

# Introduction to Random Matrix Theory

- I] General Overview and applications of RMT
  - II] Ensembles of RMT
  - III] Coulomb gas approach
  - IV] Local statistics (bulk & edge)
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## I] Intro.

$M$  :  $N \times N$  matrix w. random entries  $M_{jk}$

Real & Complex

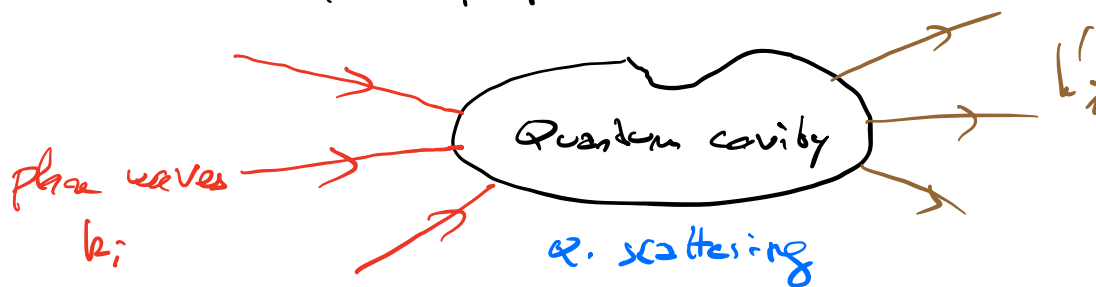
$Q$ : stat. of eigenvalues and eigenvectors  
 $N \rightarrow \infty$

In physics, introduced by Wigner ('50)

↳ study of nuclear physics to model  
the energy levels of big nuclei.

Since then many applications:

\* Mesoscopic physics:



$$\underline{k}' = S \underline{k}$$

random matrix

\* Random graph : adjacency matrix

\* nber theory (Riemann Zeta - function)

\* combinatorics

\* stochastic / Fluctuating interfaces  
Kardar-Parisi-Zhang equation

↳ related to directed polymers + disorder

In fact RMT was introduced by J. Wishart (1928)

↳ correlation in time series

$$\vec{X}_t = \begin{pmatrix} x_{1,t} \\ x_{2,t} \\ \vdots \\ x_{n,t} \end{pmatrix} \quad \text{stock price, daily temp.}$$

$$\hat{C}_{jk} = \frac{1}{T} \sum_{t=1}^T x_{jt} x_{kt}$$



random correlated matrix

→ applications in data analysis.

## 2) Ensembles of RMT

→ Matrices w. real spectrum      real symmetric  
complex Hermitian

→ Two main categories of rand. matrices:

a) Ens. with independent entries

"Wigner matrices"

$$M_{jk} = x_{jk} + i y_{jk}$$

Joint probab. density function

$$\mathcal{P}(M) = \mathcal{P}(\{M_{jk}\})$$

$M_{ij} = M_{ji} \quad i \neq j$



fixed by symmetry

$$\mathcal{P}(M) = \prod_{i=1}^N f_i(x_{ii}) \prod_{j < k} f_{jk}^{(R)}(x_{jk}) f_{jk}^{(I)}(y_{jk})$$

$$\text{Nber of degrees of freedom} = 1 + 2 + 3 + \dots + N \\ \sim \frac{N(N+1)}{2}$$

b) Rotationally invariant ensembles

$\Pi$ : real symmetric  $\Rightarrow \exists (\underline{O}, \underline{\Lambda})$

$$\underline{\Lambda} = \begin{pmatrix} \lambda_1 & & 0 \\ & \ddots & \\ 0 & & \lambda_N \end{pmatrix}, \lambda_i = \text{eigenvalues of } \Pi$$

$$\text{w. } \underline{O} \underline{O}^T = \underline{1} \text{ such that } \Pi = \underline{O} \underline{\Lambda} \underline{O}^{-1}$$

Rotationally inv.:  $\mathcal{S}(\Pi) = \mathcal{S}(\underline{O} \Pi \underline{O}^{-1})$

$\hookrightarrow$  the eigenvectors do not play a very important role.

$$\Leftrightarrow \mathcal{S}(\Pi) \propto e^{-\text{Tr } \chi(\Pi)}$$

$$P(\lambda_1, \dots, \lambda_N) \propto \prod_{i < j} |\lambda_i - \lambda_j|^\beta e^{-\sum_i \chi(\lambda_i)}$$

$\uparrow$   
 some function

Porter-Rosenzweig thm ('60) : the only

ensembles that belong to a) and b)

are the Gaussian ens. :

$$\mathcal{B}(M) = \frac{1}{Z_N} e^{-a \text{Tr}(M^2) - b \text{Tr}(M)}$$

, show that this is of type a) [Wigner]

$$\begin{aligned} \mathcal{B}(M) &= \frac{1}{Z_N} e^{-a \sum_{j,h} M_{jh} \underbrace{M_{hj}}_{= M_{jh}, j \neq h} - b \sum_j M_{jj}} \\ &= \frac{1}{Z_N} e^{-a \sum_j M_{jj}^2 - a \sum_{j \neq h} M_{jh}^2 - b \sum_j M_{jj}} \\ &= \frac{1}{Z_N} e^{-a \sum_j M_{jj}^2 - b \sum_j M_{jj} - 2a \sum_{j < h} M_{jh}^2} \\ &= \frac{1}{Z_N} \prod_j e^{-a M_{jj}^2 - b M_{jj}} \prod_{i < j} e^{-2a M_{ij}^2} \end{aligned}$$

↳ belongs to a)

NB : for  $b=0$ , the variance of the diag.

terms is twice the one of the off-diag.  
terms.

In the following: set  $b = 0$

For  $N$  real sym.: Gaussian Orthogonal Ens. <sup>GOE</sup>

$N$  complex Hermitian: Gaussian Unitary Ens. <sup>GUE</sup>

c) Joint law of eigenvalues for GOE/GUE

GOE:  $M = M^t$  and real.

$$M = \underbrace{O}_{\text{eigenvectors}} \underbrace{\Lambda}_{\text{eigenvalues}} \underbrace{O^{-1}}_{\text{eigenvectors}}, \quad \Lambda = \begin{pmatrix} \lambda_1 & & 0 \\ & \ddots & \\ 0 & & \lambda_N \end{pmatrix}$$
$$OO^t = \mathbb{1}$$

Change of variables:  $M \equiv \{M_{jk}\} \rightarrow \{\Lambda, O\}$

Rh: nber of degrees of freedom:

$$\frac{N(N+1)}{2} = N + \frac{N(N-1)}{2}$$

$$\mathcal{P}(M) \longrightarrow P(\{\lambda_1, \dots, \lambda_N\}, O)$$
$$e^{-\alpha \text{Tr}(M^2)}$$

Related by a Jacobian:

$$P(\{\lambda_1, \dots, \lambda_N\}, \underline{O}) = \mathcal{B}(N) |\det J|$$

$$J = \left\{ \frac{\partial \pi_{i\delta}}{\partial \lambda_k}, \frac{\partial \pi_{i\delta}}{\partial O_{kl}} \right\}$$

size  $\frac{N(N+1)}{2} \times \frac{N(N+1)}{2}$

$$(GOE) \quad |\det J| = \prod_{j < k} (\lambda_j - \lambda_k)$$

indep. of  $\underline{O}$

$$\Rightarrow P(\{\lambda_1, \dots, \lambda_N\}, \underline{O}) = \frac{1}{Z_N} e^{-\alpha \sum_{i=1}^N \lambda_i^2} \underbrace{\prod_{j < k} |\lambda_j - \lambda_k|^\beta}_{\text{level repulsion}}$$

for GOE  $\beta = 1$

GOE  $\beta = 2$

• Vandermonde determinant:  $x_1, x_2, \dots, x_N$

$$V_N = \begin{pmatrix} 1 & 1 & \dots & 1 \\ x_1 & x_2 & & x_N \\ x_1^2 & x_2^2 & & x_N^2 \\ \vdots & \vdots & & \vdots \\ x_1^{N-1} & x_2^{N-1} & & x_N^{N-1} \end{pmatrix}$$

$$N=2: \begin{pmatrix} 1 & 1 \\ x_1 & x_2 \end{pmatrix} = V_2 \Rightarrow \det V_2 = x_2 - x_1.$$

$$\det V_N = \prod_{i < j} (x_j - x_i)$$

Here we choose:

$$P(\lambda_1, \dots, \lambda_N) = B_N e^{-\beta \frac{N}{2} \sum_{i=1}^N \lambda_i^2} \prod_{i < j} |\lambda_i - \lambda_j|^\beta$$

III. Coulomb gas approach (Dyson '62)

$$\text{Rewriting: } \prod_{i < j} |\lambda_i - \lambda_j|^\beta = e^{\beta \sum_{i < j} \ln |\lambda_i - \lambda_j|} = e^{\frac{\beta}{2} \sum_{i \neq j} \ln |\lambda_i - \lambda_j|}$$

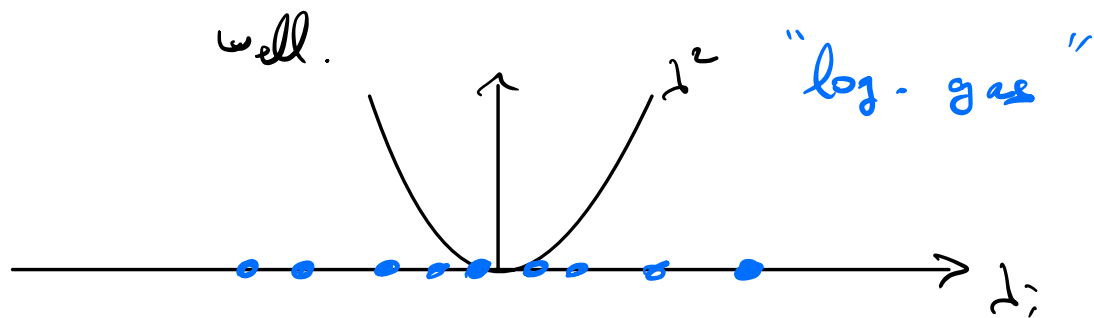
$$\Rightarrow P(\lambda_1, \dots, \lambda_N) = B_N e^{-\beta E(\{\lambda_i\})}$$

$$E(\{\lambda_i\}) = \frac{N}{2} \sum_{i=1}^N \lambda_i^2 - \frac{1}{2} \sum_{i \neq j} \ln |\lambda_i - \lambda_j|$$

repulsive

$\Rightarrow \lambda_i \equiv$  positions of charged particles  
interacting via the 2d Coulomb interaction  
confined on a line within a harmonic





$\Rightarrow$  competition between confinement and  
the repulsive interactions