

# ELECTROWEAK PHASE TRANSITIONS **S** AND HIGGS COUPLINGS

Patrick Meade

C.N. Yang Institute for Theoretical Physics  
Stony Brook University

Based mostly on

PM, H. Ramani 1807.07578

PM, S. Homiller 1808/9.xxxxxx

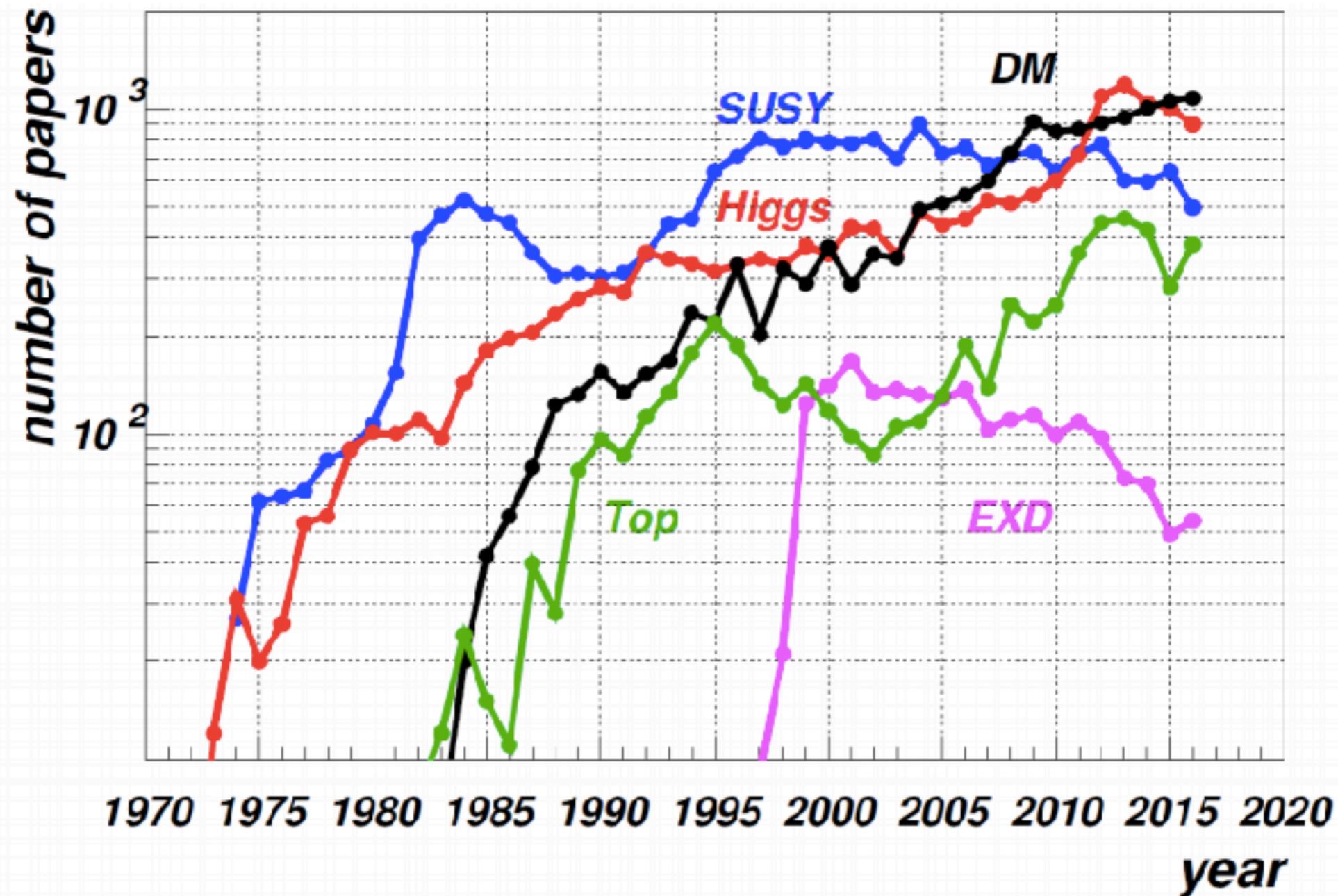
David Curtin, PM, H. Ramani 1612.00466

# MOTIVATION...

Why do I care about the early universe and the Higgs and not DM for instance?

# MOTIVATION...

## Why we are so keen to study DM?



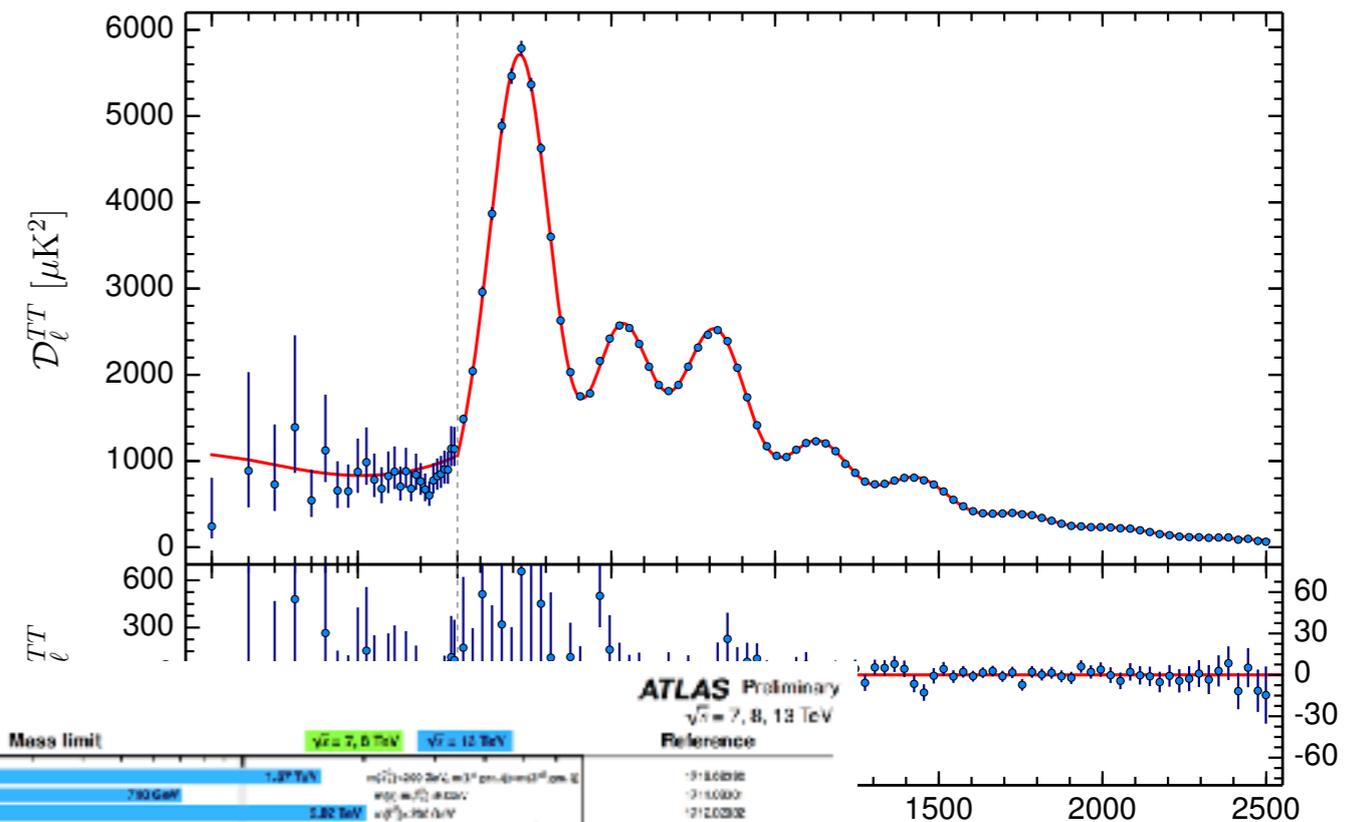
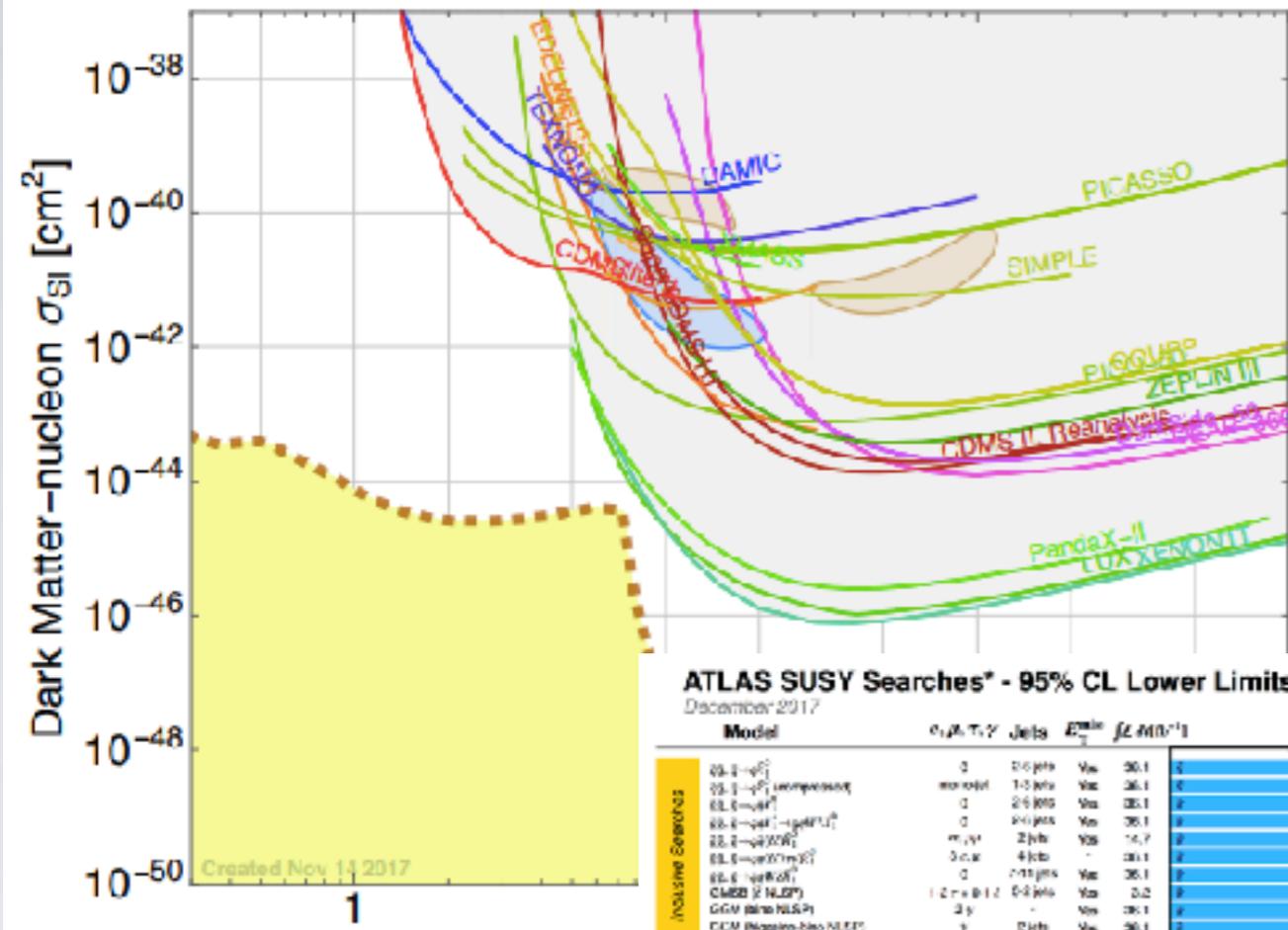
# MOTIVATION...

“It doesn’t seem like there’s anything interesting in pheno lately. Maybe Neutrinos?” L. Alvarez-Gaume (Simons Center for Geometry & Physics director at Stony Brook)

my atavistic pheno impulse is to give a panglossian view of our field:

“X, Y, and Z are being done, amazing new possibilities”

# DATA/REALITY DRIVEN VIEW:



Model	$\chi_{min}^2$	Jets	$E_{jet}^{min}$ [GeV]	Mass limit	$\chi_{min}^2$ at $\sqrt{s} = 7, 8 \text{ TeV}$	$\chi_{min}^2$ at $\sqrt{s} = 13 \text{ TeV}$	Reference		
Inclusive Searches	$\tilde{g}, \tilde{u} \rightarrow q\tilde{q}$	0	0.5 jets	Yes	30.1	1.37 TeV	$m(\tilde{g}) > 200 \text{ GeV}, m(\tilde{u}) > 100 \text{ GeV}$	1710030	
	$\tilde{g}, \tilde{d} \rightarrow q\tilde{q}$ (compressed)	non-sat	1-3 jets	Yes	36.1	730 GeV	$m(\tilde{g}) > 100 \text{ GeV}$	1710030	
	$\tilde{g}, \tilde{t} \rightarrow q\tilde{q}$	0	2-6 jets	Yes	35.1	1.82 TeV	$m(\tilde{g}) > 100 \text{ GeV}$	1710030	
	$\tilde{g}, \tilde{b} \rightarrow q\tilde{q}$	0	2-6 jets	Yes	36.1	2.31 TeV	$m(\tilde{g}) > 100 \text{ GeV}, m(\tilde{b}) > 100 \text{ GeV}$	1710030	
	$\tilde{g}, \tilde{t} \rightarrow q\tilde{q} + \tilde{g}\tilde{t}$	0	2-6 jets	Yes	36.1	1.1 TeV	$m(\tilde{g}) > 100 \text{ GeV}, m(\tilde{t}) > 100 \text{ GeV}$	1710030	
	$\tilde{g}, \tilde{t} \rightarrow q\tilde{q} + \tilde{g}\tilde{t}$	3.0e	4 jets	-	34.1	1.57 TeV	$m(\tilde{g}) > 100 \text{ GeV}$	1710030	
	$\tilde{g}, \tilde{t} \rightarrow q\tilde{q} + \tilde{g}\tilde{t}$	0	2-10 jets	Yes	36.1	1.8 TeV	$m(\tilde{g}) > 100 \text{ GeV}$	1710030	
	$\tilde{g}, \tilde{t} \rightarrow q\tilde{q} + \tilde{g}\tilde{t}$	1.2e + 0.12	0-4 jets	Yes	32.2	2.4 TeV	$m(\tilde{g}) > 100 \text{ GeV}$	1710030	
	$\tilde{g}, \tilde{t} \rightarrow q\tilde{q} + \tilde{g}\tilde{t}$	2.3e	-	Yes	36.1	2.95 TeV	$m(\tilde{g}) > 100 \text{ GeV}$	1710030	
	$\tilde{g}, \tilde{t} \rightarrow q\tilde{q} + \tilde{g}\tilde{t}$	0	0 jets	Yes	36.1	2.23 TeV	$m(\tilde{g}) > 100 \text{ GeV}, m(\tilde{t}) > 100 \text{ GeV}$	1710030	
3 $\gamma$ jet	$\tilde{g}, \tilde{g} \rightarrow 3\gamma$	0	3 $\gamma$	Yes	36.1	1.91 TeV	$m(\tilde{g}) > 100 \text{ GeV}$	1710030	
	$\tilde{g}, \tilde{g} \rightarrow 3\gamma$	0.1e	3 $\gamma$	Yes	36.1	1.60 TeV	$m(\tilde{g}) > 100 \text{ GeV}$	1710030	
	3 $\gamma$ jet, heavy direct production	$\tilde{g}, \tilde{g} \rightarrow 3\gamma$	0	3 $\gamma$	Yes	36.1	1.91 TeV	$m(\tilde{g}) > 100 \text{ GeV}$	1710030
		$\tilde{g}, \tilde{g} \rightarrow 3\gamma$	2e + 0.05	1 $\gamma$	Yes	36.1	239-760 GeV	$m(\tilde{g}) > 100 \text{ GeV}, m(\tilde{u}) > 100 \text{ GeV}$	1710030
		$\tilde{g}, \tilde{g} \rightarrow 3\gamma$	0.2e	1-3 $\gamma$	Yes	4.2(0.0)	200-700 GeV	$m(\tilde{g}) > 100 \text{ GeV}, m(\tilde{u}) > 100 \text{ GeV}$	1710030
		$\tilde{g}, \tilde{g} \rightarrow 3\gamma$	0.2e	1-3 $\gamma$	Yes	20.3(0.1)	90-190 GeV	$m(\tilde{g}) > 100 \text{ GeV}$	1710030
		$\tilde{g}, \tilde{g} \rightarrow 3\gamma$	0	non-sat	Yes	36.1	90-133 GeV	$m(\tilde{g}) > 100 \text{ GeV}$	1710030
		$\tilde{g}, \tilde{g} \rightarrow 3\gamma$	2e + 0.12	1 $\gamma$	Yes	33.0	108-133 GeV	$m(\tilde{g}) > 100 \text{ GeV}$	1710030
		$\tilde{g}, \tilde{g} \rightarrow 3\gamma$	0.1e + 0.12	1 $\gamma$	Yes	36.1	293-790 GeV	$m(\tilde{g}) > 100 \text{ GeV}$	1710030
		$\tilde{g}, \tilde{g} \rightarrow 3\gamma$	1.0e	4 $\gamma$	Yes	36.1	308-630 GeV	$m(\tilde{g}) > 100 \text{ GeV}$	1710030
LSP mass		$\tilde{g}, \tilde{g} \rightarrow 3\gamma$	2.0e	0	Yes	36.1	63-690 GeV	$m(\tilde{g}) > 100 \text{ GeV}$	1710030
		$\tilde{g}, \tilde{g} \rightarrow 3\gamma$	2.0e	0	Yes	36.1	63-690 GeV	$m(\tilde{g}) > 100 \text{ GeV}$	1710030
	$\tilde{g}, \tilde{g} \rightarrow 3\gamma$	0.1e	0	Yes	36.1	390 GeV	$m(\tilde{g}) > 100 \text{ GeV}$	1710030	
	$\tilde{g}, \tilde{g} \rightarrow 3\gamma$	0.1e	0	Yes	20.0	170 GeV	$m(\tilde{g}) > 100 \text{ GeV}$	1710030	
	$\tilde{g}, \tilde{g} \rightarrow 3\gamma$	0.1e	0	Yes	20.0	400 GeV	$m(\tilde{g}) > 100 \text{ GeV}$	1710030	
	$\tilde{g}, \tilde{g} \rightarrow 3\gamma$	1.0e + 0.12	0	Yes	20.3	1.0-179 GeV	$m(\tilde{g}) > 100 \text{ GeV}$	1710030	
	$\tilde{g}, \tilde{g} \rightarrow 3\gamma$	2.3e	0	Yes	36.1	3.83 TeV	$m(\tilde{g}) > 100 \text{ GeV}$	1710030	
	Long lived particles	$\tilde{g}, \tilde{g} \rightarrow 3\gamma$	Direct prod.	1 jet	Yes	36.1	100 GeV	$m(\tilde{g}) > 100 \text{ GeV}, m(\tilde{u}) > 100 \text{ GeV}$	1710030
		$\tilde{g}, \tilde{g} \rightarrow 3\gamma$	Direct prod.	1 jet	Yes	36.1	100 GeV	$m(\tilde{g}) > 100 \text{ GeV}, m(\tilde{u}) > 100 \text{ GeV}$	1710030
		$\tilde{g}, \tilde{g} \rightarrow 3\gamma$	Stable, stopped $\tilde{g}$ / hadron	0	1-5 jets	Yes	27.9	220 GeV	$m(\tilde{g}) > 100 \text{ GeV}, m(\tilde{u}) > 100 \text{ GeV}$
$\tilde{g}, \tilde{g} \rightarrow 3\gamma$		Stable $\tilde{g}$ / hadron	0	-	3.2	1.88 TeV	$m(\tilde{g}) > 100 \text{ GeV}$	1710030	
$\tilde{g}, \tilde{g} \rightarrow 3\gamma$		Non-stable $\tilde{g}$ / hadron	direct prod.	-	0.2	1.57 TeV	$m(\tilde{g}) > 100 \text{ GeV}, m(\tilde{u}) > 100 \text{ GeV}$	1710030	
$\tilde{g}, \tilde{g} \rightarrow 3\gamma$		Measurable $\tilde{g}$ / hadron, $\tilde{g} \rightarrow q\tilde{q}$	direct prod.	Yes	30.8	2.37 TeV	$m(\tilde{g}) > 100 \text{ GeV}, m(\tilde{u}) > 100 \text{ GeV}$	1710030	
$\tilde{g}, \tilde{g} \rightarrow 3\gamma$		CMSSM stable $\tilde{g}, \tilde{g} \rightarrow q\tilde{q}, \tilde{g} \rightarrow q\tilde{q}$	1-2 jets	-	10.1	820 GeV	$m(\tilde{g}) > 100 \text{ GeV}$	1710030	
$\tilde{g}, \tilde{g} \rightarrow 3\gamma$		CMSSM $\tilde{g} \rightarrow q\tilde{q}, \tilde{g} \rightarrow q\tilde{q}$	2 jets	Yes	20.0	610 GeV	$m(\tilde{g}) > 100 \text{ GeV}$	1710030	
$\tilde{g}, \tilde{g} \rightarrow 3\gamma$		$\tilde{g}, \tilde{g} \rightarrow q\tilde{q} + \tilde{g}\tilde{t}$	0	-	20.3	1.8 TeV	$m(\tilde{g}) > 100 \text{ GeV}$	1710030	
RPV		$\tilde{g}, \tilde{g} \rightarrow 3\gamma$	LFV prod.	0	Yes	0.2	6.8 TeV	$m(\tilde{g}) > 100 \text{ GeV}$	1710030
	$\tilde{g}, \tilde{g} \rightarrow 3\gamma$	Linear RPV CMSSM	0.1e	Yes	20.0	1.45 TeV	$m(\tilde{g}) > 100 \text{ GeV}$	1710030	
	$\tilde{g}, \tilde{g} \rightarrow 3\gamma$	$\tilde{g}, \tilde{g} \rightarrow q\tilde{q} + \tilde{g}\tilde{t}$	0	Yes	13.3	1.14 TeV	$m(\tilde{g}) > 100 \text{ GeV}, m(\tilde{u}) > 100 \text{ GeV}$	1710030	
	$\tilde{g}, \tilde{g} \rightarrow 3\gamma$	$\tilde{g}, \tilde{g} \rightarrow q\tilde{q} + \tilde{g}\tilde{t}$	0	Yes	0.19	208 GeV	$m(\tilde{g}) > 100 \text{ GeV}, m(\tilde{u}) > 100 \text{ GeV}$	1710030	
	$\tilde{g}, \tilde{g} \rightarrow 3\gamma$	$\tilde{g}, \tilde{g} \rightarrow q\tilde{q} + \tilde{g}\tilde{t}$	0	4.4e + 0.1e	36.1	1.072 TeV	$m(\tilde{g}) > 100 \text{ GeV}$	1710030	
	$\tilde{g}, \tilde{g} \rightarrow 3\gamma$	$\tilde{g}, \tilde{g} \rightarrow q\tilde{q} + \tilde{g}\tilde{t}$	1.0e	8-12 jets + 4 $\gamma$	36.1	2.1 TeV	$m(\tilde{g}) > 100 \text{ GeV}$	1710030	
	$\tilde{g}, \tilde{g} \rightarrow 3\gamma$	$\tilde{g}, \tilde{g} \rightarrow q\tilde{q} + \tilde{g}\tilde{t}$	1.0e	8-12 jets + 4 $\gamma$	36.1	1.25 TeV	$m(\tilde{g}) > 100 \text{ GeV}$	1710030	
	$\tilde{g}, \tilde{g} \rightarrow 3\gamma$	$\tilde{g}, \tilde{g} \rightarrow q\tilde{q} + \tilde{g}\tilde{t}$	0	0 jets + 2 $\gamma$	36.1	100-470 GeV	$m(\tilde{g}) > 100 \text{ GeV}$	1710030	
	$\tilde{g}, \tilde{g} \rightarrow 3\gamma$	$\tilde{g}, \tilde{g} \rightarrow q\tilde{q} + \tilde{g}\tilde{t}$	2.0e	2 $\gamma$	36.1	5.0-1.36 TeV	$m(\tilde{g}) > 100 \text{ GeV}$	1710030	
	$\tilde{g}, \tilde{g} \rightarrow 3\gamma$	Scalar charm, $\tilde{g} \rightarrow q\tilde{q}$	0	2 $\gamma$	Yes	20.3	910 GeV	$m(\tilde{g}) > 100 \text{ GeV}$	1710030

No clear sign of any deviations, or where to even test??

ow the maximum likelihood frequency-averaged th foreground and other nuisance parameters dele range  $2 \leq \ell \leq 29$ , we plot the power spectrum 4 % of the sky. The best-fit base  $\Lambda$ CDM theoret-l. Residuals with respect to this model are shown

\*Only a selection of the available mass limits on new states or phenomena is shown. Many of the limits are based on simplified models, c.f. refs. for the assumptions made.

# BY AND LARGE WE STILL HAVE ALL THE SAME PROBLEMS WE'VE HAD FOR DECADES

- Hierarchy Problem
- Dark Matter
- Matter anti-Matter asymmetry
- Neutrino Mass origin
- Strong CP problem
- Flavor
- Number of generations
- Apparent Unification of Coupling Constants
- Inflation
- Reheating
- Unification with Gravity
- Cosmological Constant Problem

EXPERIMENT TO THE RESCUE?

# THERE IS PROGRESS EXPERIMENTALLY!

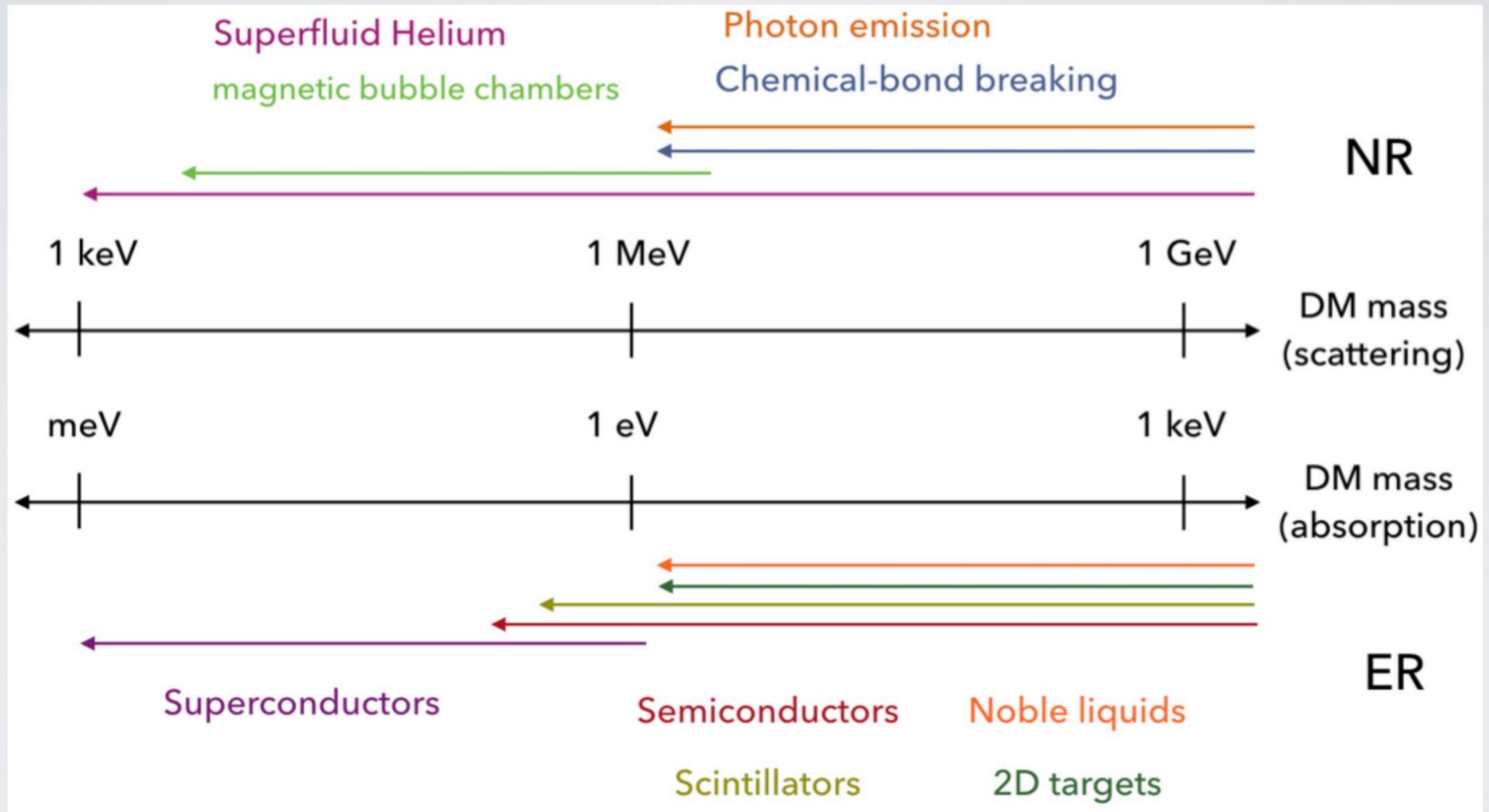


FIG. 4: Ideas to probe low-mass DM via scattering off, or absorption by, nuclei (NR) or electrons (ER).

**“Dark Sectors” abound... just ask Howie**

UN?FORTUNATELY WITH EXPERIMENTAL  
PROGRESS THERE IS **ALSO** THEORETICAL  
PROGRESS...

Wise professor tells entering graduate student 2002:

WIMPs are very motivated so it's likely:  $m_{DM} \sim m_{weak}$

# UNFORTUNATELY WITH EXPERIMENTAL PROGRESS THERE IS **ALSO** THEORETICAL PROGRESS...

Wise professor tells entering graduate student 2002:

WIMPs are very motivated so it's likely:  $m_{DM} \sim m_{weak}$

Wise professor tells entering graduate student 2018:

could be  $m_{DM} \sim 10^{-22} \text{ eV}$

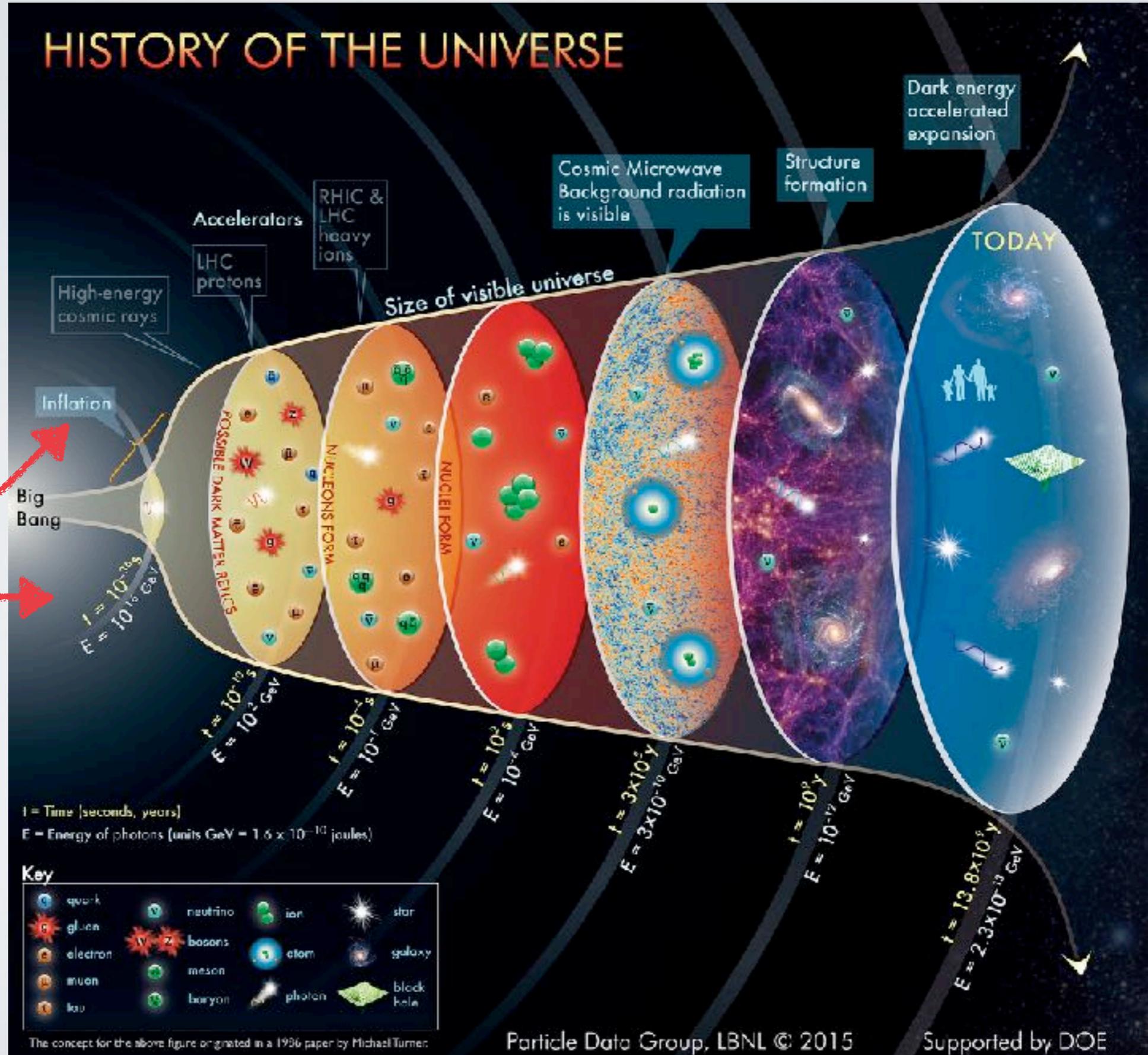
**but...** also could be **50** to **90** orders of magnitude heavier  
depending on assumptions of course

# LAMPPOST EFFECT



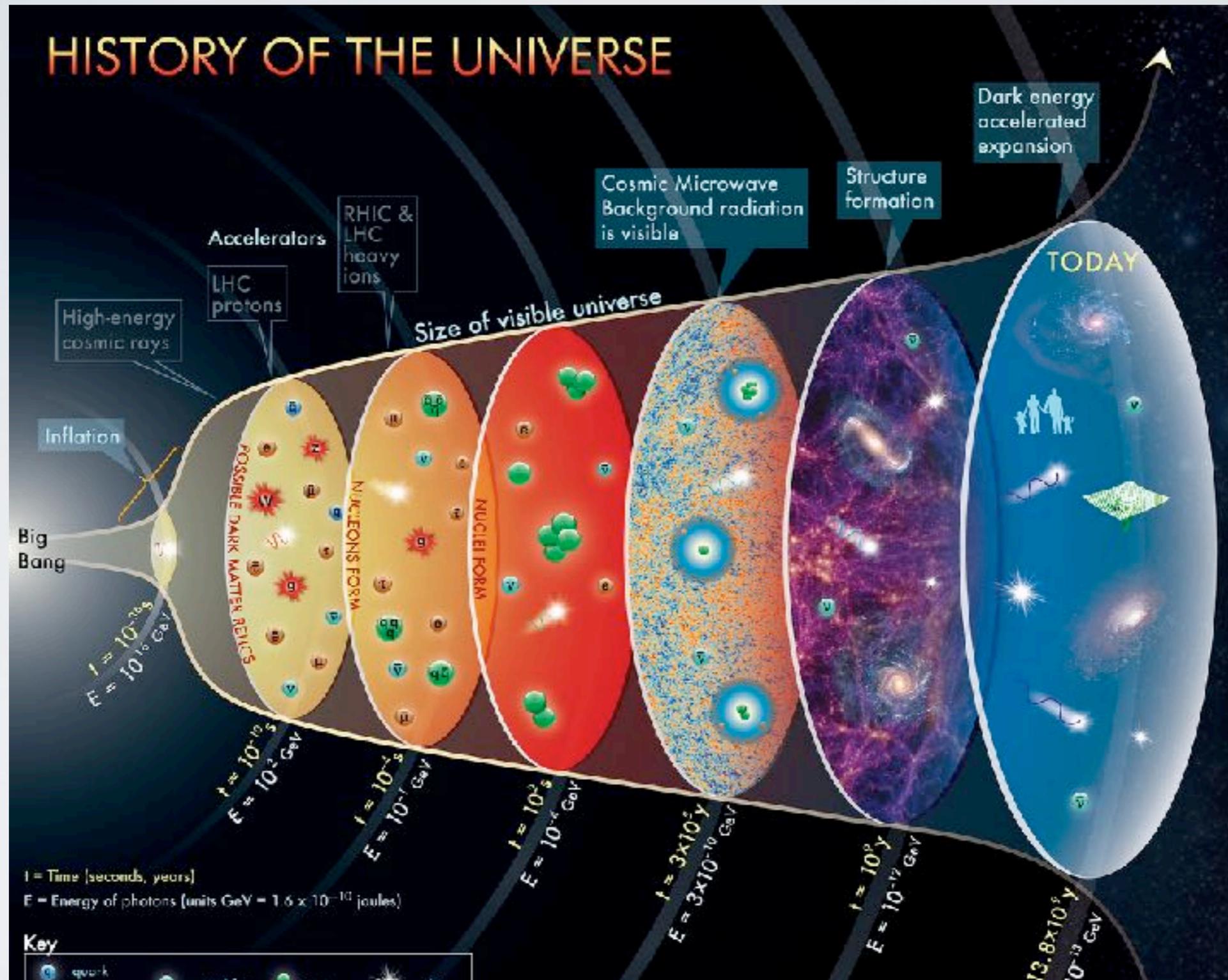
Not just limited to Dark Matter of course

# WHY NO NOBEL FOR INFLATION?



Inflation!

# WHY NO NOBEL FOR INFLATION?



We don't know **when** it happened to better than  $\sim 15$  orders of magnitude!!!

# LAMPPOST EFFECT



I'm not a wise professor nor lucky, so...

# LAMPPOST EFFECT



Is there a confluence of a lamppost **and** a theory motivation?



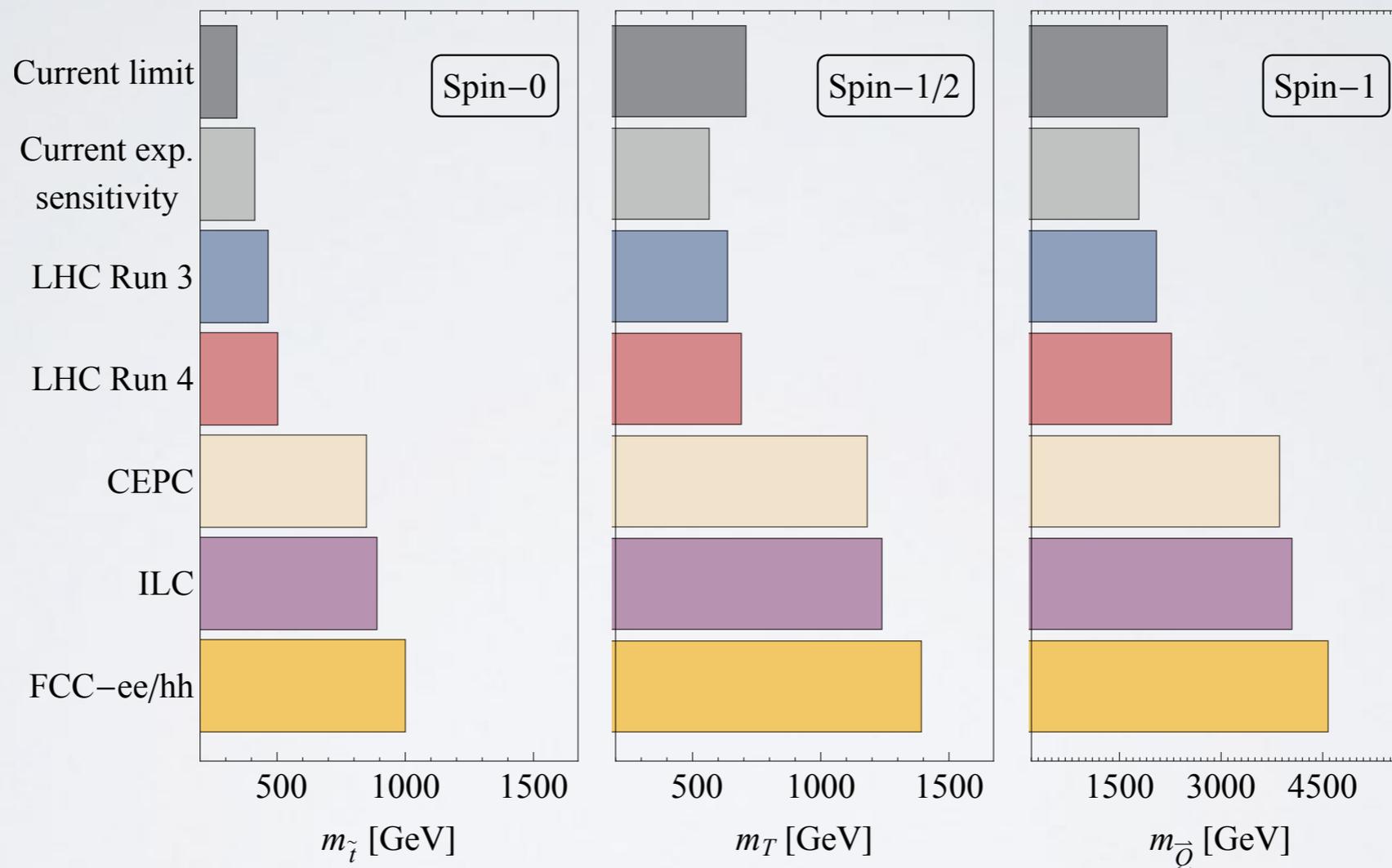
**Higgs  
Lamp**

# HIGGS LAMPPOST

- Naturalness
- Higgs Potential
- Higgs Portal to other sectors
- Cosmological History
- ...

# THIS LAMPPPOST IS LESS BRIGHT FOR SOME QUESTIONS

**Example: Colored Naturalness**



Essig, PM, Ramani, Zhong 1707.03399

**Have to get “lucky” or...**

THIS IS JUST ANOTHER LAMPPOST  
IN THE USUAL SENSE....

$$m_h^2 \sim \Lambda^2$$

$$\Lambda \sim \text{~~100~~ GeV}$$

$$\text{~~200~~ GeV}$$

$$\text{~~300~~ GeV}$$

$$\text{~~500~~ GeV}$$

$$\text{~~1000~~ GeV}$$

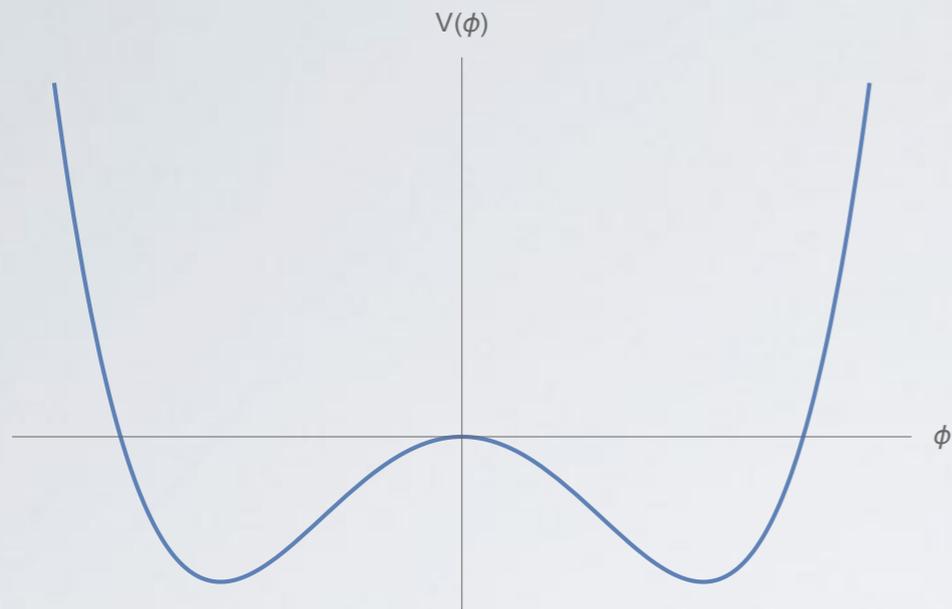
$$5000 \text{ GeV } \text{????}$$

We certainly learn something,  
but **what** it is telling us isn't as clear

# HIGGS POTENTIAL AND COSMOLOGICAL HISTORY

*Here there is a chance where  
**quantitative** measurements  
can yield **qualitative** differences!*

# HIGGS POTENTIAL



Higgs self  
interactions

OR

$$\left. \frac{\partial V(\phi)}{\partial \phi} \right|_{\phi=v} = 0$$

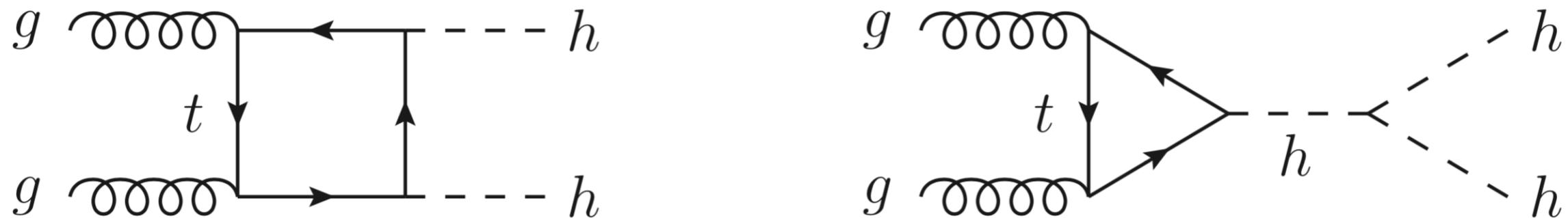
$$\left. \frac{\partial^2 V(\phi)}{\partial \phi^2} \right|_{\phi=v} = m_h^2$$

VEV and mass  
are *sufficient*

for Higgs potential, but BSM?



# NEXT UP IS THE TRIPLE HIGGS COUPLING IN THE SM...



Unfortunately it's very difficult and it interferes with itself

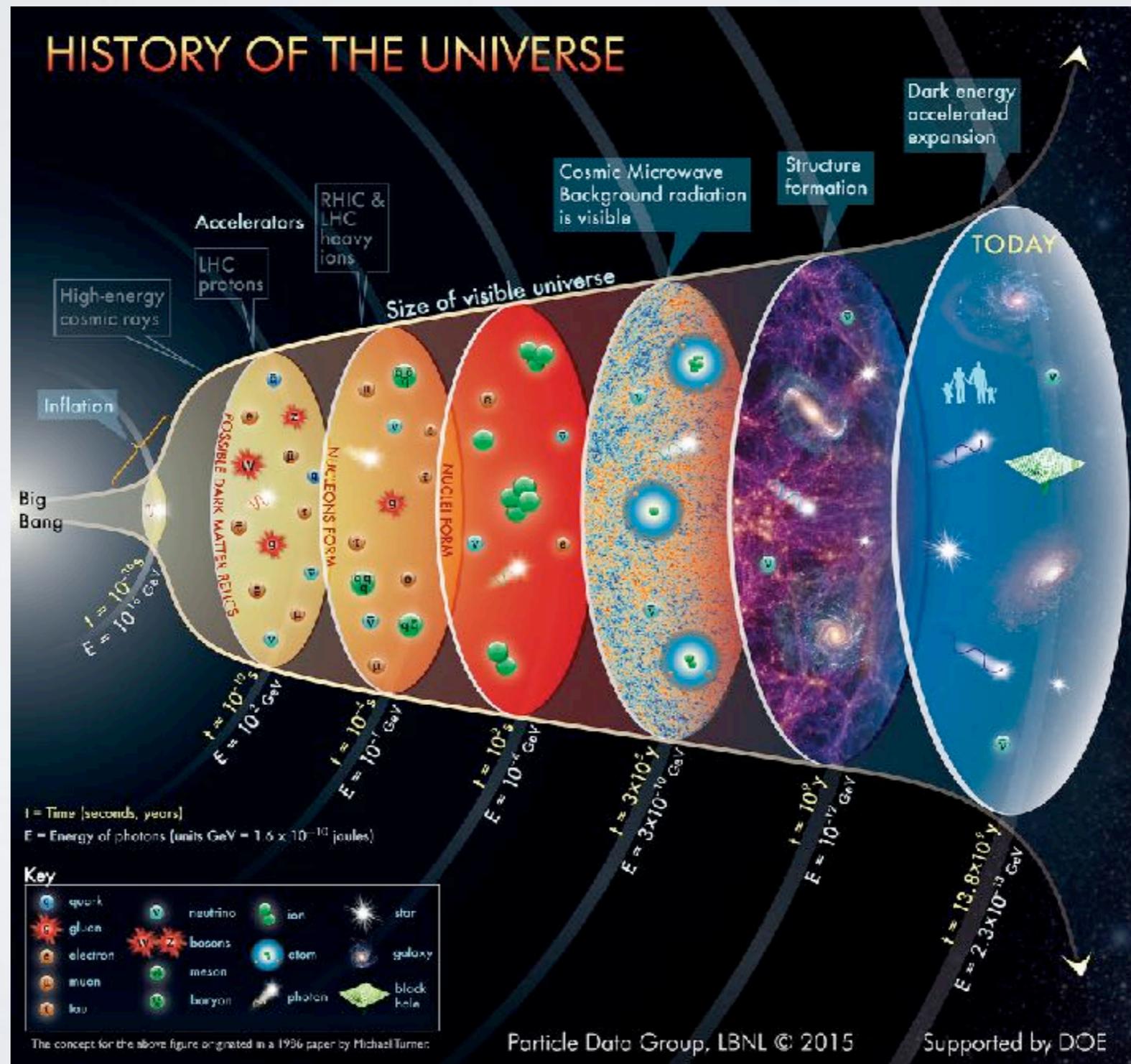
**However, just measuring the SM value  
would be seeing something qualitatively new!**

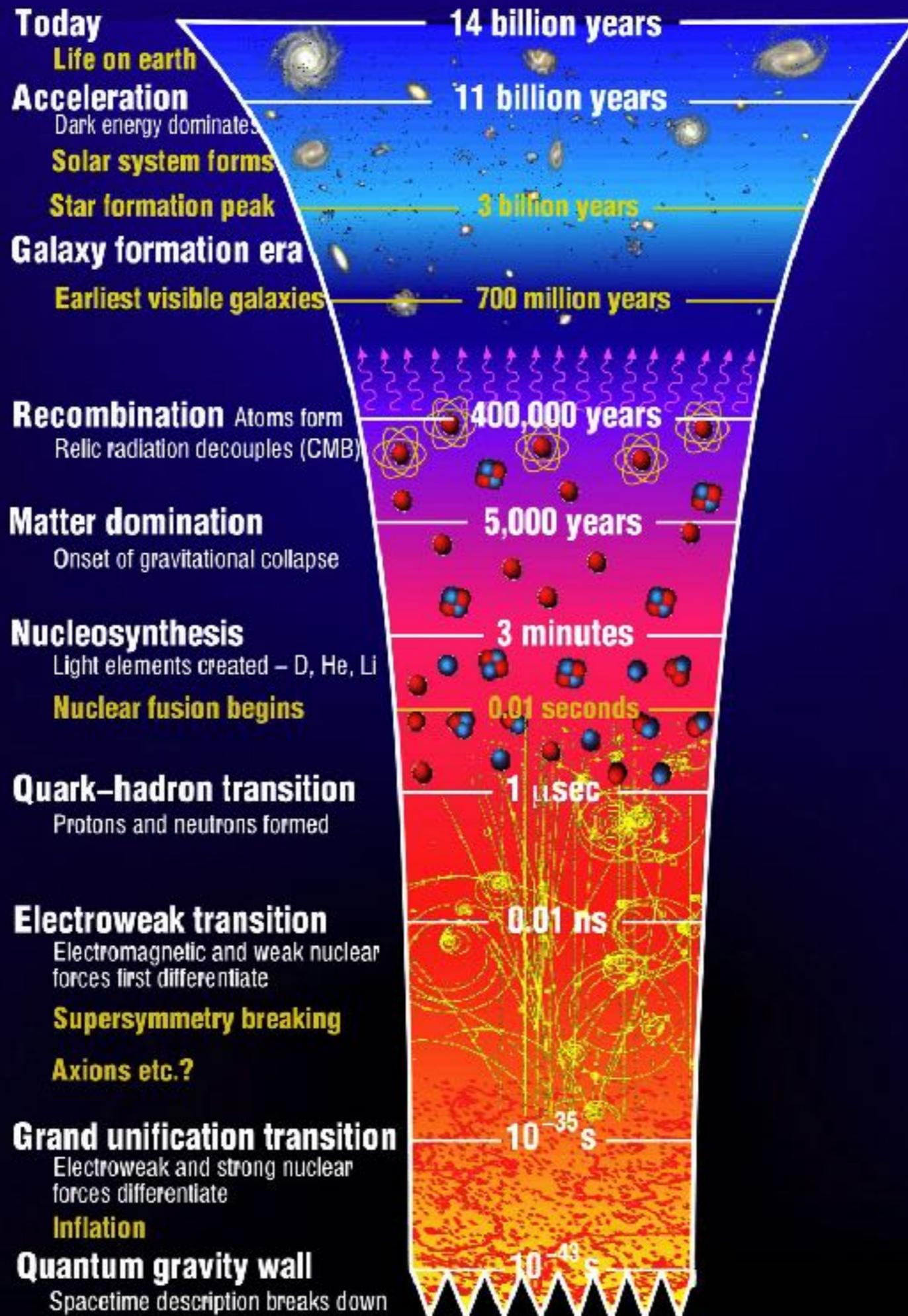
To go beyond though, is it *just* another lamppost?

Can be huge deviations...

How precisely do we *need* to measure it?

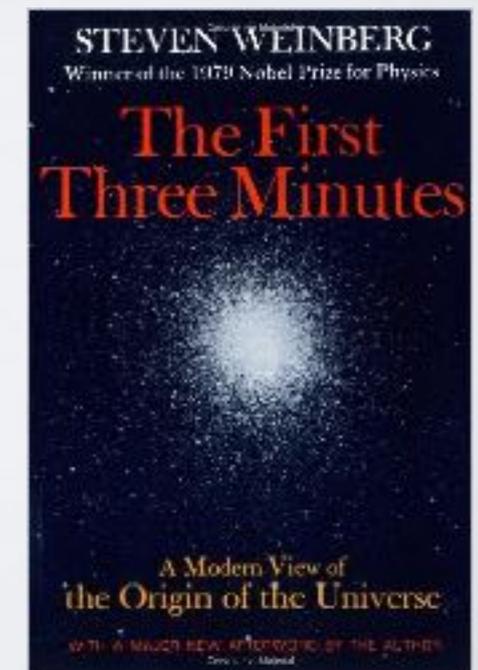
# WHAT IS OUR QUALITATIVE PICTURE OF THE COSMOLOGICAL HISTORY OF EWSB??





**Cosmology  
stuck here**

**Need particle  
physics to  
go further!**



# ELECTROWEAK PHASE TRANSITION

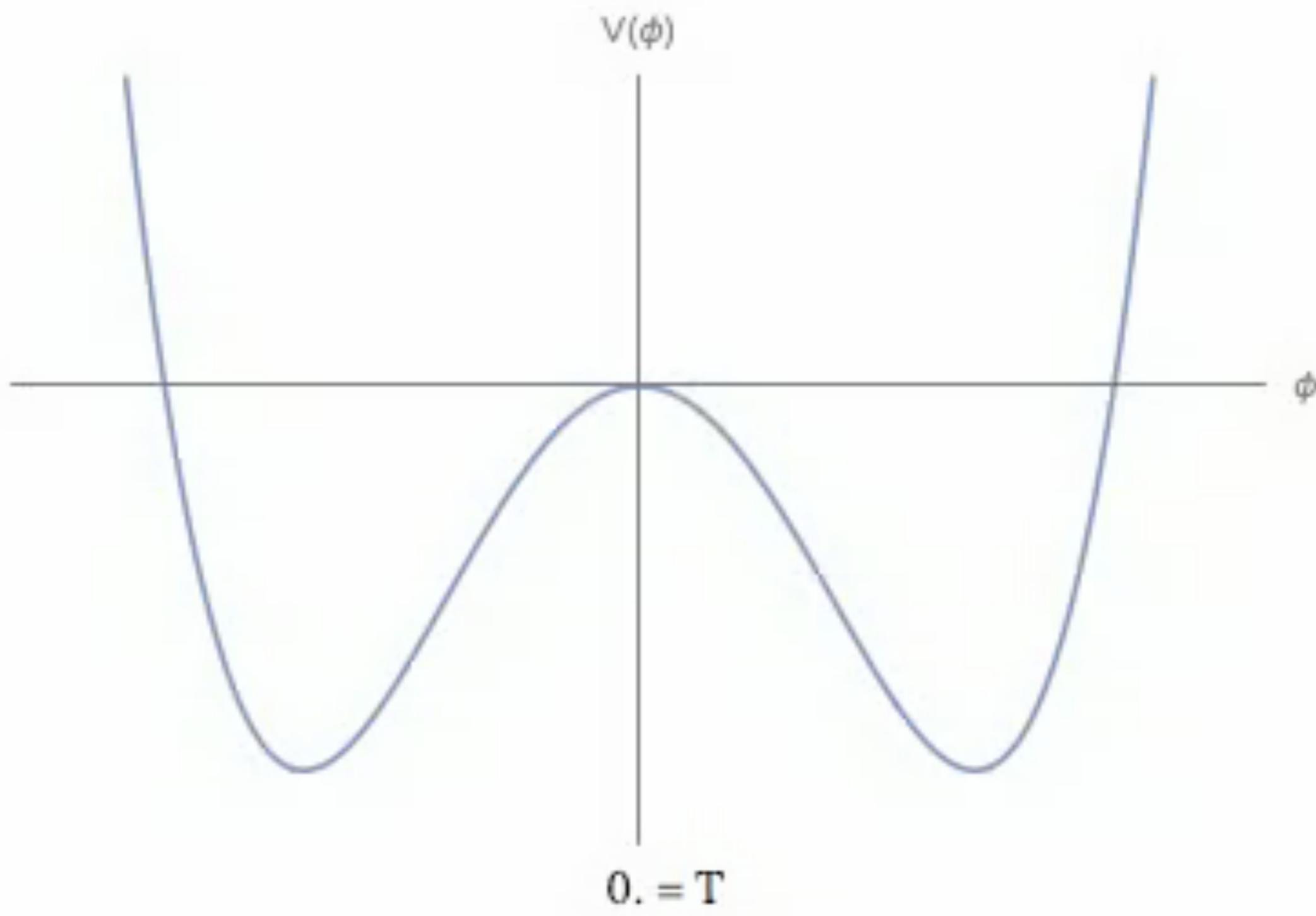
This heuristic picture of the cosmological history comes from analyzing a scalar potential at finite temperature which has been around for awhile...

# FINITE TEMPERATURE FIELD THEORY STARTS WITH SOME RUSSIANS IN 1972...



IF YOU HEAT UP A SYSTEM WITH A BROKEN SYMMETRY DOES THE SYMMETRY GET RESTORED?

WHAT'S THE CURIE TEMPERATURE OF THE UNIVERSE?



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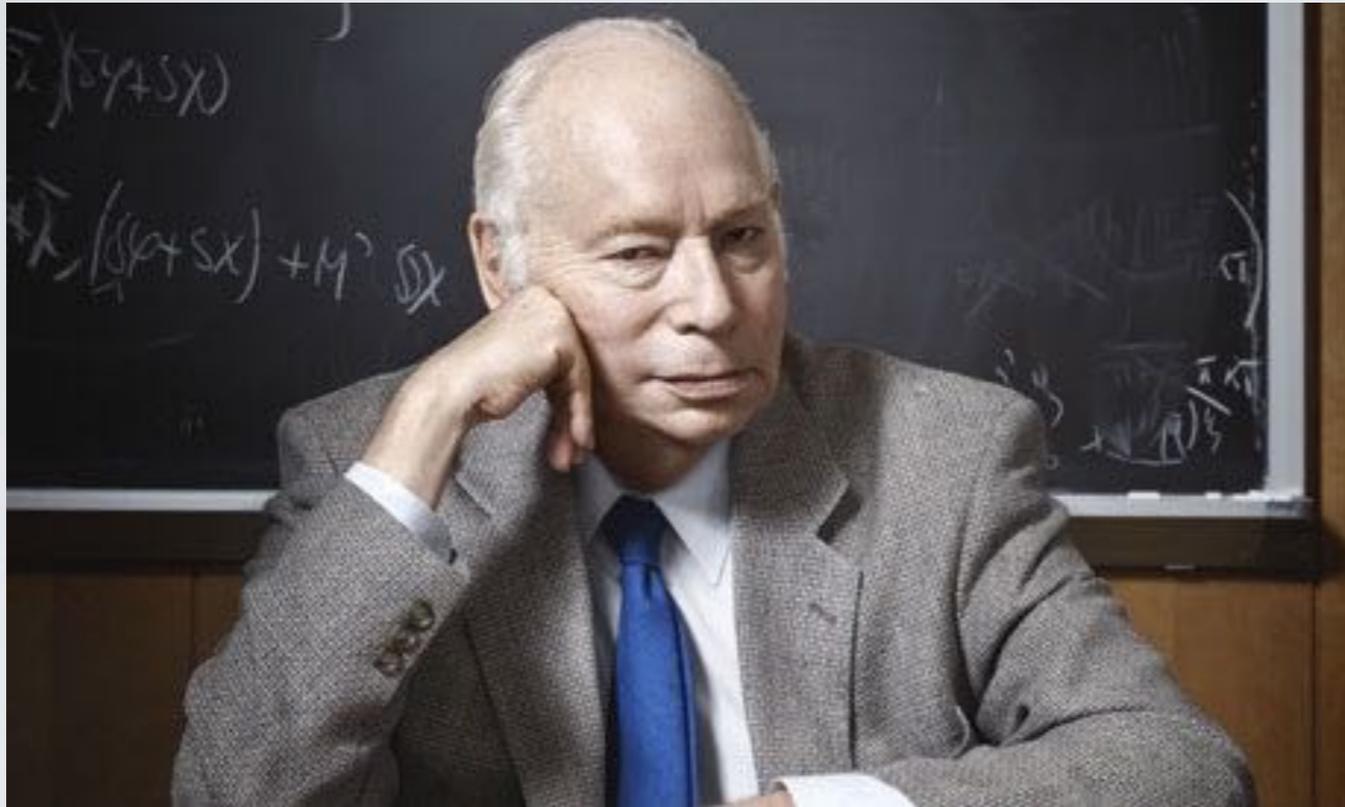
ANSWERS: YES, HMM...

There's no place  
like home,  
There's no place like home,  
There's no place like  
home...



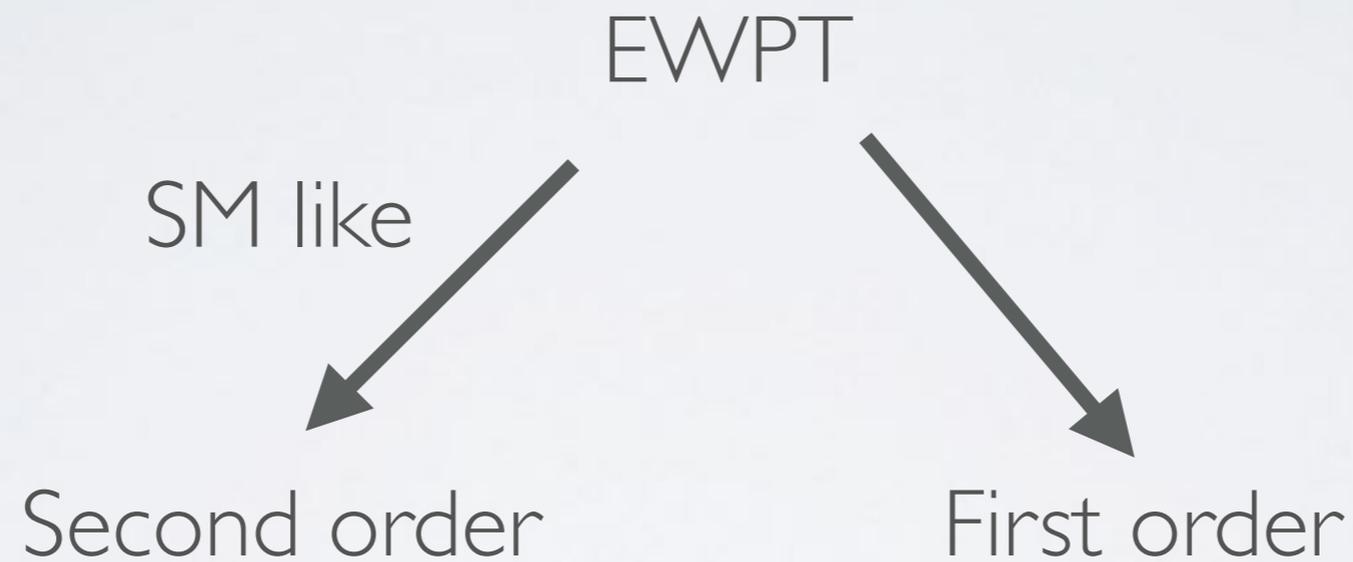
Weinberg,  
Weinberg,  
Weinberg...





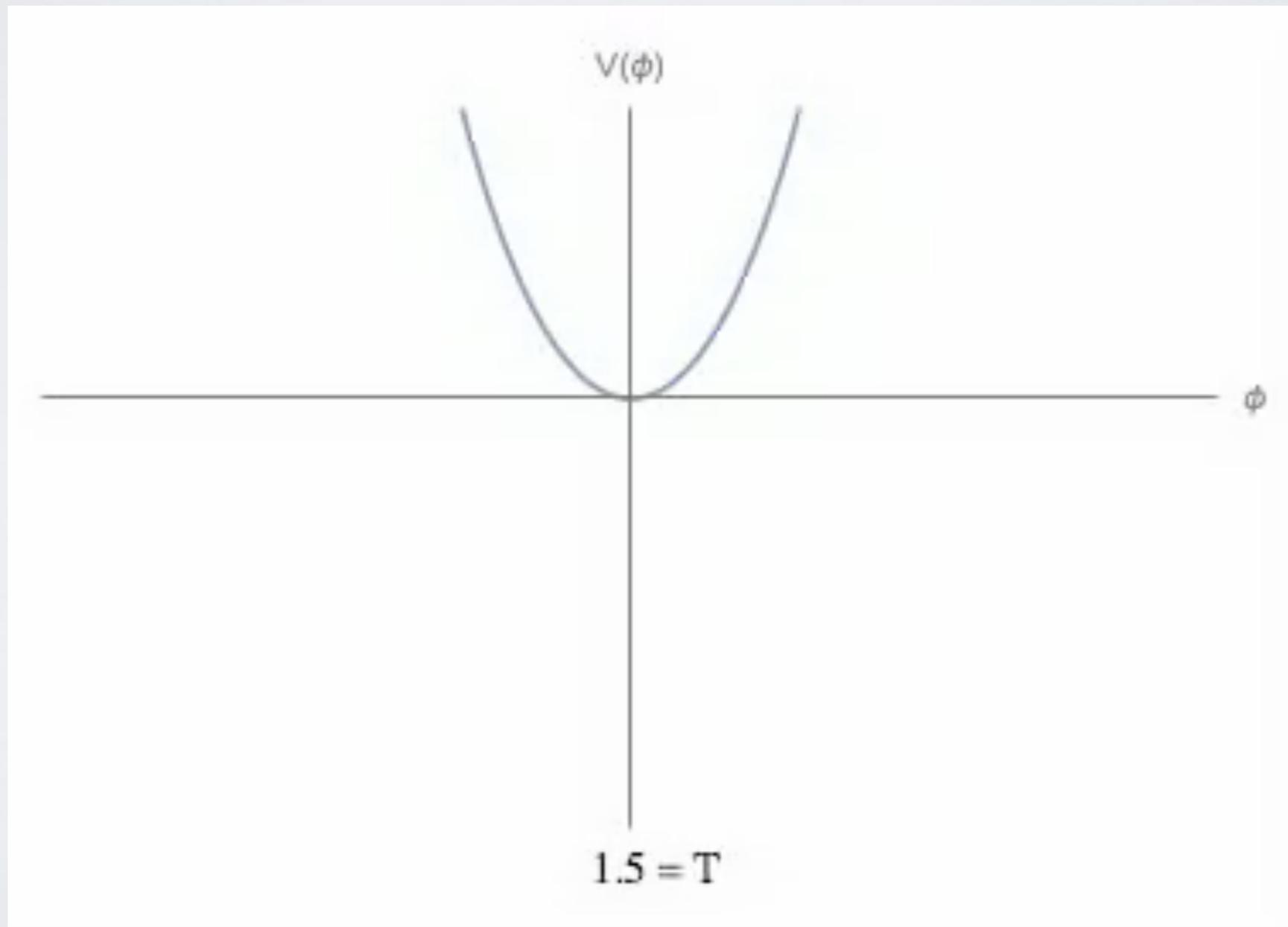
“A recent paper by Kirzhnits and Linde suggests that this is indeed the case. **However, although their title refers to a gauge theory, their analysis deals only with ordinary theories with broken global symmetries. Also, they estimate but do not actually calculate the critical temperature at which a broken symmetry is restored.**”

IF THERE'S AN EWPT HOW DO WE QUALITATIVELY DISTINGUISH?



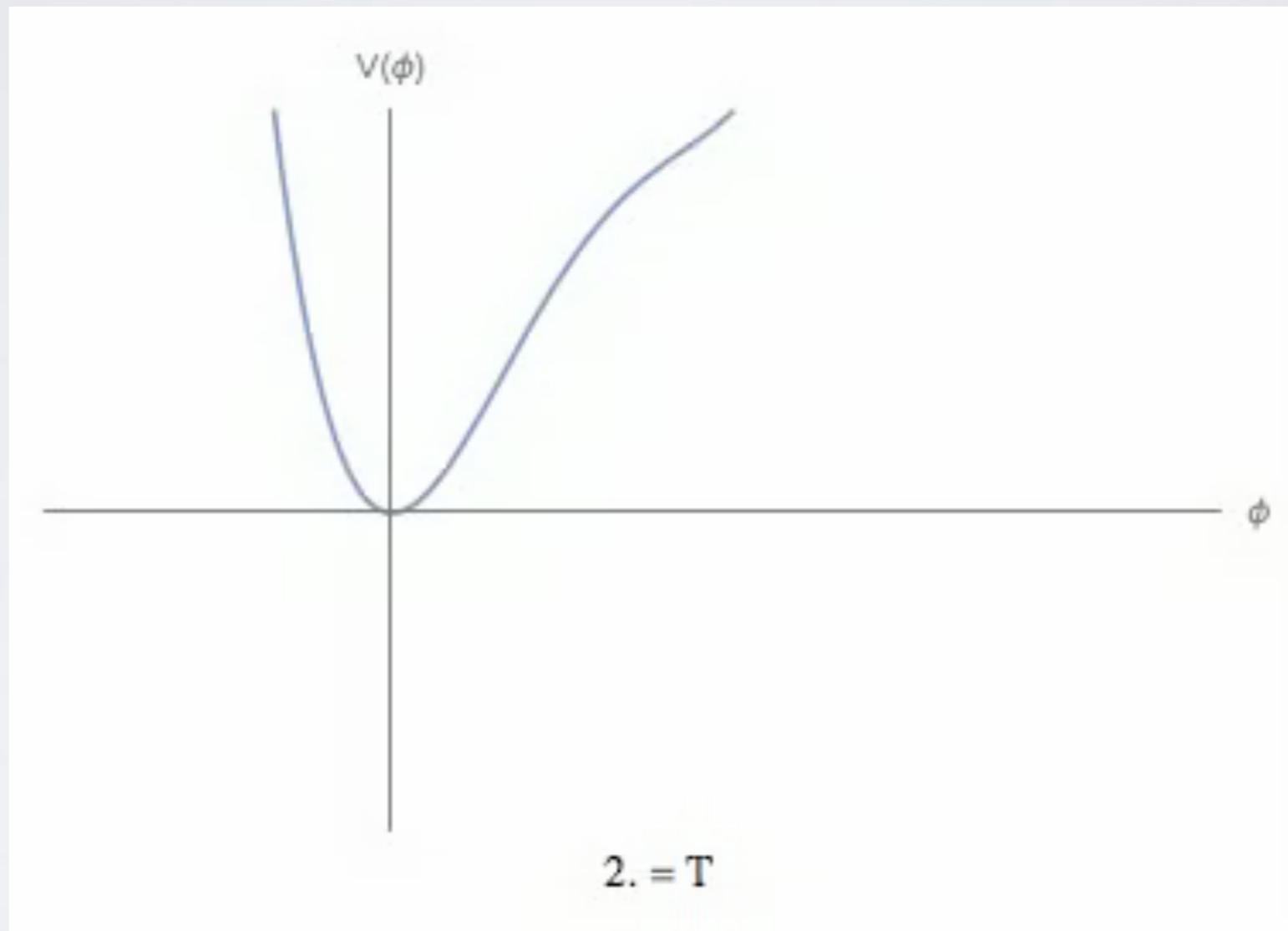
# 2ND ORDER PHASE TRANSITION

$$V(\phi, T) = D(T^2 - T_o^2)\phi^2 + \frac{\lambda(T)}{4}\phi^4$$



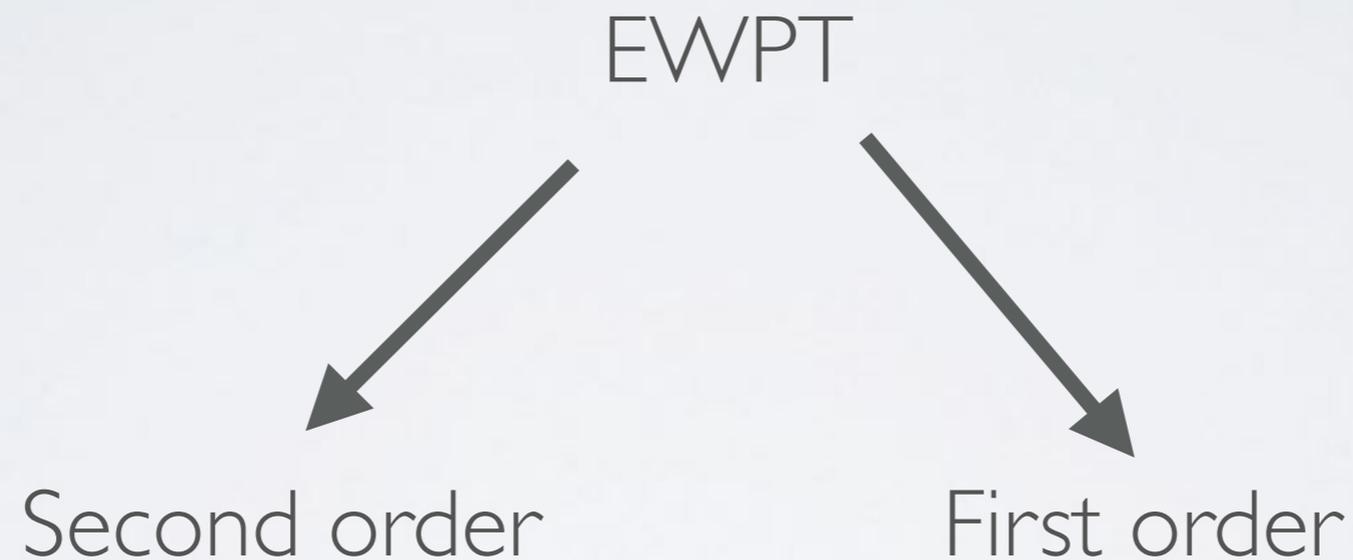
# 1ST ORDER PHASE TRANSITION

$$V(\phi, T) = D(T^2 - T_o^2)\phi^2 - ET\phi^3 + \frac{\lambda(T)}{4}\phi^4$$



A second minimum separated by a barrier!

# IF THERE'S AN EWPT HOW DO WE QUALITATIVELY DISTINGUISH?



The qualitative difference is an effective cubic at finite temperature!

Why is this so useful? Thermal Decoupling!  $e^{-\frac{m}{T}}$

# HIGGS POTENTIAL AND COSMOLOGICAL HISTORY

## **Triple Higgs + EWPT**

There **could** be a very large contribution and a FO EWPT - get lucky

There also **could** be a minimum contribution  
**within** experimental reach

# HIGGS POTENTIAL AND COSMOLOGICAL HISTORY

## Triple Higgs + EWPT

There **could** be a very large contribution and a FO EWPT - get lucky

There also **could** be a minimum contribution **within** experimental reach

### WHY?

- Has to couple to the Higgs strongly enough to affect potential
- Mass must not be too far away from EW scale!

### “No-Lose”

Can make this even sharper if you connect to Electroweak Baryogenesis

# THEORETICAL MINIMUM

## SM + SINGLET

This model has been studied numerous times for a variety of reason....

**If** the Singlet mixes with the Higgs you can see it easily via Higgs properties and has been studied quite a bit

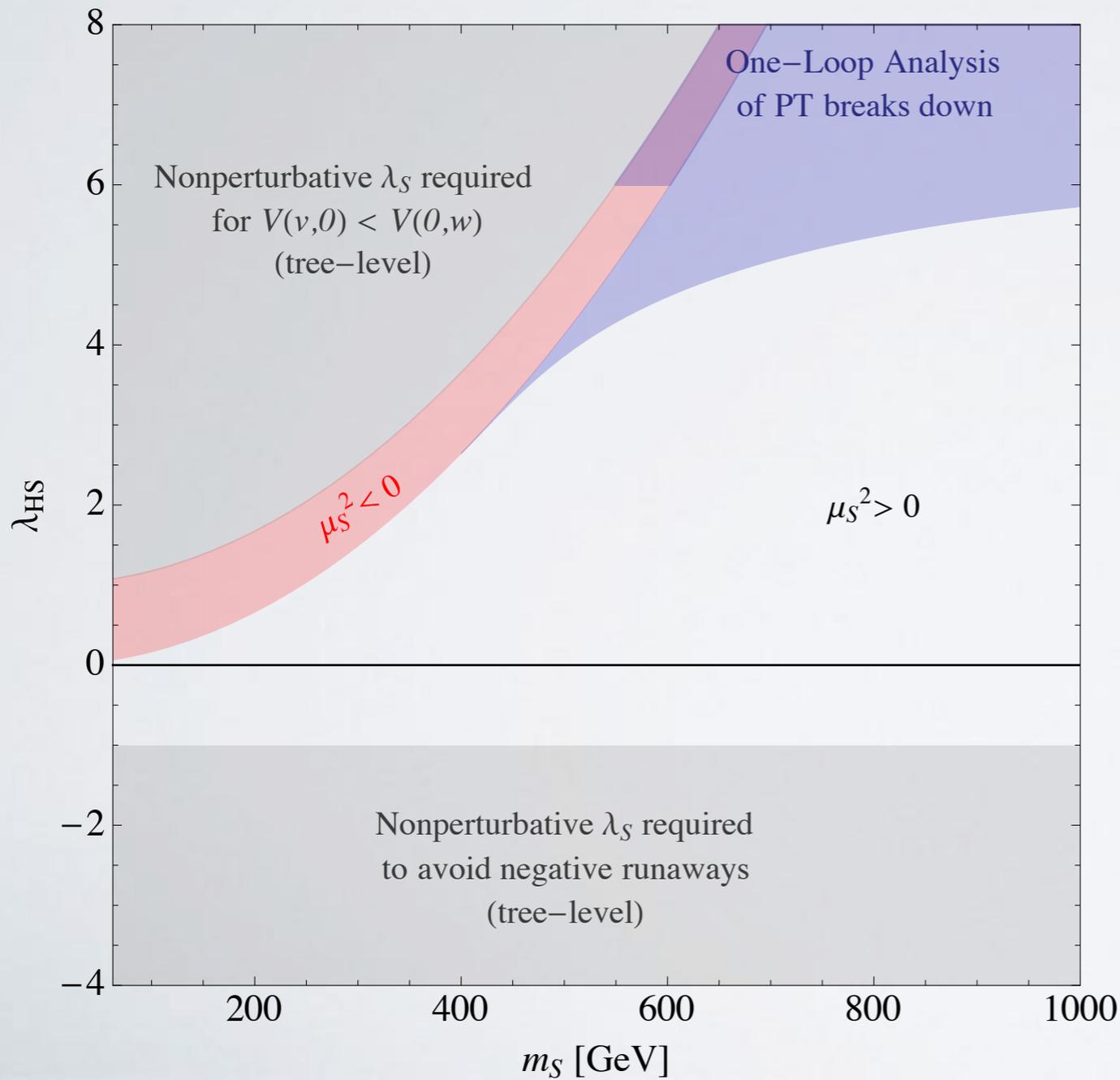
**If** the Singlet DOESN'T mix **but** its mass is less than half the Higgs mass you can see it in decays easily...

What if the singlet doesn't mix with the Higgs *and* is heavy?

# SM + SINGLET “NIGHTMARE SCENARIO”

D. Curtin, PM, T. Yu | 409.0005

$$V_0 = -\frac{1}{2}\mu^2 h^2 + \frac{1}{4}\lambda h^4 + \frac{1}{2}\mu_S^2 S^2 + \frac{1}{2}\lambda_{HS} h^2 S^2 + \frac{1}{4}\lambda_S S^4$$



$$m_S^2 = \mu_S^2 + \lambda_{HS} v^2 > 0$$

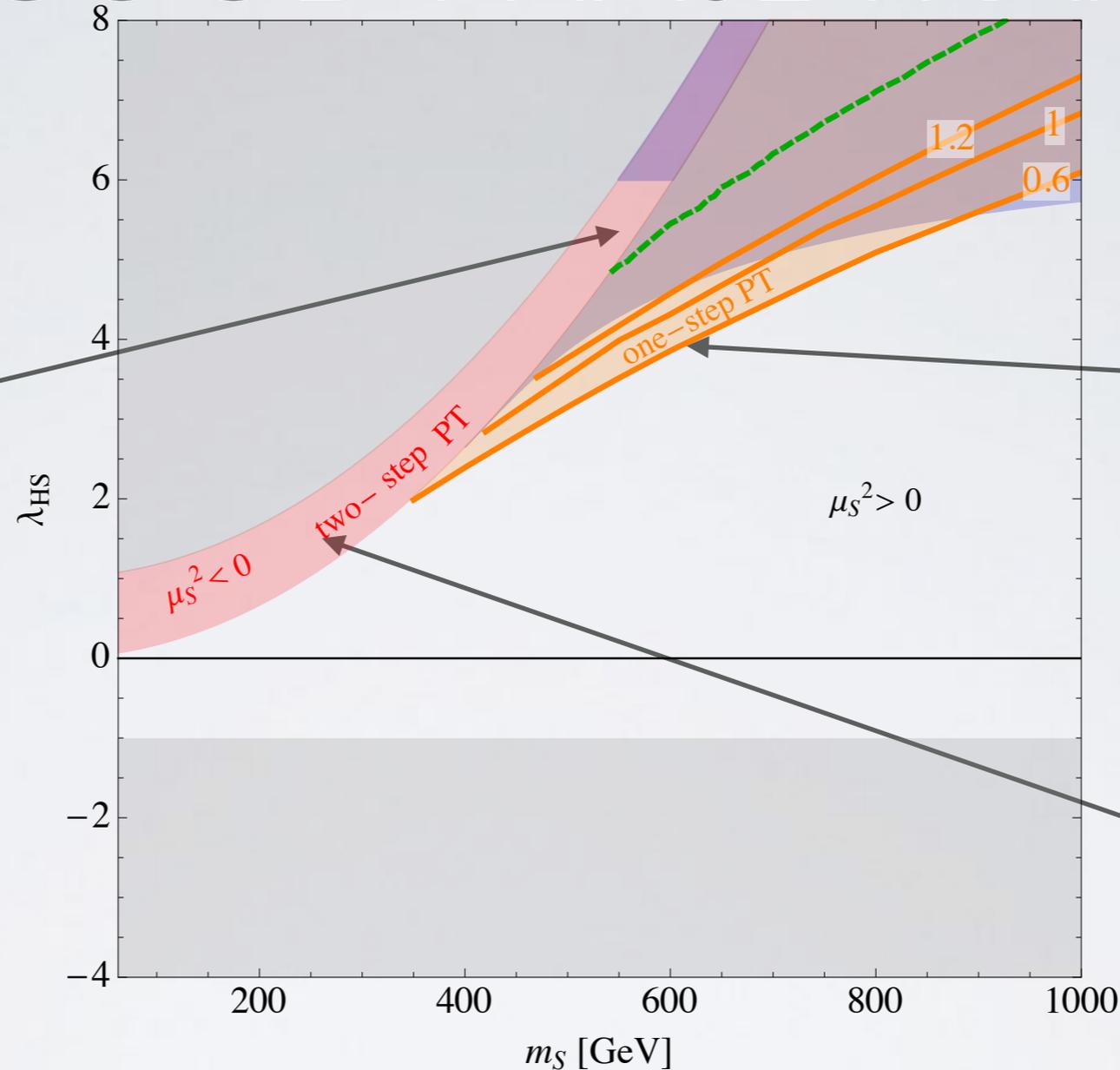
Phenomenological parameter space depends only on

$$(m_S, \lambda_{HS})$$

There are two qualitatively different regions depending on the sign of  $\mu_S^2$

# WHERE IN THE PARAMETER SPACE IS THERE A GOOD PHASE TRANSITION?

thermal  
1-step PT

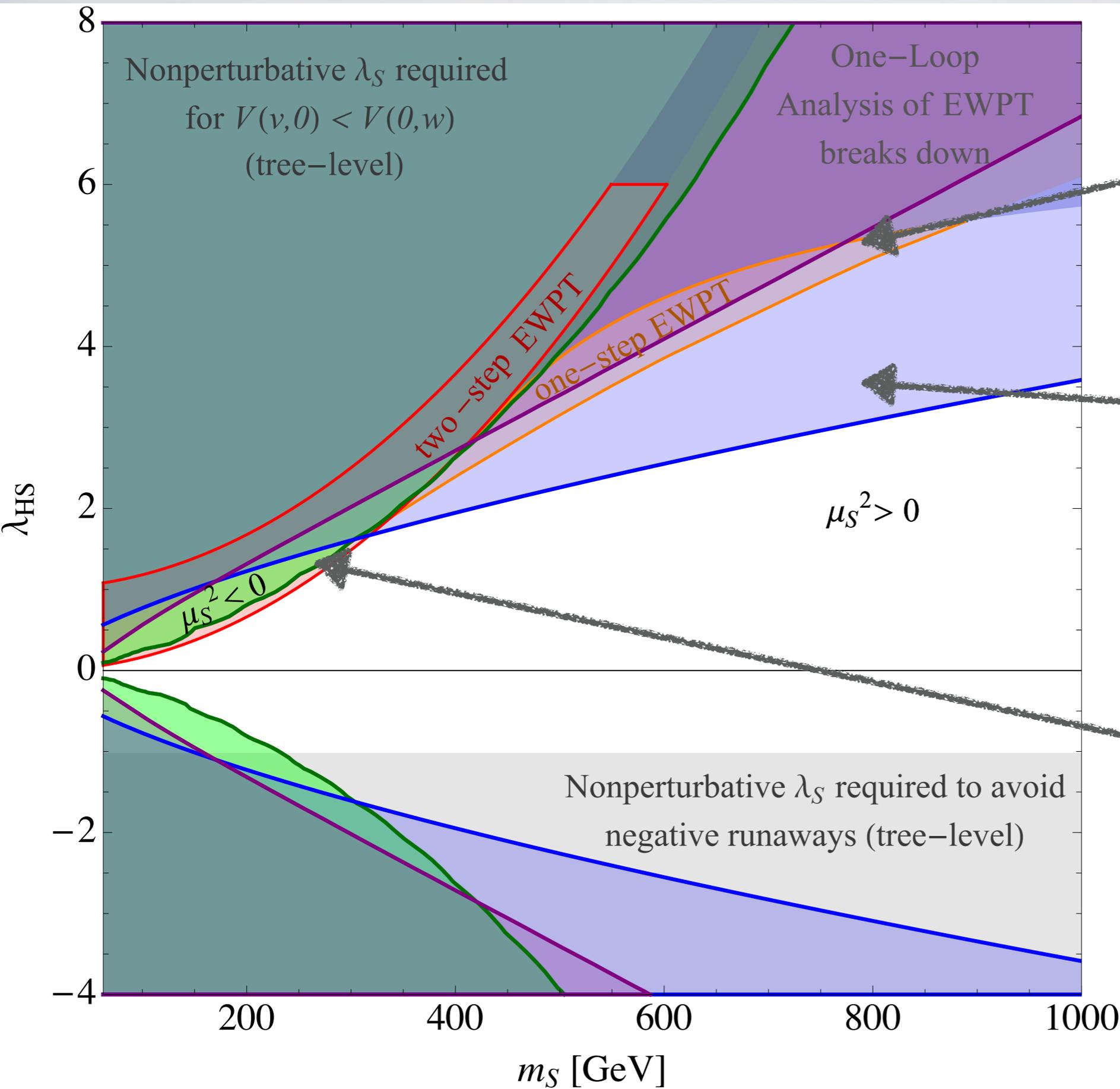


1-step PT  
zero T  
loops

2-step PT  
transition

**Figure 3.** Regions in the  $(m_S, \lambda_{HS})$  plane with viable EWBG. Red shaded region: for  $\mu_S^2 < 0$  it is possible to choose  $\lambda_S$  such that EWBG proceeds via a tree-induced strong two-step electroweak phase transition. Orange contours: value of  $v_c/T_c$  for  $\mu_S^2 > 0$ . The orange shaded region indicates  $v_c/T_c > 0.6$ , where EWBG occurs via a loop-induced strong one-step phase transition. Above the green dashed line, singlet loop corrections generate a barrier between  $h = 0$  and  $h = v$  even at  $T = 0$ , but results in the dark shaded region might not be reliable, see Section 3.1.3.

# A "NO-LOSE" THEOREM



TLEP exclusion with Zh shift

2 Sigma Exclusion with triple Higgs at 100 TeV

Collider

$$\frac{S}{\sqrt{B}} \geq 2$$

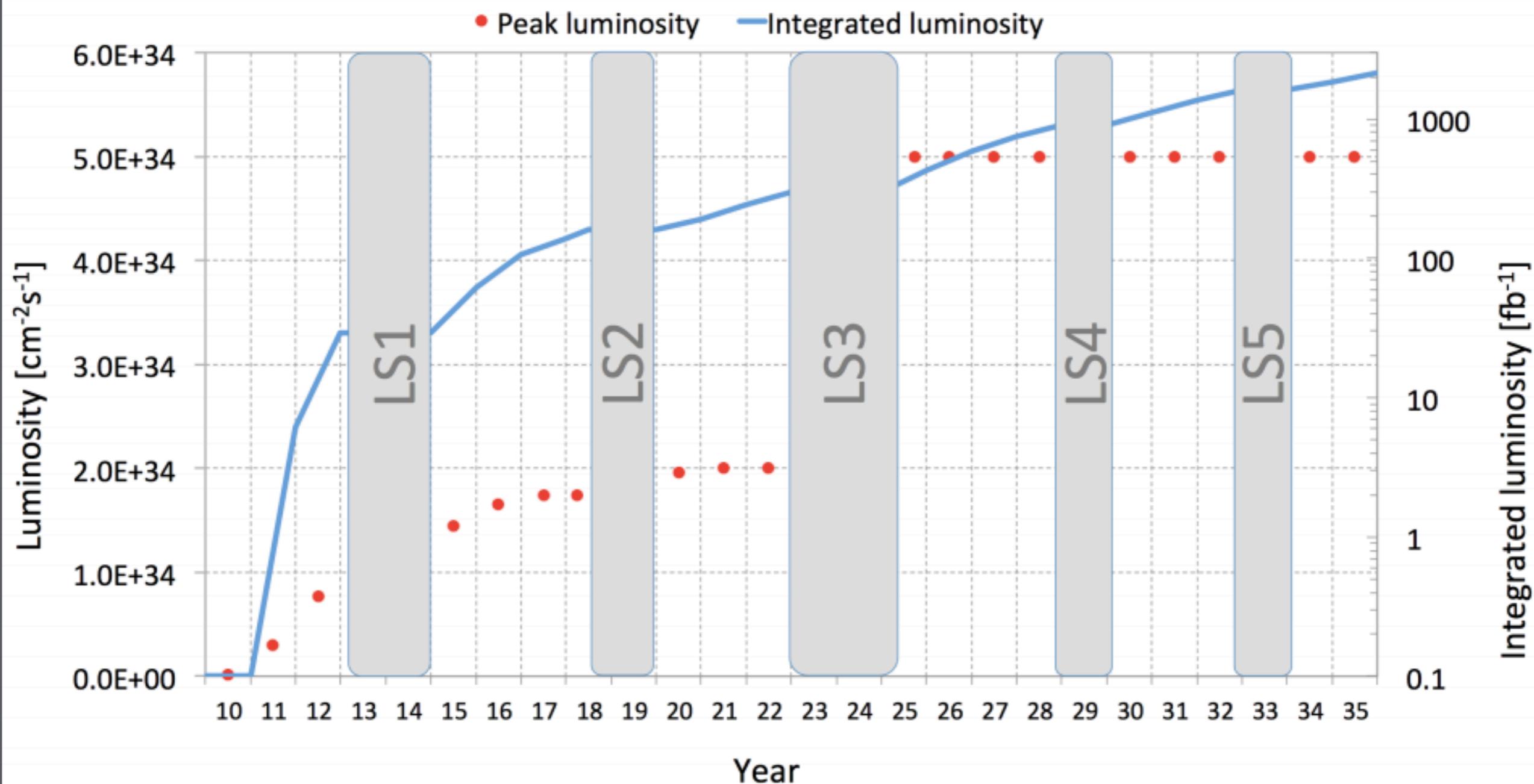
for 100 TeV 30/ab

# TRIPLE HIGGS

- Experimentally there are a number of different probes but the triple Higgs coupling does the heavy lifting
- In that study we used an assumed sensitivity of 10% on triple Higgs with 30/ab @ 100 TeV
- Strong FO EWPT typically naively has a 20-30% shift in this scenario
- A 100 TeV collider is a lot of money and a lot of time, what are the other possibilities?

# TRIPLE HIGGS MEASUREMENTS

## 2010 - 2035



# TRIPLE HIGGS MEASUREMENTS

- By 2035 we won't be able to tell the triple Higgs coupling compared to the SM better than

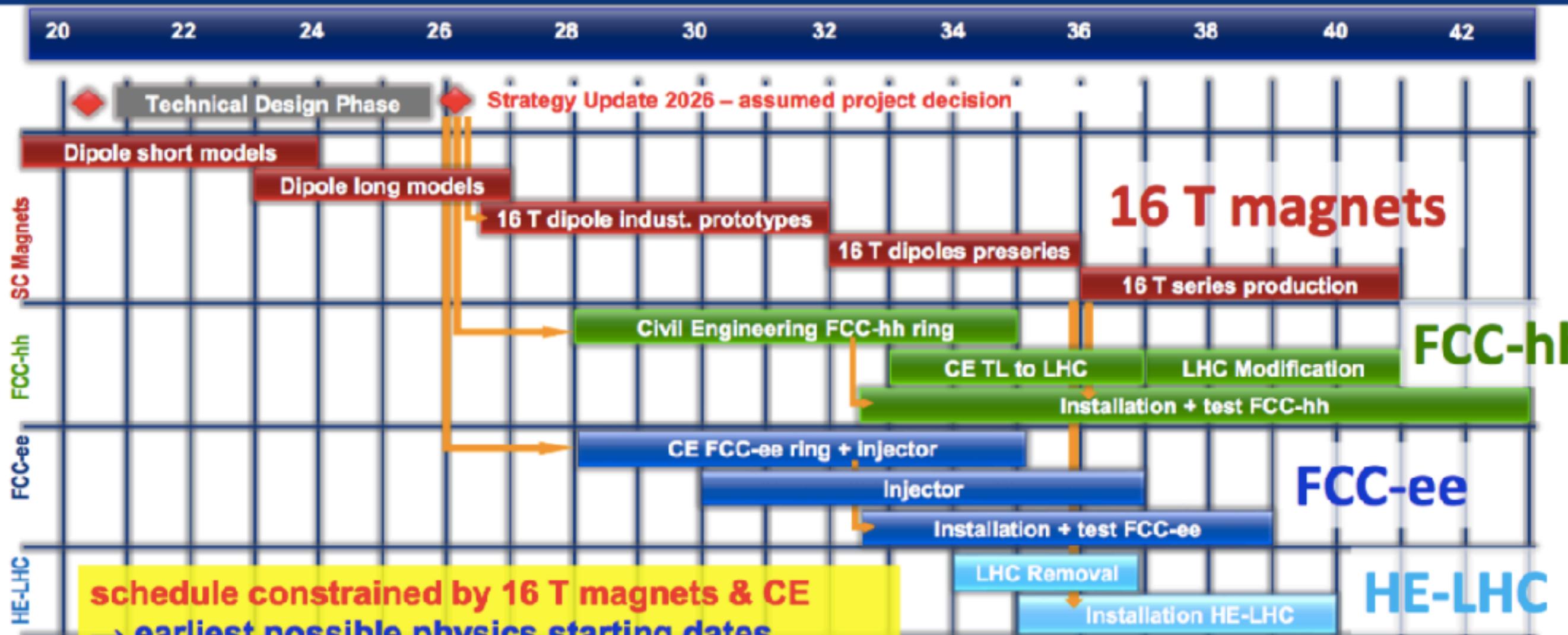
$$-0.8 < \lambda_3 < 7.7 \text{ at } 95\% \text{ C.L.}$$

Ugh...

# TRIPLE HIGGS MEASUREMENTS



## Technical Schedule for each the 3 Options



**schedule constrained by 16 T magnets & CE**  
→ earliest possible physics starting dates

- FCC-hh: 2043
- FCC-ee: 2039
- HE-LHC: 2040 (with HL-LHC stop LS5 / 2034)

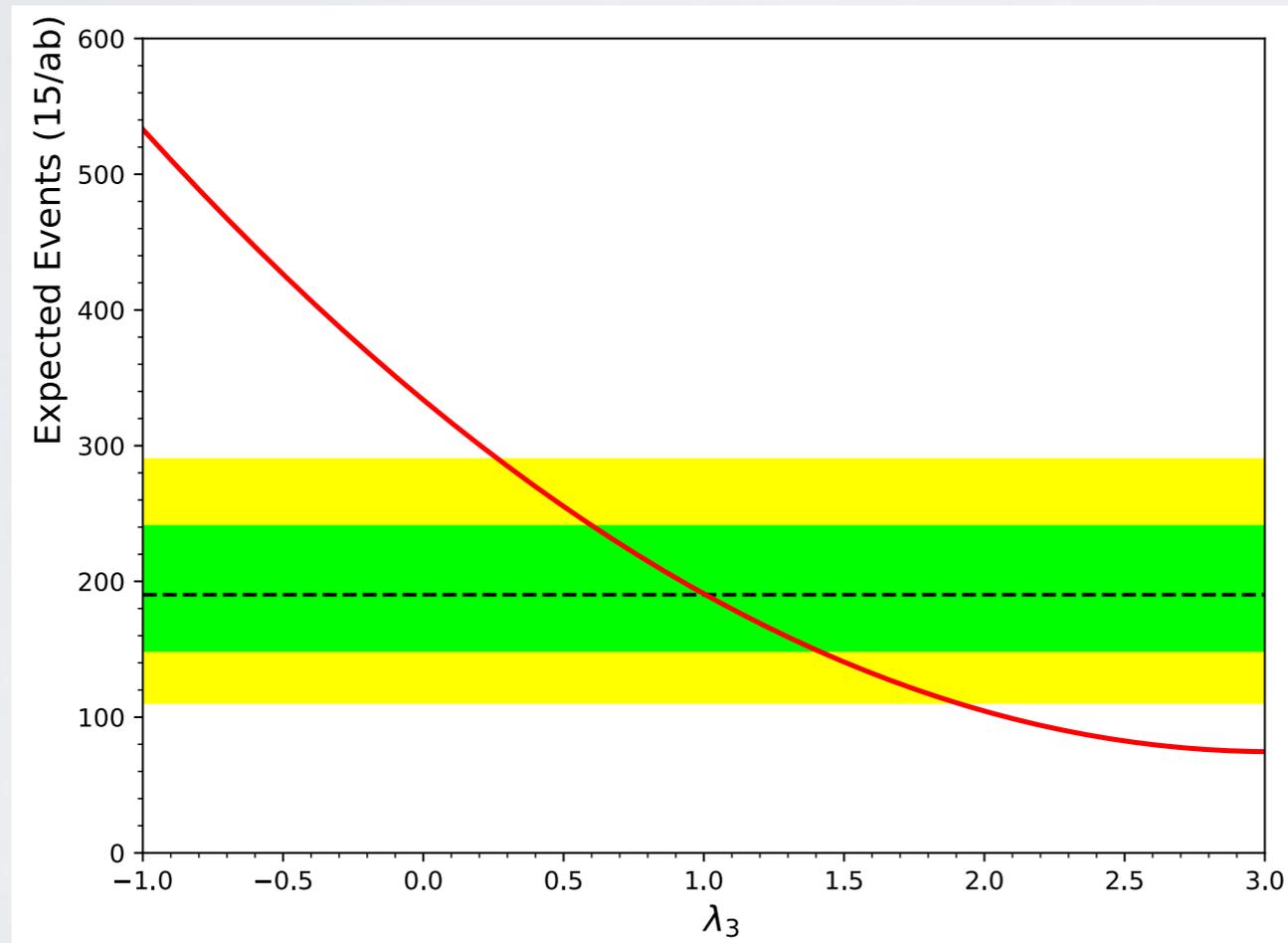
Chinese can't be faster?? although maybe politically more feasible...

# TRIPLE HIGGS MEASUREMENTS

- Refined studies show that FCC-hh could get to 1.6%-few% precision, but these are missing some backgrounds and need some work
- What about HE-LHC?

# HE-TRIPLE HIGGS MEASUREMENTS

- Han & Plehn et al claim you can get 5 sigma significance and a precision of 30%



Nevertheless HE-LHC  
does have a lot to say!

Homiller, PM to appear: a little more conservative results  
Most non-experimental studies have left out key backgrounds

NICE SIMPLE STORY...  
MEASURE TRIPLE HIGGS WELL ENOUGH  
YOU KNOW THE HISTORY OF THE  
UNIVERSE TO AN EARLIER TIME...

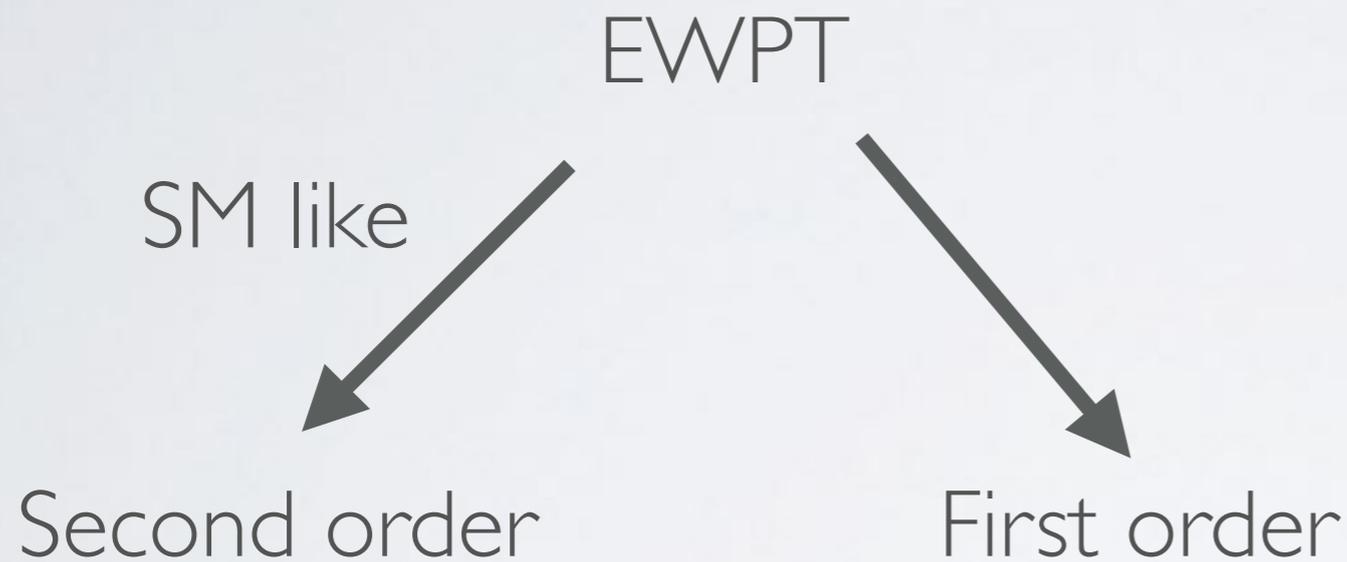
NICE SIMPLE STORY...  
MEASURE TRIPLE HIGGS WELL ENOUGH  
YOU KNOW THE HISTORY OF THE  
UNIVERSE TO AN EARLIER TIME...

**NO, NOT THAT SIMPLE!**

# THEORETICAL PROGRESS!



# EXPERIMENTAL TESTS OF COSMOLOGICAL HISTORY



## **third possibility**

The EWSB was never restored or it was delayed, or there were multiple EW phase transitions!!

Symmetry Non-Restoration

SNR phase

PM, Ramani 1807.07578

# SYMMETRY-NON RESTORATION

- Weinberg in his original finite-T paper noted counter examples
  - Rochelle salts
  - $O(N) \times O(M)$  model
- Since been verified on the lattice and with various other methods!

# VERY SIMPLE TO SEE WHERE IT COMES FROM...

$$V \sim (T^2 - \mu^2)\phi^2 + \lambda\phi^4$$

This comes from a term

$$V \supset \Pi_\phi \phi^2 \quad \Pi_\phi \sim \lambda T^2$$

In a more general theory, e.g. for the Higgs

$$\Pi_h = T^2 \left( \frac{\lambda_t^2}{4} + \frac{3g^2}{16} + \frac{g'^2}{16} + \frac{\lambda}{2} \right)$$

# NOW LET'S TAKE OUR SIMPLE SINGLET MODEL...

$$V_0 = -\frac{1}{2}\mu^2 h^2 + \frac{1}{4}\lambda h^4 + \frac{1}{2}\mu_S^2 S^2 + \frac{1}{2}\lambda_{HS} h^2 S^2 + \frac{1}{4}\lambda_S S^4$$

$$\Pi_h = T^2 \left( \frac{\lambda_t^2}{4} + \frac{3g^2}{16} + \frac{g'^2}{16} + \frac{\lambda}{2} + \frac{\lambda_{HS}}{12} \right)$$

and flip a sign...

$$\Pi_h = T^2 \left( \frac{\lambda_t^2}{4} + \frac{3g^2}{16} + \frac{g'^2}{16} + \frac{\lambda}{2} - \frac{\lambda_{HS}}{12} \right)$$

IF THE SINGLET DOMINATES WE HAVE A QUALITATIVELY DIFFERENT PICTURE...

$$V \sim -(\mu^2 + T^2)h^2 + \lambda h^4$$

The VEV **increases** with temperature!

$$\langle h \rangle \sim T$$

The EW symmetry is **never** restored in the early universe

# HOW WAS THIS MISSED?

- It's not quite so trivial, as you still have to make sure your calculation is under control and you have a good vacuum

$$\lambda_{HS}^2 \leq \lambda_s \lambda$$

So to satisfy this and dominate the thermal mass you run into non-perturbativity very quickly with the s quartic

$$\Pi_h = T^2 \left( \frac{\lambda_t^2}{4} + \frac{3g^2}{16} + \frac{g'^2}{16} + \frac{\lambda}{2} - \frac{\lambda_{HS}}{12} \right)$$

# SIMPLE TRICK - SWITCH TO $O(N)$ SINGLET

$$\Pi_h = T^2 \left( \frac{\lambda_t^2}{4} + \frac{3g^2}{16} + \frac{g'^2}{16} + \frac{\lambda}{2} - N_s \frac{\lambda_{HS}}{12} \right)$$

$$\Pi_s = T^2 \left( (N_s + 2) \frac{\lambda_s}{12} - \frac{\lambda_{HS}}{3} \right)$$

Now it can dominate the thermal mass but keep the potential stable  
for small  $\lambda_{HS}$

A rough estimate yields  $\lambda_{HS} N_s \geq 4.8$

$$\lambda_{HS}^c \equiv \lambda_{HS} N_s$$

# UNDER CONTROL...

You don't run into issues with the pure singlet quartic

$$\lambda_s \geq \left( \frac{\lambda_{HS}^c}{N_s} \right)^2 \frac{1}{\lambda}$$

$\beta$  function running of the couplings stable to high scales

# DOING THIS MORE CAREFULLY

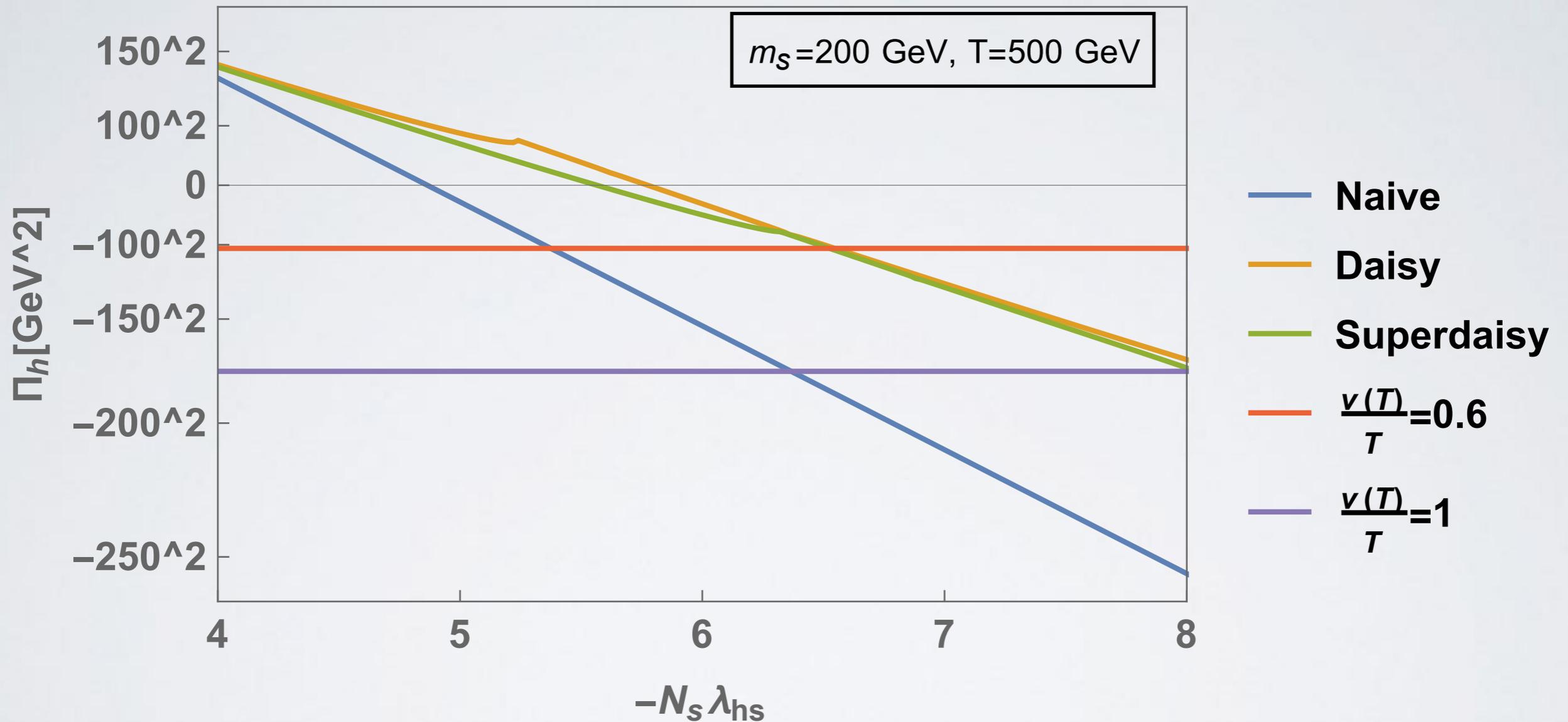
- Must take into account resummation and finite mass effects correctly - Optimized Partial Dressing

D. Curtin, PM, H. Ramani 1612.00466

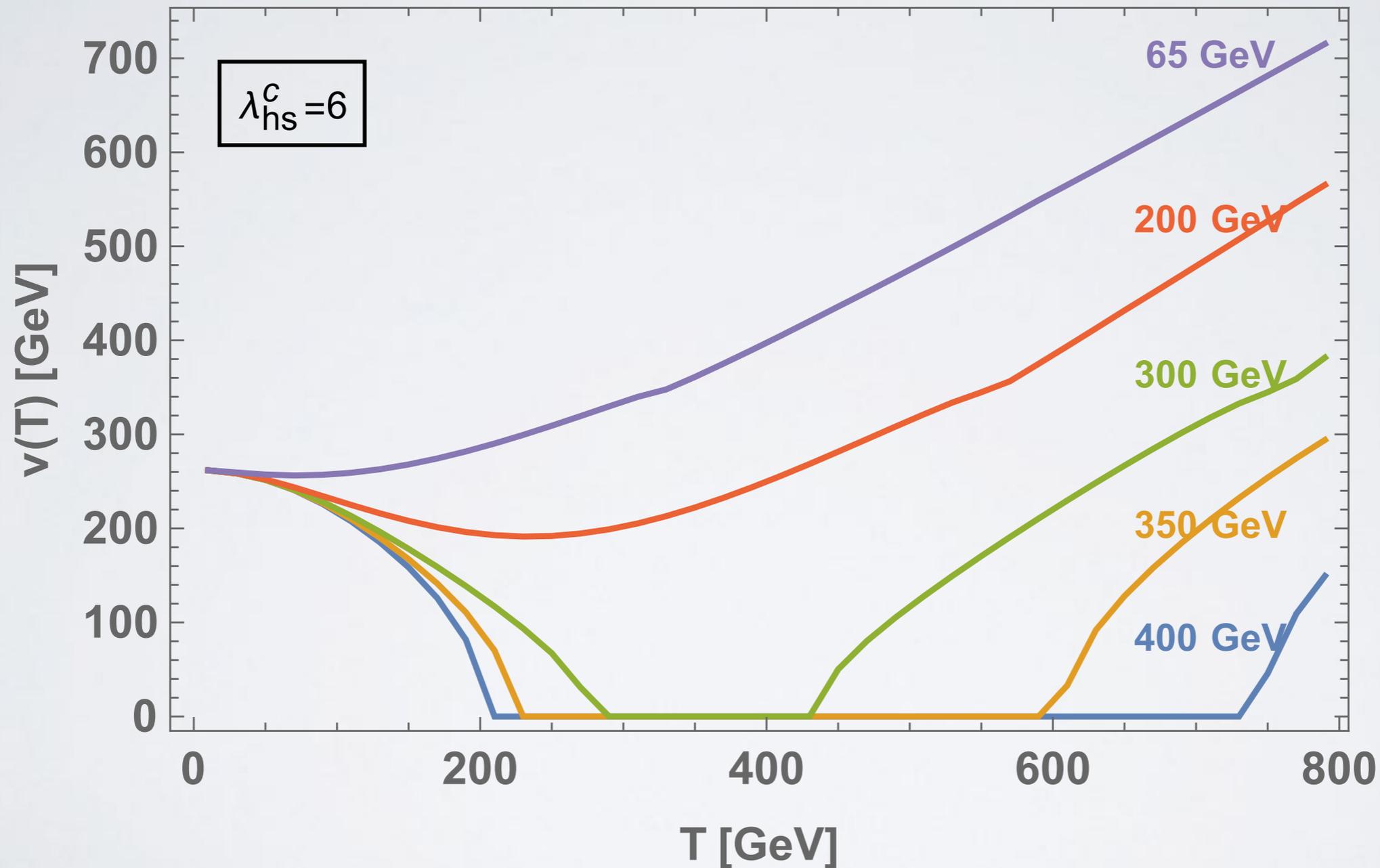
**“Unfortunately, despite the fact that one is dealing with a weakly coupled theory, many aspects of the phase transition are surprisingly complicated. Indeed, the literature contains contradictory claims and statements on almost every important question.”**

'92 Dine, Leigh, Huet, Linde ,Linde

# HIGHER ORDER EFFECTS



# VERY COOL EARLY UNIVERSE POSSIBILITIES



Depending on the Singlet Mass you can get SNR-R-SNR

# COSMOLOGICAL CHANGES

- Sphalerons are controlled by  $\frac{v(T)}{T} \equiv \kappa$ 
  - for  $\kappa \sim 1$  sphalerons are turned off
  - “GUT” Baryogenesis can work- Maximons
  - Models that use sphalerons would be dead (EWBG, some Leptogenesis) - can look SM like at low energies
    - Can also just postpone EWBG:  
see Baldes, Servant and Rattazzi, Vecchio

# COSMOLOGICAL CHANGES

- Avoid defects if you avoid phase transitions...
- Is decoupling any different? In principle yes

$$m(T) \sim g\kappa T$$

- For very large kappa, particles are non-relativistic instead of relativistic

$$\kappa \sim \frac{\sqrt{\Pi_h/\lambda}}{T}$$

Can enhance with  
running  $\lambda$

Even more interesting you can get exotic equations of state!

# LARGE N SCALING CHANGES COLLIDER OBSERVABLES AS WELL

However, let's look at the scaling for collider observables...

$$\delta_{Zh} \sim N_s \lambda_{HS}^2 \quad \sigma_{h^* \rightarrow ss} \sim N_s \lambda_{HS}^2 \quad \delta_{h^3} \sim N_s \lambda_{HS}^3$$

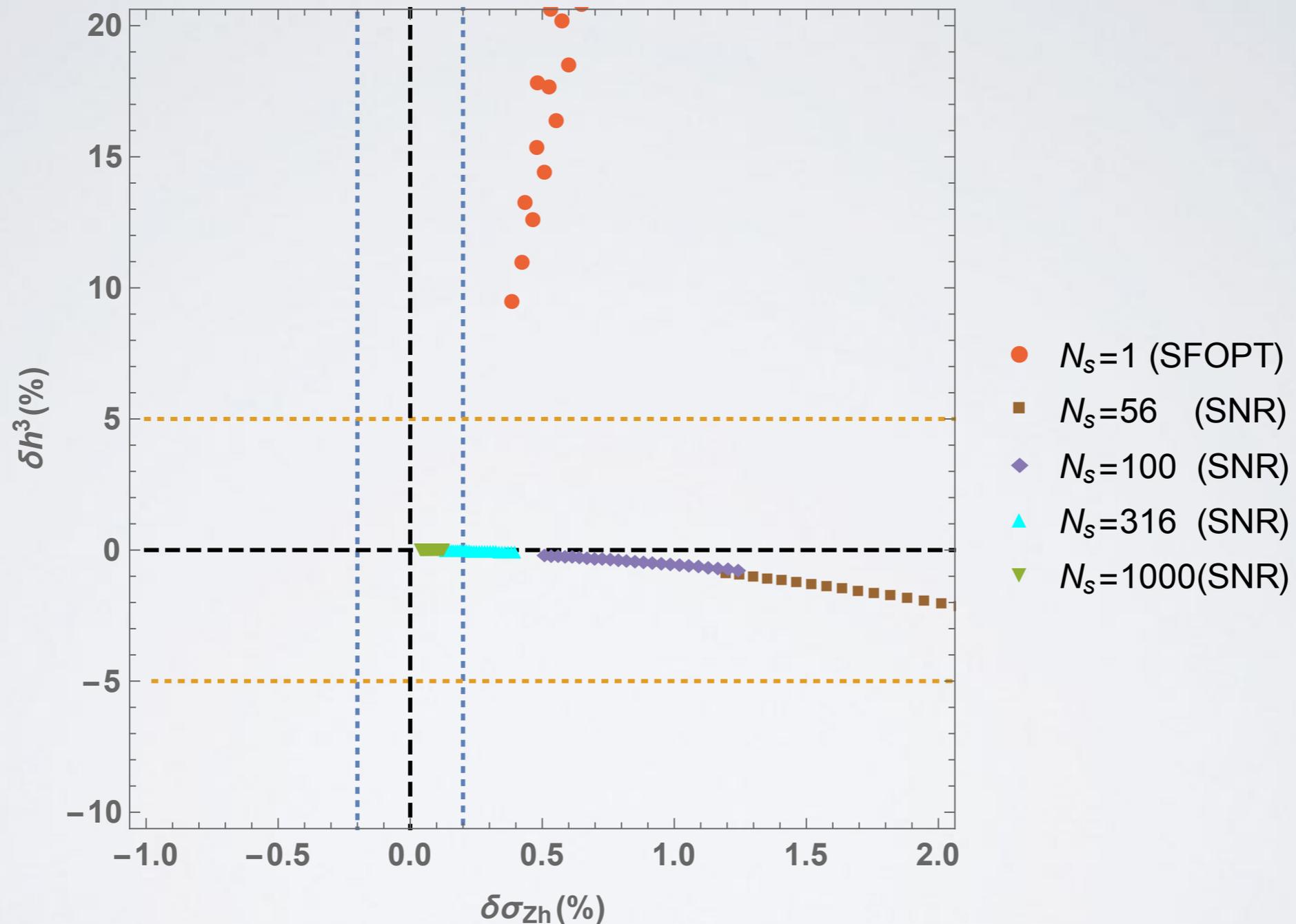
If we fix  $\lambda_{HS} N_s = \lambda_{HS}^c$

$$\delta_{Zh} \sim \frac{(\lambda_{HS}^c)^2}{N_s} \quad \sigma_{h^* \rightarrow ss} \sim \frac{(\lambda_{HS}^c)^2}{N_s} \quad \delta_{h^3} \sim \frac{(\lambda_{HS}^c)^3}{N_s^2}$$

In the scaling limit the effects disappear!

**Can we tell whether or not the early universe was in a  
SNR phase?**

# ANOTHER INTERESTING COMPLICATION



Can you confuse SNR with a strong FOPT? yes, up to triple higgs  
One would have a strong gravitational wave signal, the other wouldn't

# CONCLUSIONS

- Lots of interesting physics under the Higgs lamppost
- Need a **new** flowchart for thinking about triple Higgs couplings, but it is likely the most important measurement for understanding qualitative differences about our universe from particle perspective