

Thermo limit

$$\text{Def}^n: \frac{1}{2\pi} \frac{d\phi(k)}{dk} = \frac{1}{i} \frac{c}{k^2 + c^2} \equiv \mathcal{U}(k)$$

$$0 \leq \mathcal{U}(k) \leq \frac{2}{c} \quad k \in \mathbb{R}$$

$$\phi(k_1) - \phi(k_2) = \int_{k_2}^{k_1} dk \mathcal{U}(k) \leq \frac{2}{c} (k_1 - k_2)$$

$$\frac{2\pi}{L} |\bar{I}_j - \bar{I}_1| \geq |k_j - k_1| \geq \frac{2\pi}{L} \frac{|\bar{I}_j - \bar{I}_1|}{(1 + 2m/c)}$$

Let's view $\frac{I_j}{L} = x_j \in \mathbb{R}$

$$BE: k(x_j) + \frac{1}{L} \sum_l \phi(k(x_j) - k(x_l)) = 2\pi x_j$$

$k(x_j) \equiv k_j$ solving BE

Introduce a density $\rho(x) = \frac{1}{L} \sum_j \delta(x - \frac{I_j}{L})$

Extend defⁿ of $k(x)$ to $\mathbb{R}$ by requiring

$$k(x) + \int_{-\infty}^{\infty} dy \phi(k(x) - k(y)) \rho(y) = 2\pi x$$

Define other densities:

$$p_h(x) = \frac{1}{L} \sum_{m \in \{\tilde{I}\}} \delta(x - \frac{m}{L})$$

$\{I\}$: occupied q#
 $\{\tilde{I}\}$: unocc.

$$p_t(x) = \frac{1}{L} \sum_{m \in \{I\} \cup \{\tilde{I}\}} \delta(x - \frac{m}{L})$$

$$p(x) + p_h(x) = p_t(x)$$

$$\lim_{Ah} p_t(x) \rightarrow 1$$

$k(x)$ is a mapping $x \leftrightarrow k$
equiv, $x(k)$

Move to k space, defining densities

$$P(k) = P(x(k)) \frac{dx(k)}{dk} \quad P_h(k) = P_h(x(k)) \frac{dx(k)}{dk}$$

$$P_z(k) = \frac{dx(k)}{dk} \quad P(k) + P_h(k) = P_z(k)$$

BE: $k + \int_{-\infty}^{\infty} dk' \phi(k-k') P(k') = 2\pi x(k)$

Take $\frac{d}{dk}$ " $\Rightarrow$ $P(k) + P_h(k) = \frac{1}{2\pi} + \mathcal{C} * P(k)$

where $f * g(k) = \int_{-\infty}^{\infty} dk' f(k-k') g(k')$

Given a $\rho(k)$, eigenvalues of simple charges are

$$n = \frac{N}{L} = \int_{-\infty}^{\infty} dk \rho(k) \quad p = \frac{P}{L} = \int_{-\infty}^{\infty} dk k \rho(k)$$

$$e = \frac{E}{L} = \int_{-\infty}^{\infty} dk k^2 \rho(k)$$

The ground state: $\rho(k) = \begin{cases} \rho_F(k) & |k| < k_F \\ 0 & \text{otherwise} \end{cases}$

BE for GS:

$$\rho(k) - \int_{-k_F}^{k_F} dk' \mathcal{C}(k-k') \rho(k') = \frac{1}{2\pi} \quad |k| < k_F$$

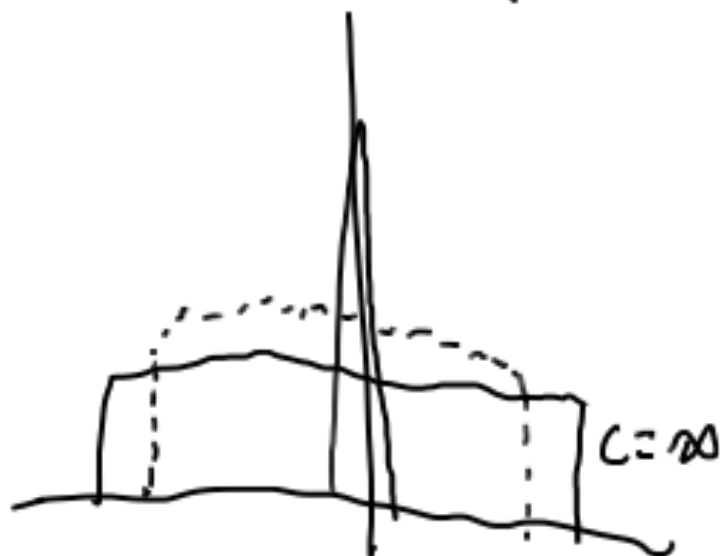
$$n = \int_{-k_F}^{k_F} dk \rho(k)$$

Only easy case: $L \rightarrow \infty$, $\ell \rightarrow 0$

so Lieb eqⁿ solved by $\rho(k) = \frac{\Theta(k_F^2 - k^2)}{2\pi}$

$$k_F = \pi n$$

Sketch:



Excitations around the GS.

Particle-like case $\bar{\Delta}^m$ (type 1)

$$Q\# : \{I_j\} = \left\{ -\frac{N}{2}, -\frac{N}{2}+1, \dots, \frac{N}{2}-1, \frac{N}{2}+m \right\}_{N+1}$$

$$\text{GS } Q\# : \{I_j^0\} = \left\{ -\frac{(N-1)}{2}, \dots, \frac{N-1}{2} \right\}_N$$

Exc state: momentum $p = \frac{2\pi}{L} m$, rep. $\{ \{k_j\}, q \}$

$$\text{GS} : \{k_j^0\}$$

Expect: $k_j - k_j^0 \equiv \frac{2\pi}{L} j$ small $j=1, \dots, N$

Subtract BE for k_j^0 from that for k_j :
 get α_j^* for d_j :

$$d_j = -\pi - \phi(k_j - q) - \sum_{l=1}^N [\phi(k_j - k_l) - \phi(k_j^0 - k_l^0)]$$

Expanding in $1/L$:

$$d_j = -\pi + \phi(k_j - q) - \frac{1}{L} \sum_{l=1}^N \frac{\partial \phi}{\partial (k_j - k_l)} \cdot (d_j - k_l) + O(1/L^2)$$

$$\Rightarrow d_j \left\{ 1 + \frac{1}{L} \sum_{l=1}^N \frac{\partial \phi}{\partial (k_j - k_l)} \right\} = -\pi - \phi(k_j - q) + \frac{1}{L} \sum_{l=1}^N \frac{\partial \phi}{\partial (k_j - k_l)} d_l$$

$$d_j \rightarrow d(k, q)$$

Introduce a $\sum$ $D(k, q) = d(k, q) \rho(k)$

get:

$$D(k, q) - \int_{-k_F}^{k_F} dk' \mathcal{P}(k-k') D(k', q) = -\frac{1}{2\pi i} \left\{ \pi + \phi(k-q) \right\}$$

$q > k_F$

Change in momentum:

$$\Delta P = q + \sum_{i=1}^N k_i - k_i^0 = q + \int_{-k_F}^{k_F} dk D(k, q)$$

$$\Delta E = q^2 - \mu + 2 \int_{-k_F}^{k_F} dk k D(k, q)$$

Simple limits:

$$C=0^+: D=0 \rightarrow \varepsilon(p) = p^2$$

$$C \rightarrow \infty: D = -1/2 \rightarrow \varepsilon(p) = p^2 + 2\pi m p$$

Hole-like exⁿ (type II)

Equilibrium thermodynamics

$$Z = \frac{1}{\Omega} \sum_{N=0}^{\infty} \sum_{\{I\}_N} e^{-(E_{\{I\}} - \mu N)/T}$$

$$Z \Rightarrow \int \mathcal{D}[p(k), p_h(k)] e^{S[p]} e^{-(E - \mu N)/T}$$

BE

In an interval $[k, k + \Delta k]$, $\frac{1}{L} \ll \Delta k \ll 1$

there are $L p(k) \Delta k$ particles

$L p_h(k) \Delta k$ holes

in $L p_t(k) \Delta k$ 'slots'

no of states described by $\rho(k)$ in the k -bin
differing by perm's in $[k, k+\Delta k]$

$$= \frac{[L(\rho+\rho_k)\Delta k]!}{[L\rho\Delta k]! [L\rho_k\Delta k]!}$$

→ get the (Yang Yang) entropy

$$S[\rho, \rho_k] = L \int_{-\infty}^{\infty} dk \left\{ (\rho + \rho_k) \ln(\rho + \rho_k) - \rho \ln \rho - \rho_k \ln \rho_k \right\}$$

$$\text{as } Z = \int_{\mathcal{G}} \mathcal{D}[p, p_h] e^{-G[p, p_h]/T} \Big|_{\text{BE}}$$

$$G[p, p_h] = L \int_{-\infty}^{\infty} dk \left[p(k) \left\{ k^2 - \mu - T \ln[p + p_h] + T \ln p \right\} + p_h(k) \left\{ -T \ln[p + p_h] + T \ln p_h \right\} \right]$$

Saddle-pt evaluation: consider variations

$$p(k) \rightarrow p(k) + \delta p(k)$$

$$p_h(k) \rightarrow p_h(k) + \delta p_h(k)$$

$$\text{get } \frac{\delta G}{L} = \int_{-\infty}^{\infty} dk \left\{ \delta\rho(k) \left\{ k^2 - \mu - T \ln[1 + \rho/\rho_0] \right\} \right. \\ \left. + \delta\rho_0(k) \left\{ -T \ln[1 + \rho/\rho_0] \right\} \right\}$$

$\delta\rho$ & $\delta\rho_0$ are related by BE:

$$\delta\rho(k) - \mathcal{U} * \delta\rho(k) = -\delta\rho_0(k)$$

$$\Rightarrow \frac{\delta G}{L} = \int_{-\infty}^{\infty} dk \delta\rho(k) \left\{ k^2 - \mu - T \ln \frac{\rho_0 + \rho}{\rho_0} - \mathcal{U} * T \ln[1 + \rho/\rho_0] \right\}$$

Saddle-pt: $\{ \} = 0 \quad \forall k$ Define $\varepsilon(k) = T \ln \frac{\rho_0 + \rho}{\rho_0}$

SP condⁿ: $\boxed{\varepsilon(k) = k^2 - \mu - \ell * T \ln[1 + e^{-\varepsilon(k)/T}]}$

The free energy itself becomes:

$$\left[g(T, \mu) = \frac{G}{L} = -T \int_{-\infty}^{\infty} \frac{dk}{2\pi} \ln[1 + e^{-\varepsilon(k)/T}] \right]$$

Simple limits:

$$\ell \rightarrow \infty: \ell \rightarrow 0, \quad \varepsilon(k) = k^2 - \mu$$

$$\ell \rightarrow 0: \ell \rightarrow \delta \quad \varepsilon(k) = T \ln \left[e^{\frac{k^2 - \mu}{T}} - 1 \right]$$

$\rightarrow$ free bosons

