

# Conformal Field Theory Problems

## Chapter 1

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**Problem 1.1.** *Using only the quantities  $\hbar$ ,  $G_N$ , and  $c$ , construct quantities that have the units of length, mass, and time. Compute the corresponding Planck length, Planck mass, and Planck time, using SI units.*

In SI units

$$\begin{aligned}G_N &= 6.67 \times 10^{-11} \text{ kg}^{-1} \text{ m}^3 \text{ s}^{-2} , \\ \hbar &= 1.05 \times 10^{-34} \text{ kg m}^2 \text{ s}^{-1} , \\ c &= 3.00 \times 10^8 \text{ m s}^{-1} .\end{aligned}$$

The quantities requested are then

$$\begin{aligned}\ell_P &= \sqrt{\frac{G_N \hbar}{c^3}} = 1.61 \times 10^{-35} \text{ m} , \\ m_P &= \sqrt{\frac{\hbar c}{G_N}} = 2.17 \times 10^{-8} \text{ kg} , \\ t_P &= \sqrt{\frac{G_N \hbar}{c^5}} = 5.37 \times 10^{-44} \text{ s} .\end{aligned}$$

**Problem 1.2.** *Another proposed source of extra physics is extra dimensions. Assume that we live not in a four dimensional world but a  $(4+p)$ -dimensional one where the extra dimensions are all extremely small circles of length  $\ell$ .*

- a) *Noting that the dimensionality of  $G_N$  is different in  $(4+p)$  dimensions, what is the new expression for the Planck energy  $E_P$  in terms of  $\hbar$ ,  $c$ , and  $G_N$ ?*
- b) *Find a relationship between  $G_N$  and the observed 4d value  $G_N^{4d}$ . Given the observed 4d value for  $G_N^{4d}$ , how small must  $\ell$  be in order to have  $E_P = 1 \text{ TeV}$ ? Are there some values of  $p$  that you can rule out?*

- a) *Coulomb's law is modified in  $p+4$  dimensions. The potential energy stored by two point masses is  $\frac{G_N m_1 m_2}{r^{1+p}}$  which means that Newton's constant has units  $[G_N] = \text{kg}^{-1} \text{m}^{3+p} \text{s}^{-2}$ , and the Planck energy should be instead*

$$E_P = \left( \frac{\hbar^{p+1} c^{p+5}}{G_N} \right)^{\frac{1}{p+2}} .$$

b) The 4d Newton's constant should be  $G_N^{4d} \sim G_N/\ell^p$ . In order for  $G_N^{4d}$  to take its usual value, we find that

$$\ell \sim \frac{\hbar^{1+\frac{1}{p}} c^{1+\frac{5}{p}}}{E_P^{1+\frac{2}{p}}} (G_N^{4d})^{-\frac{1}{p}}$$

For numbers of extra dimensions one through six,

$$\ell \sim 3 \times 10^{13} \text{ m}, 2 \text{ mm}, 1 \times 10^{-8} \text{ m}, 2.2 \times 10^{-11} \text{ m}, 5.3 \times 10^{-13} \text{ m}, 4.5 \times 10^{-14} \text{ m} .$$

It looks like we can rule out  $p = 1$  and  $2$  from this back of the envelope calculation. Note superstring theory predicts  $p = 6$ .

**Problem 1.3.** Consider an interacting scalar field

$$S = - \int d^d x ((\partial_\mu \phi)(\partial^\mu \phi) + g\phi^n) . \quad (1.1)$$

where  $n$  is a positive integer. For what pairs  $(n, d)$  can the coupling  $g$  be dimensionless?

The model  $d = 3$  and  $\phi^6$  is claimed to be in the same universality class as the tricritical Ising model. The model  $d = 4$  and  $\phi^4$  is in the same universality class as the Ising model. It is sometimes called the Wilson-Fisher theory or Landau-Ginzburg theory. Then there is  $d = 6$  and  $\phi^3$ .

**Problem 1.4.** Consider the Lagrangian for a Dirac spinor in  $d$  dimensions

$$S = - \int d^d x \left( \frac{i}{2} \bar{\psi} \gamma^\mu \partial_\mu \psi + g(\bar{\psi} \psi)^n \right) \quad (1.2)$$

What is the scaling dimension of  $\psi$ ? (You may assume the conjugate spinor  $\bar{\psi}$  has the same scaling dimension as  $\psi$ . Moreover, the gamma matrices are dimensionless.) For what  $(n, d)$  can  $g$  be made dimensionless, assuming  $n$  is a positive integer? Considering now also the scalar field of the previous problem. In what dimensions do  $\phi \bar{\psi} \psi$  and  $\phi^2 \bar{\psi} \psi$  lead to classically marginal couplings?

The scaling dimension of the free fermion is  $\Delta_\psi = \frac{d-1}{2}$ . We can make  $(\bar{\psi} \psi)^n$  dimensionless only in  $d = 2$ . These types of model are often referred to as Gross-Neveu or Thirring models. (The difference depends on how gamma matrices are sprinkled through the interaction.) The  $\phi \bar{\psi} \psi$  and  $\phi^2 \bar{\psi} \psi$  terms are sometimes called Yukawa couplings, and the associated theory a Yukawa theory.  $\phi \bar{\psi} \psi$  is classically marginal in  $d = 4$  while  $\phi^2 \bar{\psi} \psi$  is classically marginal in  $d = 3$ .

**Problem 1.5.** Start with the assumption that the supersymmetry transformation  $Q$  squares to the momentum operator  $Q^2 \sim P$  and moreover converts fermions into bosons and bosons into fermions. Try to guess how  $Q$  acts on  $\phi$  and  $\psi$ , purely based on dimensional analysis.

From  $Q^2 \sim P$ , we have that  $\Delta_Q = \frac{1}{2}$ . Since  $\Delta_\phi = \Delta_\psi + \frac{1}{2}$ , we have  $Q\phi \sim \psi$ . We also need a way to turn fermions back into bosons:  $Q\psi \sim \partial_\mu \phi$ .

**Problem 1.6.** Consider QED in  $d$  dimensions

$$S = - \int d^d x \left( \frac{1}{4} F^{\mu\nu} F_{\mu\nu} + \frac{i}{2} \bar{\psi} \gamma^\mu (\partial_\mu + ig A_\mu) \psi \right) \quad (1.3)$$

where  $F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu$ . What is the scaling dimension of  $g$ ? What is special about  $d = 4$ ?

We have that  $\Delta_{A_\mu} = \frac{d-2}{2}$ , like for a free scalar. Furthermore  $\Delta_g = d - 2\Delta_\psi - \Delta_{A_\mu} = 2 - \frac{d}{2}$ . In other words, the coupling is classically marginal in  $d = 4$  dimensions. (Quantum effects will typically lead to it running with energy scale.) In  $d > 4$ , the coupling is irrelevant, while in  $d < 4$  it is relevant.