evaluation of systematic shifts

TABLE II. Contributions from systematic effects to the amplitudes  $A_i$  (parts in 10<sup>16</sup>) at three frequencies  $\omega_i$ .

Effect	$\omega_\oplus - \Omega_\oplus$	$\omega_\oplus$	$2\omega_{\oplus}$
H-maser	<5	<5	<5
Tilt	3	3	1
Gravity	0.3	0.3	0.3
B-field	< 0.1	< 0.1	< 0.1
Temperature	<1	<1	<1
Atm. Pressure	2.3	0.3	0.4
Total	6.4	5.9	5.2

## Results

- □ Improvement by a factor of 8 for three SME parameters
- □ Non-zero at 2σ for 2 parameters but inconsistent with Müller *et al.* => a statistical coincidence, NOT a LLI violation

Müller et al. Phys. Rev. Lett. 91, 020401 (2003)

 $\square$  Better measurements with rotating oscillators (factor  $\sim 10)$ 

Stanwix et al. (2005), Herrman et al. (2005) 8

# LLI test using a Cs fountain

P. Wolf, F. Chapelet, S. Bize, A. Clairon, Phys. Rev. Lett. 96, 060801 (2006)

# SME APPLIED TO CESIUM HFS

SME shift of atomic energy levels in the local frame

$$\delta E(m_F,F) = \frac{m_F}{F} \sum_{e^-,p^+,n} \left(\beta_w \widetilde{b}_3^w + \delta_w \widetilde{d}_3^w + \kappa_w \widetilde{g}_d^w\right) + \frac{3m_F^2 - F(F+1)}{3F^2 - F(F+1)} \sum_{e^-,p^+,n} \left(\gamma_w \widetilde{c}_q^w + \lambda_w \widetilde{g}_q^w\right)$$

- $\square$   $\beta_w, \delta_w, \kappa_w, \gamma_w, \lambda_w$  are specific to the atom and the particular state
- The tilde coefficients are combinations of SME parameters
- They are in general time dependent due to atom motion wrt supposedly preferred frame
- Cs hyperfine transition in the SME

An observable which free of 1st order Zeeman effect

$$\upsilon_{+3} + \upsilon_{-3} - 2\upsilon_0 = \frac{1}{7h} K_p \widetilde{c}_q^{\ p} - \frac{9}{8} K_z^{(2)} B^2$$

 $K_{\rm p} \approx 10^{-2}$ ;  $K_{\rm e} \approx 10^{-5}$  (neglected)

# EXPERIMENTAL STRATEGY

- Alternate m<sub>F</sub> = 3 and m<sub>F</sub> = -3 measurement every second (interleaved servo-loops).
- Measure m<sub>F</sub> = 0 clock transition every 400 s (reference).
- Limited by stability of magnetic field at *τ* < 4 s.</li>
- Reduce launching height to optimize stability of observable.



#### Transforming to sun-frame SME parameters:

$$\left|\widetilde{c}_{q}^{p} = A + C_{\omega_{\oplus}}\cos(\omega_{\oplus}t) + S_{\omega_{\oplus}}\sin(\omega_{\oplus}t) + C_{2\omega_{\oplus}}\cos(2\omega_{\oplus}t) + S_{2\omega_{\oplus}}\sin(2\omega_{\oplus}t)\right|$$

- A,  $C_{i}$ ,  $S_{i}$ , are functions of the 8 proton components:  $\tilde{c}_{Q}, \tilde{c}_{X}, \tilde{c}_{Y}, \tilde{c}_{Z}, \tilde{c}_{-}, \tilde{c}_{TX}, \tilde{c}_{TY}, \tilde{c}_{TZ}$
- 3 proton components (  $\tilde{c}_{TX}$ ,  $\tilde{c}_{TY}$ ,  $\tilde{c}_{TZ}$ ) are suppressed by v<sub>⊕</sub>/c ≈ 10<sup>-4</sup>
- Search for offset, sidereal and semi-sidereal signatures in the observable

## DATA AND STATISTICAL ANALYSIS



21 days of data in April 2005, 14 days in September 2005. Least squares fit:

$$\begin{split} A = -5.3(0.04); \ C_{\omega_{\oplus}} = 0.1(0.06); \quad S_{\omega_{\oplus}} = -0.03(0.06) & \text{in mHz} \\ C_{2\omega_{\oplus}} = 0.04(0.06); \ S_{2\omega_{\oplus}} = 0.03(0.06) \end{split}$$

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# SYSTEMATICS: Residual 1st order Zeeman Shift

- Magnetic field gradients and non-identical trajectories of  $m_F = +3$  and  $m_F = -3$  atoms can lead to incomplete cancellation of  $Z^{(1)}$ .
- Confirmed by TOF difference  $\approx 158 \ \mu s$  ( $\rightarrow 623 \ \mu m$ ).
- Variation of B with launching height ≈ 0.02 pT/mm (at apogee).
  ⇒ MC simulation gives offset of only ≈ 6 µHz.
- Contrast as function of m<sub>F</sub>: 0.94, 0.93, 0.87, 0.75
- MC simulation with only vertical B gradient cannot reproduce the contrast ⇒ horizontal B gradient of ≈ 6 pT/mm (≈ 2 pT/mm from tilt measurements).
- Complete MC simulation, assuming horizontal asymmetry between trajectories is same as vertical (worst case) gives offset ≈ 25 mHz.
- Fitting sidereal and semi-sidereal variations to the TOF difference and using the above gradients we obtain no significant effect within the statistical uncertainties (≈ 0.03 mHz at both frequencies). We take this as our upper limit of the time varying part of the residual first order Zeeman.



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	(In Gev)	Parameter	$p^+$	n	$e^{-}$
8 proton paramet	ters	$\tilde{b}_X, \tilde{b}_Y$	$10^{-27}$	$10^{-31}$	$10^{-29}$
	$\frown$	Б <sub>Z</sub>			$10^{-28}$
$(c_Q) = -0.3(2.2) \times 10^{-22}$	$\widetilde{c}_{TX} = -2.7(3.0) \times 10^{-21}$	$\tilde{b_T}, \tilde{g}_T, \tilde{H}_{JT}, \tilde{d}_{\pm}$		$10^{-27}$	
$\sim$ 1.0(2.0). (10 <sup>-25</sup>	$\sim$	$\tilde{d}_{O}, \tilde{d}_{XY}, \tilde{d}_{YZ}$		$10^{-27}$	
$c_{-} = -1.8(2.8) \times 10^{-1}$	$c_{TY} = -0.2(3.0) \times 10^{-21}$	$d_X^2, d_Y$	$10^{-25}$	$10^{-29}$	$10^{-22}$
$\widetilde{c} = 0.6(1.2) \times 10^{-25}$	$\sqrt{\widetilde{a}} = \frac{1}{\sqrt{2}} 0.4(2.0) \times 10^{-21}$	$\tilde{d}_{XZ}, \tilde{d}_{Z}$			
$c_X = 0.0(1.2) \times 10^{-25}$	$C_{TZ} = -0.4(2.0) \times 10$	$\tilde{g}_{DX}, \tilde{g}_{DY}$	$10^{-25}$	$10^{-29}$	$10^{-22}$
$\tilde{c}_{y} \neq -1.9(1.2) \times 10^{-2.5}$		ÕDZ, ÕIK			
$\sim 1 1 (2 9) \times 10^{-25}$		$\tilde{g}_c$		$10^{-27}$	
$C_Z = -1.4(2.8) \times 10$		<u> ğ</u> –, <u>ğ</u> о, <u>ğ</u> ті			
$\smile$		Ĉ <sub>Q</sub>	10-22(-11)		$10^{-9}$
		$\tilde{c}_{\mathbf{X}}, \tilde{c}_{\mathbf{Y}}$	$10^{-25}$	$10^{-25}$	$10^{-19}$
		č <sub>Z</sub> , č_	10 <sup>-25</sup>	$10^{-27}$	$10^{-19}$
		$\tilde{c}_{TJ}$	10-21(-8)	• • •	$10^{-6}$

- Sensitivity to  $c_{TJ}$  reduced by a factor  $v_{\oplus}/c$  ( $\approx 10^{-4}$ ).
- Assuming no cancellation between  $c_{TJ}$  and others.
- First measurements of four components.
- Improvement by 11 and 13 orders of magnitude on previous limits (re-analysis of IS experiment, [Lane C., PRD 2005]).
- Dominated by statistical uncertainty (factor 2) except for  $c_0$ .

# LPI test: Stability of fundamental constants

S. Bize et al., J. Phys. B: At. Mol. Opt. Phys. **38**, S44 (2005) S. Bize et al., C.R. Physique **5**, 829 (2004) M. Fischer et al., Phys. Rev. Lett. **92**, 230802 (2004) H. Marion et al., Phys. Rev. Lett. **90**, 150801 (2003) Y. Sortais et al., Phys. Scripta **T95**, 50 (2001) M. Niering et al., Phys. Rev. Lett. **84**, 5496 (2000) S. Bize et al., Europhys. Lett. **45**, 558 (1999) **1** 

# COMPARISON OF Rb vs Cs HFS and H(1S-2S) vs Cs



one data point ⇔ ~1 to 2 months of measurements, with many checks of systematic shifts

$$\frac{d}{dt} \ln \left(\frac{\nu_{\rm Rb}}{\nu_{\rm Cs}}\right) = (-0.5 \pm 5.3) \times 10^{-16} \text{ yr}^{-1}$$

$$J. Prestage, et al., PRL (1995)$$

$$V. Dzuba, et al., PRL (1999)$$

$$\frac{d}{dt} \ln \left(\frac{g_{\rm Cs}}{g_{\rm Rb}} \alpha^{0.49}\right) = (0.5 \pm 5.3) \times 10^{-16} \text{ yr}^{-1}$$

With further theory, nuclear g-factors can be related to more fundamental parameters

$$\frac{d}{dt} \ln \left( \alpha^{0.49} \left[ m_q / \Lambda_{\rm QCD} \right]^{0.174} \left[ m_s / \Lambda_{\rm QCD} \right]^{0.027} \right) = (0.5 \pm 5.3) \times 10^{-16} \text{ yr}^{-1}$$

H(1S-2S) vs Cs over ~3 years (with transportable fountain at MPQ Garching)

$$\frac{\partial}{\partial t}\ln\frac{\nu_{\rm Cs}}{\nu_{\rm H}} = \frac{\partial}{\partial t} \left(\ln\frac{\mu_{\rm Cs}}{\mu_B} + (2.0 \pm 0.8)\ln\alpha\right) = (3.2 \pm 6.3) \times 10^{-15} \,\rm yr^{-1}$$

Combined with Hg<sup>+</sup> vs Cs (NIST), Yb<sup>+</sup> vs Cs (PTB), these measurements independently constrain the stability of the electroweak interaction ( $\alpha$ ) and of the strong interaction at 2x10<sup>-15</sup> per year

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# Current status and prospects in the development of LNE-SYRTE fountain ensemble

# FREQUENCY COMPARISON AT THE 10<sup>-16</sup> LEVEL

(2004)

*S. Bize et al.*, *C.R. Physique 5, 829 (2004) C. Vian et al., IEEE Trans. Instrum. Meas. 54, 833 (2005)* 

Fractional frequency instability (Allan deviation) between FO1 and FO2 fountains & fractional frequency instability of FO1 and FO2 against the CSO locked to a hydrogen maser



## UNCERTAINTY BUDGET

(2004 and improvements since then)

Systematic fractional frequency shifts for FO1 and FO2 <sup>133</sup> Cs fountains				
	<b>FO1</b> (x10 <sup>16</sup> )	<b>FO2</b> (×10 <sup>16</sup> )		
Quadratic Zeeman effect	1199.7 ± 4.5	1927.3 ± 0.3		
Blackbody radiation	-162.8 ± 2.5	-168.2 ± 2.5 ⇒	>±C	
Collisions and cavity pulling	-197.9 ± 2.4	-357.5 ± 2.0 ⇒	>±1	
Microwave spectral purity & leakage	$0.0 \pm 3.3$	0.0 ± 4.3 ⇒	>±(	
First order Doppler effect	< 3	<3 =	⇒±'	
Ramsey & Rabi pulling	< 1	< 1		
Microwave recoil	< 1.4	< 1.4	⇒<(	
Second order Doppler effect	< 0.08	< 0.08		
Background collisions	< 1	< 1		
Total uncertainty	± 7.5	± 6.5		

# FO2 SOON OPERATED AS A DUAL FOUNTAIN"Dichroic" collimatorsCs cooli

Cs cooling:  $\lambda = 852 \text{ nm}$ Rb cooling:  $\lambda = 780 \text{ nm}$ 

6 collimators for the optical molasses (+ 2 collimators for detection + 1 pusher beam)



 $\Rightarrow$  Now attached to the FO2 fountain with Rb 2DMOT



- Fountain accuracy of few 10<sup>-16</sup>
- Stability of constants in the interesting 10<sup>-17</sup> yr<sup>-1</sup> range
- Improved SME tests with dual fountain
- Stability of constants using 2 µW clocks (Rb, Cs) and 2 optical lattice clocks (Sr, Hg)
- Towards PHARAO/ACES ground segment
   Quasi-continuous operation of atomic fountain
   Improved local timescale