

# Singular Calabi-Yau threefolds: Springer resolution and quiver gauge theories

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# Outline

- F-theory and D7 brane backgrounds
- Higgs mechanism and unconventional Higgs VEVs
- Relation with ADE singularities in Calabi-Yau spaces
- Springer resolution
  - $A_n$  cases
  - Reid Pagoda
  - $D_4$
- Motivation: gauge theory of branes on ADE singularities

# D7-brane backgrounds

Consider type IIB limit of F-theory:

- Non-abelian gauge symmetries are realized stacking D7-branes on top of each other



Example:  $A_1 = SU(2)$  gauge group

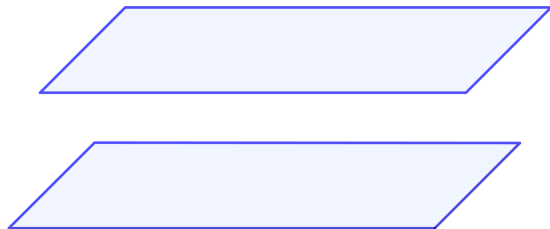
# Higgsing

Give a VEV to the Higgs field in the Cartan of  $SU(2)$ :

$$\langle \Phi \rangle = \begin{pmatrix} v & 0 \\ 0 & -v \end{pmatrix}$$

with  $v$  a **constant** on the branes worldvolume. Physically:

$$\langle \Phi \rangle = \begin{pmatrix} v & 0 \\ 0 & -v \end{pmatrix}$$



- Symmetry broken to  $U(1)$

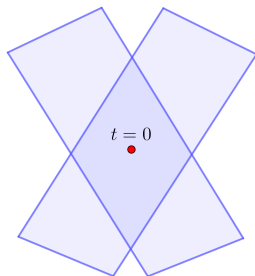
# Higgsing

Let  $\langle \Phi \rangle$  depend on the brane worldvolume coordinates, e.g.:

$$\langle \Phi \rangle = \begin{pmatrix} t & 0 \\ 0 & -t \end{pmatrix}$$

Physically:

$$\langle \Phi \rangle = \begin{pmatrix} t & 0 \\ 0 & -t \end{pmatrix}$$



- $t \neq 0$  symmetry broken to  $U(1)$
- on  $t = 0$  enhancement to  $SU(2)$

# Non-standard Higgs

Take another case:

$$\langle \Phi \rangle = \begin{pmatrix} 0 & 1 \\ t & 0 \end{pmatrix}$$

- Eigenvalues are  $\pm\sqrt{t}$ , they are not single-valued, they are **monodromic**
- Symmetry is broken to  $U(1)$  everywhere
- These backgrounds are also called **T-branes** [Cecotti, Córdova, Heckman, Vafa, 2010]

# Relation with geometry

What is the corresponding geometry? How to extract the corresponding Calabi-Yau space?

For the  $A_n$  case the procedure is known:

$$x^2 + y^2 = \det(z\mathbb{1} - \langle \Phi \rangle)$$

- Non-monodromic case:

$$\implies \text{Conifold: } x^2 + y^2 = z^2 + t^2$$

i.e. a **deformation** of the  $A_1$  singularity. A 2-cycle can be blown up.

- Monodromic case:

$$\implies x^2 + y^2 = z^2 + t$$

No 2-cycle can be blown up, due to monodromy.

What happens for more general gauge groups, i.e. Higgs fields not valued in the Cartan of some ADE gauge group?

A systematic procedure can be found using the so-called **Grothendieck-Springer resolution**



# Levi subalgebras

All the Higgs vevs (monodromic and non-monodromic) can be classified using **Levi subalgebras** of  $\mathfrak{g}$ .

Theorems by Springer and Grothendieck show that (suppressing all technical details):

*Levi subalgebras of  $\mathfrak{g}$  are in correspondence with (partial) resolutions of singular ADE Calabi-Yau threefolds.*

# Grothendieck-Springer resolution

Example: consider  $A_2 = SU(3)$  and choose a Levi subalgebra:

$$L = \begin{pmatrix} * & 0 & 0 \\ 0 & * & 0 \\ 0 & 0 & * \end{pmatrix}$$

$L$  is in correspondence with the **complete resolution** of the deformed  $SU(3)$  singularity, with equation:

$$x^2 + y^2 + z^3 - z(t_1^2 + t_2^2 + t_1 t_2) - t_1 t_2 (t_1 + t_2) = 0$$

with  $t_1(t)$  and  $t_2(t)$ .

# Grothendieck-Springer partial resolution

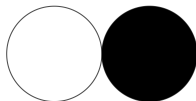
Example: if we choose a different Levi subalgebra:

$$L_p = \begin{pmatrix} * & * & 0 \\ * & * & 0 \\ 0 & 0 & * \end{pmatrix}$$

$L_p$  is in correspondence with the **partial resolution** of the deformed  $SU(3)$  singularity, in which only one two-cycle is resolved, with equation:

$$x^2 + y^2 + z^3 - z(\sigma_1^2 - \sigma_2) - \sigma_1\sigma_2 = 0$$

with  $\sigma_1(t)$  and  $\sigma_2(t)$ .

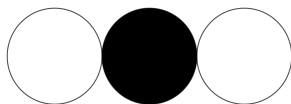


# Reid pagoda

Other example: the Reid pagoda can be seen as the (monodromic) deformation of a  $A_3$  singularity:

$$x^2 + y^2 + z^4 + t^2 = 0$$

Only the central node can be resolved.



The corresponding Levi subalgebra giving rise to this geometry is:

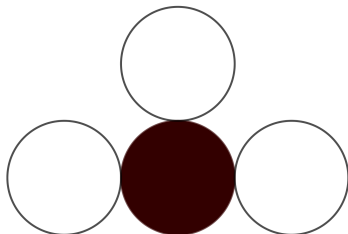
$$L_{pagoda} = \begin{pmatrix} * & * & 0 & 0 \\ * & * & 0 & 0 \\ 0 & 0 & * & * \\ 0 & 0 & * & * \end{pmatrix}$$

# A new $D_4$ singular threefold

Until now, essentially only one example of singular threefold with  $D_4$  singularity where only the central node can be resolved is known.

It is **Laufer's case** [Cachazo, Katz, Vafa, 2001], with equation:

$$x^2 + y^3 + tz^2 + t^3y = 0$$



## A new $D_4$ singular threefold

With the previously mentioned techniques we have found a new singular threefold with these properties:

$$x^2 + zy^2 - z^3 - \frac{a^3t^3}{16} + \frac{1}{8}a^2bt^4 - \frac{9}{16}a^2t^2z - \frac{1}{16}ab^2t^5 - \frac{3}{8}abt^3z - \frac{3}{2}atz^2 - \frac{1}{16}b^2t^4z - \frac{1}{2}bt^2z^2 = 0$$

with  $a$  and  $b$  constants.

The corresponding Levi subalgebra of  $SO(8)$  is:

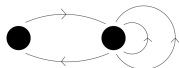
$$L_{D_4} = \begin{pmatrix} * & * & 0 & 0 & 0 & 0 & 0 & 0 \\ * & * & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & * & * & 0 & 0 & 0 & * \\ 0 & 0 & * & * & 0 & 0 & * & 0 \\ 0 & 0 & 0 & 0 & * & * & 0 & 0 \\ 0 & 0 & 0 & 0 & * & * & 0 & 0 \\ 0 & 0 & 0 & * & 0 & 0 & * & * \\ 0 & 0 & * & 0 & 0 & 0 & * & * \end{pmatrix}$$

# Gauge theory of D5 branes

Why do we care about these setups?

We can wrap D5 branes around the resolved 2-cycles:

→  $\mathcal{N} = 1$  theory with superpotential, described by a **quiver**



- Understanding these theories can be useful for large  $N$  dualities, as in [Cachazo, Katz, Vafa, 2001]

Using their techniques we have obtained a superpotential for our  $D_4$  example:

$$W = \frac{4}{3a} (X^3 + Y^3 + Z^3) - T \left( X + Y + Z + \frac{\sqrt{b}}{4} T \right)$$

with  $X, Y, Z, T$  chiral superfields.

# Recap and outlook

- Monodromic Higgses correspond to geometries where only some 2-cycles can be resolved
- The Grothendieck-Springer resolution allows to describe these geometries precisely
- Superpotentials for the theory of D5 branes wrapped on the resolved cycles can in some cases be computed using known techniques

Looking at the future:

- Consider further examples in the ADE classification
- Apply the formalism to concrete T-brane settings



**Thank you!**