

The Entropy of Black Holes from the Superconformal Index

Ziruo Zhang (SISSA)

Joint work with Francesco Benini, Edoardo Colombo, Saman Soltani, Alberto Zaffaroni ([2005.12308](#))

XVI Avogadro Meeting, Dec. 22, 2020

Summary

- In AdS/CFT, entropy of black holes
= $\log(\text{degeneracy of some states in CFT})$.

Summary

- In AdS/CFT, entropy of black holes
= $\log(\text{degeneracy of some states in CFT})$.
- We look at the dualities

$$\text{IIB on AdS}_5 \times SE_5 \leftrightarrow 4d \mathcal{N} = 1 \text{ SCFT}$$

and consider rotating, electrically charged supersymmetric black holes.

Summary

- In AdS/CFT, entropy of black holes
= $\log(\text{degeneracy of some states in CFT})$.
- We look at the dualities

$$\text{IIB on AdS}_5 \times SE_5 \leftrightarrow 4d \mathcal{N} = 1 \text{ SCFT}$$

and consider rotating, electrically charged supersymmetric black holes.

- Gave predictions for the Bekenstein-Hawking entropy from the superconformal index using the Bethe ansatz formula

Summary

- In AdS/CFT, entropy of black holes = $\log(\text{degeneracy of some states in CFT})$.
- We look at the dualities

$$\text{IIB on AdS}_5 \times SE_5 \leftrightarrow 4d \mathcal{N} = 1 \text{ SCFT}$$

and consider rotating, electrically charged supersymmetric black holes.

- Gave predictions for the Bekenstein-Hawking entropy from the superconformal index using the Bethe ansatz formula
- Successfully compared it to the near-horizon geometry of black holes in $\text{AdS}_5 \times T^{1,1}$

Black hole entropy from field theory

- In $\text{AdS}_5/\text{CFT}_4$, supersymmetric black hole \leftrightarrow ensemble of BPS states with $\{J_{1,2}, Q_{a=1, \dots, n}\}$. We study the superconformal index

$$\mathcal{I}(p, q, y) = \text{Tr}(-1)^F p^{J_1 + \frac{r}{2}} q^{J_2 + \frac{r}{2}} \prod_{a=1}^{n-1} y_a^{Q_a} .$$

Black hole entropy from field theory

- In $\text{AdS}_5/\text{CFT}_4$, supersymmetric black hole \leftrightarrow ensemble of BPS states with $\{J_{1,2}, Q_{a=1, \dots, n}\}$. We study the superconformal index

$$\mathcal{I}(p, q, y) = \text{Tr}(-1)^F p^{J_1 + \frac{r}{2}} q^{J_2 + \frac{r}{2}} \prod_{a=1}^{n-1} y_a^{Q_a} .$$

- Degeneracy of states is obtained via the Fourier transform

$$e^{S(J, Q)} = d(J, Q) = \int d\tau d\sigma \prod_{a=1}^{n-1} (d\Delta_a y_a^{-Q_a}) p^{-J_1} q^{-J_2} \mathcal{I}(p, q, y) ,$$

$$2\pi i \tau = \log p, \quad 2\pi i \sigma = \log q, \quad 2\pi i \Delta_a = \log y_a .$$

Black hole entropy from field theory

- In $\text{AdS}_5/\text{CFT}_4$, supersymmetric black hole \leftrightarrow ensemble of BPS states with $\{J_{1,2}, Q_{a=1, \dots, n}\}$. We study the superconformal index

$$\mathcal{I}(p, q, y) = \text{Tr}(-1)^F p^{J_1 + \frac{r}{2}} q^{J_2 + \frac{r}{2}} \prod_{a=1}^{n-1} y_a^{Q_a} .$$

- Degeneracy of states is obtained via the Fourier transform

$$e^{S(J, Q)} = d(J, Q) = \int d\tau d\sigma \prod_{a=1}^{n-1} (d\Delta_a y_a^{-Q_a}) p^{-J_1} q^{-J_2} \mathcal{I}(p, q, y) ,$$

$$2\pi i \tau = \log p, \quad 2\pi i \sigma = \log q, \quad 2\pi i \Delta_a = \log y_a .$$

- $S(J, Q)$ at $\mathcal{O}(N^2)$ is the integrand at the saddle point, i.e. the extremal value wrt. τ, σ, Δ_a of the entropy function

$$\mathcal{S} = \log \mathcal{I}(\tau, \sigma, \Delta) - 2\pi i \left(\sum_{a=1}^n \Delta_a Q_a + \tau J_1 + \sigma J_2 \right) .$$

Bethe ansatz formula

- For angular fugacities $p = e^{2\pi i a \omega}$, $q = e^{2\pi i b \omega}$, $\gcd\{a, b\} = 1$, the superconformal index can be expressed as [Benini and Milan, 2018]

$$\begin{aligned}\mathcal{I}(p, q, y) &= \kappa \oint_{\mathbb{T}^{\text{rk}(G)}} \mathcal{Z}(u; p, q, y) \\ &= \kappa \sum_{\hat{u} \in \text{BAEs}} \sum_{\{m_i\}=1}^{ab} \mathcal{Z}(\hat{u} + m\omega; p, q, y) H^{-1}.\end{aligned}$$

Bethe ansatz formula

- For angular fugacities $p = e^{2\pi i a \omega}$, $q = e^{2\pi i b \omega}$, $\gcd\{a, b\} = 1$, the superconformal index can be expressed as [Benini and Milan, 2018]

$$\begin{aligned}\mathcal{I}(p, q, y) &= \kappa \oint_{\mathbb{T}^{\text{rk}(G)}} \mathcal{Z}(u; p, q, y) \\ &= \kappa \sum_{\hat{u} \in \text{BAEs}} \sum_{\{m_i\}=1}^{ab} \mathcal{Z}(\hat{u} + m\omega; p, q, y) H^{-1}.\end{aligned}$$

- The Bethe ansatz equations are defined on a torus with modulus ω and have the form

$$Q_i(u; \Delta, \omega) = 1, \quad i = 1, \dots, \text{rk}(G).$$

Prediction for the entropy

- Conjecture that one term in the sum is dual to the contribution of BHs: \hat{u} evenly distributed on the B-cycle, and $m_i = i \bmod ab$.

Prediction for the entropy

- Conjecture that one term in the sum is dual to the contribution of BHs: \hat{u} evenly distributed on the B-cycle, and $m_i = i \bmod ab$.
- For this contribution, we computed $\log \mathcal{Z}$ at leading order for toric quivers and found

$$\log \mathcal{Z} = -i\pi N^2 \sum_{a,b,c=1}^n \frac{C_{abc} \Delta_a \Delta_b \Delta_c}{6\tau\sigma}$$

C_{abc} are 't Hooft anomaly coefficients.

Prediction for the entropy

- Conjecture that one term in the sum is dual to the contribution of BHs: \hat{u} evenly distributed on the B-cycle, and $m_i = i \bmod ab$.
- For this contribution, we computed $\log \mathcal{Z}$ at leading order for toric quivers and found

$$\log \mathcal{Z} = -i\pi N^2 \sum_{a,b,c=1}^n \frac{C_{abc} \Delta_a \Delta_b \Delta_c}{6\tau\sigma}$$

C_{abc} are 't Hooft anomaly coefficients.

- Reproduces the entropy function for known black holes [Gutowski and Reall, 2004] in $\text{AdS}_5 \times S^5$ [Hosseini et al., 2017] but **does this work in other cases?**

Supergravity checks

For $SE_5 \neq S^5$:

- Supergravity duals contain hypermultiplets and full supersymmetric black hole solutions are not known.

Supergravity checks

For $SE_5 \neq S^5$:

- Supergravity duals contain hypermultiplets and full supersymmetric black hole solutions are not known.
- Simple check: truncating to minimal gauged supergravity, entropy of Gutowski-Reall black hole with single R-charge is reproduced by the entropy function.

Supergravity checks

For $SE_5 \neq S^5$:

- Supergravity duals contain hypermultiplets and full supersymmetric black hole solutions are not known.
- Simple check: truncating to minimal gauged supergravity, entropy of Gutowski-Reall black hole with single R-charge is reproduced by the entropy function.
- Otherwise, try to apply the strategy [Hosseini et al., 2017] .

The strategy

- Rotating charged black holes are topologically AdS_2 fibered over S^3 .

The strategy

- Rotating charged black holes are topologically AdS_2 fibered over S^3 .
- Reducing along the Hopf fiber of S^3 , generically get rotating charged black hole in 4d with the same entropy.

The strategy

- Rotating charged black holes are topologically AdS_2 fibered over S^3 .
- Reducing along the Hopf fiber of S^3 , generically get rotating charged black hole in 4d with the same entropy.
- $J_1 - J_2$ is the 4d angular momentum \implies black holes with $J_1 = J_2$ become spherically symmetric in 4d.

The strategy

- Rotating charged black holes are topologically AdS_2 fibered over S^3 .
- Reducing along the Hopf fiber of S^3 , generically get rotating charged black hole in 4d with the same entropy.
- $J_1 - J_2$ is the 4d angular momentum \implies black holes with $J_1 = J_2$ become spherically symmetric in 4d.
- Restricting to this case, the entropy is determined by horizon data, via the attractor mechanism, which is also an extremization problem.

The conifold

- List of truncations of IIB on SE_5 is limited, the simplest example is a truncation on $T^{1,1}$ given in [Cassani and Faedo, 2011].

The conifold

- List of truncations of IIB on SE_5 is limited, the simplest example is a truncation on $T^{1,1}$ given in [Cassani and Faedo, 2011].
- Only $n_V = 2$ vector multiplets dual to the field theory R charge and baryonic charge were kept. $n_H = 2$.

The conifold

- List of truncations of IIB on SE_5 is limited, the simplest example is a truncation on $T^{1,1}$ given in [Cassani and Faedo, 2011].
- Only $n_V = 2$ vector multiplets dual to the field theory R charge and baryonic charge were kept. $n_H = 2$.
- Restricting to black holes carrying these two charges and performing the dimensional reduction, we found that indeed

entropy function \equiv attractor mechanism

The conifold

- List of truncations of IIB on SE_5 is limited, the simplest example is a truncation on $T^{1,1}$ given in [Cassani and Faedo, 2011].
- Only $n_V = 2$ vector multiplets dual to the field theory R charge and baryonic charge were kept. $n_H = 2$.
- Restricting to black holes carrying these two charges and performing the dimensional reduction, we found that indeed

entropy function \equiv attractor mechanism

- Charges were matched using the AdS_5/CFT_4 dictionary and then reduced.

Outlook

- Construction of rotating black holes with generic electric charges in truncations on SE_5 .

Outlook

- Construction of rotating black holes with generic electric charges in truncations on SE_5 .
- Is there an attractor mechanism in 5d? Or a rotating attractor mechanism in 4d?

Outlook

- Construction of rotating black holes with generic electric charges in truncations on SE_5 .
- Is there an attractor mechanism in 5d? Or a rotating attractor mechanism in 4d?
- Can we compute subleading corrections to the Bekenstein-Hawking entropy?

Outlook

- Construction of rotating black holes with generic electric charges in truncations on SE_5 .
- Is there an attractor mechanism in 5d? Or a rotating attractor mechanism in 4d?
- Can we compute subleading corrections to the Bekenstein-Hawking entropy?
- Can we go beyond extremality and SUSY?