

SCHRÖDINGER CONNECTION WITH SELF DUAL NON-METRICITY VECTOR IN $2+1$ DIMENSIONS

Based on: S. Klemm, L. Rovera, 2008.12740

1918: Weyl attempts to unify EM with gravity geometrically $\rightarrow$ Connection with non-metricity $\leadsto$ Weyl vector $\leftrightarrow$ gauge pot.

- Physical inconsistencies ("2nd clock")

1940s: Schrödinger $\rightarrow$ Symmetric connection that involves non-metricity but preserves the length of vectors under parallel transport! $\rightarrow$ No "2nd clock"

Schrodinger connection

$$\hat{\Gamma}^{\lambda}_{\mu\nu} = \tilde{\Gamma}^{\lambda}_{\mu\nu} + g^{\lambda\rho} Z_{\rho\mu\nu}$$

$\hookrightarrow LC$

$$Z_{\lambda\mu\nu} = Z_{\lambda\nu\mu}, \quad Z(\lambda\mu\nu) = 0$$

Now, consider a generic conn. $\Gamma^{\lambda}_{\mu\nu} = \tilde{\Gamma}^{\lambda}_{\mu\nu} + N^{\lambda}_{\mu\nu}$

$N^{\lambda}_{\mu\nu}$: deflection ($Q_{\lambda\mu\nu}$) - contorsion ($T_{\lambda\mu\nu}$)

$$\leadsto \text{if } \Gamma^\lambda_{\mu\nu} = 0 \text{ and } N(\lambda_{\mu\nu}) = 0$$

$$\Rightarrow \Gamma^\lambda_{\mu\nu} = \tilde{\Gamma}^\lambda_{\mu\nu} - g^{\lambda\rho} Q_{\rho\mu\nu}, \quad Q(\lambda_{\mu\nu}) = 0$$

$$\text{Taking } Z\lambda_{\mu\nu} = -Q\lambda_{\mu\nu} \Rightarrow \Gamma = \hat{\Gamma}, \text{ Schr.}$$

$$\Rightarrow \hat{\Gamma}^\lambda_{\mu\nu} = \tilde{\Gamma}^\lambda_{\mu\nu} - g^{\lambda\rho} Q_{\rho\mu\nu}$$

more physical
than Weyl's
connection!

Non-metricity decomposition in 3-dim.

$$Q_{\lambda\mu\nu} = \frac{2}{5} Q_{\lambda} g_{\mu\nu} - \frac{1}{5} \tilde{Q}_{\lambda} g_{\mu\nu} + \frac{3}{5} g_{\lambda(\nu} \tilde{Q}_{\mu)} \\ - \frac{1}{5} g_{\lambda(\nu} Q_{\mu)} + \Omega_{\lambda\mu\nu}$$

→ traceless

$$Q_{\lambda} = Q_{\lambda\mu}{}^{\mu} \quad \tilde{Q}_{\lambda} = Q^{\mu}{}_{\mu\lambda}$$

Our model (3 dim., 2+1)

$$S = \frac{1}{2\kappa^2} \int d^3x \left(\sqrt{-g} \mathcal{I}(R) + \frac{1}{2\mu} \epsilon^{\mu\nu\rho} Q_\rho \hat{R}_{\nu\mu} \right) \\ + \int d^3x \epsilon^{\mu\nu\rho} \sum_{\nu\sigma} T_{g\mu}{}^\sigma \quad \begin{array}{l} \text{CS term} \\ \downarrow \\ \text{LM} \end{array}$$

$$R = g^{\mu\nu} R_{\mu\nu}(\Gamma)$$

$$\hat{R}_{\mu\nu} = R^\lambda{}_{\lambda\mu\nu} = \partial_{[\mu} Q_{\nu]}$$

$$\epsilon^{\mu\nu\rho} = \sqrt{-g} \varepsilon^{\mu\nu\rho}$$

E.O.M.:

$$\delta \mathcal{L}_{\mu\nu} : T_{\rho\sigma}{}^\lambda = 0$$

$$\delta g^{\mu\nu} \quad f'(R) R_{(\mu\nu)} - \frac{1}{2} f(R) g_{\mu\nu} = 0$$

$$\text{Trace} : \frac{f}{2f'} = \frac{R}{3} \quad \rightarrow \quad f(R) = C R^{3/2}$$

back $\downarrow$

$$R_{(\mu\nu)} - \frac{R}{3} g_{\mu\nu} = 0$$

$$\text{inv. } g_{\mu\nu} \rightarrow e^{2\Omega} g_{\mu\nu}$$

$$R = C_k, \text{ const.}$$

(would be EW if
Weyl non-metricity)

$$\delta \Gamma^\lambda_{\mu\nu} : P_{\lambda\mu\nu} + \delta_\lambda^\nu \frac{\partial^\mu f'}{f'} - g_{\mu\nu} \frac{\partial_\lambda f'}{f'} \\ + \frac{2}{\mu f'} \varepsilon^{\nu\rho\sigma} \delta_\lambda^\mu \hat{R}_{\rho\sigma} + \frac{2\kappa^2}{f'} \varepsilon^{\mu\nu\rho} \zeta_{\rho\lambda} = 0$$

$$\Rightarrow \Omega_{\lambda\mu\nu} = 0, \quad \zeta_{\rho\sigma} = \frac{3}{\kappa^2 \mu} \hat{R}_{\rho\sigma}$$

$$\hat{R}_{\rho\sigma} = \frac{\mu f'}{20} \varepsilon^{\mu\lambda} g_{\rho\mu} g_{\sigma\nu} (Q_\lambda - 3\tilde{Q}_\lambda)$$

$$\text{And } \frac{\partial \mathcal{L}}{\partial \tilde{Q}_\lambda} = \frac{1}{10} (Q_\lambda + 2 \tilde{Q}_\lambda)$$

$$\Rightarrow \Gamma^\lambda_{\mu\nu} = \tilde{\Gamma}^\lambda_{\mu\nu} - \frac{3}{10} g_{\mu\nu} Q^\lambda + \frac{2}{5} \delta^\lambda_{(\mu} Q_{\nu)} \\ + \frac{2}{5} g_{\mu\nu} \tilde{Q}^\lambda - \frac{1}{5} \delta^\lambda_{(\mu} \tilde{Q}_{\nu)}$$

- For $C=1$: Generalized monopole eq.

$$\star (d\Sigma + h\Sigma) = dh$$

$$\Sigma = 3\mu\sqrt{R}/8 \leftarrow$$

$$\hookrightarrow h_\lambda = -Q_\lambda/6$$

And $\partial_\mu \ln R = \frac{1}{5} (Q_\mu + 2 \tilde{Q}_\mu)$

- If Σ const. (achieved by $\Sigma \rightarrow e^{-\Omega \Sigma}$)

$\Rightarrow d\ln = \star \ln \Sigma$

Self-duality in
3 dimensions!

In this case $\tilde{Q}_\mu = -\frac{1}{2} Q_\mu$

- Setting Σ const = Sol. $R = C_0 \Rightarrow f(R)$ const
and $f'(R) = C_0$

- Consider $f'(R) = C_0$

$$\partial_\lambda f' = 0$$

$$R(\mu\nu) = \frac{C_0}{3} g_{\mu\nu}$$

$$\text{and } \xi_{\rho\sigma} = \frac{3C_0}{8\kappa^2} \epsilon_{\rho\sigma\tau} Q^\tau$$

$$\hat{R}_{\rho\sigma} = \frac{C_0}{8} \epsilon_{\rho\sigma\tau} Q^\tau$$

↓ can be dualized as



$$-\frac{\omega_\mu}{4} Q_\alpha = \varepsilon_{\alpha\rho\sigma} \partial^\rho Q^\sigma$$

which implies

$$\tilde{\nabla}^\mu Q_\mu = 0$$

In particular we have $Q(\lambda_{\mu\nu}) = 0$

$$\text{and } Q_{\lambda\mu} = \frac{1}{2} Q_\lambda g_{\mu\nu} - \frac{1}{2} Q_\mu g_{\lambda\nu} - \frac{1}{4} Q_\nu g_{\lambda\mu}$$

$$\Rightarrow \Gamma^\lambda_{\mu\nu} = \tilde{\Gamma}^\lambda_{\mu\nu} - g^{\lambda\rho} Q_{\rho\mu\nu}$$

Schrödinger
with selfdual Q_λ

Self-duality in odd dim and inh. Maxwell eqs.

Townsend - Pilch - van Nieuwenhuizen :

dim. $m = 4k - 1$, $k = 1, 2, 3, \dots$

"Square root" of Proca eq. for a massive

AS tensor field $\rightarrow$ Self dual field

- Our model : $k = 1$

- We have:

$$Q_\mu = \frac{1}{2M} \epsilon_{\mu\nu\rho} F^{\nu\rho}, \quad F_{\mu\nu} = 2\partial_{[\mu} Q_{\nu]} = 2\hat{R}_{\mu\nu}$$

$$M = -C_0 \mu / 4$$

Dual form: $F_{\alpha\beta} = M \epsilon_{\mu\beta\alpha} Q^\mu$

$\Rightarrow \tilde{\nabla}^\mu Q_\mu = 0$ and covariant Proca eq.

$$\tilde{\nabla}^\mu F_{\mu\nu} - M^2 Q_\nu = 0$$

$$\rightarrow \partial^\mu \bar{F}_{\mu\nu} - \frac{1}{2} g^{\alpha\beta} \partial^\mu g_{\alpha\beta} F_{\mu\nu} - M^2 Q_\nu = 0$$

$\Rightarrow Q_\mu$: massive photon gauge field

- Furthermore, cov. Proca = inhom. EM wave eq.
following from inhom. Maxwell eqs

$\rightarrow Q_\mu$: massless and "source of itself"

See this explicitly:

- Wave eq. in our model $(\square - M^2) Q^\mu = 0$

$$\square \equiv \tilde{\nabla}_\alpha \tilde{\nabla}^\alpha$$

- Ymhom EM wave eq for Q^μ in the gauge

$$\tilde{\nabla}^\mu Q_\mu = 0 : \quad \square Q^\mu = \mu_0 J^\mu \Leftrightarrow \square Q^\mu = \tilde{\nabla}_\rho F^{\rho\mu}$$

μ_0 : vacuum permeability

J^μ : current

• Using S-d in our model & $(\square - M^2) Q^\mu = 0$:

$$\square Q^\mu = \frac{M}{2} \epsilon^{\mu\rho\sigma} F_{\rho\sigma} = \tilde{\nabla}_\rho (M \epsilon^{\mu\rho\sigma} Q_\sigma) = \tilde{\nabla}_\rho F^{\rho\mu}$$

$$\Rightarrow \square Q^\mu = \tilde{\nabla}_\rho F^{\rho\mu} \text{ in our model}$$

→ We can write $\square Q^\mu = \mu_0 J^\mu$,

$$\underline{J^\mu} \equiv \frac{1}{\mu_0} \tilde{\nabla}_\rho F^{\rho\mu} = \frac{M}{2\mu_0} \epsilon^{\mu\rho\sigma} F_{\rho\sigma} = \underline{\frac{M^2}{\mu_0} Q^\mu}$$

- The source current is covariantly cons. $\tilde{\nabla}_\mu J^\mu = 0$

↳ Cannot def globally cons. charge,
since we need a local cons. law for this

⇒ Take Proca: $\partial^\mu F_{\mu\nu} - \frac{1}{2} g^{\alpha\beta} \partial^\mu g_{\alpha\beta} F_{\mu\nu} - M^2 Q_\nu = 0$

→ Total current
(source + self)

$$j^\mu = J^\mu + \frac{1}{2\mu_0} g^{\alpha\beta} \partial_\beta g_{\alpha\beta} F^{\mu\rho}$$

locally conserved: $\partial_\mu j^\mu = 0$

Conclusions

- Our Proca, massive photon $\rightsquigarrow$ ^{interpreted as} inh. Maxwell eqs., with corresponding wave eq.
 - $\rightarrow$ Photon "Source of itself" due to self-duality
- Gauge inv. of connection and inh. Maxwell eqs follows from self-duality

$\Rightarrow Q_\mu$ "photon"

- In our framework

Inh. Maxwell eqs.

from a purely
gravitational setup!

- CS and Logr. mult. break proj. inv. of S

$\Rightarrow$ Consistent matter couplings allowed

- Dynamical d.o.f. : 1 for Q_μ ($\partial^\mu Q_\mu = 0$ + S-d.)

→ Extension to higher dimensions?

THANK

you !