

LECTURE III

2 D - CFT



RECAP

GAUKE FIXING DIFF X KEYL

$$S_{\text{BS}} = S_{\text{matter}}(x) + S_{\text{ghost}}(b_{1c})$$

$$\text{BRST SYMMETRY} \quad \delta_B = [Q_B + \bar{Q}_B, \cdot]$$

$$Q_B^2 = 0 \iff C_{\text{matter}} + C_{\text{ghost}} = 0 \quad (0=26)$$

PHYSICAL STATE CONDITIONS

WORLD-SHEET

SPACE-TIME

"PHYSICAL STATE"

$$Q_B |\Psi\rangle = 0$$

EQUATION OF MOTION

EQUIVALENCE RELATION

$$|\Psi\rangle \sim |\Psi\rangle + Q_B |1\rangle$$

GAUKE INVARIANCE

$$(b_0 |\Psi\rangle = 0 \text{ GAUKE FIXNt})$$

$$S = \frac{1}{2\pi} \int d^2\omega \left(\frac{\partial X^a \partial X^a}{d^2} + b \bar{\partial} c + \bar{b} \partial \bar{c} \right)$$

CONFORMAL SYMMETRY

$$\omega \rightarrow f(\omega) \quad \bar{\omega} \rightarrow \bar{f}(\bar{\omega})$$

• GAUSS REDUNDANCY

$$d\omega d\bar{\omega} \rightarrow |f'(\omega)|^2 d\omega d\bar{\omega}$$

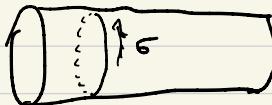
• WEYL SYMMETRY (FADDEEV)

$$d\omega d\bar{\omega} \xrightarrow{\text{WEYL}} |f'(\omega)|^{-2} d\omega d\bar{\omega}$$

→ CONFORMAL SYMMETRY IS A GAUSS SYMMETRY

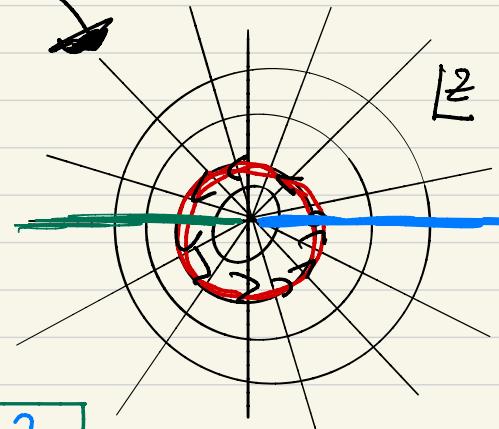
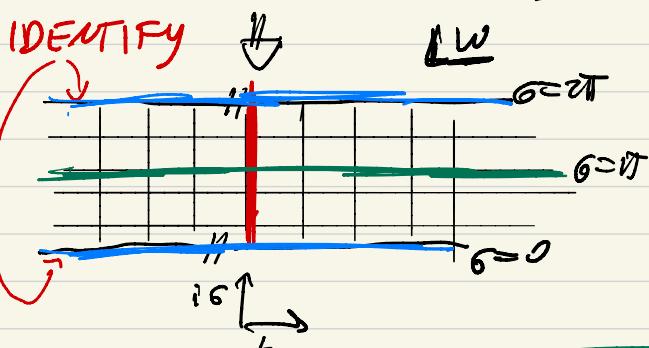
→ NEED ONLY CONFORMAL INVARIANT
"OBJECTS"

CYLINDER \longrightarrow COMPLEX PLANE



CYLINDER

IDENTIFY



$$Z = e^{\omega}$$

$$\omega = \ln z$$

$$X(\omega, \bar{w}) = X(\ln z, \ln \bar{z})$$

$$\partial X(\omega) = \left(\frac{d z}{d \omega} \right) \partial X(z)$$

$$\partial X d\omega = \partial_z X dz$$

$$b_{\omega\omega}(\omega) = \left(\frac{d z}{d \omega} \right)^2 b(z)$$

$$b(\omega) d\omega^2 = b(z) dz^2$$

$$c^\omega(\omega) = \left(\frac{d z}{d \omega} \right)^{-1} c(z)$$

$$c(\omega) d\omega^{-1} = c(z) dz^{-1}$$

PRIMARY FIELD OF WEIGHT (k, \bar{k})

$\phi(\omega, \bar{\omega}) (\mathrm{d}\omega)^k (\mathrm{d}\bar{\omega})^{\bar{k}}$ is invariant

$$\phi(\omega, \bar{\omega}) = \left(\frac{\mathrm{d}\omega}{\mathrm{d}z} \right)^k \left(\frac{\mathrm{d}\bar{\omega}}{\mathrm{d}\bar{z}} \right)^{\bar{k}} \phi(z, \bar{z})$$

$$X \rightarrow (1, 0)$$

$$\bar{X} \rightarrow (0, 1)$$

$$b \rightarrow (2, 0)$$

$$\bar{b} \rightarrow (0, 2)$$

$$c \rightarrow (-1, 0)$$

$$\bar{c} \rightarrow (0, -1)$$

$$\phi_{\text{cyc}}^{(k)}(w) = \sum_r \phi_r e^{-rw}$$

$$\phi_{\text{cyc}}^{(w)}(z) = \left(\frac{d w}{d z}\right)^k \phi_{\text{cyc}}^{(w)}(w(z)) =$$

$$\boxed{\phi_{\text{cyc}}^{(k)}(z) = \sum_r \phi_r z^{-r-k}}$$

Ex

$$j^{\mu}(z) = i \sqrt{\frac{2}{d!}} \partial X^{\mu} = \sum_{m \in \mathbb{Z}} d_m^{\mu} z^{m-1} \quad \alpha_0^{\mu} = \sqrt{\frac{2}{d!}} p^{\mu}$$

$$c(z) = \sum_{m \in \mathbb{Z}} c_m z^{-m+1}$$

$$b(z) = \sum_{n \in \mathbb{Z}} b_n z^{-n-2}$$

INFINITESIMAL CONFORMAL TRANSFORMATIONS

$$z \rightarrow z + \epsilon(z) = (1 + \underline{\epsilon(z)} \partial_z) z$$

$$\epsilon(z) = \sum_{m \in \mathbb{Z}} \epsilon_m z^{-m+1}$$

$$z \rightarrow \boxed{e^{\epsilon(z) \partial_z} z = f(z)}$$

$$\epsilon(z) \partial_z = - \sum_{m \in \mathbb{Z}} \epsilon_m l_{-m} = \delta_\epsilon$$

$$l_{-m} = -z^{m+1} \partial_z$$

EXERCISE 1

$$[l_m, l_n] = (m-n) l_{m+n}$$

EXERCISE 2

$$[\delta_{\epsilon_1}, \delta_{\epsilon_2}] = \int_{\epsilon_1} \partial \epsilon_2 - \epsilon_2 \partial \epsilon_1$$

PRIORITARY FIELD UNDER KEP- CORF- TRANSF

$$\phi(z) \rightarrow \phi'(z') = \left(\frac{dz}{dz'}\right)^4 \phi(z)$$

$$z'(z) = z + \epsilon(z) \quad \frac{\partial z'}{\partial z} = 1 + \partial \epsilon$$

$$\phi'(z + \epsilon(z)) = (1 - 4\partial \epsilon) \phi(z)$$

$$\begin{aligned}\phi'(z) &= (1 - 4\partial \epsilon) \phi(z - \epsilon) = \\ &= \phi(z) + \delta_\epsilon \phi(z)\end{aligned}$$

$$\delta_\epsilon \phi(z) = -(\epsilon \partial \phi + 4\partial \epsilon)(z)$$

RADIAL ORDERING

cyclic order = t-ordering $\omega = t + i\delta$



$$z = e^\omega$$

$|z| \sim \text{cyclic order}$

$$R(\phi_1(z)\phi_2(\omega)) = \begin{cases} \phi_1(z)\phi_2(\omega) & |z| > |\omega| \\ (-)^{\epsilon} \phi_2(\omega)\phi_1(z) & |z| < |\omega| \end{cases}$$

ORTHOGONALITY OF ϕ_1 and ϕ_2

~~GRAPED~~

EQUAL RADII COMMUTATOR

$$[\phi_1(z), \phi_2(\omega)]_{|z|=|\omega|} =$$

$$= \lim_{\epsilon \rightarrow 0^+} \left. \phi_1(z)\phi_2(\omega) \right|_{|z|=|\omega|+\epsilon} - \left. \phi_2(\omega)\phi_1(z) \right|_{|z|=|\omega|-\epsilon}$$

ORTHOGONALITY

ENERGY-MOMENTUM TENSOR

$$T_{\alpha\beta} = \frac{1}{4\pi} \delta_{\alpha\beta} S$$

$$T_{zz} = T(z)$$

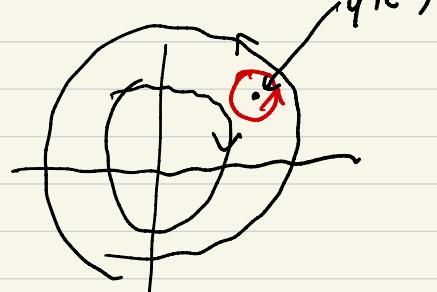
$$T_{z\bar{z}}, T_{\bar{z}\bar{z}}, \cancel{T_{\bar{z}\bar{z}}}$$

$$\bar{T}_{\bar{z}\bar{z}} = \bar{T}(\bar{z})$$

Diff : $\partial_{\bar{z}} T_{z\bar{z}} = 0 \quad \partial_z T_{\bar{z}\bar{z}} = 0$

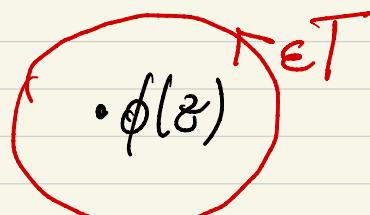
Wyl : $T_{\alpha}^{\alpha} = 0 \rightarrow T_{z\bar{z}} = 0$

$$\delta_{\varepsilon} = -[T_{\varepsilon}, \cdot]$$



$$T_{\varepsilon} = \oint_0 \frac{dw}{2\pi i} \varepsilon(w) T(w)$$

$$\delta_{\varepsilon} \phi(z) = - \oint_{\mathcal{C}} \frac{dw}{2\pi i} \varepsilon(w) T(w) \phi(z)$$



⇒ ONLY SENSIBLE TO

THE SHORT-DISTANCE BEHAVIOR OF $\underbrace{T(w)}_{\phi(z)}$

OPERATOR PRODUCT EXPANSION

$$\phi_i^{(k_i)}(z) \phi_j^{(l_j)}(w) = \sum_{\kappa} C_{ij}^{\kappa} \frac{\phi_{\kappa}^{(h_{\kappa})}(w)}{(z-w)^{h_{i+j}-h_{\kappa}}}$$

STRUCTURE CONSTANTS

$$T(\epsilon) \phi^{(n)}(z) , \quad \underline{\phi^{(4)}} \text{ e PRIMARY}$$

$$\delta_{\epsilon} \phi(z) = - \oint_{\gamma} \frac{dw}{2\pi i} \epsilon(w) T(w) \phi(z) = \\ = - (\epsilon \partial \phi + h \partial \epsilon \phi)(z)$$

$$T(w) \phi^{(n)}(z) = \frac{h \phi(z)}{(w-z)^2} + \frac{\partial \phi(z)}{w-z} + \text{regular}$$

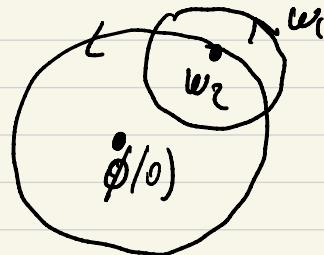
↳ DEFINITION OF A PRIMARY FIELD

$$T(\varepsilon) T(\omega) = ?$$

$$[\delta_{\varepsilon_1}, \delta_{\varepsilon_2}] = \delta_{\varepsilon_1 \partial \varepsilon_2 - \varepsilon_2 \partial \varepsilon_1} \quad \phi^{(4)}(0)$$

$$\delta_{\varepsilon} = \oint \frac{d\omega}{2\pi i} \varepsilon(\omega) T(\omega)$$

$$\oint_0 \frac{d\omega_2}{2\pi i} \oint_{w_2} \frac{d\omega_1}{2\pi i} \varepsilon_1(\omega_1) \varepsilon_2(\omega_2) T(\omega_1) T(\omega_2) \phi(0)$$



$$T(\omega_1) T(\omega_2) = \frac{\#1}{(\omega_1 - \omega_2)^4} + \frac{c T(\omega_1)}{(\omega_1 - \omega_2)^2} + \frac{\partial T(\omega_2)}{\omega_1 - \omega_2} + \text{reg.}$$

$$\boxed{\# = \frac{c}{2}}$$

$c \rightarrow$ CENTRAL CHARGE
OF THE CFT

\hookrightarrow DATA OF THE CFT

$$T(z)T(w) = \frac{\%}{(z-w)^4} + \frac{zT(w)}{(z-w)^2} + \frac{\partial T(w)}{z-w} + \text{res.}$$

$$T(z) = \sum_{m \in \mathbb{Z}} L_m z^{-m-2}$$

$$L_m = \oint_0 \frac{dz}{2\pi i} z^{m+1} T(z)$$

$$[L_n, L_m] = (n-m) L_{n+m} + \frac{c}{12} \underbrace{(n^3 - n)}_{\downarrow} S_{n+m}$$

$$L_{-1}, L_0, L_1 \rightarrow SL(2, \mathbb{C}) \quad \text{O } n= \pm 1, 0$$

\Rightarrow QUASI-PRIMARY FIELD

($T \mapsto$ quasi-primary)

EXERCISE

$$[L_m, \phi^{(h)}(z)] = z^m (z\partial + (m+1)h) \phi(z)$$

$$[L_m, \phi_m] = (m(h-1) - m) \phi_{m+m}$$

HILBERT SPACE OF A CFT

① $SL(2|\mathbb{C})$ -invariant vacuum

$$|0\rangle \quad \lim_{w \rightarrow 0} T(w)|0\rangle = \text{WELL-DEFINED}$$

$$\lim_{w \rightarrow 0} \# \frac{1}{w^m}|0\rangle \quad m > 0$$

$$\#|0\rangle = 0$$

$$L_m |0\rangle = 0 \quad m \geq -1$$

$$L_m^+ = L_{-m} \rightarrow \langle 0| L_m = 0 \quad m \leq 1$$

L_{-1}, L_0, L_1 annihilate $|0\rangle$
 $\langle 0|$

$|0\rangle, \langle 0|$ $SL(2|\mathbb{C})$ invariant
 VACUUM

$$\textcircled{2} \quad \phi^{(n)}(z) \Big|_{z=0} |0\rangle \quad \text{well-defined}$$

$$\begin{aligned} \phi_m |0\rangle &= 0 & m \geq -h+1 \\ \langle 0| \phi_m &= 0 & m \leq h-1 \end{aligned} \quad \rightarrow \quad c_{-1}, c_0, c_1 \\ \text{zero modes} \end{math>$$

\textcircled{3} STATE-OPERATOR CORRESPONDENCE

$$\phi(z) \rightarrow |\phi\rangle = \lim_{z \rightarrow 0} \phi(z) |0\rangle_{SL(2, \mathbb{C})}$$

$$\phi(z) = e^{zL_1} \phi(0) e^{-zL_1} \quad L_1 = \partial_z = l_{-1}$$

$$\begin{aligned} \phi(z) |0\rangle_{SL(2, \mathbb{C})} &= e^{zL_1} \phi(0) e^{-zL_1} |0\rangle_{SL(2, \mathbb{C})} = \\ &= e^{zL_1} \phi(0) |0\rangle = e^{zL_1} |\phi\rangle \end{aligned}$$

$$|\phi\rangle \rightarrow e^{zL_1} |\phi\rangle = \phi(z) |0\rangle$$

PRIMARIES AND DESCENDENTS

$$\phi^{(h)}(z) \rightarrow |\phi^{(h)}\rangle = \phi^{(0)}|0\rangle_{SL(2|C)}$$

$$T-\phi \text{ opE} \rightarrow \begin{cases} L_0 |\phi^{(h)}\rangle = h |\phi\rangle \\ L_m |\phi^{(h)}\rangle = 0 \quad m > 0 \end{cases}$$

$L_{-m} \quad m > 0 \rightarrow$ RAISING OPERATORS ACTING
ON "VACUA" GIVEN BY
PRIMARY STATES

Span $\{L_{-k_1}, \dots, L_{-k_n} | \phi^{(h)}\rangle\} \rightarrow$ VERMA MODULE

\Downarrow

DESCENDENTS OF $\phi^{(h)}$

$L_0 = h + \sum_{i=1}^m k_i$

BOSONIC STRING CFT

Matter and ghost

2- DECOUPLED CFT's

$$X: T^{(X)} = -\frac{1}{d!} : \partial X_\mu \partial X^\mu :$$

$$bc: T^{(bc)} = -z : b \partial c : (z) - : \partial bc : (z)$$

$\therefore \rightarrow$ normal ordering on $|0\rangle_{SL(2, \mathbb{C})}$

\rightarrow "bc" CFT is universal

\rightarrow X-CFT is more "occidental"

\hookrightarrow A chosen background of
(bosonic) string theory

$$A(z)B(w) = \underbrace{A(z)B(w)}_{\text{CONTRACTION}} + :A(z)B(w):$$

CONTRACTION

\Downarrow
THE SINGULAR PART
OF THE OPE

WORM-ORDERED
PRODUCT

\Downarrow
ROUND PART

$$:AB:(w) = \lim_{z \rightarrow w} :A(z)B(w): = \oint \frac{dz}{2\pi i} \frac{A(z)B(w)}{z-w}$$

\rightarrow IN CASE OF FREE CFT'S (W $\in X, b_c$)

\rightarrow EQUIVALENT TO OSCILLATOR WORMS ORDERED
ON $|0\rangle_{SCC}$

6HOST CFT

$$S = \frac{1}{2\pi} \int d^2z (b\bar{\partial}c + \bar{b}\partial\bar{c})$$

$b(z)$ ($z_1, 0$)

$$C(\mathcal{E}) \quad (-1, 0)$$

$$b(z) c(w) = \frac{1}{z-w}$$

$$[b_m, c_n] = \delta_{m+n}$$

$$T^{(bc)} = -2 \cdot b \partial c : (z) - : \partial b c : (z)$$

EXERCISES

$$T_{(8)}^{bc} T_{(\omega)}^{bc} = \frac{-13}{(2-\omega)^4} + \frac{2T(\omega)}{(2-\omega)^2} + \frac{\partial T(\omega)}{2-\omega} \quad (= -26)$$

$$j_{f^k}(z) = - :bc:(z) = \sum_{m \in \mathbb{Z}} j_m^{(k)} z^{-m-1}$$

$$Q_{\text{fgh}} = j_0^{(fgh)} = \oint \frac{dz}{2\pi i} j_{\text{fgh}}(z)$$

$$T(z) \underset{z \rightarrow \omega}{\overset{\text{lim}}{\longrightarrow}} j_{ph}(\omega) = \frac{3}{(z-\omega)^3} + \frac{j_{ph}(\omega)}{(z-\omega)^2} + \frac{\partial j_{ph}(\omega)}{z-\omega}$$

Priamaries

$c, c\partial c, c\partial c\partial^2 c, \dots$ (right # $\#_{\text{gh}}$)

$b, b\partial b, b\partial b\partial^2 b, \dots$ (left-right # $\#_{\text{gh}}$)

$$|0\rangle_{SL(2)\mathbb{C}} \longleftrightarrow 1$$

$$|\downarrow\rangle = c(0)|0\rangle \longleftrightarrow c(\varepsilon)$$

$$|\uparrow\rangle = \partial c(c(0)|0\rangle) \underset{c_0 c_1 |0\rangle}{\longleftrightarrow} \partial c c(\varepsilon)$$

$$|0'\rangle = \frac{1}{2} \partial^2 c \partial c(c(0)|0\rangle) \underset{c_-(c_0 c_1 |0\rangle)}{\longleftrightarrow} \frac{1}{2} \partial^2 c \partial c c(\varepsilon)$$

EXERCISE

$$\cdot \langle 0|c_{-1}c_0c_1|0\rangle = 1$$

$$\cdot c(\varepsilon) = \sum_{n \in \mathbb{Z}} c_n \varepsilon^{-n+1}$$

$$\textcircled{1} \quad \langle 0|\frac{1}{2}\partial^2 c \partial c c(\varepsilon)|0\rangle = \langle 0|\frac{1}{2}\partial^2 c \partial c c(\varepsilon)|0\rangle = 1$$

$$\textcircled{2} \quad \langle c(\varepsilon_1)c(\varepsilon_2)c(\varepsilon_3) \rangle = \varepsilon_1 \varepsilon_2 \varepsilon_3 \\ z_{ij} = \varepsilon_i - \varepsilon_j$$

BEST-CHARGE

$$Q_B = \sum_{m \in \mathbb{Z}} c_m l_m^{(m)} + \frac{1}{z} : c_m l_m^{(gh)} : =$$

$$= \oint \frac{dz}{2\pi i} \underbrace{\left(c T^{(m)}(z) + \frac{1}{z} : c T^{(gh)} : (z) \right)}_{\tilde{J}_B(z)}$$

$$j_B(z) = \tilde{J}_B(z) + \frac{3}{2} \partial^c c \rightarrow \begin{matrix} \text{PRIMARY OF} \\ h=1 \end{matrix}$$

$$\textcircled{1} [Q_B, X^m(z)] = \oint \frac{d\omega}{2\pi i} j_B(\omega) X^m(z) \\ = c \partial X^m$$

$$\textcircled{2} [Q_B, c(z)] = c \partial c(z)$$

$$\textcircled{3} [Q_B, b(z)] = T^x(z) + T^{(gh)}(z) = T^{(\text{tot})}(z)$$

$$j_B(z) \underset{\substack{\longleftarrow \\ \text{---}}}{} j_B(w) = \dots + \frac{c^{(m)} - 26}{12} \frac{1}{z-w} \overset{\rightharpoonup}{\partial} c c(w) + \dots$$

$$Q^2 \rightarrow \oint \frac{d\omega}{2\pi i} \underset{\substack{\longleftarrow \\ \text{---}}}{} j_B(\omega) \underset{\substack{\longleftarrow \\ \text{---}}}{} j_B(z) = 0 \quad c^{(m)} = 26$$

$j_B(z)$ is BRST-EXACT

$$j_B(z) = - [Q_B, j_{gh}(z)]$$

↓

$$[Q_{gh}, Q_B] = + Q_B$$

COHOMOLOGY OF Q_B

$$\begin{array}{ll}
 f_h=0 & 1 \\
 \boxed{f_h=1} & \subset V_{h=1}^{(m)}(z) \\
 f_h=2 & \subset \partial V_{h=1}^{(m)}(z) \\
 f_h=3 & \subset \partial^2 V_{h=1}^{(m)}(z)
 \end{array}
 \rightarrow \text{PHYSICAL FIELDS}$$

$$\int \underbrace{\phi_h \partial X}_{V_{\text{gauge}}} \underbrace{V_1 \dots V_m}_{e^{-S[\phi_h, X]}} \rightarrow \text{GAUSS INVARIANT OPERATORS}$$

$f_h=1$ elements of the cohomology are
the first set of GAUSS-INVARIANT OPERATORS

$$C V_1^{(m)}(z) \xrightarrow[Q_B = 0]{} \text{NON-INTEGRATED}\\ \text{VERTEX OPERATORS} \\ (\text{CHIRAL PART})$$

FULL CLOSED STRING THEORY

$$C \in V_{(1,1)}^{(m)}(z, \bar{z}) \rightarrow \frac{Q_B}{Q_S} = 0 \quad [Q_B + \bar{Q}_B = 0]$$

→ ANOTHER SET OF STATE-INVARIANT OPERATORS

$$CV_1^{(m)} \xrightarrow{\text{---}} \int dz \underbrace{V_1^{(m)}(z)}_{(-1)^z} \quad Q = cT + \frac{1}{2}cT^4$$

$$\int dz V_1^{(m)}(z) \Rightarrow \boxed{\text{DIMENSION } \ell = 0}$$

\Downarrow → INTEGRATED
RENEX OPERATOR

$$[Q_B, V_1^{(m)}(z)] = \oint \frac{dw}{2\pi i} CT \underbrace{V_1^{(m)}(w)}_{(w-z)^2} V_1^{(m)}(z) =$$

$$= \oint \frac{dw}{2\pi i} c(w) \left(\frac{V_1^{(m)}(z)}{(w-z)^2} + \frac{\partial V_1^{(m)}(z)}{w-z} \right) =$$

$$= \partial(cV_1^{(m)})$$

$$[Q_B, \mathcal{D}] = 0 = \int dz \partial(cV_1^{(m)}) = 0$$

| UP TO
POSSIBLE
BOUNDARY
TERMS

$$\textcircled{1} \quad C V_i^{(m)}(z) \rightarrow \text{new integrals}$$

$V, O, f_k = 1$

$$\textcircled{2} \quad \int dz V_i^{(m)}(z) \rightarrow \text{INTEGRALS}$$

$V=0, f_k = 0$

\Rightarrow STEWART'S AMPLITUDES USE
BOTH

$$\int p_b p_c \langle \mu, b \rangle V_i \dots V_j e^{b_{\alpha\beta}(p, c)^{\alpha\beta}}$$