Complexity of mixed Gaussian states from Fisher information geometry

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G.D.G, E.Tonni, JHEP **2020**,101 (2020) G.D.G, E.Tonni, JHEP **2021**,22 (2021)

Cortona Young 2021

9 June 2021





Circuit Complexity

Reference state $\ket{\psi_{\mathrm{R}}}$

Target state $|\psi_{\mathrm{T}}
angle$

Quantum circuit:

$$|\psi_{\mathrm{T}}\rangle = U |\psi_{\mathrm{R}}\rangle = Q_{i_{D}} \dots Q_{i_{2}} Q_{i_{1}} |\psi_{\mathrm{R}}\rangle$$

Set of elementary (unitary) gates: $\{Q_i\}_{i=1,...}$

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Complexity: minimum number of allowed gates that is needed to construct the target state starting from the assigned reference state

Holographic Complexity

Holographic entanglement entropy is not able to capture the late-time dynamics behind the horizon of a black hole



Complexity in quantum field theories still to be understood

Gaussian states in harmonic lattices

We compute the complexity for lattice models with a known continuum limit

$$\widehat{H} = \sum_{i=1}^{N} \left(\frac{1}{2m} \hat{p}_i^2 + \frac{m\omega^2}{2} \hat{q}_i^2 \right) + \sum_{\langle i,j \rangle} \frac{\kappa}{2} (\hat{q}_i - \hat{q}_j)^2 \qquad [\hat{q}_i, \hat{p}_j] = \mathrm{i}\delta_{ij}$$

- The continuum limit is a Klein-Gordon field theory
- In the continuum, when $\omega \rightarrow 0$, we have a conformal field theory

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- The continuum limit is a (1+1)-dimensional Klein-Gordon field theory
- In the continuum, when $\omega \rightarrow 0$, we have a c = 1 conformal field theory

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We restrict our attention to the Gaussian states (with vanishing first moments)

$$\hat{\rho} \longrightarrow \gamma$$
 γ is the covariance matrix

The entries of γ are given by $\langle \{\hat{q}_i, \hat{q}_j\} \rangle$, $\langle \{\hat{p}_i, \hat{p}_j\} \rangle$ and $\langle \{\hat{q}_i, \hat{p}_j\} \rangle$ and it satisfies the uncertainty principle

$$\gamma + i \frac{J}{2} \ge 0$$

Complexity as geodesic distance [Nielsen, '06]

The **quantum circuit** connecting two **mixed** states $\hat{\rho}_{\rm R}$ and $\hat{\rho}_{\rm T}$ can be seen as a **curve** connecting two points in a manifold of quantum mixed states



 $\hat{
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Optimal circuit (circuit with the minimum number of gates) \longrightarrow Geodesics Complexity \longrightarrow Length of the geodesics

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Compute the complexity becomes a geometric problem with ambiguities due the choice of the **cost function**

Complexity of pure states: [Jefferson, Myers, '17; Hackl, Myers, '18;

Chapman, Eisert, Hackl, Heller, Jefferson, Marrochio, Myers, '19]

Complexity from Fisher-Rao distance

Assumption: we restrict the space of all the states we can reach through a quantum circuit to the set of bosonic Gaussian **mixed** state

Reference state $\gamma_{
m R}$

Target state γ_{T}

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The complexity is given by the Fisher-Rao distance [Rao, '45; GDG, Tonni, '20]:

$$d(\gamma_{\mathrm{R}},\gamma_{\mathrm{T}}) \equiv \sqrt{\mathrm{Tr}\{[\log(\gamma_{\mathrm{T}}\gamma_{\mathrm{R}}^{-1})]^2\}} \qquad \mathcal{C} = rac{1}{2\sqrt{2}} d(\gamma_{\mathrm{R}},\gamma_{\mathrm{T}})$$

Explicit expressions for spectrum and basis complexity in [GDG, Tonni, '20]

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Subsystem complexity at equilibrium [GDG, Tonni, '20]

$$|\psi\rangle \in \mathcal{H}_A \otimes \mathcal{H}_B \ \Rightarrow \ \hat{\rho} = |\psi\rangle \langle \psi| \ \Rightarrow \ \hat{\rho}_A = \mathrm{Tr}_B \hat{\rho} \ \rightarrow \ \gamma_A$$

$$\mathcal{C}_{\mathcal{A}} = \frac{1}{2\sqrt{2}} \sqrt{\mathrm{Tr}\left\{[\log(\gamma_{\mathcal{A},\mathrm{T}} \gamma_{\mathcal{A},\mathrm{R}}^{-1})]^2\right\}}$$

Subsystem complexity in harmonic lattices [GDG, Tonni, '20] Subsystem complexity at equilibrium [GDG, Tonni, '20]

$$|\psi\rangle \in \mathcal{H}_A \otimes \mathcal{H}_B \implies \hat{\rho} = |\psi\rangle\langle\psi| \implies \hat{\rho}_A = \mathrm{Tr}_B \hat{\rho} \longrightarrow \gamma_A$$

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Subsystem complexity in harmonic lattices [GDG, Tonni, '20]

A is an interval in infinite line





Global quantum quench

$$\hat{H} = \sum_{i=1}^{N} \left(\frac{1}{2m} \, \hat{p}_i^2 + \frac{m\omega^2}{2} \, \hat{q}_i^2 + \frac{\kappa}{2} (\hat{q}_i - \hat{q}_{i-1})^2 \right) = \hat{H}(\omega, \kappa, m)$$

Global quench protocol [Calabrese, Cardy,'05]

t = 0

$$|\Psi_0
angle$$
 ground state of $\widehat{H}(\omega_0,\kappa,m)$

$$[\hat{H}(\omega_0,\kappa,m),\hat{H}(\omega,\kappa,m)] \neq 0 \implies$$

$$|\Psi(t)
angle = e^{-i \hat{H}(\omega,\kappa,m)t} |\Psi_0
angle$$

t > 0

non-trivial time evolution (preserving Gaussianity of $|\Psi_0\rangle)$

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non-trivial time evolution (preserving Gaussianity of $|\Psi_0\rangle$)

• When $\omega = 0 \longrightarrow CFT$ description in the continuum [Calabrese, Cardy, '05]

 Vaidya spacetimes are employed as the gravitational duals of global quantum quenches in the CFT on their boundary

Temporal evolution of the subsystem complexity [GDG, Tonni, '21]

$$\gamma_{A,R} \rightarrow$$
 initial state $(t = 0)$

 $\gamma_{A,\mathrm{T}} \rightarrow \text{state at time } t \text{ after}$ the quench

$$\Delta S_A = S_A(t) - S_A(0)$$







Temporal evolution of the subsystem complexity [GDG, Tonni, '21]

3.0

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- Different initial growth
- Saturation when t > L + d: relaxation to the generalised Gibbs ensamble



Qualitative comparison with holographic results

[GDG, Tonni, '21]



Holographic subsystem complexity: CV proposal



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Qualitative comparison with holographic results



[GDG-Tonni,'21]

Holographic subsystem complexity: CV proposal

- Similar qualitative behaviours when 0 < t < L/2: initial linear growth and local maximum
- Different behaviours when t > L/2



Conclusions

- Complexity for circuits made by mixed states exploiting Fisher-Rao distance
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Future perspectives

- Local quench
- Non-Gaussian states
- Complexity of fermionic mixed states

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