GRAVITATIONAL WAVE – BACKGROUNDS –



GGI LECTURES ON THE THEORY OF FUNDAMENTAL INTERACTIONS – 2022 Program (3rd week)

OUTLINE



2) GWs from Inflation

- 3) GWs from Preheating
- 4) GWs from Phase Transitions
- 5) GWs from Cosmic Defects

OUTLINE



2) GWs from Inflation

- 3) GWs from Preheating
- 4) GWs from Phase Transitions
- 5) GWs from Cosmic Defects

A PRIMER ON GRAVITATIONAL WAVES

General Relativity (GR)



 $G_{\mu
u} = \frac{1}{m_p^2} T_{\mu
u}$ geometry matter

$$[m_p = (8\pi G)^{-1/2} = 2.44 \cdot 10^{18} \,\text{GeV}]$$

$$ds^2 = g_{\mu\nu}(x)dx^{\mu}dx^{\nu}$$

DIFF:
$$x^{\mu} \to x'^{\mu}(x)$$

symmetry





$$\begin{array}{ll} \underset{\mu\nu}{\text{metric}} & & \uparrow \\ G_{\mu\nu} \equiv \mathcal{D}[g_{\alpha\beta}] & = & m_p^{-2}T_{\mu\nu}(\text{Matt, Rad, Top.Defects, DarkEnergy, }...) \\ & \downarrow & \\ \text{source} \\ \text{2nd order, non-Linear} \end{array}$$







Cosmological Principle Background metric and matter Homogeneous & Isotropy

FLRW expanding Universe ! $ds^{2} = g_{\mu\nu}dx^{\mu}dx^{\nu} = a^{2}(t)\left(-dt^{2} + d\mathbf{x}^{2}\right)$



How do we define GWs ?



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How do we define GWs ?

$g_{\alpha\beta} = \bar{g}_{\alpha\beta} + \delta g_{\alpha\beta}$ Let's continue this approach...

I hope you took a good load of coffee ('cause you are gonna need it)



Definition of GWs 1st approach

1st approach to GWs



1st approach to GWs

1st approach to GWs

 $\begin{array}{l} \text{Minkowski} \\ f \\ g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu}(x) \\ (\left| h_{\mu\nu} \right| \ll 1 \end{array} \right) \begin{array}{l} \text{fixed} \\ \text{frame} \\ \end{array}$

1st approach to GWs

$$\begin{array}{l} \text{Minkowski} \\ f \\ g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu}(x) \\ (\left| h_{\mu\nu} \right| \ll 1 \end{array} \right) \\ \end{array}$$

DIFF:
$$x^{\mu} \rightarrow x'^{\mu}(x)$$

symmetry?

1st approach to GWs

Minkowski

$$g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu}(x)$$
fixed
($|h_{\mu\nu}| \ll 1$)

DIFF:
$$x^{\mu} \rightarrow x'^{\mu}(x)$$

Minkowski

1st approach to GWs $g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu}(x) \qquad \text{fixed} \\ (|h_{\mu\nu}| \ll 1) \qquad (|h_{\mu\nu}| \ll 1) \qquad (|h_{\mu\nu}| \ll 1) \qquad (|h_{\mu\nu}| \ll 1) \qquad (|h_{\mu\nu}| \gg 1) \qquad (|h_{\mu\nu}(x)| \lesssim |h_{\mu\nu}|) \qquad \text{residual} \\ (|\partial_{\mu}\xi_{\nu}(x)| \lesssim |h_{\mu\nu}|) \qquad (|h_{\mu\nu}(x) - h_{\mu\nu}(x) - \partial_{\mu}\xi_{\nu}|) \qquad (|h_{\mu\nu}(x) \to h_{\mu\nu}(x) - h_{\mu\nu}(x) - \partial_{\mu}\xi_{\nu}|) \qquad (|h_{\mu\nu}(x) \to h_{\mu\nu}(x) - h_{\mu\nu}(x) - \partial_{\mu}\xi_{\nu}|) \qquad (|h_{\mu\nu}(x) \to h_{\mu\nu}(x) - h_{\mu\nu}(x$



Notation:
$$\begin{cases} \partial_{(\mu}\xi_{\nu)} \equiv \partial_{\mu}\xi_{\nu} + \partial_{\nu}\xi_{\mu} \\ \partial_{[\mu}\xi_{\nu]} \equiv \partial_{\mu}\xi_{\nu} - \partial_{\nu}\xi_{\mu} \end{cases}$$

Minkowski

1st approach to GWs $g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu}(x) \qquad \text{fixed} \\ (|h_{\mu\nu}| \ll 1) \qquad (|h_{\mu\nu}| \ll 1) \qquad (|h_{\mu\nu}| \ll 1) \qquad (|h_{\mu\nu}| \ll 1) \qquad (|h_{\mu\nu}| \gg 1) \qquad (|h_{\mu\nu}(x)| \lesssim |h_{\mu\nu}|) \qquad \text{residual} \\ (|\partial_{\mu}\xi_{\nu}(x)| \lesssim |h_{\mu\nu}|) \qquad (|h_{\mu\nu}(x) - h_{\mu\nu}(x) - \partial_{\mu}\xi_{\nu}|) \qquad (|h_{\mu\nu}(x) \to h_{\mu\nu}(x) - h_{\mu\nu}(x) - \partial_{\mu}\xi_{\nu}|) \qquad (|h_{\mu\nu}(x) \to h_{\mu\nu}(x) - h_{\mu\nu}(x) - \partial_{\mu}\xi_{\nu}|) \qquad (|h_{\mu\nu}(x) \to h_{\mu\nu}(x) - h_{\mu\nu}(x$

1st approach to GWs

$$\begin{array}{l} {\rm Minkowski} \\ g_{\mu\nu} = \overset{\mbox{\boldmath\uparrow}}{\eta_{\mu\nu}} + h_{\mu\nu}(x) & {\rm fixed} \\ & (|h_{\mu\nu}| \ll 1 \) \end{array}$$

Let's expand Einstein Equations !

1st approach to GWs

$$\begin{array}{l} \text{Minkowski} \\ f \\ g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu}(x) \\ (\left| h_{\mu\nu} \right| \ll 1 \end{array} \right) \begin{array}{l} \text{fixed} \\ \text{frame} \\ \end{array}$$

$$\bar{h}_{\mu\nu} \equiv h_{\mu\nu} - \frac{1}{2}h\eta_{\mu\nu}$$






























1st approach to GWs

 $\begin{array}{l} \text{Minkowski} \\ f \\ g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu}(x) \\ (\left| h_{\mu\nu} \right| \ll 1 \end{array} \right) \\ \end{array}$

Is that all ?

1st approach to GWs

$$\begin{array}{l} \text{Minkowski} \\ g_{\mu\nu} = \stackrel{\uparrow}{\eta_{\mu\nu}} + h_{\mu\nu}(x) & \stackrel{\text{fixed}}{\underset{(|h_{\mu\nu}| \ll 1)}{\text{frame}}} \end{array}$$

Is that all ? Not really ...

$$\begin{array}{l} \text{Minkowski} \\ f \\ g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu}(x) \\ (\left| h_{\mu\nu} \right| \ll 1 \end{array} \right) \begin{array}{l} \text{fixed} \\ \text{frame} \\ \end{array}$$

$$x'^{\mu} = x^{\mu} + \xi^{\mu}(x)$$

with $\partial_{\alpha}\partial^{\alpha}\xi_{\mu} = 0$
(further residual gauge)

$$\begin{array}{l} \text{Minkowski} \\ f \\ g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu}(x) \\ (\left| h_{\mu\nu} \right| \ll 1 \end{array} \right) \begin{array}{l} \text{fixed} \\ \text{frame} \\ \end{array}$$

$$\begin{aligned} x'^{\mu} &= x^{\mu} + \xi^{\mu}(x) \\ \text{with } \partial_{\alpha} \partial^{\alpha} \xi_{\mu} &= 0 \\ \text{(further residual gauge)} \\ (\partial^{\mu} \bar{h}_{\mu\nu} = 0 \quad \rightarrow \quad \partial'^{\mu} \bar{h}'_{\mu\nu} = 0) \\ \text{(Lorentz preserving)} \end{aligned}$$

$$\begin{array}{l} \text{Minkowski} \\ f \\ g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu}(x) \\ (\left| h_{\mu\nu} \right| \ll 1 \end{array} \right) \begin{array}{l} \text{fixed} \\ \text{frame} \\ \end{array}$$

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$$x'^{\mu} = x^{\mu} + \xi^{\mu}(x)$$

with $\partial_{\alpha}\partial^{\alpha}\xi_{\mu} = 0$
(further residual gauge)

$$\label{eq:F} \begin{split} \mathbf{F} \ T_{\mu\nu} &= 0 \\ \mathbf{Outside} \\ \mathbf{Source} \end{split}$$



Minkowski $g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu}(x) \quad \text{fixed} \\ (|h_{\mu\nu}| \ll 1) \quad \text{frame}$ 1st approach to GWs $h^{0\mu} = 0, \qquad h^i_i = 0, \qquad \partial_j h_{ij} = 0$ $x'^{\mu} = x^{\mu} + \xi^{\mu}(x)$ (transversewith $\partial_{\alpha}\partial^{\alpha}\xi_{\mu}=0$ traceless $\partial_{\mu}\partial^{\mu}h_{ij} = 0$ IF $T_{\mu\nu} = 0$ (further residual gauge) gauge) Outside (6 - 4 = 2 d.o.f.)Source

Minkowski $g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu}(x) \qquad \text{fixed} \\ (|h_{\mu\nu}| \ll 1) \qquad \text{frame}$ 1st approach to GWs $h^{0\mu} = 0, \qquad h^i_i = 0, \qquad \partial_j h_{ij} = 0$ $x'^{\mu} = x^{\mu} + \xi^{\mu}(x)$ $\partial_{\alpha}\partial^{\alpha}\bar{h}_{\mu\nu} = -\frac{2}{m_{p}^{2}}T_{\mu\nu}$ (transv with $\partial_{\alpha}\partial^{\alpha}\xi_{\mu}=0$ **IF** $T_{\mu\nu} \neq 0$ (further residual gauge) gauge) Inside 6 - 4 = 2 d.o.f.? Source !

1st approach to GWs $g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu}(x)$ fixed $(|h_{\mu\nu}| \ll 1)$





(TT gauge: 6 - 4 = 2 d.o.f.)

$$h^{0\mu} = 0$$
, $h^i_i = 0$, $\partial_j h_{ij} = 0$
Outside
Source

(TT gauge: 6 - 4 = 2 d.o.f.)

$$h^{0\mu} = 0$$
, $h^i_i = 0$, $\partial_j h_{ij} = 0$
Outside
Source

$$\partial_{\mu}\partial^{\mu}h_{ij} = 0$$
 Wave Eq. \rightarrow Gravitational Waves !

1st approach to GWs

(TT gauge: 6 - 4 = 2 d.o.f.)

$$h^{0\mu} = 0$$
, $h^i_i = 0$, $\partial_j h_{ij} = 0$
Outside
Source

$$\partial_{\mu}\partial^{\mu}h_{ij} = 0$$
 Wave Eq. \rightarrow Gravitational Waves !

can GW be 'gauged away' ?

1st approach to GWs

(TT gauge: 6 - 4 = 2 d.o.f.)

$$h^{0\mu} = 0$$
, $h^i_i = 0$, $\partial_j h_{ij} = 0$
Outside
Source

$$\partial_{\mu}\partial^{\mu}h_{ij} = 0$$
 Wave Eq. \rightarrow Gravitational Waves !

1st approach to GWs

(11 gauge: 6 - 4 = 2 d.O.T.)

$$h^{0\mu} = 0, \quad h^i_i = 0, \quad \partial_j h_{ij} = 0$$

Outside
Source

$$\partial_{\mu}\partial^{\mu}h_{ij} = 0$$
 Wave Eq. \rightarrow Gravitational Waves !



1st approach to GWs

(TT gauge: 6 - 4 = 2 d.o.f.)

$$h^{0\mu} = 0$$
, $h^i_i = 0$, $\partial_j h_{ij} = 0$
Outside
Source

$$\partial_{\mu}\partial^{\mu}h_{ij} = 0$$
 Wave Eq. \rightarrow Gravitational Waves !



1st approach to GWs

(TT gauge: 6 - 4 = 2 d.o.f.)

$$h^{0\mu} = 0$$
, $h^i_i = 0$, $\partial_j h_{ij} = 0$
Outside
Source

$$\partial_{\mu}\partial^{\mu}h_{ij} = 0$$
 Wave Eq. \rightarrow Gravitational Waves !

can GW be 'gauged away'? No !

2 *dof* = 2 polarizations

$$h_{ab}(t, \mathbf{x}) = \int_{-\infty}^{\infty} df \int d\hat{n} \, h_{ab}(f, \hat{n}) e^{-2\pi i f(t - \hat{n}\mathbf{x})}$$
(plane wave)
transverse plane

1st approach to GWs

(TT gauge: 6 - 4 = 2 d.o.f.)

$$h^{0\mu} = 0$$
, $h^i_i = 0$, $\partial_j h_{ij} = 0$
Outside
Source

$$\partial_{\mu}\partial^{\mu}h_{ij} = 0$$
 Wave Eq. \rightarrow Gravitational Waves !

can GW be 'gauged away'? No !

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2 dof = 2 polarizations
$$h_{ab}(t, \mathbf{x}) = \int_{-\infty}^{\infty} df \int d\hat{n} h_{ab}(f, \hat{n}) e^{-2\pi i f(t-\hat{n}\mathbf{x})}$$
 (plane wave)
transverse plane $h_{ab}(f, \hat{n}) = \sum_{A=+,\mathbf{x}} h_A(f, \hat{n}) \epsilon_{ab}^{(A)}(\hat{n}) = \begin{pmatrix} h_+ & h_x & 0\\ h_x & -h_+ & 0\\ 0 & 0 & 0 \end{pmatrix}$ Transverse-Traceless (2 dof)

1st approach to GWs

(TT gauge: 6 - 4 = 2 d.o.f.)

$$h^{0\mu} = 0$$
, $h^i_i = 0$, $\partial_j h_{ij} = 0$
Outside
Source

$$\partial_{\mu}\partial^{\mu}h_{ij} = 0$$
 Wave Eq. \rightarrow Gravitational Waves !



Definition of GWs 2nd approach

2nd approach to GWs

(gauge invariant def.)

$$\begin{aligned} & \text{Minkowski} \\ & \uparrow \\ g_{\mu\nu} = \eta_{\mu\nu} + \delta g_{\mu\nu} \qquad (|\delta g_{\mu\nu}| \ll 1) \end{aligned}$$

 $g_{\mu\nu} = \eta_{\mu\nu}$ $\delta g_{\mu\nu}$ Mink**b**wski $T_{\mu}g_{\mu\nu} = \eta_{\mu\nu} + \delta g_{\mu\nu} \qquad (|\delta g_{\mu\nu}| \ll 1)$ 2nd approach to GWs (gauge invariant def.) (svt decomposition) $\delta g_{00} = -2\phi,$ s: scalar $\delta g_{0i} = \delta g_{i0} = (\partial_i B + S_i),$ v: vector $\delta g_{ij} = \delta g_{ji} = -2\psi \delta_{ij} + (\partial_i \partial_j - \frac{1}{3}\delta_{ij}\nabla^2)E + \partial_i F_j + \partial_j F_i + h_{ij},$ t: tensor

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 $T_{00}=\rho,$

T

$$T_{0i}=T_{i0}=\partial_i u+u_i,$$

Gravitational Wave Definition $\delta g_{\mu\nu}$ $\delta g_{\mu\nu}$ $\delta g_{\mu\nu}$ Minkbyyski

 $T_{\mu\nu} = \eta_{\mu\nu} + \delta g_{\mu\nu} \qquad (|\delta g_{\mu\nu}| \ll 1)$

2nd approach to GWs

(gauge invariant def.)

$$\begin{split} \delta g_{00} &= -2\phi, \\ \delta g_{00} &= -2\phi, \\ \delta g_{0i} &= \delta g_{i0} \equiv (\partial_i B \pm \hat{s}_i), \\ \delta g_{0i} &= \delta g_{i0} \equiv (\partial_i B \pm \hat{s}_i), \\ \delta \delta g_{ij} &= \delta g_{ji} = -2\psi \delta_{ij} + (\partial_i \partial_j - \frac{1^1}{3^3} \delta_{ij} \nabla^2) E + \partial_i F_j + \partial_j F_i + h_{ij}, \\ \delta \delta g_{ij} &= \delta g_{ji} = -2\psi \delta_{ij} + (\partial_i \partial_j - \frac{1^1}{3^3} \delta_{ij} \nabla^2) E + \partial_i F_j + \partial_j F_i + h_{ij}, \\ T T_{00} &= \rho, \\ T T_{00} &= \partial_i w + u_i, \end{split}$$

$$T_{ij}^{T} = T_{ji} = p \,\delta_{ij} + (\partial_i \partial_j = \frac{\mathfrak{h}}{\mathfrak{H}} \delta_{ij} \nabla^2) \sigma + \partial_i v_j + \partial_j v_i + \Pi_{ij}.$$

$$\begin{split} & \delta g_{00} = -2\phi, \\ & \delta g_{0i} = \delta g_{i0} = (\partial_i B + S_i), \\ & \delta g_{0i} = \delta g_{i0} = (\partial_i B + S_i), \\ & \delta g_{0i} = \delta g_{i0} = (\partial_i B + S_i), \\ & \delta \delta g_{ij} = \delta g_{ji} = -\overline{2}\psi \delta_{ij} + (\partial_i \partial_j - \frac{1}{3}\delta_{ij} \nabla^2)E + \partial_i F_j + \partial_j F_i + h_{ij}, \\ & \delta \delta g_{ij} = \delta g_{ji} = -\overline{2}\psi \delta_{ij} + (\partial_i \partial_j - \frac{1}{3}\delta_{ij} \nabla^2)E + \partial_i F_j + \partial_j F_i + h_{ij}, \\ & \mathcal{T}_{I0i} = \mathcal{T}_{I0} = \partial_{I} u + u \mu_i, \\ & \mathcal{T}_{Ij} = T_{ji} = p \, \delta_{ij} + (\partial_i \partial_j - \frac{1}{3} \frac{1}{3} \overline{\delta}_{ij} \nabla^2)\sigma + \partial_i v_j + \partial_j v_i + \Pi_{ij}. \end{split}$$

 $\delta g_{\mu
u} \over \delta g_{\mu
u}$

 $T_{\mu
u}$ $T_{\mu
u}$

 v_i

 Π_{ij} v_i

 \prod_{ij}

δg_{0i} G δg_{i0} $\overline{Vitatt} \delta i$ hal Wave Definition





δg_{0i} G δg_{i0} $\overline{Vitatton}$ hal Wave Definition





δg_{0i} G δg_{i0} $\overline{Vitatton}$ hal Wave Definition





δg_{0i} G δg_{i0} $\overline{Vitatt} \delta i$ hal Wave Definition





δg_{0i} G δg_{i0} $\overline{Vitatt} \delta i$ hal Wave Definition





δg_{0i} G δg_{i0} $\overline{Vitatton}$ hal Wave Definition





δg_{0i} G δg_{i0} $\nabla i t^{2i} t^{3i} t^{3i}$





$$\begin{split}
\delta g_{00} &= -2\phi, \\
\delta g_{0i} &= \delta g_{i0} = (\partial_i B + S_i), \\
\delta g_{0i} &= \delta g_{i0} = (\partial_i B + S_i), \\
\delta g_{0i} &= \delta g_{i0} = (\partial_i B + S_i), \\
\delta g_{0i} &= \delta g_{i0} = (\partial_i B + S_i), \\
\delta g_{ij} &= \delta g_{ji} = -2\psi \delta_{ij} + (\partial_i \partial_j - \frac{1}{3} \frac{1}{3} \delta_{ij} \nabla^2) E + \partial_i F_j + \partial_j F_i + h_{ij}, \\
\hline
\mathbf{T}_{\mathbf{T}_{0i}} &= \mathbf{T}_{i0} = \partial_i u + u \mu_i, \\
\mathbf{T}_{ij} &= \mathbf{T}_{ji} = p \, \delta_{ij} + (\partial_i \partial_j - \frac{1}{3} \frac{1}{3} \overline{\delta}_{ij} \nabla^2) \sigma + \partial_i v_j + \partial_j v_i + \Pi_{ij}.
\end{split}$$
16 degrees of freedom
16 degrees of freedom

In order NOT to over-count degrees of freedom

$$\delta g_{\mu
u} \over \delta g_{\mu
u}$$

 $T_{\mu
u}$ $T_{\mu\nu}$

 v_i \prod_{ij} \mathcal{V}_i \prod_{ij}

$$\delta g_{0i} = T \delta g_{i0} = (\partial_i B + S_i), \quad -\frac{1}{3} \nabla$$

$$\delta g_{ij} = \delta g_{ji} = -2\psi \delta_{ij} + (\partial_i \partial_j - \frac{1}{3} \delta_{ij} \nabla^2) E + \partial_i F_j + \partial_j F_i + h_{ij}, \quad 16 \text{ degrees}$$
of freedom
$$T_{00} \equiv \rho, \quad (s_{y_i t} E/p \text{-tensor components})_{\mu\nu}$$

$$T_{0ii} = T_{i0} = \partial_{ij} \mu \mu + \mu \mu_i, \quad 16 \text{ degrees}$$

$$T_{ij} = T_{ji} = p \delta_{ij} + (\partial_i \partial_j - \frac{1}{3} \frac{1}{3} \overline{\delta}_{ij} \nabla^2) \sigma + \partial_i v_j + \partial_j v_i + \Pi_{ij}. \quad 16 \text{ degrees}$$

In order NOT to over-count degrees of freedom

$$\partial_{i}S_{i} = 0 \text{ (1 constraint)}, \quad \partial_{i}F_{i} = 0 \text{ (1 constraint)}, \\ \partial_{i}h_{ij} = 0 \text{ (3 constraints)} \stackrel{\delta g}{}_{\beta} \stackrel{\mu\nu}{}_{\mu\nu} h_{ii} = 0 \text{ (1 constraint)} \end{cases} \begin{cases} \text{Metric} \\ T_{\mu\nu} \text{ perturbations} \\ T_{\mu\nu} \text{ perturbations} \\ T_{\mu\nu} \text{ perturbations} \\ T_{\mu\nu} \text{ perturbations} \end{cases}$$
Gravitational Wave Definition $\frac{\delta g_{00}}{\delta g_{00}} = \frac{-2\phi}{\overline{T}_{0i}} = \partial_i u + u_i,$ (svt metric perturbations) $\delta g_{0i} = \delta g_{i0} = (\partial_i B + S_i), \qquad -\frac{1}{3} \nabla$ 16 degrees

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of freedom

$$\begin{split} & \delta_{g_{ij}} = \delta g_{ji} = -2\psi \delta_{ij} + (\partial_i \partial_j - \frac{1}{3} \delta_{ij} \nabla^2) E + \partial_i F_j + \partial_j F_i + h_{ij}, \end{split}, \quad \text{Or freedom} \\ & \mathcal{T}_{T_{00}} \equiv \beta_2 \\ & \mathcal{T}_{00i} = \mathcal{T}_{i0} = \partial_{ti} u + u \mu_i, \\ & \mathcal{T}_{ij} = T_{ji} = p \, \delta_{ij} + (\partial_i \partial_j - \frac{1}{3} \frac{1}{3} \overline{\delta}_{ij} \nabla^2) \sigma + \partial_i v_j + \partial_j v_i + \Pi_{ij}. \end{split}$$

In order NOT to over-count degrees of freedom

$$\begin{aligned} \partial_i S_i &= 0 \ (1 \ \text{constraint}), \quad \partial_i F_i = 0 \ (1 \ \text{constraint}), \\ \partial_i h_{ij} &= 0 \ (3 \ \text{constraints}) \overset{\delta g}{\overset{\mu \nu}{\overset{\mu \nu}{\delta g}_{\mu \nu}} h_{ii} = 0 \ (1 \ \text{constraint}) \end{aligned} \right\} \begin{array}{l} \text{Metric} \\ T_{\mu \nu} \ \text{perturbations} \\ T_{\mu \nu} \end{array} \\ \partial_i u_i &= 0 \ (1 \ \text{constraint}), \quad \partial_i v_i = 0 \ (1 \ \text{constraint}), \end{aligned} \right\} \begin{array}{l} \prod_{ij} \overset{v_i}{\overset{\tau}{\underset{ij$$

Gravitational Wave Definition $\frac{\delta g_{00}}{\delta g_{00}} = \frac{-2\phi}{-2\phi}, \quad \exists u + u_i, \quad (\text{svt metric perturbations})$ $\delta g_{0i} = \delta g_{i0} = (\partial_i B + S_i), \qquad -\frac{1}{3} \nabla$ $\delta g_{0i} = T \delta g_{i0} = (\partial_i B + S_i), \qquad -\frac{1}{3} \nabla$ $\delta g_{0i} = -\frac{1}{2} \sqrt{2} \delta_{ii} + (\partial_i \partial_i - \frac{1}{2} \delta_{ii} \nabla^2) E + \partial_i F_i + \partial_i F_i + h$ 16 degrees of freedom

$$\delta g_{ij} = \delta g_{ji} = -2\psi \delta_{ij} + (\partial_i \partial_j - \frac{1}{3} \delta_{ij} \nabla^2)E + \partial_i F_j + \partial_j F_i + h_{ij},$$

$$T_{I_{00i}} \equiv \beta,$$

$$T_{I_{00i}} \equiv T_{i0} = \partial_{i} \mu + \mu \mu_i,$$

$$T_{ij} = T_{ji} = p \, \delta_{ij} + (\partial_i \partial_j - \frac{1}{3} \frac{1}{3} \overline{\delta}_{ij} \nabla^2)\sigma + \partial_i v_j + \partial_j v_i + \Pi_{ij}.$$

$$I_{ij} = T_{i0} = \beta \delta_{ij} + (\partial_i \partial_j - \frac{1}{3} \frac{1}{3} \overline{\delta}_{ij} \nabla^2)\sigma + \partial_i v_j + \partial_j v_i + \Pi_{ij}.$$

In order NOT to over-count degrees of freedom

$$\partial_{i}S_{i} = 0 \text{ (1 constraint)}, \quad \partial_{i}F_{i} = 0 \text{ (1 constraint)}, \\ \partial_{i}h_{ij} = 0 \text{ (3 constraints)}_{\delta g_{\mu\nu}}^{\mu\nu}h_{ii} = 0 \text{ (1 constraint)} \end{cases} \begin{cases} 6 \text{ constraints for} \\ T_{\mu\nu} \text{ perturbations} \\ T_{\mu\nu} \text{ perturbations} \end{cases}$$

 $\partial_i u_i = 0$ (1 constraint), $\partial_i v_i = 0$ (1 constraint), $\partial_i \Pi_{ii} = 0$ (3 constraints), $\Pi_{ii} = 0$ (1 constraint), 6 constraints for E/p tensor components

constraints for



In order NOT to over-count degrees of freedom

$$\partial_i S_i = 0 \text{ (1 constraint)}, \quad \partial_i F_i = 0 \text{ (1 constraint)}, \\ \partial_i h_{ij} = 0 \text{ (3 constraints)} \\ \delta_g \mu\nu h_{ii} = 0 \text{ (1 constraint)} \\ \partial_i u_i = 0 \text{ (1 constraint)}, \quad \partial_i v_i = 0 \text{ (1 constraint)}, \\ \partial_i v_i = 0 \text{ (1 constraint)}, \quad \partial_i v_i = 0 \text{ (1 constraint)}, \\ \partial_i v_i = 0 \text{ (1 constraint)}, \quad \partial_i v_i = 0 \text{ (1 constraint)}, \\ \partial_i v_i = 0 \text{ (1 constraint)}, \quad \partial_i v_i = 0 \text{ (1 constraint)}, \\ \partial_i v_i = 0 \text{ (1 constraint)}, \quad \partial_i v_i = 0 \text{ (1 constraint)}, \\ \partial_i v_i = 0 \text{ (1 constraint)}, \quad \partial_i v_i = 0 \text{ (1 constraint)}, \\ \partial_i v_i = 0 \text{ (1 constraint)}, \quad \partial_i v_i = 0 \text{ (1 constraint)}, \\ \partial_i v_i = 0 \text{ (1 constraint)}, \quad \partial_i v_i = 0 \text{ (1 constraint)}, \\ \partial_i v_i = 0 \text{ (1 constraint)}, \quad \partial_i v_i = 0 \text{ (1 constraint)}, \\ \partial_i v_i = 0 \text{ (1 constraint)}, \quad \partial_i v_i = 0 \text{ (1 constraint)}, \\ \partial_i v_i = 0 \text{ (1 constraint)}, \quad \partial_i v_i = 0 \text{ (1 constraint)}, \\ \partial_i v_i = 0 \text{ (1 constraint)}, \quad \partial_i v_i = 0 \text{ (1 constraint)}, \\ \partial_i v_i = 0 \text{ (1 constraint)}, \quad \partial_i v_i = 0 \text{ (1 constraint)}, \\ \partial_i v_i = 0 \text{ (1 constraint)}, \quad \partial_i v_i = 0 \text{ (1 constraint)}, \\ \partial_i v_i = 0 \text{ (1 constraint)}, \quad \partial_i v_i = 0 \text{ (1 constraint)}, \\ \partial_i v_i = 0 \text{ (1 constraint)}, \quad \partial_i v_i = 0 \text{ (1 constraint)}, \\ \partial_i v_i = 0 \text{ (1 constraint)}, \quad \partial_i v_i = 0 \text{ (1 constraint)}, \\ \partial_i v_i = 0 \text{ (1 constraint)}, \quad \partial_i v_i = 0 \text{ (1 constraint)}, \\ \partial_i v_i = 0 \text{ (1 constraint)}, \quad \partial_i v_i = 0 \text{ (1 constraint)}, \\ \partial_i v_i = 0 \text{ (1 constraint)}, \quad \partial_i v_i = 0 \text{ (1 constraint)}, \\ \partial_i v_i = 0 \text{ (1 constraint)}, \quad \partial_i v_i = 0 \text{ (1 constraint)}$$

 $\partial_i \Pi_{ij} = 0$ (3 constraints), $\Pi_{ii} = 0$ (1 constraint),

Constraints for E/p

$\begin{array}{c} & T_{\mu\nu} \\ \hline \mathbf{Gravitational Wave Definition} \\ \hline \mathbf{Gravitational Wave Definition } \\ \hline \mathbf{Gravitational Wave Definitional Wave Definition \\ \hline \mathbf{Gravitational Wave Definitional Wave Definitional Wave Definition \\ \hline \mathbf{Gravitational Wave Definitional Wave Definitional Wave Definitional Wave Definitional Wave Definition \\ \hline \mathbf{Gravitational Wave Definitional Wave Definitional Wave Definitional Wave Definitional Wave Definition \\ \hline \mathbf{Gravitational Wave Definitional Wave Definitional Wave Definitional Wave Definitional Wave Definition \\ \hline \mathbf{Gravitational Wave Definitional Wave Defin$

$$T_{T_{0}} \equiv \beta;$$

$$T_{0} \equiv T_{i0} = \partial_{i}u + u\mu_{i},$$

$$T_{0} = T_{i0} = \partial_{i}u + u\mu_{i},$$

$$T_{ij} = T_{ji} = p \,\delta_{ij} + (\partial_{i}\partial_{j} - \frac{1}{3} \frac{1}{3} \overline{\delta}_{ij} \nabla^{2})\sigma + \partial_{i}v_{j} + \partial_{j}v_{i} + \Pi_{ij}.$$

$$T_{ij} = T_{ji} = p \,\delta_{ij} + (\partial_{i}\partial_{j} - \frac{1}{3} \frac{1}{3} \overline{\delta}_{ij} \nabla^{2})\sigma + \partial_{i}v_{j} + \partial_{j}v_{i} + \Pi_{ij}.$$

In order NOT to over-count degrees of freedom

$$\partial_i S_i = 0 \ (1 \ \text{constraint}), \quad \partial_i F_i = 0 \ (1 \ \text{constraint}), \\ \partial_i h_{ij} = 0 \ (3 \ \text{constraints}) \int_{\delta g_{\mu\nu}}^{\delta g_{\mu\nu}} h_{ii} = 0 \ (1 \ \text{constraint})$$

 $\partial_i u_i = 0$ (1 constraint), $\partial_i v_i = 0$ (1 constraint), $\partial_i \Pi_{ij} = 0$ (3 constraints), $\Pi_{ii} = 0$ (1 constraint),

tensor components



 $\delta g_{\mu
u} \delta g_{\mu
u}$

 $T_{\mu\nu}$ $T_{\mu\nu}$

 \mathcal{U}_i

 Π_{ii}

 v_i

 Π_{ij}



Physical Constraints

$$\partial^{\mu}T_{\mu\nu}=0$$

$$\delta g_{\mu
u} \over \delta g_{\mu
u}$$

 $T_{\mu\nu}$ $T_{\mu\nu}$

 \mathcal{U}_i

 \prod_{ij}

 v_i

 \prod_{ij}

Gravitational Wave Definition $^{T_{\mu\nu}}$

$$\begin{array}{cccc} \rho, u_{i}^{b} & g_{i} p_{i}^{c} \sigma, \overline{q}_{i}^{b}, \Pi_{ij} \longrightarrow 0 & S_{i}, F_{i} & (\text{svt metric perturbations}) & 3 \times 3 \\ \delta g_{0i} &= \delta g_{i0} &= (\partial_{i}B + S_{i}), \\ 3 \times 3 & \delta g_{0i} &= \delta g_{i0} &= (\partial_{i}B + S_{i}), \\ 3 \times 3 & \delta g_{0i} &= \delta g_{i0} &= (\partial_{i}B + S_{i}), \\ 3 \times 3 & & 10 \text{ degrees} \\ \delta g_{ij} &= \delta g_{ii} &= (\partial_{i}B + S_{i}), \\ \delta g_{ij} &= \delta g_{ij} &= (\partial_{i}B + S_{i}), \\ \delta g_{ij} &= \delta g_{ij} &= (\partial_{i}B + S_{i}), \\ \delta g_{ij} &= \delta g_{ij} &= (\partial_{i}B + S_{i}), \\ \delta g_{ij} &= \delta g_{ij} &= (\partial_{i}B + S_{i}), \\ \delta g_{ij} &= \delta g_{ij} &= (\partial_{i}B + S_{i}), \\ \delta g_{ij} &= \delta g_{ij} &= (\partial_{i}B + S_{i}), \\ \delta g_{ij} &= \delta g_{ij} &= (\partial_{i}B + S_{i}), \\ \delta g_{ij} &= \delta g_{ij} &= (\partial_{i}B + S_{i}), \\ \delta g_{ij} &= \delta g_{ij} &= (\partial_{i}B + S_{i}), \\ \delta g_{ij} &= \delta g_{ij} &= (\partial_{i}B + S_{i}), \\ \delta g_{ij} &= \delta g_{ij} &= (\partial_{i}B + S_{i}), \\ \delta g_{ij} &= \delta g_{ij} &= (\partial_{i}B + S_{i}), \\ \delta g_{ij} &= \delta g_{ij} &= (\partial_{i}B + S_{i}), \\ T_{0} &= \delta g_{ij} &= (\partial_{i}B + S_{i}), \\ \delta g_{ij} &= \delta g_{ij} &= (\partial_{i}B + S_{i}), \\ \delta g_{ij} &= \delta g_{ij} &= (\partial_{i}B + S_{i}), \\ \delta g_{ij} &= \delta g_{ij} &= (\partial_{i}B + S_{i}), \\ \delta g_{ij} &= \delta g_{ij} &= (\partial_{i}B + S_{i}), \\ \delta g_{ij} &= \delta g_{ij} &= (\partial_{i}B + S_{i}), \\ \delta g_{ij} &= (S_{ij} + (\partial_{i}\partial_{j} - \frac{1}{3} \frac{1}{3} \overline{S}_{ij} \overline{S}_{ij} \nabla^{2}) + (\partial_{i}\partial_{j} - \partial_{i} + \partial_{i}\nabla_{j} + \partial_{j}\nabla_{i} + \Pi_{ij}), \\ \delta g_{ij} &= \delta g_{ij} &= (\partial_{i}\partial_{j} - \frac{1}{3} \frac{1}{3} \overline{S}_{ij} \overline{S}_{ij} \nabla^{2}) + (\partial_{i}\partial_{j} - \partial_{i} + \partial_{i}\nabla_{j} + \partial_{i}\nabla_{j} + \partial_{j}\nabla_{i} + \Pi_{ij}), \\ \partial^{\mu} T_{\mu\nu} &= 0 \\ \\ \delta g_{\mu\nu} &= \delta g_{\mu\nu} &=$$

Phys Constraints

Gravitational Wave Definition $^{T_{\mu\nu}}$

Physical Constraints $\partial^{\mu}T_{\mu\nu} = 0 \implies \begin{cases} \delta g_{\mu\nu} \\ \nabla^{2}\sigma = \frac{3}{2}(\dot{u} - p) \text{ (1 constraint),} \\ \Pi_{ij} \end{cases} \begin{bmatrix} 4 \text{ constraints} \\ (\text{due to E/p} \\ \text{conservation}) \\ \Pi_{ij} \end{bmatrix}$



Physical Constraints

$$\partial^{\mu}T_{\mu
u}=0$$

$$\delta g_{\mu
u} \over \delta g_{\mu
u}$$

 $T_{\mu\nu}$ $T_{\mu\nu}$

 \mathcal{U}_i

 \prod_{ij}

 v_i

 \prod_{ij}

$$\begin{split} & \delta g_{00} = -2\phi, \\ & \delta g_{0i} = \delta g_{i0} = (\partial_i B + S_i), \\ & \delta g_{0i} = \delta g_{i0} = (\partial_i B + S_i), \\ & \delta g_{0i} = \delta g_{i0} = (\partial_i B + S_i), \\ & \delta g_{0i} = \delta g_{ii} = -2\psi \delta_{ij} + (\partial_i \partial_j - \frac{1}{3} \delta_{ij} \nabla^2) E + \partial_i F_j + \partial_j F_i + h_{ij}, \\ & \delta g_{ij} = \delta g_{ii} = -2\psi \delta_{ij} + (\partial_i \partial_j - \frac{1}{3} \delta_{ij} \nabla^2) E + \partial_i F_j + \partial_j F_i + h_{ij}, \\ & T_{0ii} = T_{a0} = \partial_i u + u \mu_i, \\ & T_{0ii} = T_{a0} = \partial_i u + u \mu_i, \\ & T_{ij} = T_{ji} = p \, \delta_{ij} + (\partial_i \partial_j - \frac{1}{3} \frac{1}{3} \overline{\delta}_{ij} \nabla^2) \sigma + \partial_i v_j + \partial_j v_i + \Pi_{ij}. \\ & \rho, u, u_i, p, \sigma, v_i, \Pi_{ij} \\ & \rho, u_i, p, \Pi_{ij} \\ & O^{\mu} G_{\mu\nu} = 0 \quad \Longrightarrow \quad \begin{bmatrix} \delta g_{\mu\nu} & T_{\mu\nu} \\ \delta g_{\mu\nu} & T_{\mu\nu} \\ & \Pi_{ij} & U_i \\ & \Pi_{ij} &$$

nts.

 v_i

,

,

nts.

, **Gravitational Wave Definition**



$$\begin{array}{c} \underset{2}{\overset{3}{2}} - (\overbrace{\mathbf{Gravitational Wave Definition}^{T_{\mu\nu}}} \\ \overbrace{\mathbf{Gravitational Wave Definition}^{T_{\mu\nu}}} \\ \underset{2}{\overset{\delta g_{n} \lor i}{\mathcal{G}_{n}, u, u_{i}, p, \sigma, v_{i}, \Pi_{ij}}} \\ \underset{2}{\overset{\delta g_{n} \lor i}{\mathcal{G}_{n}, u_{i}, p, \sigma, v_{i}, \Pi_{ij}}} \\ \underset{2}{\overset{\delta g_{n} \lor i}{\mathcal{G}_{n}, u_{i}, p, \sigma, v_{i}, \Pi_{ij}}} \\ \underset{2}{\overset{\delta g_{n} \lor i}{\mathcal{G}_{n}, u_{i}, p, \sigma, v_{i}, \Pi_{ij}}} \\ \underset{2}{\overset{\delta g_{n} \lor i}{\mathcal{G}_{n}, u_{i}, p, \sigma, v_{i}, \Pi_{ij}}} \\ \underset{2}{\overset{\delta g_{n} \lor i}{\mathcal{G}_{n}, u_{i}, p, \sigma, v_{i}, \Pi_{ij}}} \\ \underset{2}{\overset{\delta g_{n} \lor i}{\mathcal{G}_{n}, u_{i}, p, \sigma, v_{i}, \Pi_{ij}}} \\ \underset{2}{\overset{\delta g_{n} \lor i}{\mathcal{G}_{n}, u_{i}, p, \sigma, v_{i}, \Pi_{ij}}} \\ \underset{2}{\overset{\delta g_{n} \lor i}{\mathcal{G}_{n}, u_{i}, p, \sigma, v_{i}, \Pi_{ij}}} \\ \underset{2}{\overset{\delta g_{n} \lor i}{\mathcal{G}_{n}, u_{i}, p, \sigma, v_{i}, \Pi_{ij}}} \\ \underset{2}{\overset{\delta g_{n} \lor i}{\mathcal{G}_{n}, v_{i}, v_{i}}} \\ \underset{2}{\overset{\delta g_{n} \lor i}{\mathcal{G}_{n}, v_{i}, v_{i}}} \\ \underset{2}{\overset{\delta g_{n} \lor i}{\mathcal{G}_{n}, v_{i}, v_{i}}} \\ \underset{2}{\overset{\delta g_{n} \lor i}{\mathcal{G}_{n}, v_{i}, v_{i}, v_{i}}} \\ \underset{2}{\overset{\delta g_{n} \lor i}{\mathcal{G}_{n}, v_{i}, v_{i}, v_{i}}} \\ \underset{2}{\overset{\delta g_{n} \lor i}{\mathcal{G}_{n}, v_{i}, v_{i}}} \\ \underset{2}{\overset{\delta g_{n} \lor i}{\mathcal{G}_{n}, v_{i}, v_{i}}} \\ \underset{2}{\overset{\delta g_{n} \lor i}{\mathcal{G}_{n}, v_{i}} \\ \underset{2}{\overset{\delta g_{n} \lor i}{\mathcal{G}_{n}, v_{i}}} \\ \underset{2}{\overset{\delta g_{n} \lor i}{\mathcal{G}_{n}, v_{i}} \\ \underset{2}{\overset{\delta g_{n} \lor i}{\mathcal{G}_{n}, v_{i}}} \\ \underset{2}{\overset{\delta g_{n} \lor i}{\mathcal{G}_{n}, v_{i}} \\ \underset{2}{\overset{\delta g_{n} \lor i}{\mathcal{G}_{n}, v_{i}} \\ \underset{2}{\overset{\delta g_{n} \lor$$

$$\begin{array}{c} \underset{2}{\overset{3}{2}} - (\overbrace{\mathbf{Gravitational Wave Definition}}^{T_{\mu\nu}} \\ \overbrace{\mathbf{Gravitational Wave Definition}}^{T_{\mu\nu}} \\ \underset{2}{\overset{\delta g_{n} \vee g_{1} \vee g_{2} \vee g_{1} \vee g_{2} \vee$$

Gravitation Gravitation $T_{\mu\nu} \rho, u_i, p, \Pi_{ij}$

$$\begin{array}{c} \delta g_{00} = -2\phi, \\ \delta g_{0i} = \delta g_{i0} = (\partial_i B + S_i), \\ \delta g_{0i} = \delta g_{i0} = (\partial_i B + S_i), \\ \delta g_{0i} = \delta g_{i0} = (\partial_i B + S_i), \\ \delta g_{0i} = \delta g_{i0} = (\partial_i B + S_i), \\ \delta g_{ii} = \delta g_{ii} = -2\psi \delta_{ij} + (\partial_i \partial_j - \frac{1}{3} \delta_{ij} \nabla^2) E + \partial_i F_j + \partial_j F_i + h_{ij}, \\ \end{array}$$

$$\begin{array}{c} 6 \text{ degrees} \\ \text{of freedom} \\ \hline f_{00} \equiv \rho, \\ \hline f_{0i} = T_{ji} = p \, \delta_{ij} + (\partial_i \partial_j - \frac{1}{3} \frac{1}{3} \overline{\delta}_{ij} \nabla^2) \sigma + \partial_i v_j + \delta_{ij} \partial_{ij} v_i + \Pi_{gjuv} \\ - \partial_{\mu} \xi_{\nu} - \partial_{\nu} \xi_{\mu} \\ \hline f_{0i} = T_{ji} = p \, \delta_{ij} + (\partial_i \partial_j - \frac{1}{3} \frac{1}{3} \overline{\delta}_{ij} \nabla^2) \sigma + \partial_i v_j + \delta_{ij} \partial_{ij} v_i + \Pi_{gjuv} \\ - \partial_{\mu} \xi_{\nu} - \partial_{\nu} \xi_{\mu} \\ \hline f_{0i} = 0, \\$$

$$\begin{array}{c} \delta g_{00} = -2\phi, & (svt metric perturbations) & \rightarrow - , \\ \delta g_{0i} = \delta g_{i0} = (\partial_i B + S_i), & (svt metric perturbations) & \rightarrow - , \\ \delta g_{0i} = \delta g_{i0} = (\partial_i B + S_i), & \psi \rightarrow - \frac{1}{3} \nabla & \rightarrow - 6 \text{ degrees} \\ \delta g_{ij} = \delta g_{ij} = -2\psi \delta_{ij} + (\partial_i \partial_j - \frac{1}{3} \delta_j \nabla^2) E + \partial_i F_j + \partial_j F_j + h_{ij}, \\ \delta g_{ij} = \delta g_{ji} = -2\psi \delta_{ij} + (\partial_i \partial_j - \frac{1}{3} \delta_j \nabla^2) E + \partial_i F_j + \partial_j F_i + h_{ij}, \\ \delta g_{ij} = \delta g_{ij} = -2\psi \delta_{ij} + (\partial_i \partial_j - \frac{1}{3} \delta_j \nabla^2) E + \partial_i F_j + \partial_j F_i + h_{ij}, \\ \delta g_{ij} = \delta_{ij} = 0, \\ \delta g_{ij} = \delta_{ij} = 0, \\ \delta g_{ij} = \delta_{ij} = 0, \\ \delta g_{ij} = \delta_{ij} + (\partial_i \partial_j - \frac{1}{3} |\frac{1}{3} \nabla_{ij} \nabla^2 \otimes \sigma + \partial_i \nabla_i \Phi_j \otimes \rho + \frac{1}{3} \nabla_i \nabla_i \Phi_j \otimes \rho + \frac{1}{3} \nabla_i \nabla_i \Phi_j \otimes \rho + \frac{1}{3} \nabla_$$

(4

$$\begin{array}{c} \delta g_{00} = -2\phi, & (\text{svt metric perturbations}) & \rightarrow - , \\ \delta g_{0i} = \delta g_{i0} = (\partial_i B + S_i), & \psi \rightarrow - \frac{1}{3} \nabla & \rightarrow - 6 \text{ degrees} \\ \delta g_{0i} = \delta g_{i0} = (\partial_i B + S_i), & \psi \rightarrow - \frac{1}{3} \nabla & \rightarrow - 6 \text{ degrees} \\ \delta g_{0i} = \delta g_{i0} = (\partial_i B + S_i), & \psi \rightarrow - \frac{1}{3} \nabla & \rightarrow - 6 \text{ degrees} \\ \delta g_{0i} = \delta g_{i0} = (\partial_i B + S_i), & \psi \rightarrow - \frac{1}{3} \nabla & \rightarrow - 6 \text{ degrees} \\ \delta g_{0i} = \delta g_{0i} = -2\psi \delta_{ij} + (\partial_i \partial_i - \frac{1}{3} \delta_j \nabla^2) E + \partial_i F_j + \partial_j F_i + h_{ij}, & F_i - 2d_i, \\ S_i \rightarrow S_i \rightarrow d_i, & F_i \rightarrow T_i - 2d_i, & F_i \rightarrow T_i - 2d_i, \\ \hline T_{00} \equiv \beta; & (\text{svt E/p-tensor Components}) \\ d_{i0}, d_i & T_{ij} = T_{ji} = p \delta_{ij} + (\partial_i \partial_j^T - \frac{1}{3} |\frac{1}{3} \nabla_{ij} \nabla^2 \otimes \sigma + \partial_i \nabla_i \Phi_{ij} - (f, \mathbf{x}) \\ T_{ij} = T_{ji} = p \delta_{ij} + (\partial_i \partial_j^T - \frac{1}{3} |\frac{1}{3} \nabla_{ij} \nabla^2 \otimes \sigma + \partial_i \nabla_i \Phi_{ij} - \partial$$

(4

44 (= 6 = 2) Gauge Invariant ! $\Phi \equiv -\phi + \dot{B} = \frac{1}{2}\ddot{E}, \quad (1)$ $\Theta \equiv -2\psi - \frac{1}{3}\nabla^{2}E, \quad (1)$ $\Sigma_{i} \equiv S_{i} - \frac{1}{2}\dot{F}_{i}, \quad (\partial_{i}\Sigma_{i} = 0) \quad (2)$ $h_{ij} \equiv h_{ij}, \quad (h_{ii} = \partial_{i}h_{ij} = 0) \quad (2)$

6 gauge invariant degrees of freedom

 $\Phi; \Theta; \Sigma_i^i$

 \sum

Mague I Gravitational Wavē¹^bDefinition









 Θ, Φ, Σ_i



 Θ, Φ, Σ_i



$$h_{ij}, \ (h_{ii} = \partial_i h_{ij} = 0)$$

transverse & traceless (tensor dof)

Only radiative (~ propagating wave Eq.) gauge invariant degrees of freedom !

$$G \stackrel{1}{r} G \text{ gauge invariant d.o.f.} \qquad \begin{array}{l} \text{Gauge Invariant}\\ (\text{perturbed})\\ \text{Einstein Eqs.} \end{array}$$

$$\nabla^2 \Theta = -\frac{1}{m_p^2} \rho, (1) \qquad \nabla^2 \Phi = \frac{1}{2m_p^2} \left(\rho + 3p - 3\dot{u}\right) (1)$$

$$\nabla^2 \Sigma_i = -\frac{2}{m_p^2} S_i, (2) \qquad \Box h_{ij} = -\frac{2}{m_p^2} \Pi_{ij}. \qquad (2)$$

$$h_{ij}, \ (h_{ii} = \partial_i h_{ij} = 0)$$

transverse & traceless (tensor dof)

Only radiative (~ propagating wave Eq.) gauge invariant degrees of freedom !

 \sum_{i}

Gravitational Waves (GWs) are TT *d.o.f.* metric perturbational, independently of system of reference

Cosmological Backgrounds of Gravitational Waves

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Abstract. Gravitational waves (GWs) have a great potential to probe cosmology. We review early universe sources that can lead to cosmological backgrounds of GWs. We begin by presenting proper definitions of GWs in flat space-time and in a

•
$$g_{\mu\nu}(x) = \eta_{\mu\nu} + h_{\mu\nu}(x), \qquad |h_{\mu\nu}(x)| \ll 1.$$

*Notice that under a Lorentz transformation $x'_{\mu} = \Lambda_{\mu}{}^{\nu}x_{\nu}$, $g'_{\mu\nu}(x') = \Lambda_{\mu}{}^{\alpha}\Lambda_{\nu}{}^{\beta}g_{\alpha\beta}(x)$, preservation of Eq. (3) requires $|\Lambda_{\mu}{}^{\alpha}\Lambda_{\nu}{}^{\beta}h_{\alpha\beta}(x)| \ll 1$, so that it remains true that $|h'_{\mu\nu}(x')| \ll 1$. Rotations do not spoil the condition $|h_{\mu\nu}(x)| \ll 1$, but boosts could, and therefore must be restricted to those that do not spoil such condition. As $h_{\mu\nu}(x)$ is invariant under constant displacements $x'^{\mu} \longrightarrow x^{\mu} + a^{\mu}$, linearised gravity Eq. (3) is also invariant under Poincaré transformations.

Definition of GWs 3rd approach

3rd approach to GWs

(for a FLRW space-time)

$$g_{\mu\nu}(x) = \underline{\bar{g}}_{\mu\nu}(x) + \delta g_{\mu\nu}(x), \quad |\delta g_{\mu\nu}| \ll 1$$
(FLRW)

3rd approach to GWs $g_{\mu\nu}(x) = \overline{g}_{\mu\nu}(x) + \delta g_{\mu\nu}(x), \quad |\delta g_{\mu\nu}| \ll 1$ (for a FLRW space-time) (FLRW)

Flat-FLRW: $ds^2 = \bar{g}_{\mu\nu} dx^{\mu} dx^{\nu}$ $(t \rightarrow \text{Conformal time})$ = $a^2(t)(-dt^2 + d\mathbf{x} \cdot d\mathbf{x})$ = $a^2(t)\eta_{\mu\nu} dx^{\mu} dx^{\nu}$

3rd approach to GWs $g_{\mu\nu}(x) = \overline{g}_{\mu\nu}(x) + \delta g_{\mu\nu}(x), \quad |\delta g_{\mu\nu}| \ll 1$ (for a FLRW space-time) (FLRW)

Flat-FLRW: $ds^2 = a^2(t)\eta_{\mu\nu}dx^{\mu}dx^{\nu}$ ($t \rightarrow$ Conformal time)

3rd approach to GWs $g_{\mu\nu}(x) = \overline{g}_{\mu\nu}(x) + \delta g_{\mu\nu}(x), |\delta g_{\mu\nu}| \ll 1$ (for a FLRW space-time) (FLRW)

Flat-FLRW: $ds^2 = a^2(t)\eta_{\mu\nu}dx^{\mu}dx^{\nu}$ ($t \rightarrow$ Conformal time)

Flat-FLRW + GWs :
$$ds^2 = a^2(t)[\eta_{\mu\nu} + h_{\mu\nu}]dx^{\mu}dx^{\nu}$$

where $h_{0\mu} = 0$, $h_{ii} = 0$, $\partial_i h_{ij} = 0$ Traceless (TT)
d.o.f.

3rd approach to GWs $g_{\mu\nu}(x) = \overline{g}_{\mu\nu}(x) + \delta g_{\mu\nu}(x), |\delta g_{\mu\nu}| \ll 1$ (for a FLRW space-time) (FLRW)

Flat-FLRW: $ds^2 = a^2(t)\eta_{\mu\nu}dx^{\mu}dx^{\nu}$ ($t \rightarrow$ Conformal time)

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$$ds^2 = a^2(t)[\eta_{\mu\nu} + h_{\mu\nu}]dx^{\mu}dx^{\nu}$$

where $h_{0\mu} = 0$, $h_{ii} = 0$, $\partial_i h_{ij} = 0$ Traceless (TT)
d.o.f.

Conformal Transf.: $ds^2 = \tilde{g}_{\mu\nu}(x)dx^{\mu}dx^{\nu} = \Omega^2(x)g_{\mu\nu}(x)dx^{\mu}dx^{\nu}$

3rd approach to GWs $g_{\mu\nu}(x) = \overline{g}_{\mu\nu}(x) + \delta g_{\mu\nu}(x), |\delta g_{\mu\nu}| \ll 1$ (for a FLRW space-time) (FLRW)

Flat-FLRW: $ds^2 = a^2(t)\eta_{\mu\nu}dx^{\mu}dx^{\nu}$ ($t \rightarrow$ Conformal time)

Flat-FLRW + GWs :
$$ds^2 = a^2(t)[\eta_{\mu\nu} + h_{\mu\nu}]dx^{\mu}dx^{\nu}$$

where $h_{0\mu} = 0$, $h_{ii} = 0$, $\partial_i h_{ij} = 0$ Transverse-
Traceless (TT)
d.o.f.

Conformal Transf.:
$$ds^2 = \tilde{g}_{\mu\nu}(x)dx^{\mu}dx^{\nu} = \underbrace{\Omega^2(x)g_{\mu\nu}(x)dx^{\mu}dx^{\nu}}_{a^2(t)}\underbrace{\eta_{\mu\nu} + h_{\mu\nu}}_{a^2(t)}$$

3rd approach to GWs $g_{\mu\nu}(x) = \overline{g}_{\mu\nu}(x) + \delta g_{\mu\nu}(x), |\delta g_{\mu\nu}| \ll 1$ (for a FLRW space-time) (FLRW)

Flat-FLRW: $ds^2 = a^2(t)\eta_{\mu\nu}dx^{\mu}dx^{\nu}$ ($t \rightarrow$ Conformal time)

Flat-FLRW + GWs :
$$ds^2 = a^2(t)[\eta_{\mu\nu} + h_{\mu\nu}]dx^{\mu}dx^{\nu}$$

where $h_{0\mu} = 0$, $h_{ii} = 0$, $\partial_i h_{ij} = 0$ Transverse-
Traceless (TT)
d.o.f.
Conformal Transf.: $ds^2 = \tilde{g}_{\mu\nu}(x)dx^{\mu}dx^{\nu} = \Omega^2(x)g_{\mu\nu}(x)dx^{\mu}dx^{\nu}$

$$ds^{2} = \tilde{g}_{\mu\nu}(x)dx^{\mu}dx^{\nu} = \underbrace{\Omega^{2}(x)g_{\mu\nu}(x)dx^{\mu}dx^{\nu}}_{\substack{a^{2}(t) [\eta_{\mu\nu} + h_{\mu\nu}] \\ = \tilde{g}_{\mu\nu}(x)}}$$

Einstein Eqs:
$$G_{\mu\nu} \equiv R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} = m_p^{-2}T_{\mu\nu}$$

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Question: $\tilde{R}_{\mu\nu}[\tilde{g}_{**} \equiv \Omega^2(x)g_{**}]$?

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Note:
$$\tilde{R}_{\mu\nu}[\tilde{g}_{**}] \equiv (\partial_{\lambda}\tilde{\Gamma}^{\lambda}_{\mu\nu} - \partial_{\nu}\tilde{\Gamma}^{\lambda}_{\mu\lambda}) + (\tilde{\Gamma}^{\alpha}_{\alpha\lambda}\tilde{\Gamma}^{\lambda}_{\mu\nu} - \tilde{\Gamma}^{\alpha}_{\nu\lambda}\tilde{\Gamma}^{\lambda}_{\mu\alpha})$$

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Notation:
$$\begin{cases} \partial_{(\mu}\xi_{\nu)} \equiv \partial_{\mu}\xi_{\nu} + \partial_{\nu}\xi_{\mu} \\ \partial_{[\mu}\xi_{\nu]} \equiv \partial_{\mu}\xi_{\nu} - \partial_{\nu}\xi_{\mu} \end{cases}$$

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 $\tilde{\Gamma}^{\lambda}_{\mu\nu}[\tilde{g}_{**} \equiv \Omega^2 g_{**}] = \Gamma^{\lambda}_{\mu\nu}[g_{**}] + \delta\Gamma^{\lambda}_{\mu\nu}[g_{**}, \omega]; \quad \omega \equiv \log(\Omega)$ \downarrow^{\downarrow} where $\delta\Gamma^{\lambda}_{\mu\nu} = g^{\lambda\sigma}(\partial_{(\mu}\omega \cdot g_{\sigma\nu)} - g_{\mu\nu}\partial_{\sigma}\omega)$

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 $\tilde{R}_{\mu\nu} \equiv \partial_{[\lambda}(\Gamma^{\lambda}_{\mu\nu]} + \delta\Gamma^{\lambda}_{\mu\nu]}) + (\Gamma^{\alpha}_{[\alpha\lambda} + \delta\Gamma^{\alpha}_{[\alpha\lambda]})(\Gamma^{\lambda}_{\mu\nu]} + \delta\Gamma^{\lambda}_{\mu\nu]})$

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$$\begin{split} \tilde{R}_{\mu\nu} &\equiv \partial_{[\lambda} (\Gamma^{\lambda}_{\mu\nu]} + \delta \Gamma^{\lambda}_{\mu\nu]}) + (\Gamma^{\alpha}_{[\alpha\lambda} + \delta \Gamma^{\alpha}_{[\alpha\lambda]}) (\Gamma^{\lambda}_{\mu\nu]} + \delta \Gamma^{\lambda}_{\mu\nu]}) \\ &= R_{\mu\nu} [g_{**}] + \delta R_{\mu\nu} \end{split}$$

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How does it look $\delta R_{\mu\nu}$?

Then:
$$\tilde{R}_{\mu\nu}[\tilde{g}_{**}] \equiv R_{\mu\nu}[g_{**}] + \delta R_{\mu\nu}$$

 $\left[\tilde{\Gamma}^{\lambda}_{\mu\nu} = \Gamma^{\lambda}_{\mu\nu}[g_{**}] + \delta \Gamma^{\lambda}_{\mu\nu}[g_{**}, \omega]\right]$
 $\delta R_{\mu\nu} = \partial_{[\lambda}\delta\Gamma^{\lambda}_{\mu\nu]} + \delta\Gamma^{\lambda}_{[\lambda\sigma}\Gamma^{\sigma}_{\mu\nu]} + \Gamma^{\lambda}_{[\lambda\sigma}\delta\Gamma^{\sigma}_{\mu\nu]} + \delta\Gamma^{\lambda}_{[\lambda\sigma}\delta\Gamma^{\sigma}_{\mu\nu]}$
where $\delta\Gamma^{\lambda}_{\mu\nu} = \omega_{(\mu}\delta^{\lambda}_{\ \nu)} - g_{\mu\nu}\omega^{\lambda}; \quad \omega_{\mu} \equiv \omega_{,\mu}; \quad \omega^{\mu} \equiv \omega^{,\mu}$
 $\delta R_{\mu\nu}[g_{**}, \omega] \equiv A\omega_{\mu}\omega_{\mu} + B\omega_{\mu;\nu} + Cg_{\mu\nu}\omega_{\alpha}\omega^{\alpha} + Dg_{\mu\nu}(\omega^{\alpha})_{;\alpha}$
 $(A, B, C, D \ constants)$
It can only
take this form !

Then:
$$\tilde{R}_{\mu\nu}[\tilde{g}_{**}] \equiv R_{\mu\nu}[g_{**}] + \delta R_{\mu\nu}$$

 $\left[\tilde{\Gamma}^{\lambda}_{\mu\nu} = \Gamma^{\lambda}_{\mu\nu}[g_{**}] + \delta \Gamma^{\lambda}_{\mu\nu}[g_{**}, \omega]\right]$
 $\delta R_{\mu\nu} = \partial_{[\lambda}\delta\Gamma^{\lambda}_{\mu\nu]} + \delta\Gamma^{\lambda}_{[\lambda\sigma}\Gamma^{\sigma}_{\mu\nu]} + \Gamma^{\lambda}_{[\lambda\sigma}\delta\Gamma^{\sigma}_{\mu\nu]} + \delta\Gamma^{\lambda}_{[\lambda\sigma}\delta\Gamma^{\sigma}_{\mu\nu]}$
where $\delta\Gamma^{\lambda}_{\mu\nu} = \omega_{(\mu}\delta^{\lambda}_{\ \nu)} - g_{\mu\nu}\omega^{\lambda}; \quad \omega_{\mu} \equiv \omega_{,\mu}; \quad \omega^{\mu} \equiv \omega^{,\mu}$
 $\delta R_{\mu\nu}[g_{**}, \omega] \equiv A\omega_{\mu}\omega_{\mu} + B\omega_{\mu;\nu} + Cg_{\mu\nu}\omega_{\alpha}\omega^{\alpha} + Dg_{\mu\nu}(\omega^{\alpha})_{;\alpha}$
After some Calculation... $A = +2, B = -2, C = -2, D = -1$

Then:
$$\tilde{R}_{\mu\nu}[\tilde{g}_{**}] \equiv R_{\mu\nu}[g_{**}] + \delta R_{\mu\nu}$$

 $\left[\delta R_{\mu\nu} = 2\omega_{\mu}\omega_{\mu} - 2\omega_{\mu;\nu} - 2g_{\mu\nu}\omega_{\alpha}\omega^{\alpha} - g_{\mu\nu}(\omega^{\alpha})_{;\alpha} \right] ; \quad \omega \equiv \log a(t)$
 $\omega_{\mu} \equiv \omega_{,\mu} = \partial_{\mu}\omega \quad \omega^{\mu} \equiv \omega^{,\mu} = g^{\mu\nu}\partial_{\nu}\omega \quad \omega_{\mu;\nu} = \omega_{\mu,\nu} - \Gamma^{\lambda}_{\mu\nu}\omega_{\lambda} \quad \omega^{\alpha}{}_{;\alpha} = \omega^{\alpha}_{\alpha} + \Gamma^{\alpha}_{\alpha\beta}\omega^{\beta}$

Then:
$$\tilde{R}_{\mu\nu}[\tilde{g}_{**}] \equiv R_{\mu\nu}[g_{**}] + \mathcal{D}_{\mu\nu}\omega$$

 $\left[\mathcal{D}_{\mu\nu}\omega = 2\omega_{\mu}\omega_{\mu} - 2\omega_{\mu;\nu} - 2g_{\mu\nu}\omega_{\alpha}\omega^{\alpha} - g_{\mu\nu}(\omega^{\alpha})_{;\alpha}\right]$; $\omega \equiv \log a(t)$
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Then: $\tilde{R}_{\mu\nu}[\tilde{g}_{**}] \equiv R_{\mu\nu}[\eta_{**} + h_{**}] + \mathcal{D}_{\mu\nu}\omega$ $\left[\mathcal{D}_{\mu\nu}\omega = 2\omega_{\mu}\omega_{\mu} - 2\omega_{\mu;\nu} - 2g_{\mu\nu}\omega_{\alpha}\omega^{\alpha} - g_{\mu\nu}(\omega^{\alpha})_{;\alpha}\right]; \quad \omega \equiv \log a(t)$



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 $\omega_{\mu} \equiv \omega_{,\mu} = \partial_{\mu}\omega$ $\omega^{\mu} \equiv \omega^{,\mu} = g^{\mu\nu}\partial_{\nu}\omega$ $\omega_{\mu;\nu} = \omega_{\mu,\nu} - \Gamma^{\lambda}_{\mu\nu}\omega_{\lambda}$ $\omega^{\alpha}_{;\alpha} = \omega^{\alpha}_{\alpha} + \Gamma^{\alpha}_{\alpha\beta}\omega^{\beta}$

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 $2\omega_{\mu}\omega_{\nu} = 2\mathcal{H}^{2}\delta_{\mu0}\delta_{\nu0}, \quad \mathcal{H} \equiv a'/a$
 $-2\omega_{\mu;\nu} = 2(\mathcal{H}^{2} - a''/a)\delta_{\mu0}\delta_{\nu0} + 2\mathcal{H}\Gamma^{0}_{\mu\nu}[\eta_{**} + h_{**}]$

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 $\Gamma^{\alpha}_{\mu\nu}[\eta_{**} + h_{**}] = \frac{1}{2}\left(\eta^{\alpha\beta} - h^{\alpha\beta} + h^{\alpha}_{\gamma}h^{\gamma\beta} + \dots\right)\left(\underbrace{\partial_{(\mu}h_{\beta\nu)} - \partial_{\beta}h_{\mu\nu}}{\mathcal{O}(h_{**})}\right)$

Then:
$$\tilde{R}_{\mu\nu}[\tilde{g}_{**}] \equiv R_{\mu\nu}[\eta_{**} + h_{**}] + \mathcal{D}_{\mu\nu}\omega$$

 $\left[\mathcal{D}_{\mu\nu}\omega = 2\omega_{\mu}\omega_{\mu} - 2\omega_{\mu;\nu} - 2g_{\mu\nu}\omega_{\alpha}\omega^{\alpha} - g_{\mu\nu}(\omega^{\alpha})_{;\alpha}\right]; \quad \omega \equiv \log a(t)$
 $2\omega_{\mu}\omega_{\nu} = 2\mathcal{H}^{2}\delta_{\mu0}\delta_{\nu0}, \quad \mathcal{H} \equiv a'/a$
 $-2\omega_{\mu;\nu} = 2(\mathcal{H}^{2} - a''/a)\delta_{\mu0}\delta_{\nu0} + 2\mathcal{H}\Gamma^{0}_{\mu\nu}[\eta_{**} + h_{**}]$
 $-2g_{\mu\nu}\omega_{\alpha}\omega^{\alpha} = +2(\eta_{\mu\nu} + h_{\mu\nu})\mathcal{H}^{2}$
 $-g_{\mu\nu}(\omega^{\alpha})_{;\alpha} = (\eta_{\mu\nu} + h_{\mu\nu})(\mathcal{H}' + \mathcal{H}\Gamma^{\alpha}_{\alpha0}[\eta_{**} + h_{**}])$
 $\Gamma^{\alpha}_{\mu\nu}[\eta_{**} + h_{**}] = \frac{1}{2}\left(\eta^{\alpha\beta} - h^{\alpha\beta} + h^{\alpha}_{\gamma}h^{\gamma\beta} + \dots\right)\left(\partial_{(\mu}h_{\beta\nu)} - \partial_{\beta}h_{\mu\nu}\right) = \Gamma^{(1)}_{\mu\nu} + \Gamma^{(2)}_{\mu\nu} + \dots$

$$\begin{aligned} \text{Then:} \quad \tilde{R}_{\mu\nu}[\tilde{g}_{**}] &\equiv R_{\mu\nu}[\eta_{**} + h_{**}] + \mathcal{D}_{\mu\nu}\omega \\ & \left[\mathcal{D}_{\mu\nu}\omega = 2\omega_{\mu}\omega_{\mu} - 2\omega_{\mu;\nu} - 2g_{\mu\nu}\omega_{\alpha}\omega^{\alpha} - g_{\mu\nu}(\omega^{\alpha})_{;\alpha} \right] \; ; \; \; \omega \equiv \log a(t) \\ & \left[2\omega_{\mu}\omega_{\nu} = 2\mathcal{H}^{2}\delta_{\mu0}\delta_{\nu0} , \quad \mathcal{H} \equiv a'/a \\ & -2\omega_{\mu;\nu} = 2(\mathcal{H}^{2} - a''/a)\delta_{\mu0}\delta_{\nu0} + 2\mathcal{H}\Gamma^{0}_{\mu\nu}[\eta_{**} + h_{**}] \\ & -2g_{\mu\nu}\omega_{\alpha}\omega^{\alpha} = + 2(\eta_{\mu\nu} + h_{\mu\nu})\mathcal{H}^{2} \\ & -g_{\mu\nu}(\omega^{\alpha})_{;\alpha} = (\eta_{\mu\nu} + h_{\mu\nu})\left(\mathcal{H}' + \mathcal{H}\Gamma^{\alpha}_{\alpha0}[\eta_{**} + h_{**}]\right) \\ & \Gamma^{\alpha}_{\mu\nu}[\eta_{**} + h_{**}] \equiv \Gamma^{(1)}_{\mu\nu} + \Gamma^{(2)}_{\mu\nu} + \dots \quad \begin{cases} \Gamma^{(1)}_{\mu\nu} \equiv +\frac{1}{2}\eta^{\alpha\beta}\left(\partial_{(\mu}h_{\beta\nu)} - \partial_{\beta}h_{\mu\nu}\right) \\ \Gamma^{\alpha}_{\mu\nu} \equiv -\frac{1}{2}h^{\alpha\beta}\left(\partial_{(\mu}h_{\beta\nu)} - \partial_{\beta}h_{\mu\nu}\right) \end{cases} \end{aligned}$$

$$\begin{aligned} \text{Then:} \quad \tilde{R}_{\mu\nu}[\tilde{g}_{**}] &\equiv R_{\mu\nu}[\eta_{**} + h_{**}] + \mathcal{D}_{\mu\nu}\omega \\ & \left[\mathcal{D}_{\mu\nu}\omega = 2\omega_{\mu}\omega_{\mu} - 2\omega_{\mu;\nu} - 2g_{\mu\nu}\omega_{\alpha}\omega^{\alpha} - g_{\mu\nu}(\omega^{\alpha})_{;\alpha} \right] \; ; \; \; \omega \equiv \log a(t) \\ & \left[2\omega_{\mu}\omega_{\nu} = 2\mathcal{H}^{2}\delta_{\mu0}\delta_{\nu0} , \quad \mathcal{H} \equiv a'/a \\ & -2\omega_{\mu;\nu} = 2(\mathcal{H}^{2} - a''/a)\delta_{\mu0}\delta_{\nu0} + 2\mathcal{H}(\Gamma^{0}_{\mu\nu} + \Gamma^{0}_{\mu\nu})] \\ & -2g_{\mu\nu}\omega_{\alpha}\omega^{\alpha} = + 2(\eta_{\mu\nu} + h_{\mu\nu})\mathcal{H}^{2} \\ & -g_{\mu\nu}(\omega^{\alpha})_{;\alpha} = (\eta_{\mu\nu} + h_{\mu\nu})\left(\mathcal{H}' + \mathcal{H}(\Gamma^{\alpha}_{\alpha0} + \Gamma^{\alpha}_{\alpha0})\right) \\ & \left[\Gamma^{\alpha}_{\mu\nu} [\eta_{**} + h_{**}] \equiv \Gamma^{(1)}_{\mu\nu} + \Gamma^{(2)}_{\mu\nu} + \dots \right] \\ & \left\{ \begin{array}{c} \Gamma^{(1)}_{\mu\nu} = \frac{1}{2}\eta^{\alpha\beta} \left(\partial_{(\mu}h_{\beta\nu)} - \partial_{\beta}h_{\mu\nu} \right) \\ \Gamma^{\alpha}_{\mu\nu} = -\frac{1}{2}h^{\alpha\beta} \left(\partial_{(\mu}h_{\beta\nu)} - \partial_{\beta}h_{\mu\nu} \right) \end{array} \right. \end{aligned}$$

$$\begin{aligned} \text{Then:} \quad \tilde{R}_{\mu\nu} &\equiv R_{\mu\nu} [\eta_{**} + h_{**}] + (\mathcal{D}_{\mu\nu}\omega)^{(0)} + (\mathcal{D}_{\mu\nu}\omega)^{(1)} + (\mathcal{D}_{\mu\nu}\omega)^{(2)} \\ &\left[\mathcal{D}_{\mu\nu}\omega = 2\omega_{\mu}\omega_{\mu} - 2\omega_{\mu;\nu} - 2g_{\mu\nu}\omega_{\alpha}\omega^{\alpha} - g_{\mu\nu}(\omega^{\alpha})_{;\alpha} \right] \; ; \; \; \omega \equiv \log a(t) \\ \\ &\left[2\omega_{\mu}\omega_{\nu} = 2\mathcal{H}^{2}\delta_{\mu0}\delta_{\nu0}, \quad \mathcal{H} \equiv a'/a & (1) & (2) \\ &-2\omega_{\mu;\nu} = 2(\mathcal{H}^{2} - a''/a)\delta_{\mu0}\delta_{\nu0} + 2\mathcal{H}(\Gamma^{0}_{\mu\nu} + \Gamma^{0}_{\mu\nu}) \right] \\ &-2g_{\mu\nu}\omega_{\alpha}\omega^{\alpha} = + 2(\eta_{\mu\nu} + h_{\mu\nu})\mathcal{H}^{2} \\ &-g_{\mu\nu}(\omega^{\alpha})_{;\alpha} = (\eta_{\mu\nu} + h_{\mu\nu})\left(\mathcal{H}' + \mathcal{H}(\Gamma^{\alpha}_{\alpha0} + \Gamma^{\alpha}_{\alpha0})\right) \\ \\ &\left[\Gamma^{\alpha}_{\mu\nu} [\eta_{**} + h_{**}] \equiv \Gamma^{(1)}_{\mu\nu} + \Gamma^{(2)}_{\mu\nu} + \dots \right] \\ &\left[\Gamma^{(2)}_{\mu\nu} = -\frac{1}{2}h^{\alpha\beta} \left(\partial_{(\mu}h_{\beta\nu)} - \partial_{\beta}h_{\mu\nu} \right) \\ \end{array} \right] \end{aligned}$$

Then:
$$\tilde{R}_{\mu\nu} \equiv R_{\mu\nu} [\eta_{**} + h_{**}] + (\mathscr{D}_{\mu\nu}\omega)^{(0)} + (\mathscr{D}_{\mu\nu}\omega)^{(1)} + (\mathscr{D}_{\mu\nu}\omega)^{(2)}$$

 $\left[\mathscr{D}_{\mu\nu}\omega = 2\omega_{\mu}\omega_{\mu} - 2\omega_{\mu;\nu} - 2g_{\mu\nu}\omega_{\alpha}\omega^{\alpha} - g_{\mu\nu}(\omega^{\alpha})_{;\alpha} \right] ; \quad \omega \equiv \log a(t)$



Then:
$$\tilde{R}_{\mu\nu} \equiv \underbrace{R_{\mu\nu} [\eta_{**} + h_{**}]}_{?} + (\mathscr{D}_{\mu\nu}\omega)^{(0)} + (\mathscr{D}_{\mu\nu}\omega)^{(1)} + (\mathscr{D}_{\mu\nu}\omega)^{(2)}$$

$$\Gamma^{\alpha}_{\mu\nu}[\eta_{**} + h_{**}] \equiv \Gamma^{(1)}_{\mu\nu} + \Gamma^{(2)}_{\mu\nu} + \dots \left\{ \begin{array}{c} \Gamma^{(1)}_{\mu\nu} \equiv +\frac{1}{2}\eta^{\alpha\beta} \left(\partial_{(\mu}h_{\beta\nu)} - \partial_{\beta}h_{\mu\nu}\right) \\ \Gamma^{(2)}_{\mu\nu} \equiv -\frac{1}{2}h^{\alpha\beta} \left(\partial_{(\mu}h_{\beta\nu)} - \partial_{\beta}h_{\mu\nu}\right) \end{array} \right.$$

Then:
$$\tilde{R}_{\mu\nu} \equiv R_{\mu\nu} [\eta_{**} + h_{**}] + (\mathscr{D}_{\mu\nu}\omega)^{(0)} + (\mathscr{D}_{\mu\nu}\omega)^{(1)} + (\mathscr{D}_{\mu\nu}\omega)^{(2)}$$
?

$$\Gamma^{\alpha}_{\mu\nu}[\eta_{**} + h_{**}] \equiv \Gamma^{(1)}_{\mu\nu} + \Gamma^{\alpha}_{\mu\nu} + \dots \begin{cases} \Gamma^{(1)}_{\mu\nu} \equiv +\frac{1}{2}\eta^{\alpha\beta} \left(\partial_{(\mu}h_{\beta\nu)} - \partial_{\beta}h_{\mu\nu}\right) \\ \Gamma^{(2)}_{\mu\nu} \equiv -\frac{1}{2}h^{\alpha\beta} \left(\partial_{(\mu}h_{\beta\nu)} - \partial_{\beta}h_{\mu\nu}\right) \end{cases}$$
$$R_{\mu\nu}[\eta_{**} + h_{**}] \equiv \partial_{[\lambda}\Gamma^{\lambda}_{\mu\nu]} + \Gamma^{\alpha}_{[\alpha\lambda}\Gamma^{\lambda}_{\mu\nu]}$$

Then:
$$\tilde{R}_{\mu\nu} \equiv \underbrace{R_{\mu\nu} [\eta_{**} + h_{**}]}_{?} + (\mathscr{D}_{\mu\nu}\omega)^{(0)} + (\mathscr{D}_{\mu\nu}\omega)^{(1)} + (\mathscr{D}_{\mu\nu}\omega)^{(2)}$$

$$\Gamma^{\alpha}_{\mu\nu}[\eta_{**} + h_{**}] \equiv \Gamma^{(1)}_{\mu\nu} + \Gamma^{\alpha}_{\mu\nu} + \dots \left\{ \begin{array}{c} \Gamma^{(1)}_{\mu\nu} \equiv +\frac{1}{2}\eta^{\alpha\beta} \left(\partial_{(\mu}h_{\beta\nu)} - \partial_{\beta}h_{\mu\nu}\right) \\ \Gamma^{(2)}_{\mu\nu} \equiv -\frac{1}{2}h^{\alpha\beta} \left(\partial_{(\mu}h_{\beta\nu)} - \partial_{\beta}h_{\mu\nu}\right) \end{array} \right. \\ R_{\mu\nu}[\eta_{**} + h_{**}] \equiv \partial_{[\lambda}\Gamma^{\lambda}_{\mu\nu]} + \Gamma^{\alpha}_{[\alpha\lambda}\Gamma^{\lambda}_{\mu\nu]} \\ \Gamma^{(1)}_{[\alpha\lambda]} = \partial_{[\lambda}(\Gamma^{\lambda}_{\mu\nu]} + \Gamma^{\lambda}_{\mu\nu]} + \dots) + (\Gamma^{(1)}_{[\alpha\lambda]} + \Gamma^{\alpha}_{[\alpha\lambda]} + \dots)(\Gamma^{\lambda}_{\mu\nu]} + \Gamma^{\lambda}_{\mu\nu]} + \dots)$$

Then:
$$\tilde{R}_{\mu\nu} \equiv \underbrace{R_{\mu\nu} [\eta_{**} + h_{**}]}_{?} + (\mathscr{D}_{\mu\nu}\omega)^{(0)} + (\mathscr{D}_{\mu\nu}\omega)^{(1)} + (\mathscr{D}_{\mu\nu}\omega)^{(2)}$$

$$\Gamma^{\alpha}_{\mu\nu}[\eta_{**} + h_{**}] \equiv \Gamma^{(1)}_{\mu\nu} + \Gamma^{\alpha}_{\mu\nu} + \dots \left\{ \begin{array}{c} \Gamma^{(1)}_{\mu\nu} \equiv +\frac{1}{2}\eta^{\alpha\beta} \left(\partial_{(\mu}h_{\beta\nu)} - \partial_{\beta}h_{\mu\nu}\right) \\ \Gamma^{(2)}_{\mu\nu} \equiv -\frac{1}{2}h^{\alpha\beta} \left(\partial_{(\mu}h_{\beta\nu)} - \partial_{\beta}h_{\mu\nu}\right) \end{array} \right. \\ R_{\mu\nu}[\eta_{**} + h_{**}] \equiv \partial_{[\lambda}\Gamma^{\lambda}_{\mu\nu]} + \partial_{[\lambda}\Gamma^{\lambda}_{\mu\nu]} + \Gamma^{(1)}_{[\alpha\lambda}\Gamma^{\lambda}_{\mu\nu]} + \dots$$
Then:
$$\tilde{R}_{\mu\nu} \equiv \underbrace{R_{\mu\nu} [\eta_{**} + h_{**}]}_{?} + (\mathscr{D}_{\mu\nu}\omega)^{(0)} + (\mathscr{D}_{\mu\nu}\omega)^{(1)} + (\mathscr{D}_{\mu\nu}\omega)^{(2)}$$

$$\Gamma^{\alpha}_{\mu\nu}[\eta_{**} + h_{**}] \equiv \Gamma^{(1)}_{\mu\nu} + \Gamma^{(2)}_{\mu\nu} + \dots \begin{cases} \Gamma^{(1)}_{\mu\nu} \equiv +\frac{1}{2}\eta^{\alpha\beta} \left(\partial_{(\mu}h_{\beta\nu)} - \partial_{\beta}h_{\mu\nu}\right) \\ \Gamma^{(2)}_{\mu\nu} \equiv -\frac{1}{2}h^{\alpha\beta} \left(\partial_{(\mu}h_{\beta\nu)} - \partial_{\beta}h_{\mu\nu}\right) \end{cases}$$

$$R_{\mu\nu}[\eta_{**} + h_{**}] \equiv \underbrace{\partial_{[\lambda}\Gamma^{\lambda}_{\mu\nu]}}_{(1)} + \underbrace{\partial_{[\lambda}\Gamma^{\lambda}_{\mu\nu]}}_{R_{\mu\nu}} + \underbrace{\Gamma^{(1)}_{[\alpha\lambda}\Gamma^{\lambda}_{\mu\nu]}}_{\delta R_{\mu\nu}} + \dots$$

Then:
$$\tilde{R}_{\mu\nu} \equiv \underbrace{R_{\mu\nu} [\eta_{**} + h_{**}]}_{0 + \delta R_{\mu\nu}} + (\mathscr{D}_{\mu\nu}\omega)^{(0)} + (\mathscr{D}_{\mu\nu}\omega)^{(1)} + (\mathscr{D}_{\mu\nu}\omega)^{(2)}$$

$$\Gamma^{\alpha}_{\mu\nu}[\eta_{**} + h_{**}] \equiv \Gamma^{(1)}_{\mu\nu} + \Gamma^{(2)}_{\mu\nu} + \dots \begin{cases} \Gamma^{(1)}_{\mu\nu} \equiv +\frac{1}{2}\eta^{\alpha\beta} \left(\partial_{(\mu}h_{\beta\nu)} - \partial_{\beta}h_{\mu\nu}\right) \\ \Gamma^{(2)}_{\mu\nu} \equiv -\frac{1}{2}h^{\alpha\beta} \left(\partial_{(\mu}h_{\beta\nu)} - \partial_{\beta}h_{\mu\nu}\right) \end{cases}$$
$$R_{\mu\nu}[\eta_{**} + h_{**}] \equiv \underbrace{\partial_{[\lambda}\Gamma^{\lambda}_{\mu\nu]}}_{(1)} + \underbrace{\partial_{[\lambda}\Gamma^{\lambda}_{\mu\nu]}}_{(\lambda} + \Gamma^{(1)}_{[\alpha\lambda}\Gamma^{(1)}_{\mu\nu]} + \dots \\ \underbrace{\partial_{[\lambda}\Gamma^{\lambda}_{\mu\nu]}}_{(2)} + \underbrace{\partial_{[\lambda}\Gamma^{\lambda}_{\mu\nu]}}_{(2)} + \underbrace{\partial_{[\lambda}\Gamma^{\lambda}_{\mu\nu]}}_{(2)} + \underbrace{\partial_{[\lambda}\Gamma^{\lambda}_{\mu\nu]}}_{(2)} + \dots \\ \underbrace{\partial_{[\lambda}\Gamma^{\lambda}_{\mu\nu]}}_{(2)} + \underbrace{\partial_{[\lambda}\Gamma^{\lambda}_{\mu\mu}}_{(2)} + \underbrace{\partial_{[\lambda}\Gamma$$

Then: $\tilde{R}_{\mu\nu} \equiv \delta R^{(1)}_{\mu\nu} + \delta R^{(2)}_{\mu\nu} + (\mathscr{D}_{\mu\nu}\omega)^{(0)} + (\mathscr{D}_{\mu\nu}\omega)^{(1)} + (\mathscr{D}_{\mu\nu}\omega)^{(2)}$



Then:
$$\tilde{R}_{\mu\nu} \equiv (\mathscr{D}_{\mu\nu}\omega)^{(0)} + \left(\delta R^{(1)}_{\mu\nu} + (\mathscr{D}_{\mu\nu}\omega)^{(1)}\right) + \left(\delta R^{(2)}_{\mu\nu} + (\mathscr{D}_{\mu\nu}\omega)^{(2)}\right)$$



Then:
$$\tilde{R}_{\mu\nu} \equiv \underbrace{(\mathscr{D}_{\mu\nu}\omega)^{(0)}}_{\mathcal{O}(h^0_{**})} + \underbrace{\left(\underbrace{\delta R_{\mu\nu}^{(1)} + (\mathscr{D}_{\mu\nu}\omega)^{(1)}}_{\mathcal{O}(h_{**})}\right) + \underbrace{\left(\underbrace{\delta R_{\mu\nu}^{(2)} + (\mathscr{D}_{\mu\nu}\omega)^{(2)}}_{\mathcal{O}(h_{**})}\right)}_{\mathcal{O}(h_{**})}$$





Then:
$$\tilde{R}_{\mu\nu} \equiv \tilde{R}_{\mu\nu}^{(0)} + \tilde{R}_{\mu\nu}^{(1)} + \tilde{R}_{\mu\nu}^{(2)}$$



Then:
$$\tilde{R}_{\mu\nu} \equiv \tilde{\tilde{R}}_{\mu\nu} + \tilde{\tilde{R}}_{\mu\nu} + \tilde{\tilde{R}}_{\mu\nu}^{(2)}$$



Then:
$$\tilde{R}_{\mu\nu} \equiv \tilde{R}_{\mu\nu}^{(0)} + \tilde{R}_{\mu\nu}^{(1)} + \tilde{R}_{\mu\nu}^{(2)}$$

$$\begin{aligned} \stackrel{(0)}{\tilde{R}}_{\mu\nu} &= 2(2\mathscr{H}^2 - a''/a)\delta_{\mu0}\delta_{\nu0} + (\mathscr{H}^2 + a''/a)\eta_{\mu\nu} \text{ [Background]} \\ \stackrel{(1)}{\tilde{R}}_{\mu\nu} &= \delta_{i\mu}\delta_{j\nu} \Big(-\frac{1}{2}\eta^{\alpha\beta}\partial_{\alpha}\partial_{\beta}h_{ij} + \mathscr{H}h'_{ij} + (\mathscr{H}^2 + a''/a)h_{ij} \Big) \end{aligned}$$

Then:
$$\tilde{R}_{\mu\nu} \equiv \tilde{R}_{\mu\nu}^{(0)} + \tilde{R}_{\mu\nu}^{(1)} + \tilde{R}_{\mu\nu}^{(2)}$$

$$\begin{split} & \stackrel{(0)}{\tilde{R}}_{\mu\nu} = 2(2\mathscr{H}^2 - a''/a)\delta_{\mu0}\delta_{\nu0} + (\mathscr{H}^2 + a''/a)\eta_{\mu\nu} \text{ [Background]} \\ & \stackrel{(1)}{\tilde{R}}_{\mu\nu} = \delta_{i\mu}\delta_{j\nu} \Big(-\frac{1}{2}\eta^{\alpha\beta}\partial_{\alpha}\partial_{\beta}h_{ij} + \mathscr{H}h'_{ij} + (\mathscr{H}^2 + a''/a)h_{ij} \Big) \\ & \stackrel{(2)}{\tilde{R}}_{\mu\nu} = -\frac{1}{2}\mathscr{H}\eta_{\mu\nu}h^{\alpha\beta}h'_{\alpha\beta} + \delta^{(2)}_{R\mu\nu} \\ & \stackrel{(2)}{\to_{[\lambda}\Gamma^{\lambda}_{\mu\nu]} + \Gamma^{(1)}_{[\alpha\lambda}\Gamma^{\lambda}_{\mu\nu]}} \end{split}$$

Then:
$$\tilde{R}_{\mu\nu} \equiv \tilde{R}_{\mu\nu}^{(0)} + \tilde{R}_{\mu\nu}^{(1)} + \tilde{R}_{\mu\nu}^{(2)}$$

Let's forget for the moment of second order parts ...

$$\begin{split} & \stackrel{(0)}{\tilde{R}}_{\mu\nu} = 2(2\mathscr{H}^2 - a''/a)\delta_{\mu0}\delta_{\nu0} + (\mathscr{H}^2 + a''/a)\eta_{\mu\nu} \text{ [Background]} \\ & \stackrel{(1)}{\tilde{R}}_{\mu\nu} = \delta_{i\mu}\delta_{j\nu}\Big(-\frac{1}{2}\eta^{\alpha\beta}\partial_{\alpha}\partial_{\beta}h_{ij} + \mathscr{H}h'_{ij} + (\mathscr{H}^2 + a''/a)h_{ij}\Big) \\ & \stackrel{(2)}{\tilde{R}}_{\mu\nu} = -\frac{1}{2}\mathscr{H}\eta_{\mu\nu}h^{\alpha\beta}h'_{\alpha\beta} + \delta\stackrel{(2)}{\tilde{R}}_{\mu\nu} \\ & \stackrel{(2)}{\tilde{R}}_{i\mu\nu} = -\frac{1}{2}\mathscr{H}\eta_{\mu\nu}h^{\alpha\beta}h'_{\alpha\beta} + \stackrel{(2)}{\tilde{R}}_{\mu\nu} \\ & \stackrel{(2)}{\tilde{R}}_{i\mu\nu]} + \stackrel{(1)}{\tilde{\Gamma}}_{i\alpha\lambda}^{(1)}\Gamma^{\lambda}_{\mu\nu]} \end{split}$$

Then:
$$\tilde{R}_{\mu\nu} \equiv \tilde{R}_{\mu\nu}^{(0)} + \tilde{R}_{\mu\nu}^{(1)}$$

Let's forget for the moment of second order parts ...

$$\begin{aligned} \stackrel{(0)}{\tilde{R}}_{\mu\nu} &= 2(2\mathscr{H}^2 - a''/a)\delta_{\mu0}\delta_{\nu0} + (\mathscr{H}^2 + a''/a)\eta_{\mu\nu} \text{ [Background]} \\ \stackrel{(1)}{\tilde{R}}_{\mu\nu} &= \delta_{i\mu}\delta_{j\nu} \Big(-\frac{1}{2}\eta^{\alpha\beta}\partial_{\alpha}\partial_{\beta}h_{ij} + \mathscr{H}h'_{ij} + (\mathscr{H}^2 + a''/a)h_{ij} \Big) \end{aligned}$$

Then:
$$\tilde{R}_{\mu\nu} \equiv \tilde{\tilde{R}}_{\mu\nu} + \tilde{\tilde{R}}_{\mu\nu}^{(1)}$$

.....

Let's focus on the Einstein Equations

$$\begin{aligned} \stackrel{(0)}{\tilde{R}}_{\mu\nu} &= 2(2\mathscr{H}^2 - a''/a)\delta_{\mu0}\delta_{\nu0} + (\mathscr{H}^2 + a''/a)\eta_{\mu\nu} \text{ [Background]} \\ \stackrel{(1)}{\tilde{R}}_{\mu\nu} &= \delta_{i\mu}\delta_{j\nu} \Big(-\frac{1}{2}\eta^{\alpha\beta}\partial_{\alpha}\partial_{\beta}h_{ij} + \mathscr{H}h'_{ij} + (\mathscr{H}^2 + a''/a)h_{ij} \Big) \end{aligned}$$

Then:
$$\tilde{R}_{\mu\nu} \equiv \overset{(0)}{\tilde{R}}_{\mu\nu} + \overset{(1)}{\tilde{R}}_{\mu\nu}$$
; $m_p^2 \tilde{R}_{\mu\nu} = S_{\mu\nu}$
[$S_{\mu\nu} \equiv T_{\mu\nu} - \frac{1}{2} g_{\mu\nu} g^{\alpha\beta} T_{\alpha\beta}$]

$$\begin{aligned} \stackrel{(0)}{\tilde{R}}_{\mu\nu} &= 2(2\mathscr{H}^2 - a''/a)\delta_{\mu0}\delta_{\nu0} + (\mathscr{H}^2 + a''/a)\eta_{\mu\nu} \text{ [Background]} \\ \stackrel{(1)}{\tilde{R}}_{\mu\nu} &= \delta_{i\mu}\delta_{j\nu} \Big(-\frac{1}{2}\eta^{\alpha\beta}\partial_{\alpha}\partial_{\beta}h_{ij} + \mathscr{H}h'_{ij} + (\mathscr{H}^2 + a''/a)h_{ij} \Big) \end{aligned}$$

Then:
$$\tilde{R}_{\mu\nu} \equiv \underbrace{\tilde{R}_{\mu\nu}}_{\mu\nu} + \underbrace{\tilde{R}_{\mu\nu}}_{\mu\nu}$$
; $m_p^2 \tilde{R}_{\mu\nu} = S_{\mu\nu}$

$$\begin{aligned} \stackrel{(0)}{\tilde{R}}_{\mu\nu} &= 2(2\mathscr{H}^2 - a''/a)\delta_{\mu0}\delta_{\nu0} + (\mathscr{H}^2 + a''/a)\eta_{\mu\nu} \text{ [Background]} \\ \stackrel{(1)}{\tilde{R}}_{\mu\nu} &= \delta_{i\mu}\delta_{j\nu} \Big(-\frac{1}{2}\eta^{\alpha\beta}\partial_{\alpha}\partial_{\beta}h_{ij} + \mathscr{H}h'_{ij} + (\mathscr{H}^2 + a''/a)h_{ij} \Big) \end{aligned}$$

Then:
$$\tilde{R}_{\mu\nu} \equiv \underbrace{\tilde{R}_{\mu\nu}}_{(0)} + \underbrace{\tilde{R}_{\mu\nu}}_{(1)}$$
; $m_p^2 \tilde{R}_{\mu\nu} = S_{\mu\nu}$
 $\underbrace{\tilde{R}_{\mu\nu}}_{(0)} = \underbrace{S_{\mu\nu}}_{(0)} + \underbrace{S_{\mu\nu}}_{(1)}$

$$\begin{split} & \stackrel{(0)}{\tilde{R}}_{\mu\nu} = 2(2\mathscr{H}^2 - a''/a)\delta_{\mu0}\delta_{\nu0} + (\mathscr{H}^2 + a''/a)\eta_{\mu\nu} \text{ [Background]} \\ & \stackrel{(1)}{\tilde{R}}_{\mu\nu} = \delta_{i\mu}\delta_{j\nu} \Big(-\frac{1}{2}\eta^{\alpha\beta}\partial_{\alpha}\partial_{\beta}h_{ij} + \mathscr{H}h'_{ij} + (\mathscr{H}^2 + a''/a)h_{ij} \Big) \end{split}$$

$$\begin{aligned} & \stackrel{(0)}{\tilde{R}}_{\mu\nu} = 2(2\mathscr{H}^2 - a''/a)\delta_{\mu 0}\delta_{\nu 0} + (\mathscr{H}^2 + a''/a)\eta_{\mu\nu} \text{ [Background]} \\ & \stackrel{(1)}{\tilde{R}}_{\mu\nu} = \delta_{i\mu}\delta_{j\nu} \Big(-\frac{1}{2}\eta^{\alpha\beta}\partial_{\alpha}\partial_{\beta}h_{ij} + \mathscr{H}h'_{ij} + (\mathscr{H}^2 + a''/a)h_{ij} \Big) \end{aligned}$$



$$\begin{split} S_{\mu\nu} &= (\rho + p)u_{\mu}u_{\nu} + \frac{1}{2}(\rho - p)\tilde{g}_{\mu\nu} + \Pi_{ij} \quad ; \quad u_{\mu} \equiv (a, 0, 0, 0) \\ &= (\rho + p)a^{2}\delta_{\mu0}\delta_{\mu0} + \frac{1}{2}(\rho - p)a^{2}\eta_{\mu\nu} + \frac{1}{2}(\rho - p)a^{2}h_{\mu\nu} + \Pi_{ij} \\ &= \underbrace{ \overset{(0)}{S_{\mu\nu}}}_{S_{\mu\nu}} + \underbrace{ \overset{(1)}{S_{\mu\nu}}}_{S_{\mu\nu}} \end{split}$$

$$\begin{split} S_{\mu\nu} &= (\rho + p)u_{\mu}u_{\nu} + \frac{1}{2}(\rho - p)\tilde{g}_{\mu\nu} + \Pi_{ij} \quad ; \quad u_{\mu} \equiv (a, 0, 0, 0) \\ &= (\rho + p)a^{2}\delta_{\mu0}\delta_{\mu0} + \frac{1}{2}(\rho - p)a^{2}\eta_{\mu\nu} + \frac{1}{2}(\rho - p)a^{2}h_{\mu\nu} + \Pi_{ij} \\ &= \underbrace{(0)}_{S_{\mu\nu}} \qquad + \underbrace{(1)}_{S_{\mu\nu}} + \underbrace{(1)}_{S_{\mu\nu}} \\ &= \underbrace{(1)}_{I_{ij}} = \Pi_{ij}^{(S)} + \Pi_{ij}^{(V)} + \Pi_{ij}^{(T)}] \end{split}$$

$$m_{p}^{2} \begin{pmatrix} {}^{(0)}_{\tilde{R}_{\mu\nu}} + {}^{(1)}_{\tilde{R}_{\mu\nu}} \end{pmatrix} = \underbrace{(\rho + p)a^{2}\delta_{\mu0}\delta_{\mu0} + \frac{1}{2}(\rho - p)a^{2}\eta_{\mu\nu}}_{(0)} + \underbrace{\frac{1}{2}(\rho - p)a^{2}h_{\mu\nu} + \Pi_{ij}^{(T)}}_{(1)} + \underbrace{\frac{1}{2}(\rho - p)a^{2}h_{\mu\nu} + \Pi_{ij}^{(T)}}_{S_{\mu\nu}} + \underbrace{\frac{1}{2}(\rho - p)a^{2}h_{\mu\nu} + \underbrace{\frac{1}{2}(\rho -$$



Background: $m_p^2 \tilde{R}_{\mu\nu} = \overset{(0)}{S}_{\mu\nu}$

$$(\mu,\nu) = (0,0): (\mathscr{H}^2 - a''/a) = \frac{a^2}{6m_p^2}(\rho+3p)$$
 (I)

$$(\mu,\nu) = (i,i): (\mathscr{H}^2 + a''/a) = \frac{a^2}{2m_p^2}(\rho-p)$$
 (II)

Background:
$$m_p^2 \tilde{R}_{\mu\nu} = \overset{(0)}{S}_{\mu\nu}$$

0

P

(I) + (II):
$$\mathscr{H}^2 = \frac{a^2}{3m_p^2}\rho$$

(II) - (I): $\frac{a''}{a} = \frac{a^2}{6m_p^2}(\rho - 3p)$





 $\langle \mathbf{n} \rangle$



Then:
$$\tilde{R}_{\mu\nu} \equiv \underbrace{\tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu}}_{(0)}$$
; $m_p^2 \tilde{R}_{\mu\nu} = S_{\mu\nu} \longrightarrow \begin{cases} m_p^2 \tilde{R}_{\mu\nu} = S_{\mu\nu} \\ m_p^2 \tilde{R}_{\mu\nu} = S_{\mu\nu} \end{cases}$

First Order:
$$m_p^2 \tilde{R}_{\mu\nu} = \overset{(1)}{S}_{\mu\nu}$$

 $\underbrace{h_{ij}'' - \nabla^2 h_{ij} + 2\mathscr{H} h_{ij}'}_{\text{wave operator}} + \underbrace{2(\mathscr{H}^2 + a''/a)h_{ij}}_{\text{mass term?}} = \frac{2}{m_p^2} \Pi_{ij}^{(\text{T})} + \frac{a^2(\rho - p)}{m_p^2} h_{ij}$

Then:
$$\tilde{R}_{\mu\nu} \equiv \underbrace{\tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu}}_{(0)}$$
; $m_p^2 \tilde{R}_{\mu\nu} = S_{\mu\nu} \longrightarrow \begin{cases} m_p^2 \tilde{R}_{\mu\nu} = S_{\mu\nu} \\ m_p^2 \tilde{R}_{\mu\nu} = S_{\mu\nu} \end{cases}$

First Order:
$$m_p^2 \tilde{R}_{\mu\nu} = \overset{(1)}{S}_{\mu\nu}$$

 $\underbrace{h_{ij}'' - \nabla^2 h_{ij} + 2\mathcal{H} h_{ij}'}_{\text{wave operator}} + \underbrace{2(\mathcal{H}^2 + a'' + a)h_{ij}}_{\text{mass term?}} = \frac{2}{m_p^2} \Pi_{ij}^{(T)} + \frac{a^2(\rho - p)}{m_p^2} h_{ij}$
wave operator mass term?

 (\mathbf{n})



 (\mathbf{n})



Then:
$$\tilde{R}_{\mu\nu} \equiv \tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu}^{(1)}$$

[Friedmann [GW Eq.
Equations] motion]



Then:
$$\tilde{R}_{\mu\nu} \equiv \tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu}$$

[Friedmann [GW Eq. ?

$$\begin{split} \stackrel{(2)}{\tilde{R}}_{\mu\nu} &= -\frac{1}{2} \mathscr{H} \eta_{\mu\nu} h^{\alpha\beta} h_{\alpha\beta}' + \delta \stackrel{(2)}{R}_{\mu\nu} \quad ; \quad \delta \stackrel{(2)}{R}_{\mu\nu} \equiv \partial_{[\lambda} \Gamma^{\lambda}_{\mu\nu]} + \Gamma^{\alpha}_{[\alpha\lambda} \Gamma^{\lambda}_{\mu\nu]} \\ & \left\{ \stackrel{(1)}{\Gamma^{\alpha}_{\mu\nu}} \equiv +\frac{1}{2} \eta^{\alpha\beta} \left(\partial_{(\mu} h_{\beta\nu)} - \partial_{\beta} h_{\mu\nu} \right) \right. \\ \left. \stackrel{(2)}{\Gamma^{\alpha}_{\mu\nu}} \equiv -\frac{1}{2} h^{\alpha\beta} \left(\partial_{(\mu} h_{\beta\nu)} - \partial_{\beta} h_{\mu\nu} \right) \right\} \end{split}$$

Then:
$$\tilde{R}_{\mu\nu} \equiv \tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu}$$

[Friedmann [GW Eq. ?

$$S_{\rm HE} = \frac{m_p^2}{2} \int d^4 x \sqrt{-\tilde{g}} \tilde{R}$$

Then:
$$\tilde{R}_{\mu\nu} \equiv \tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu}$$

[Friedmann [GW Eq. ?

$$S_{\rm HE} = \frac{m_p^2}{2} \int d^4 x \sqrt{-\tilde{g}} \tilde{R}$$
$$= \frac{m_p^2}{2} \int d^4 x \sqrt{-\tilde{g}} \tilde{g}^{\mu\nu} \tilde{R}_{\mu\nu}$$
$$= \frac{m_p^2}{2} \int d^4 x \tilde{f}^{\mu\nu} \tilde{R}_{\mu\nu}$$

Then:
$$\tilde{R}_{\mu\nu} \equiv \tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu}$$

[Friedmann [GW Eq. ?

$$S_{\rm HE} = \frac{m_p^2}{2} \int d^4x \, \tilde{f}^{\mu\nu} \tilde{R}_{\mu\nu}$$

Then:
$$\tilde{R}_{\mu\nu} \equiv \tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu}$$

[Friedmann [GW Eq. ?

$$S_{\rm HE} = \frac{m_p^2}{2} \int d^4x \, \tilde{f}^{\mu\nu} \tilde{R}_{\mu\nu}$$
$$\tilde{f}^{\mu\nu} \equiv \sqrt{-\tilde{g}} \, \tilde{g}^{\mu\nu} = a(t)^4 \left(1 - \frac{1}{4}h_{\mu\nu}h^{\mu\nu}\right) a^{-2} \left(\eta^{\mu\nu} - h^{\mu\nu} + h^{\mu}_{\ \alpha}h^{\alpha\nu}\right)$$
Then:
$$\tilde{R}_{\mu\nu} \equiv \tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu}$$

[Friedmann [GW Eq. ?

$$S_{\rm HE} = \frac{m_p^2}{2} \int d^4x \, \tilde{f}^{\mu\nu} \tilde{R}_{\mu\nu}$$

$$\tilde{f}^{\mu\nu} \equiv \sqrt{-\tilde{g}} \, \tilde{g}^{\mu\nu} = a(t)^4 \left(1 - \frac{1}{4}h_{\mu\nu}h^{\mu\nu}\right) a^{-2} \left(\eta^{\mu\nu} - h^{\mu\nu} + h^{\mu}_{\ \alpha}h^{\alpha\nu}\right)$$

$$= \underbrace{a(t)^2 \eta^{\mu\nu}}_{(0)} - \underbrace{a(t)^2 h^{\mu\nu}}_{(1)} + \underbrace{h^{\mu}_{\ \alpha}h^{\alpha\nu} - \frac{1}{4}\eta^{\mu\nu}h_{\alpha\beta}h^{\alpha\beta}}_{(2)} + \underbrace{\tilde{f}^{\mu\nu}}_{(2)}$$

Then:
$$\tilde{R}_{\mu\nu} \equiv \tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu}$$

[Friedmann [GW Eq. ?

$$\begin{split} S_{\rm HE} &= \frac{m_p^2}{2} \int d^4 x \, \tilde{f}^{\mu\nu} \tilde{R}_{\mu\nu} \\ &= \frac{m_p^2}{2} \int d^4 x \left(\tilde{f}^{\mu\nu} + \tilde{f}^{\mu\nu} + \tilde{f}^{\mu\nu} \right) \begin{pmatrix} 0 & 0 & 0 \\ \tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu} \end{pmatrix} \\ &= S_{\rm HE}^{(0)} + S_{\rm HE}^{(1)} + S_{\rm HE}^{(2)} \end{split}$$

Then:
$$\tilde{R}_{\mu\nu} \equiv \tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu}$$

[Friedmann [GW Eq. ?

$$\begin{split} S_{\rm HE} &= S_{\rm HE}^{(0)} + S_{\rm HE}^{(1)} + S_{\rm HE}^{(2)} \\ S_{\rm HE}^{(0)} &= \frac{m_p^2}{2} \int d^4 x \, \tilde{f}^{\mu\nu} \tilde{R}_{\mu\nu} \\ S_{\rm HE}^{(1)} &= \frac{m_p^2}{2} \int d^4 x \left(\tilde{f}^{(0)} \, {}^{(0)} \, {}^{(1)} + {}^{(1)} \, {}^{(0)} \, {}^{(0)} \right) \\ S_{\rm HE}^{(2)} &= \frac{m_p^2}{2} \int d^4 x \left(\tilde{f}^{\mu\nu} \tilde{R}_{\mu\nu} + \tilde{f}^{\mu\nu} \tilde{R}_{\mu\nu} + \tilde{f}^{\mu\nu} \tilde{R}_{\mu\nu} \right) \\ S_{\rm HE}^{(2)} &= \frac{m_p^2}{2} \int d^4 x \left(\tilde{f}^{\mu\nu} \tilde{R}_{\mu\nu} + \tilde{f}^{\mu\nu} \tilde{R}_{\mu\nu} + \tilde{f}^{\mu\nu} \tilde{R}_{\mu\nu} \right) \end{split}$$

Then:
$$\tilde{R}_{\mu\nu} \equiv \tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu}$$

[Friedmann [GW Eq. ?

$$S_{\rm HE} = S_{\rm HE}^{(0)} + S_{\rm HE}^{(1)} + S_{\rm HE}^{(2)}$$

$$S_{\rm HE} = 3m_p^2 \int d^4 x \ a(t)a''(t)$$

$$S_{\rm HE}^{(1)} = 0$$

$$S_{\rm HE}^{(2)} = -\frac{m_p^2}{4} \int d^4 x \ a^2(t) \left(\frac{1}{2}\eta^{\mu\nu}\partial_{\mu}h_{ij}\partial_{\nu}h_{ij} + 2\mathscr{H}h_{ij}h'_{ij} + 3\frac{a''}{a}h_{ij}h_{ij}\right)$$

Then:
$$\tilde{R}_{\mu\nu} \equiv \tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu}$$

[Friedmann [GW Eq. ?

$$S_{\rm HE}^{(2)} \equiv -\frac{m_p^2}{4} \int d^4x \, a^2(t) \left(\frac{1}{2}\eta^{\mu\nu}\partial_\mu h_{ij}\partial_\nu h_{ij} + 2\mathscr{H}h_{ij}h_{ij}' + 3\frac{a''}{a}h_{ij}h_{ij}\right)$$

Then:
$$\tilde{R}_{\mu\nu} \equiv \tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu}$$

[Friedmann [GW Eq. ?

$$S_{\rm HE}^{(2)} \equiv -\frac{m_p^2}{4} \int d^4x \, a^2(t) \left(\frac{1}{2}\eta^{\mu\nu}\partial_{\mu}h_{ij}\partial_{\nu}h_{ij} + 2\mathcal{H}h_{ij}h'_{ij} + 3\frac{a''}{a}h_{ij}h_{ij}\right)$$

Consistency check: Find Eq.'s of motion of h_{ii}



Then:
$$\tilde{R}_{\mu\nu} \equiv \tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu}$$

[Friedmann [GW Eq. ?

$$\begin{split} S_{\rm HE}^{(2)} &\equiv -\frac{m_p^2}{4} \int d^4x \, a^2(t) \Big(\frac{1}{2} \eta^{\mu\nu} \partial_\mu h_{ij} \partial_\nu h_{ij} + 2 \mathscr{H} h_{ij} h_{ij}' + 3 \frac{a''}{a} h_{ij} h_{ij} \Big) \\ \delta S_{\rm HE}^{(2)} &\equiv -\frac{m_p^2}{4} \int d^4x \, a^2 \Big(\frac{h_{ij}'' + 2 \mathscr{H} h_{ij}' - \nabla^2 h_{ij}}{wave operator} + \frac{2(\mathscr{H}' + a''/a) h_{ij}}{-\frac{2a^2 p}{m_p^2}} \Big) \delta h_{ij} \end{split}$$

Then:
$$\tilde{R}_{\mu\nu} \equiv \tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu}$$

[Friedmann [GW Eq. ?

 $\delta S_{\rm HE}^{(2)} \equiv -\frac{m_p^2}{4} \left[d^4 x \, a^2 \left(h_{ij}'' + 2\mathcal{H} h_{ij}' - \nabla^2 h_{ij} \right) \delta h_{ij} + \frac{1}{2} \left[d^4 x \, a^4 p \, h_{ij} \delta h_{ij} \right] \right]$

Then:
$$\tilde{R}_{\mu\nu} \equiv \tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu}$$

[Friedmann [GW Eq. ?

$$\delta S_{\rm HE}^{(2)} \equiv -\frac{m_p^2}{4} \int d^4x \, a^2 \Big(h_{ij}'' + 2\mathcal{H} h_{ij}' - \nabla^2 h_{ij} \Big) \delta h_{ij} + \frac{1}{2} \int d^4x \, a^4p \, h_{ij} \delta h_{ij}$$
$$S_{\rm m} \equiv \int d^4x \, \sqrt{-\tilde{g}} \, \mathscr{L}_{\rm m} \quad \text{(matter sector)}$$

Then:
$$\tilde{R}_{\mu\nu} \equiv \tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu}$$

[Friedmann [GW Eq. ?

$$\delta S_{\rm HE}^{(2)} \equiv -\frac{m_p^2}{4} \int d^4x \, a^2 \Big(h_{ij}'' + 2\mathcal{H} h_{ij}' - \nabla^2 h_{ij} \Big) \delta h_{ij} + \frac{1}{2} \int d^4x \, a^4p \, h_{ij} \delta h_{ij}$$
$$S_{\rm m} \equiv \int d^4x \, \sqrt{-\tilde{g}} \, \mathscr{L}_{\rm m} = S_{\rm m}^{(0)} + S_{\rm m}^{(1)} + S_{\rm m}^{(2)} + \dots$$

Then:
$$\tilde{R}_{\mu\nu} \equiv \tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu}$$

[Friedmann [GW Eq. ?

$$\delta S_{\text{HE}}^{(2)} \equiv -\frac{m_p^2}{4} \int d^4x \, a^2 \Big(h_{ij}'' + 2\mathscr{H} h_{ij}' - \nabla^2 h_{ij} \Big) \delta h_{ij} + \frac{1}{2} \int d^4x \, a^4p \, h_{ij} \delta h_{ij}$$

$$S_{\text{m}} \equiv \int d^4x \, \sqrt{-\tilde{g}} \, \mathscr{L}_{\text{m}} = S_{\text{m}}^{(0)} + S_{\text{m}}^{(1)} + S_{\text{m}}^{(2)} + \dots$$

$$\delta S_{\text{m}}^{(2)} \equiv -\frac{1}{2} \int d^4x \, \sqrt{-\tilde{g}} \, T_{\mu\nu} \, \delta \tilde{g}^{\mu\nu} - \frac{1}{4} \int d^4x \, \sqrt{-\tilde{g}} \, \Big(\frac{1}{\sqrt{-\tilde{g}}} \, \frac{\delta(\sqrt{-\tilde{g}} \, T_{\mu\nu})}{\delta \tilde{g}^{\alpha\beta}} \Big) \delta \tilde{g}^{\mu\nu} \delta \tilde{g}^{\alpha\beta}$$

Then:
$$\tilde{R}_{\mu\nu} \equiv \tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu}$$

[Friedmann [GW Eq. ?

$$\delta S_{\text{HE}}^{(2)} \equiv -\frac{m_p^2}{4} \int d^4 x \, a^2 \Big(h_{ij}'' + 2\mathscr{H} h_{ij}' - \nabla^2 h_{ij} \Big) \delta h_{ij} + \frac{1}{2} \int d^4 x \, a^4 p \, h_{ij} \delta h_{ij}$$

$$S_{\text{m}} \equiv \int d^4 x \, \sqrt{-\tilde{g}} \, \mathscr{L}_{\text{m}} = S_{\text{m}}^{(0)} + S_{\text{m}}^{(1)} + S_{\text{m}}^{(2)} + \dots$$

$$\delta S_{\text{m}}^{(2)} \equiv -\frac{1}{2} \int d^4 x \, \sqrt{-\tilde{g}} \, T_{\mu\nu} \, \delta \tilde{g}^{\mu\nu} - \frac{1}{4} \int d^4 x \, \sqrt{-\tilde{g}} \, \Big(\frac{1}{\sqrt{-\tilde{g}}} \frac{\delta(\sqrt{-\tilde{g}} \, T_{\mu\nu})}{\delta \tilde{g}^{\alpha\beta}} \Big) \delta \tilde{g}^{\mu\nu} \delta \tilde{g}^{\alpha\beta}$$

$$\frac{1}{2} \int d^4 x \, a^2(t) \, \Pi_{ij}^{(\text{T})} \, h_{ij} - \frac{1}{4} \int d^4 x \, a^4(t) \, p \, h_{ij} h_{ij}$$

Then:
$$\tilde{R}_{\mu\nu} \equiv \tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu}$$

[Friedmann [GW Eq. ?

$$\delta S_{\rm HE}^{(2)} \equiv -\frac{m_p^2}{4} \int d^4x \, a^2 \Big(h_{ij}'' + 2\mathscr{H} h_{ij}' - \nabla^2 h_{ij} \Big) \delta h_{ij} + \frac{1}{2} \int d^4x \, a^4p \, h_{ij} \delta h_{ij}$$

$$S_{\rm m}^{(2)} \equiv \frac{1}{2} \int d^4x \, a^2(t) \, \Pi_{ij}^{({\rm T})} h_{ij} - \frac{1}{4} \int d^4x \, a^4(t) \, p \, h_{ij} h_{ij}$$

Then:
$$\tilde{R}_{\mu\nu} \equiv \tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu}$$

[Friedmann [GW Eq. ?

$$\delta S_{\rm HE}^{(2)} \equiv -\frac{m_p^2}{4} \int d^4x \, a^2 \Big(h_{ij}'' + 2\mathscr{H} h_{ij}' - \nabla^2 h_{ij} \Big) \delta h_{ij} + \frac{1}{2} \int d^4x \, a^4p \, h_{ij} \delta h_{ij}$$

$$\delta S_{\rm m}^{(2)} = \frac{1}{2} \int d^4x \, a^2(t) \, \Pi_{ij}^{(T)} \, \delta h_{ij} - \frac{1}{2} \int d^4x \, a^4(t) \, p \, h_{ij} \, \delta h_{ij}$$

Then:
$$\tilde{R}_{\mu\nu} \equiv \tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu}$$

[Friedmann [GW Eq. ?

$$\delta S_{\text{HE}}^{(2)} \equiv -\frac{m_p^2}{4} \int d^4 x \, a^2 \Big(h_{ij}'' + 2\mathcal{H} h_{ij}' - \nabla^2 h_{ij} \Big) \delta h_{ij} + \frac{1}{2} \int d^4 x \, a^4 p \, h_{ij} \delta h_{ij}$$

$$\delta S_{\text{m}}^{(2)} = \frac{1}{2} \int d^4 x \, a^2(t) \, \Pi_{ij}^{(\text{T})} \, \delta h_{ij} - \frac{1}{2} \int d^4 x \, a^4(t) \, p \, h_{ij} \, \delta h_{ij}$$

$$\delta S_{\text{m}}^{(2)} + \delta S_{\text{HE}}^{(2)} = 0$$

Then:
$$\tilde{R}_{\mu\nu} \equiv \tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu}$$

[Friedmann [GW Eq. ?

$$\delta S_{\rm HE}^{(2)} = -\frac{m_p^2}{4} \int d^4 x \, a^2 \Big(h_{ij}'' + 2\mathcal{H} h_{ij}' - \nabla^2 h_{ij} \Big) \delta h_{ij} + \frac{1}{2} \int d^4 x \, a^4 p \, h_{ij} \delta h_{ij}$$

$$\delta S_{\rm m}^{(2)} = \frac{1}{2} \int d^4 x \, a^2(t) \, \Pi_{ij}^{({\rm T})} \, \delta h_{ij} - \frac{1}{2} \int d^4 x \, a^4(t) \, p \, h_{ij} \, \delta h_{ij}$$

$$\delta S_{\rm m}^{(2)} + \delta S_{\rm HE}^{(2)} = 0 = \int d^4 x \, a^2 \Big[-\frac{m_p^2}{4} \Big(h_{ij}'' + 2\mathcal{H} h_{ij}' - \nabla^2 h_{ij} \Big) + \frac{1}{2} \Pi_{ij}^{({\rm T})} \Big] \delta h_{ij}$$

Then:
$$\tilde{R}_{\mu\nu} \equiv \tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu}$$

[Friedmann [GW Eq. ?

$$\delta S_{\text{HE}}^{(2)} = -\frac{m_p^2}{4} \int d^4 x \, a^2 \Big(h_{ij}'' + 2 \mathscr{R} h_{ij}' - \nabla^2 h_{ij} \Big) \delta h_{ij} + \frac{1}{2} \int d^4 x \, a^4 p \, h_{ij} \delta h_{ij}$$

$$\delta S_{\text{m}}^{(2)} = \frac{1}{2} \int d^4 x \, a^2(t) \, \Pi_{ij}^{(\text{T})} \, \delta h_{ij} - \frac{1}{2} \int d^4 x \, a^4(t) \, p \, h_{ij} \, \delta h_{ij}$$

$$\delta S_{\text{m}}^{(2)} + \delta S_{\text{HE}}^{(2)} = -\frac{m_p^2}{4} \int d^4 x \, a^2 \Big(\frac{h_{ij}'' + 2 \mathscr{R} h_{ij}' - \nabla^2 h_{ij}}{\text{wave operator}} - \frac{2}{m_p^2} \Pi_{ij}^{(\text{T})} \Big) \delta h_{ij} = 0$$
Source

Then:
$$\tilde{R}_{\mu\nu} \equiv \tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu}$$

[Friedmann [GW Eq. ?

$$\delta S_{\text{HE}}^{(2)} \equiv -\frac{m_p^2}{4} \int d^4x \, a^2 \Big(h_{ij}'' + 2\mathcal{H} h_{ij}' - \nabla^2 h_{ij} \Big) \delta h_{ij} + \frac{1}{2} \int d^4x \, a^4 p \, h_{ij} \delta h_{ij}$$

$$\delta S_{\text{m}}^{(2)} = \frac{1}{2} \int d^4x \, a^2(t) \, \Pi_{ij}^{(\text{T})} \, \delta h_{ij} - \frac{1}{2} \int d^4x \, a^4(t) \, p \, h_{ij} \, \delta h_{ij}$$

$$\delta S_{\text{m}}^{(2)} + \delta S_{\text{HE}}^{(2)} = -\frac{m_p^2}{4} \int d^4x \, a^2 \Big(h_{ij}'' + 2\mathcal{H} h_{ij}' - \nabla^2 h_{ij} - \frac{2}{m_p^2} \Pi_{ij}^{(\text{T})} \Big) \delta h_{ij} = 0$$
Correct Eq. of motion !

Then:
$$\tilde{R}_{\mu\nu} \equiv \tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu} + \tilde{R}_{\mu\nu}$$

[Friedmann [GW Eq. ?



Then:
$$\tilde{R}_{\mu\nu} \equiv \overset{(0)}{\tilde{R}}_{\mu\nu} + \overset{(1)}{\tilde{R}}_{\mu\nu} + \overset{(2)}{\tilde{R}}_{\mu\nu}$$

[Friedmann [GW Eq. $\mathcal{O}(h_{**}^2) \rightarrow$ GW's Energy-momentum ?



Then:
$$\tilde{R}_{\mu\nu} \equiv \begin{pmatrix} 0 \\ R \\ \mu\nu \end{pmatrix} + \begin{pmatrix} 1 \\ R \\ \mu\nu \end{pmatrix} + \begin{pmatrix} 2 \\ R \\ \mu\nu \end{pmatrix}$$

[Friedmann [GW Eq. $\mathcal{O}(h_{**}^2) \rightarrow \text{GW's Energy-momentum }?$
 $(2) = (2) + (2)$
 $S_{\text{tot}} \equiv S_{\text{m}} + S_{\text{HE}}$

Then:
$$\tilde{R}_{\mu\nu} \equiv \begin{pmatrix} 0 \\ \tilde{R}_{\mu\nu} \end{pmatrix} + \begin{pmatrix} 1 \\ \tilde{R}_{\mu\nu} \end{pmatrix} + \begin{pmatrix} 2 \\ \tilde{R}_{\mu\nu} \end{pmatrix}$$

[Friedmann [GW Eq. $\mathcal{O}(h_{**}^2) \rightarrow \mathbf{GW's Energy-momentum ?}$
 $S_{tot}^{(2)} \equiv S_m^{(2)} + S_{HE}^{(2)}$
 $= -\frac{m_p^2}{4} \int d^4x \sqrt{-g} \left[-\frac{1}{2} g^{\mu\nu} \partial_{\mu} h_{ij} \partial_{\nu} h_{ij} + 4 \mathcal{H} h_{ij} g^{0\mu} \partial_{\mu} h_{ij} + \frac{1}{a^2} \left(\mathcal{H}^2 + \frac{a''}{a} \right) h_{ij} h_{ij} - \frac{2}{a^2 m_p^2} h_{ij} \Pi_{ij}^{(T)} \right]$

Then:
$$\tilde{R}_{\mu\nu} \equiv \begin{pmatrix} 0 \\ \tilde{R}_{\mu\nu} \end{pmatrix} + \begin{pmatrix} 1 \\ \tilde{R}_{\mu\nu} \end{pmatrix} + \begin{pmatrix} 2 \\ \tilde{R}_{\mu\nu} \end{pmatrix}$$

[Friedmann [GW Eq. $\mathcal{O}(h_{**}^2) \rightarrow$ GW's Energy-momentum?
 $\int_{tot}^{(2)} = \int d^4x \sqrt{-g} \mathcal{L}(h_{ij}, \partial_{\mu}h_{ij}) \quad ; \quad g_{\mu\nu} \equiv a^2(t)\eta_{\mu\nu}$ [FLRW]

Then:
$$\tilde{R}_{\mu\nu} \equiv \begin{pmatrix} 0 \\ R \\ \mu\nu \end{pmatrix} + \begin{pmatrix} 1 \\ R \\ \mu\nu \end{pmatrix} + \begin{pmatrix} 2 \\ R \\ \mu\nu \end{pmatrix}$$

[Friedmann [GW Eq. $O(h_{**}^2) \rightarrow GW$'s Energy-momentum?
 $\int_{tot}^{(2)} = \int d^4x \sqrt{-g} \mathcal{L}(h_{ij}, \partial_{\mu}h_{ij}) \quad ; \quad g_{\mu\nu} \equiv a^2(t)\eta_{\mu\nu}$ [FLRW]
Noether's Theorem: $T_{\mu\nu} \equiv -\frac{\partial \mathcal{L}}{\partial(\partial^{\mu}h_{ij})}\partial_{\nu}h_{ij} + g_{\mu\nu}\mathcal{L}$

Then:
$$\tilde{R}_{\mu\nu} \equiv \begin{pmatrix} 0 \\ R \\ \mu\nu \end{pmatrix} + \begin{pmatrix} 1 \\ R \\ \mu\nu \end{pmatrix} + \begin{pmatrix} 2 \\ R \\ \mu\nu \end{pmatrix}$$

[Friedmann [GW Eq. $\mathcal{O}(h_{**}^2) \rightarrow GW$'s Energy-momentum?
 $\int_{tot}^{(2)} \equiv \int d^4x \sqrt{-g} \mathcal{L}(h_{ij}, \partial_{\mu}h_{ij}) \quad ; \quad g_{\mu\nu} \equiv a^2(t)\eta_{\mu\nu} \quad [FLRW]$
Noether's Theorem: $T_{\mu\nu} \equiv -\frac{\partial \mathcal{L}}{\partial(\partial^{\mu}h_{ij})}\partial_{\nu}h_{ij} + g_{\mu\nu}\mathcal{L} + \partial_{\lambda}f^{\lambda\mu\nu}$

Then:
$$\tilde{R}_{\mu\nu} \equiv \begin{pmatrix} 0 \\ R \\ \mu\nu \end{pmatrix} + \begin{pmatrix} 1 \\ R \\ \mu\nu \end{pmatrix} + \begin{pmatrix} 2 \\ R \\ \mu\nu \end{pmatrix}$$

[Friedmann [GW Eq. $O(h_{**}^2) \rightarrow GW$'s Energy-momentum ?
 $\int_{\text{Equations}}^{(2)} \left[\int d^4x \sqrt{-g} \mathcal{L}(h_{ij}, \partial_{\mu}h_{ij}) \right]; g_{\mu\nu} \equiv a^2(t)\eta_{\mu\nu}$ [FLRW]
Noether's Theorem: $T_{\mu\nu} \equiv \left\langle -\frac{\partial \mathcal{L}}{\partial(\partial^{\mu}h_{ij})} \partial_{\nu}h_{ij} + g_{\mu\nu}\mathcal{L} + \partial_{\lambda}f^{\lambda\mu\nu} \right\rangle$

Then:
$$\tilde{R}_{\mu\nu} \equiv \begin{pmatrix} 0 \\ R \\ \mu\nu \end{pmatrix} + \begin{pmatrix} 1 \\ R \\ \mu\nu \end{pmatrix} + \begin{pmatrix} 2 \\ R \\ \mu\nu \end{pmatrix}$$

[Friedmann [GW Eq. $O(h_{**}^2) \rightarrow GW$'s Energy-momentum ?
 $\int_{\text{Equations}}^{(2)} \left[\int d^4x \sqrt{-g} \mathcal{L}(h_{ij}, \partial_{\mu}h_{ij}) \right]; g_{\mu\nu} \equiv a^2(t)\eta_{\mu\nu}$ [FLRW]
Noether's Theorem: $T_{\mu\nu} \equiv \left\langle -\frac{\partial \mathcal{L}}{\partial(\partial^{\mu}h_{ij})} \partial_{\nu}h_{ij} + g_{\mu\nu}\mathcal{L} + \partial_{\nu}\mathcal{L}^{\mu\nu} \right\rangle$

Then:
$$\tilde{R}_{\mu\nu} \equiv \begin{pmatrix} 0 \\ \tilde{R}_{\mu\nu} \end{pmatrix} + \begin{pmatrix} 1 \\ \tilde{R}_{\mu\nu} \end{pmatrix} + \begin{pmatrix} 2 \\ \tilde{R}_{\mu\nu} \end{pmatrix}$$

[Friedmann [GW Eq. $\mathcal{O}(h_{**}^2) \rightarrow$ GW's Energy-momentum ?
 $\int_{\text{Equations]}}^{(2)} \int d^4x \sqrt{-g} \mathcal{L}(h_{ij}, \partial_{\mu}h_{ij}) \quad ; \quad g_{\mu\nu} \equiv a^2(t)\eta_{\mu\nu} \quad \text{[FLRW]}$
Noether's Theorem: $\bar{T}_{\mu\nu} \equiv \left\langle -\frac{\partial \mathcal{L}}{\partial(\partial^{\mu}h_{ij})} \partial_{\nu}h_{ij} + g_{\mu\nu}\mathcal{L} \right\rangle$

[Volume averaging over $V \gg \lambda^3$]

Then:
$$\tilde{R}_{\mu\nu} \equiv \begin{pmatrix} 0 \\ \mu\nu \end{pmatrix} + \begin{pmatrix} 1 \\ \mu\nu \end{pmatrix} + \begin{pmatrix} 2 \\ \mu\nu \end{pmatrix}$$

[Friedmann [GW Eq. $O(h_{**}^2) \rightarrow GW$'s Energy-momentum?
 $\int_{tot}^{(2)} \equiv \int d^4x \sqrt{-g} \mathscr{L}(h_{ij}, \partial_{\mu}h_{ij}) \quad ; \quad g_{\mu\nu} \equiv a^2(t)\eta_{\mu\nu}$ [FLRW]
Noether's Theorem: $\bar{T}_{\mu\nu} \equiv \left\langle -\frac{\partial\mathscr{L}}{\partial(\partial^{\mu}h_{ij})} \partial_{\nu}h_{ij} + g_{\mu\nu}\mathscr{L} \right\rangle$
 $\rho_{GW} = a^{-2}\bar{T}_{00} \equiv \left\langle \frac{m_p^2}{4} \left(\frac{1}{2a^2}(h_{ij}')^2 + \frac{1}{2a^2}(\nabla h_{ij})^2 + \left(\mathscr{H}^2 + \frac{a''}{a}\right)h_{ij}^2 \right) - \frac{1}{2a^2}\Pi_{ij}^{(T)}h_{ij} \right\rangle$

Then:
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[Friedmann [GW Eq. $\mathcal{O}(h_{**}^2) \rightarrow GW$'s Energy-momentum ?
 $\begin{pmatrix} 2 \\ S_{tot} \end{bmatrix} = \int d^4x \sqrt{-g} \mathcal{L}(h_{ij}, \partial_{\mu}h_{ij}) \quad ; \quad g_{\mu\nu} \equiv a^2(t)\eta_{\mu\nu} \quad [FLRW]$
Noether's Theorem: $\bar{T}_{\mu\nu} \equiv \left\langle -\frac{\partial \mathcal{L}}{\partial(\partial^{\mu}h_{ij})} \partial_{\nu}h_{ij} + g_{\mu\nu}\mathcal{L} \right\rangle$
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Kinetic Gradient Interaction

Then:
$$\tilde{R}_{\mu\nu} \equiv \begin{pmatrix} 0 \\ R \\ \mu\nu \end{pmatrix} + \begin{pmatrix} 1 \\ R \\ \mu\nu \end{pmatrix} + \begin{pmatrix} 2 \\ R \\ \mu\nu \end{pmatrix}$$

[Friedmann [GW Eq. $\mathcal{O}(h_{**}^2) \rightarrow \text{GW's Energy-momentum ?}$
 $\int_{\text{formations}}^{(2)} \left[\int d^4x \sqrt{-g} \mathcal{L}(h_{ij}, \partial_{\mu}h_{ij}) \right] ; g_{\mu\nu} \equiv a^2(t)\eta_{\mu\nu}$ [FLRW]
Noether's Theorem: $\bar{T}_{\mu\nu} \equiv \left\langle -\frac{\partial \mathcal{L}}{\partial(\partial^{\mu}h_{ij})} \partial_{\nu}h_{ij} + g_{\mu\nu}\mathcal{L} \right\rangle$
 $\rho_{\text{GW}} = a^{-2}\bar{T}_{00} \equiv \left\langle \frac{m_p^2}{4} \left(\frac{1}{2a^2}(h_{ij}')^2 + \frac{1}{2a^2}(\nabla h_{ij})^2 + \left(\mathcal{H}^2 + \frac{a''}{a} \right)h_{ij}^2 \right) - \frac{1}{2a^2}\Pi_{ij}^{(T)}h_{ij} \right\rangle$
Kinetic Gradient Caution! differs in the literature

Then:
$$\tilde{R}_{\mu\nu} \equiv \begin{pmatrix} 0 \\ R \\ \mu\nu \end{pmatrix} + \begin{pmatrix} 1 \\ R \\ \mu\nu \end{pmatrix} + \tilde{R}_{\mu\nu}$$

[Friedmann [GW Eq. $\mathcal{O}(h_{***}^2) \rightarrow$ GW's Energy-momentum ?

$$\rho_{GW} = \left\langle \frac{m_p^2}{4} \left(\frac{1}{2a^2} (h_{ij}')^2 + \frac{1}{2a^2} (\nabla h_{ij})^2 + \left(\mathscr{H}^2 + \frac{a''}{a} \right) h_{ij}^2 \right) - \frac{1}{2a^2} \Pi_{ij}^{(T)} h_{ij} \right\rangle$$

Then:
$$\tilde{R}_{\mu\nu} \equiv \overset{(0)}{\tilde{R}}_{\mu\nu} + \overset{(1)}{\tilde{R}}_{\mu\nu} + \overset{(2)}{\tilde{R}}_{\mu\nu}$$

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$$\rho_{\rm GW} = \left\langle \frac{m_p^2}{4} \left(\frac{1}{2a^2} (h_{ij}')^2 + \frac{1}{2a^2} (\nabla h_{ij})^2 + \left(\mathcal{H}^2 + \frac{a''}{a} \right) h_{ij}^2 \right) - \frac{1}{2a^2} \Pi_{ij}^{(\rm T)} h_{ij} \right\rangle$$

Sub-horizon : $(k \gg \mathcal{H})$

 $\sim k^2 h^2 \gg \sim \mathcal{H}^2 h^2$



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 $\rho_{\mathrm{GW}} = \left\langle \frac{m_p^2}{4} \left(\frac{1}{2a^2} (h_{ij}')^2 + \frac{1}{2a^2} (\nabla h_{ij})^2 + \left(\mathcal{H}^2 + \frac{a''}{a} \right) h_{ij}^2 \right) - \frac{1}{2a^2} \mathbf{H}_{ij}^{(1)} h_{ij} \right\rangle$
Sub-horizon : $\sim k^2 h^2 \gg \sim \mathcal{H}^2 h^2$

 $\Pi_{ij} \rightarrow 0$

Free fields : (after emission)

Then:
$$\tilde{R}_{\mu\nu} \equiv \begin{pmatrix} 0 \\ \tilde{R}_{\mu\nu} \end{pmatrix} + \begin{pmatrix} 1 \\ \tilde{R}_{\mu\nu} \end{pmatrix} + \begin{pmatrix} 2 \\ \tilde{R}_{\mu\nu} \end{pmatrix}$$

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 $\rho_{GW} = \left\langle \frac{m_p^2}{4} \left(\frac{1}{2a^2} (h_{ij}')^2 + \frac{1}{2a^2} (\nabla h_{ij})^2 + \left(\mathcal{H}^2 + \frac{a''}{a} \right) h_{ij}^2 \right) - \frac{1}{2a^2} \prod_{ij}^{(1)} h_{ij} \right\rangle$
Sub-horizon : $\sim k^2 h^2 \gg \sim \mathcal{H}^2 h^2$
 $(k \gg \mathcal{H})$
Free fields : $\frac{1}{2a^2} (h_{ij}')^2 = \frac{1}{2a^2} (\nabla h_{ij})^2$
 $\Pi_{ij} \rightarrow 0$
Then:
$$\tilde{R}_{\mu\nu} \equiv \begin{pmatrix} 0 \\ \tilde{R}_{\mu\nu} \end{pmatrix} + \begin{pmatrix} 1 \\ \tilde{R}_{\mu\nu} \end{pmatrix} + \begin{pmatrix} 2 \\ \tilde{R}_{\mu\nu} \end{pmatrix}$$

[Friedmann [GW Eq. $\mathcal{O}(h_{**}^2) \rightarrow GW$'s Energy-momentum ?
 $\mathcal{O}(m_{**}^2) \rightarrow GW$'s Energy density
 $\mathcal{O}_{GW} = \left\langle \frac{m_p^2}{4} \left(\frac{1}{2a^2} (h_{ij}')^2 + \frac{1}{2a^2} (\nabla h_{ij})^2 \right) \right\rangle$
Energy density
carried by
Grav. Waves
Sub-horizon & Free fields
 $(k \gg \mathcal{H})$ (after emission)
[Volume averaging over $V \gg \lambda^3$]

Then:
$$\tilde{R}_{\mu\nu} \equiv \begin{pmatrix} 0 \\ \tilde{R}_{\mu\nu} \end{pmatrix} + \begin{pmatrix} 1 \\ \tilde{R}_{\mu\nu} \end{pmatrix} + \begin{pmatrix} 2 \\ \tilde{R}_{\mu\nu} \end{pmatrix}$$

[Friedmann [GW Eq. $\mathcal{O}(h_{**}^2) \rightarrow GW$'s Energy-momentum ?

$$\int \rho_{GW} = \left\langle \frac{m_p^2}{4} \left(\frac{1}{2a^2} (h_{ij}')^2 + \frac{1}{2a^2} (\nabla h_{ij})^2 \right) \right\rangle$$
Energy density carried by Grav. Waves
Sub-horizon & Free fields
 $(k \gg \mathcal{H})$ (after emission)
[Volume averaging over $V \gg \lambda^3$]

Then:
$$\tilde{R}_{\mu\nu} \equiv \begin{pmatrix} 0 \\ \tilde{R}_{\mu\nu} \end{pmatrix} + \begin{pmatrix} 1 \\ \tilde{R}_{\mu\nu} \end{pmatrix} + \begin{pmatrix} 2 \\ \tilde{R}_{\mu\nu} \end{pmatrix}$$

[Friedmann [GW Eq. $\mathcal{O}(h_{**}^2) \rightarrow \mathbf{GW}$'s Energy-momentum ?
 $\mathcal{P}_{\mathrm{GW}} = \frac{m_p^2}{4a^2} \left\langle (h_{ij}')^2 \right\rangle$
Energy density carried by Grav. Waves
Sub-horizon & Free fields
 $(k \gg \mathcal{H})$ (after emission)
[Volume averaging over $V \gg \lambda^3$]

Then:
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[Friedmann Equations] [GW Eq. $\mathcal{O}(h_{**}^2) \rightarrow$ GW's Energy-momentum ?
Energy density carried by
Gravitational Waves $\rho_{GW} = \frac{m_p^2}{4a^2} \langle h'_{ij} h'_{ij} \rangle_{V \gg \lambda^3}$
Sub-horizon & Free fields (after emission)







Definition of GWs 4th approach

4th approach to GWs

(for a curved space-time)

 $g_{\mu\nu}(x) = \bar{g}_{\mu\nu}(x) + \delta g_{\mu\nu}(x), \quad |\delta g_{\mu\nu}| \ll 1$ (separation not well defined)

4th approach to GWs

(for a curved space-time)

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4th approach to GWs

(for a curved space-time)

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(for a curved space-time)

 $g_{\mu\nu}(x) = \bar{g}_{\mu\nu}(x) + \delta g_{\mu\nu}(x), \quad |h_{\mu\nu}| \ll 1$ (separation not well defined)

$$R_{\mu\nu} = \frac{1}{m_p^2} \left(T_{\mu\nu} - \frac{1}{2} g_{\mu\nu} T \right)$$

4th approach to GWs

(for a curved space-time)

 $g_{\mu\nu}(x) = \bar{g}_{\mu\nu}(x) + \delta g_{\mu\nu}(x), \quad |h_{\mu\nu}| \ll 1$ (separation not well defined)

$$R_{\mu\nu} = \frac{1}{m_p^2} \left(T_{\mu\nu} - \frac{1}{2} g_{\mu\nu} T \right) \longrightarrow R_{\mu\nu} = \bar{R}_{\mu\nu} + R_{\mu\nu}^{(1)} + R_{\mu\nu}^{(2)} + \dots ,$$

(background) $\mathcal{O}(\delta g) \quad \mathcal{O}(\delta g^2)$

4th approach to GWs

(for a curved space-time)

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More subtle problem! <u>Solution</u>: Separation of scales !

$$R_{\mu\nu} = \frac{1}{m_p^2} \left(T_{\mu\nu} - \frac{1}{2} g_{\mu\nu} T \right) \longrightarrow R_{\mu\nu} = \bar{R}_{\mu\nu} + R_{\mu\nu}^{(1)} + R_{\mu\nu}^{(2)} + \dots,$$

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Low Freq. / Long Scale: $\bar{R}_{\mu\nu} = -[R^{(2)}_{\mu\nu}]^{\text{Low}} + \frac{1}{m_p^2} \left(T_{\mu\nu} - \frac{1}{2}g_{\mu\nu}T\right)^{\text{Low}}$

4th approach to GWs

(for a curved space-time)

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High Freq. / Short Scale: $R_{\mu\nu}^{(1)} = -[R_{\mu\nu}^{(2)}]^{\text{High}} + \frac{1}{m_p^2} \left(T_{\mu\nu} - \frac{1}{2}g_{\mu\nu}T\right)^{\text{High}}$

4th approach to GWs

(for a curved space-time)

 $g_{\mu\nu}(x) = \bar{g}_{\mu\nu}(x) + \delta g_{\mu\nu}(x), \quad |h_{\mu\nu}| \ll 1$ (separation not well defined)

Low Freq. / Long Scale:
$$\bar{R}_{\mu\nu} = -[R^{(2)}_{\mu\nu}]^{\text{Low}} + \frac{1}{m_p^2} \left(T_{\mu\nu} - \frac{1}{2}g_{\mu\nu}T\right)^{\text{Low}}$$

Low Freq. / Long Scale: $\bar{R}_{\mu\nu} = -[R^{(2)}_{\mu\nu}]^{\text{Low}} + \frac{1}{m_p^2} \left(T_{\mu\nu} - \frac{1}{2}g_{\mu\nu}T\right)^{\text{Low}}$



Low Freq. / Long Scale: $\bar{R}_{\mu\nu} = -\langle R^{(2)}_{\mu\nu} \rangle + \frac{1}{m_p^2} \langle T_{\mu\nu} - \frac{1}{2}g_{\mu\nu}T \rangle$ average)



Low Freq. / Long Scale: $\bar{R}_{\mu\nu} = -\langle R^{(2)}_{\mu\nu} \rangle + \frac{1}{m_p^2} \langle T_{\mu\nu} - \frac{1}{2}g_{\mu\nu}T \rangle$ average)







Low Freq. / Long Scale:

$$\bar{R}_{\mu\nu} = \frac{1}{m_p^2} \left(t_{\mu\nu} - \frac{1}{2} \bar{g}_{\mu\nu} t \right) + \frac{1}{m_p^2} \left(\bar{T}_{\mu\nu} - \frac{1}{2} \bar{g}_{\mu\nu} \bar{T} \right)$$

$$t_{\mu\nu} = -\frac{1}{m_p^2} \left\langle R_{\mu\nu}^{(2)} - \frac{1}{2} \bar{g}_{\mu\nu} R^{(2)} \right\rangle \qquad \langle T_{\mu\nu} - \frac{1}{2} g_{\mu\nu} T \rangle = \bar{T}^{\mu\nu} - \frac{1}{2} \bar{g}_{\mu\nu} \bar{T}$$



Low Freq. / Long Scale:

$$\bar{R}_{\mu\nu} = \frac{1}{m_p^2} \left(t_{\mu\nu} - \frac{1}{2} \bar{g}_{\mu\nu} t \right) + \frac{1}{m_p^2} \left(\bar{T}_{\mu\nu} - \frac{1}{2} \bar{g}_{\mu\nu} \bar{T} \right)$$

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$$\langle R^{(2)}_{\mu\nu} \rangle = -\frac{1}{4} \langle \partial_{\mu} \delta g_{\alpha\beta} \, \partial_{\nu} \delta g^{\alpha\beta} \rangle \longrightarrow t_{\mu\nu} = \frac{m_p^2}{4} \langle \partial_{\mu} \delta g_{\alpha\beta} \, \partial_{\nu} \delta g^{\alpha\beta} \rangle$$

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$$\bar{R}_{\mu\nu} = \frac{1}{m_p^2} \left(t_{\mu\nu} - \frac{1}{2} \bar{g}_{\mu\nu} t \right) + \frac{1}{m_p^2} \left(\bar{T}_{\mu\nu} - \frac{1}{2} \bar{g}_{\mu\nu} \bar{T} \right)$$

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It can be shown that only TT *dof* contribute to < ... >

Low Freq. / Long Scale:
$$\bar{R}_{\mu\nu} = \frac{1}{m_p^2} \left(t_{\mu\nu} - \frac{1}{2} \bar{g}_{\mu\nu} t \right) + \frac{1}{m_p^2} \left(\bar{T}_{\mu\nu} - \frac{1}{2} \bar{g}_{\mu\nu} \bar{T} \right)$$
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It can be shown that only TT *dof* contribute to < ... >

$$t_{\mu\nu} = \frac{m_p^2}{4} \left\langle \partial_\mu \delta g_{ij}^{\rm TT} \, \partial_\nu \delta g_{ij}^{\rm TT} \right\rangle$$

GW energy-momentum tensor

Low Freq. / Long Scale:

$$\bar{R}_{\mu\nu} = \frac{1}{m_p^2} \left(t_{\mu\nu} - \frac{1}{2} \bar{g}_{\mu\nu} t \right) + \frac{1}{m_p^2} \left(\bar{T}_{\mu\nu} - \frac{1}{2} \bar{g}_{\mu\nu} \bar{T} \right)$$

$$t_{\mu\nu} = -\frac{1}{m_p^2} \left\langle R_{\mu\nu}^{(2)} - \frac{1}{2} \bar{g}_{\mu\nu} R^{(2)} \right\rangle \qquad \langle T_{\mu\nu} - \frac{1}{2} g_{\mu\nu} T \rangle = \bar{T}^{\mu\nu} - \frac{1}{2} \bar{g}_{\mu\nu} \bar{T}$$

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It can be shown that only TT *dof* contribute to < ... >

GW energy-momentum tensor

GW energy density

What about the High Freq. / Short Scale?

$$R_{\mu\nu}^{(1)} = -[R_{\mu\nu}^{(2)}]^{\text{High}} + \frac{1}{m_p^2} \left(T_{\mu\nu} - \frac{1}{2} g_{\mu\nu} T \right)^{\text{High}}$$

What about the
High Freq. / Short Scale?
$$R^{(1)}_{\mu\nu} = -[R^{(2)}_{\mu\nu}]^{\text{High}} + \frac{1}{m_p^2} \left(T_{\mu\nu} - \frac{1}{2}g_{\mu\nu}T\right)^{\text{High}}$$

$$\frac{|R_{\mu}^{(2)}|^{\text{High}}}{|R_{\mu}^{(1)}|} \sim \mathcal{O}\left(\frac{\lambda_{*}}{L_{B}}\right) \longrightarrow |R_{\mu}^{(2)}|^{\text{High}} \text{ negligible}$$

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$$R^{(1)}_{\mu\nu} = -[R^{(2)}_{\mu\nu}]^{\text{High}} + \frac{1}{m_p^2} \left(T_{\mu\nu} - \frac{1}{2}g_{\mu\nu}T\right)^{\text{High}}$$

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(-)

$$\begin{split} R^{(1)}_{\mu\nu} &= \bar{g}^{\alpha\beta} \left(D_{\alpha} D_{(\mu} \delta g_{\nu)\beta} - D_{\mu} D_{\nu} \delta g_{\alpha\beta} - D_{\alpha} D_{\beta} \delta g_{\mu\nu} \right) \\ D_{\mu} \overline{\delta g}_{\mu\nu} &= 0 \quad (\overline{\delta g}_{\mu\nu} \equiv \delta g_{\mu\nu} - \frac{1}{2} \bar{g}_{\mu\nu} \bar{g}^{\alpha\beta} \delta g_{\alpha\beta}) \quad \underset{\substack{\text{lorentz} \\ \text{gauge}}}{\text{lorentz}} \end{split}$$

What about the High Freq. / Short Scale?
$$R^{(1)}_{\mu\nu} = -[R^{(2)}_{\mu\nu}]^{\text{High}} + \frac{1}{m_p^2} \left(T_{\mu\nu} - \frac{1}{2}g_{\mu\nu}T\right)^{\text{High}}$$

$$\frac{|R_{\mu}^{(2)}|^{\text{High}}}{|R_{\mu}^{(1)}|} \sim \mathcal{O}\left(\frac{\lambda_{*}}{L_{B}}\right) \longrightarrow |R_{\mu}^{(2)}|^{\text{High}} \text{ negligible}$$

$$\begin{split} R^{(1)}_{\mu\nu} &= \bar{g}^{\alpha\beta} \left(D_{\alpha} D_{(\mu} \delta g_{\nu)\beta} - D_{\mu} D_{\nu} \delta g_{\alpha\beta} - D_{\alpha} D_{\beta} \delta g_{\mu\nu} \right) \\ D_{\mu} \overline{\delta g}_{\mu\nu} &= 0 \quad (\ \overline{\delta g}_{\mu\nu} \equiv \delta g_{\mu\nu} - \frac{1}{2} \bar{g}_{\mu\nu} \bar{g}^{\alpha\beta} \delta g_{\alpha\beta} \) \qquad \begin{array}{c} \text{Lorentz} \\ \text{gauge} \end{array}$$

vacuum $D_{\alpha}D^{\alpha}\overline{\delta g}_{\mu\nu} = 0$ Propagation of GWs in curved space-time

What about the High Freq. / Short Scale?
$$R^{(1)}_{\mu\nu} = -[R^{(2)}_{\mu\nu}]^{\text{High}} + \frac{1}{m_p^2} \left(T_{\mu\nu} - \frac{1}{2}g_{\mu\nu}T\right)^{\text{High}}$$

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vacuum $D_{\alpha}D^{\alpha}\delta g_{ij}^{\mathrm{TT}} = 0$ Propagation of GWs in curved space-time ($D_{i}\delta g_{ij}^{\mathrm{TT}} = \bar{g}^{ij}\delta g_{ij}^{\mathrm{TT}} = 0$)

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Creation of GWs in curved space-time TT dof = truly radiative ! [no gauge choice]

GW Propagation/Creation in Cosmology

FLRW: $ds^2 = a^2(-d\eta^2 + (\delta_{ij} + h_{ij})dx^i dx^j), \quad \text{TT}: \begin{cases} h_{ii} = 0 \\ h_{ij}, j = 0 \end{cases}$

GW Propagation/Creation in Cosmology

FLRW:
$$ds^2 = a^2(-d\eta^2 + (\delta_{ij} + h_{ij})dx^i dx^j),$$
 TT: $\begin{cases} h_{ii} = 0 \\ h_{ij,j} = 0 \end{cases}$

Creation of GWs in curved space-time

Source: Anisotropic Stress

Eom:
$$h_{ij}'' + 2\mathcal{H}h_{ij}' - \nabla^2 h_{ij} = 16\pi G \Pi_{ij}^{\mathrm{TT}}$$

$$\Pi_{ij} = T_{ij} - \left\langle T_{ij} \right\rangle_{\rm FRW}$$
GW Propagation/Creation in Cosmology

FLRW:
$$ds^2 = a^2(-d\eta^2 + (\delta_{ij} + h_{ij})dx^i dx^j), \quad \text{TT}: \begin{cases} h_{ii} = 0 \\ h_{ij}, j = 0 \end{cases}$$

Creation of GWs in curved space-time

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.

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GW Source(s): (SCALARS , VECTOR , FERMIONS) $\Pi_{ij}^{TT} \propto \{\partial_i \chi^a \partial_j \chi^a\}^{TT}, \{E_i E_j + B_i B_j\}^{TT}, \{\bar{\psi} \gamma_i D_j \psi\}^{TT}$

Cosmic History

BiGGER size, **SMALLER Temp**



TODAY [Galaxies, Clusters, ...] (13.700 Million years)

FIRST GALAXIES (500 Millions years)

ATOMS CREATION (300.000-400.000 years)

ATOMIC NUCLEI CREATION (3 minutes !)

FIRST SECOND of the UNIVERSE !

SMALLER SIZE, LARGER Temperature

GWs: probe of the early Universe

WEAKNESS of GRAVITY:

ADVANTAGE: GW DECOUPLE upon Production **DISADVANTAGE: DIFFICULT DETECTION**

O ADVANTAGE: GW \rightarrow Probe for Early Universe $\rightarrow \left\{ \begin{array}{l} \mathbf{Decouple} \rightarrow \mathrm{Spectral} \ \mathrm{Form} \ \mathrm{Retained} \\ \mathbf{Specific} \ \mathbf{HEP} \ \Leftrightarrow \ \mathrm{Specific} \ \mathrm{GW} \end{array} \right.$

Physical Processes: Inflation
Reheating
Phase Transitions
Cosmic Defects

The Early Universe



GWs: probe of the early Universe

OUTLINE



2) GWs from Inflation

Early Universe 3) GWs from Preheating

4) GWs from Phase Transitions

5) GWs from Cosmic Defects

