

# BRST-Lagrangian Double Copy of Yang-Mills Theory

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21 December 2020

Based on joint work with L Borsten, B Jurčo, H Kim, C Saemann, M Wolf  
2007.13803, 2101.?????

# Introduction: gravity = gauge $\times$ gauge

- Naive idea: “ $A_\mu \otimes \bar{A}_\nu = g_{\mu\nu} \oplus B_{\mu\nu} \oplus \phi$ ”
- First concrete incarnation: KLT relations
- A purely field theoretic approach: BCJ color–kinematic duality and double copy prescription

# Flash review of BCJ duality and double-copy

- Every Lagrangian field theory is equivalent to a theory with only cubic interactions
- $n$ -points  $L$ -loops YM amplitude as sums of trivalent graphs

$$\mathcal{A}_{n,L}^{\text{YM}} = \sum_i \int \prod_{l=1}^L d^d p_l \frac{1}{S_i} \frac{C_i N_i}{D_i}$$

- $i$  ranges over all trivalent  $L$ -loops graphs
- $C_i$ : color factor, composed of gauge group structure constants
- $N_i$ : kinematic factor, composed of Lorentz-invariant contractions of polarisations and momenta

# Flash review of BCJ duality and double-copy

## BCJ color–kinematic duality

There is a choice of kinematic factors such that  $N_i$ s obey the same algebraic relations (e.g., Jacobi identity) of the correspondent  $C_i$

- True at tree-level, conjectured for loop-level
- Gravity amplitudes can be represented as sum over trivalent graphs, too

## Yang–Mills double copy

If **BCJ duality holds true**, replacing the color factor with a copy of the kinematic factor in  $\mathcal{A}_{n,L}^{\text{YM}}$  produces a  $\mathcal{N} = 0$  supergravity amplitude

- All-loop statement, the problem is then to validate BCJ duality at loop level

- Until now, on-shell scattering amplitude approach
- An off-shell Lagrangian realization of color-kinematic duality and double-copy could solve the all-loop conundrum

## Our approach

Double-copy YM Lagrangian and BRST operator to obtain a theory equivalent to  $\mathcal{N} = 0$  supergravity

- There exist a non-local YM Lagrangian with manifest tree-level BCJ duality for on-shell physical gluons (Tolotti, Weinzierl, '13)
- We can insert auxiliary fields to make it local, and strictify to a Lagrangian with only cubic interaction

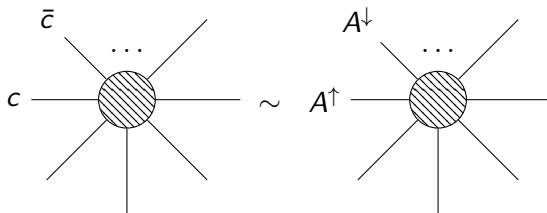
$$\mathcal{L} = \frac{1}{2} \Phi^{\alpha i} g_{\alpha\beta} G_{ij} \square \Phi^{\beta j} + \frac{1}{3!} \Phi^{\alpha i} f_{\alpha\beta\gamma} F_{ijk} \Phi^{\beta j} \Phi^{\gamma k} + \text{ghosts} + \text{gauge fixing}$$

- We want to double-copy the BRST-extended field space
- BCJ duality is satisfied only on-shell for physical gluons: eventual BCJ violations due to unphysical gluons and ghosts!

- We compensate for these eventual BCJ violations with suitable field redefinitions
- We start with 4-gluon tree amplitude
- Unphysical gluons ruin BCJ duality
- We compensate the violation of BCJ duality adding terms to the gauge-fixing fermion and shifting the Nakanishi–Lautrup fields

# BRST-Lagrangian perspective on double copy

- Next, gluon-gluon-ghost-antighost tree amplitude
- Thanks to the Ward identity, we can write this amplitude in terms of 4-gluon tree amplitudes: we redefine the ghost sector to compensate for BCJ violations



## Recursive construction

For every tree amplitude with  $n$  legs and  $k$  pairs ghost-antighost, recursively redefine gauge-fixing sector and ghost sector



- We obtain a cubic action that manifest on-shell tree-level BCJ duality for the BRST-extended field space:

$$\mathcal{L} = \frac{1}{2} \Phi^{\alpha i} g_{\alpha\beta} G_{ij} \square \Phi^{\beta j} + \frac{1}{3!} \Phi^{\alpha i} f_{\alpha\beta\gamma} F_{ijk} \Phi^{\beta j} \Phi^{\gamma k}$$

- BRST operator

$$(Q\Phi)^{\alpha i} = \delta_{\beta}^{\alpha} q_{\beta}^i \Phi^{\beta j} + \frac{1}{2} f_{\beta\gamma}^{\alpha} Q_{\beta\gamma}^i \Phi^{\beta j} \Phi^{\gamma k} + \frac{1}{3!} f_{\beta\gamma\delta}^{\alpha} Q_{\beta\gamma\delta}^i \Phi^{\beta j} \Phi^{\gamma k} \Phi^{\delta l}$$

- Double-copy (both Lagrangian and  $Q$ )

$$\begin{aligned} \mathcal{L}_{\text{dc}} &= \frac{1}{2} \Phi^{i'i} G_{i'j'} G_{ij} \square \Phi^{j'j} + \frac{1}{3!} \Phi^{i'i} F_{i'j'k'} F_{ijk} \Phi^{j'j} \Phi^{k'k} \\ (Q_{\text{dc}}\Phi)^{i'i} &= \dots \end{aligned}$$

- If  $F_{ijk}$  satisfies the same algebraic properties of  $f_{\alpha\beta\gamma}$ , then  $Q_{\text{dc}}^2 = 0$  and  $Q_{\text{dc}}S_{\text{dc}} = 0$
- We have  $Q_{\text{dc}}^2 = 0$  and  $Q_{\text{dc}}S_{\text{dc}} = 0$  on-shell, and that's enough: it ensures the correct Ward identities and a consistent quantization
- $\mathcal{L}_{\text{dc}}$  is related to  $\mathcal{N} = 0$  supergravity by local field redefinitions
- We obtain a double-copy action  $\mathcal{L}_{\text{dc}}$  that is (perturbatively) quantum equivalent to  $\mathcal{N} = 0$  supergravity

- Further clarification of the algebraic structure of Lagrangian double-copy in terms of homotopy algebras (*paper to appear, hopefully January*)
- Extend our approach to supersymmetric theories
- Homotopic description of open–closed string duality?

Thank you for listening!