

# STRINGS AT THE END OF THE SWAMPLAND

#### Stefano Lanza

based on arXiv: **2104.05726** with **Fernando Marchesano, Luca Martucci, Irene Valenzuela** 

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# PLAN OF THE TALK

#### Introduction:

the Swampland Program

Swampland distance conjecture: infinite field distances and EFT inconsistencies

#### A novel perspective:

Axionic strings as tools to explore the moduli space

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# INCLUDING GRAVITY IN EFTS



Can the EFT be embedded into a proper quantum theory of gravity? Or: Does the EFT have a quantum gravity origin?

# SWAMPLAND CONJECTURES

Only some them are embedded within a **quantum gravity completion**:



The Swampland program aims at identifying the borders via "Swampland Conjectures"

# THE WEB OF SWAMPLAND CONJECTURES

Many consistency criteria, named *Swampland Conjectures*, have been identified:



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Many consistency criteria, named *Swampland Conjectures*, have been identified:



# THE SWAMPLAND DISTANCE CONJECTURE

### SINGULARITIES AND INFINITE FIELD DISTANCES

Consider a generic *d*-dimensional EFT, valid up to the cutoff  $\Lambda$ :

$$S_d = \int \left(\frac{1}{2}R - \mathcal{G}_{ij}(\phi)\partial\phi^i \cdot \partial\phi^j + \ldots\right) * 1$$

 $\phi^i$  span the manifold  $\,\mathcal{M}\,$  and  $\,\mathcal{G}_{ij}(\phi)$  is the field space metric. where



Geodesic distance between the points P and Q:

$$d(P,Q) = \int_{\gamma} \sqrt{\mathcal{G}_{ij}(\phi)\dot{\phi}^i \dot{\phi}^j} \, \mathrm{d}\sigma$$

Some singular points are located at infinite distance

 $d(P,Q) \to \infty$ 

# THE SWAMPLAND DISTANCE CONJECTURE

[Ooguri, Vafa, 2006]

The **Swampland Distance Conjecture** states that, in any EFT consistent with quantum gravity:

 $\exists$  an infinite tower of states that becomes exponentially light at any infinite field distance limit. Namely:  $M(Q) \sim M(P) e^{-\lambda \operatorname{d}(P,Q)}$ 

in terms of the geodesic distance d(P,Q) and with  $\lambda$  an O(1)-parameter.



<u>Example</u>: Consider a 10D string theory compactified over an  $S^1$ . In 9D, a modulus  $\rho$ , the *radion*, is present

$$S_9 = \int \left(\frac{1}{2}R + \frac{1}{2\rho^2}\partial\rho \cdot \partial\rho + \ldots\right)$$

The distance conjecture is realized towards decompactification limits:

 $ho \to \infty$  at  $d(
ho_0, 
ho) \to \infty$   $\longrightarrow$   $\exists$  Infinite tower of KK modes  $m_{\rm KK}^2 \sim \frac{n}{
ho^2} \sim m_{{\rm KK},0}^2 e^{-2d(P,P_0)}$ 

# THE SWAMPLAND DISTANCE CONJECTURE

[Ooguri, Vafa, 2006]

The **Swampland Distance Conjecture** states that, in any EFT consistent with quantum gravity:

 $\exists$  an infinite tower of states that becomes exponentially light at *any* infinite field distance limit. Namely:

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Many top-down arguments support the conjecture:

- Tower of D3-particles in general N=2 settings [Grimm, Li, Palti, Valenzuela 2018-2020]
- Emergent string conjecture [Lee, Lerche, Weigand 2019]
- Towers of strings and membranes
   [Font, Herraez, Ibanez 2019]

# THE SWAMPLAND DISTANCE CONJECTURE

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How can the Distance Conjecture be realized at the EFT level?Or:Is there a mechanism which drives the fields to the field space boundary?

# A NOVEL PERSPECTIVE: INTRODUCING AXIONIC STRINGS

# AXIONIC STRINGS

Focus on 4D EFTs. The objects that realize the Distance Conjecture at the EFT level are

### **Fundamental Axionic Strings**

Their distinguishing features are:

- strict codimension two;
- electrically coupled to gauge two-forms via:

$$e^i \int_{\mathcal{S}} \mathcal{B}_{2\,i}$$

magnetically coupled to the dual axions, such that, encircling the string:

$$a^i \rightarrow a^i + e^i$$

• they are semiclassical, i.e. within an EFT with cutoff  $\Lambda$ ,

$$\mathcal{T}_{
m str}\gtrsim\Lambda^2$$



# AXIONIC STRINGS

Focus on 4D EFTs. The objects that realize the Distance Conjecture at the EFT level are

### **Fundamental Axionic Strings**

Axionic strings can be included in EFTs as unresolved, semiclassical objects via the action

Nambu-Goto term

Chern-Simons term

$$S_{\rm str} = -\int_{\mathcal{S}} \mathrm{d}^2 \xi \sqrt{-h} \,\mathcal{T}_{\rm str} + e^i \int_{\mathcal{S}} \mathcal{B}_{2\,i}$$

String tension (field dependent)

Electric charges/axion windings

V



## AXIONIC STRINGS: WHAT ARE THEY?





## AXIONIC STRINGS IN N=1 SUPERGRAVITY

Consider a set of chiral multiplets  $\Phi^i$  is



with the Kähler potential

$$K(\phi, \bar{\phi}) \equiv K(s)$$
  
Invariant under axionic shifts  
$$a^i \to a^i + c^i , \qquad c^i \in \mathbb{R}$$

• The bosonic components of an N=1 supersymmetric action of a set of chiral multiplets  $\Phi^i$  is

$$S = \int M_{\rm P}^2 \left( \frac{1}{2} R - 2\mathcal{G}_{ij} \partial s^i \cdot \partial s^j - 2\mathcal{G}_{ij} \partial a^i \cdot \partial a^j \right) * 1$$

with the field space metric

$$\mathcal{G}_{ij}(s) = \frac{1}{2} \frac{\partial^2 K}{\partial s^i \partial s^j}$$

# AXIONIC STRINGS IN N=1 SUPERGRAVITY



The dual action expressed in terms of a set of linear multiplets is:

$$S = \frac{1}{2} \int \left( M_{\mathrm{P}}^2 R * 1 - M_{\mathrm{P}}^2 \mathcal{G}^{ij} \mathrm{d}\ell_i \wedge * \mathrm{d}\ell_j - \frac{1}{M_{\mathrm{P}}^2} \mathcal{G}^{ij} \mathcal{H}_{3\,i} \wedge * \mathcal{H}_{3\,j} \right)$$

…to which we can couple a BPS string

$$S_{\rm str} = -\int_{\mathcal{S}} \mathrm{d}^2 \xi \sqrt{-h} \,\mathcal{T}_{\rm str} + e^i \int_{\mathcal{S}} \mathcal{B}_{2\,i}$$

The string tension is fully fixed by supersymmetry  $\mathcal{T}_{\rm str} = M_{\rm P}^2 e^i \ell_i(s) = -\frac{1}{2} M_{\rm P}^2 e^i \frac{\partial K}{\partial s^i}$ 

# AXIONIC STRINGS IN N=1 SUPERGRAVITY



The dual action expressed in terms of a set of linear multiplets is:

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Solve the equations of motion of bulk + string ⇒ String backreaction

 $a^i \rightarrow a^i + e^i$ 

 $r_{\Lambda} = \frac{1}{\Lambda}$ 

The backreaction of the string onto the fields is



 $\gamma_{\mathbf{e}} \equiv \{ \bar{\boldsymbol{t}} = \bar{\boldsymbol{a}}_0 + \mathrm{i}\bar{\boldsymbol{s}}(\sigma) \} \subset \mathcal{M}, \qquad \bar{\boldsymbol{s}}(\sigma) = \bar{\boldsymbol{s}}_0 + \sigma \mathbf{e}, \qquad \bar{\boldsymbol{a}}_0, \bar{\boldsymbol{s}}_0 = \mathrm{const.}$ 

The backreaction of the string onto the fields is

$$s(r) = s_0 - \frac{e}{2\pi} \log \frac{r}{r_0}, \qquad a = \frac{\theta}{2\pi} e + \text{const}$$

Approaching the string core

 $s \to \infty$ 

#### For Kähler potentials such as

$$K = -n\log s$$

these are infinite field distance limits

 $d(s,s_0)\to\infty$ 

*Infinite distances reached toward the core* 



The backreaction of the string onto the fields is

$$s(r) = s_0 - \frac{e}{2\pi} \log \frac{r}{r_0}, \qquad a = \frac{\theta}{2\pi} e + \text{const}$$

Approaching the string core, the string appears tensionless

$$\mathcal{T}_{\rm str} = M_{\rm P}^2 e \,\ell(r) = \frac{M_{\rm P}^2 e \,n}{2s_0 - \frac{e}{\pi} \log \frac{r}{r_0}} \xrightarrow{r \to 0} 0$$

Define the energy scale  $\Lambda = r^{-1}$ Consider the string tension as EFT coupling

 $\blacksquare$  ...and we regard the backreaction as an RG flow

[Polchinski et al., 2015, Goldberger, Wise, 2001]  
$$\mathcal{T}_{\rm str}^{\rm eff}(\Lambda) = \frac{\mathcal{T}_{\rm str}^0}{1 + \frac{\mathcal{T}_{\rm str}^0}{2\pi M_{\rm P}^2} \log(\Lambda r_0)}$$

#### Weak coupling closer to the string!

Changing energy scales, following the RG flow, different, distant regions can be explored

 $r_{\Lambda} = \frac{1}{\Lambda}$ 

 $\overline{S}$ 

 $s^i(r)$ 

 $\mathcal{M}_{\Lambda}$ 

 $a^i \rightarrow a^i + \epsilon$ 

The backreaction of the string onto the fields is

$$s(r) = s_0 - \frac{e}{2\pi} \log \frac{r}{r_0}, \qquad a = \frac{\theta}{2\pi} e + \text{const}$$

Approaching the string core, the string appears tensionless

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As the string becomes tensionless, an <u>infinite</u> <u>tower of states</u> appears

$$m \sim \mathcal{T}_{\mathrm{str}}^{1/2} \sim (\mathcal{T}_{\mathrm{str}}^0)^{1/2} \exp\left(-\frac{\gamma}{2}\mathrm{d}(s,s_0)\right) \to 0$$

with 
$$\gamma = \mathcal{Q}_{\mathrm{str}} / \mathcal{T}_{\mathrm{str}}$$
 and  $\mathcal{Q}_{\mathrm{str}} = \mathcal{G}_{ij} e^i e^j$ 

The EFT breaks down approaching the string!

 $\mathcal{M}_\Lambda$ 

 $s^i(r)$ 

 $r_{\Lambda} = \frac{1}{\Lambda}$ 

 $\overline{S}$ 

 $a^i \rightarrow a^i + e$ 

 $d \rightarrow \infty$ 

The backreaction of the string onto the fields is

$$s(r) = s_0 - \frac{e}{2\pi} \log \frac{r}{r_0}, \qquad a = \frac{\theta}{2\pi} e + \text{const}$$

Approaching the string core, or moving along the RG flow

- Boundaries of moduli space reached
- Infinite field distances
- Infinite tower of massless states

$$m \sim \mathcal{T}_{\mathrm{str}}^{1/2} \sim (\mathcal{T}_{\mathrm{str}}^0)^{1/2} \exp\left(-\frac{\gamma}{2}\mathrm{d}(s,s_0)\right) \to 0$$

We are realising the distance conjecture at the EFT level!

![](_page_22_Picture_9.jpeg)

# THE DISTANT AXIONIC STRING CONJECTURE

#### Distant Axionic String Conjecture (DASC): All infinite distance limits of a 4d EFT can be realised as an RG flow endpoint of a fundamental axionic string.

![](_page_23_Figure_2.jpeg)

For finite cutoff  $\Lambda$ , the moduli space is only partially explorable.

Infinite field distances are explorable via the string RG flow.

# THE DISTANT AXIONIC STRING CONJECTURE

Distant Axionic String Conjecture (DASC):

All infinite distance limits of a 4d EFT can be realised as an RG flow endpoint of a fundamental axionic string.

Infinite tower of states becomes massless

 $\Rightarrow$  EFT breaking down

In Stringy EFTs: KK/winding modes, D0 branes, ... (see SL, Marchesano, Martucci, Valenzuela '21 for the detailed analysis)

Infinite field distances are explorable via the string RG flow.

 $\mathcal{M}$ 

 $\rightarrow a^i + e^i$ 

 $\mathcal{S}$ 

# THE DISTANT AXIONIC STRING CONJECTURE

#### Distant Axionic String Conjecture (DASC): All infinite distance limits of a 4d EFT can be realised as an RG flow endpoint of a fundamental axionic string.

#### Supported by:

#### **Cut-off Asymptotics:**

Along an asymptotic limit specified by the RG flow of an EFT string, its tension goes to zero. The maximal EFT cut-off  $m_*$  then scales like:

$$m_*^2 \simeq M_{\rm P}^2 A \left(\frac{\mathcal{T}_{\rm str}}{M_{\rm P}^2}\right)^u$$

with the scaling weight  $w \in \mathbb{N}$ . (In stringy EFTs w = 1,2,3)

![](_page_25_Figure_7.jpeg)

Infinite field distances are explorable via the string RG flow.

# CONCLUSIONS

# CONCLUSIONS AND FUTURE OUTLOOK

With Axionic strings we can explore the moduli space

- They realise the Distance Conjecture at the EFT level;
- Their existence may be experimentally tested;
- They draw the correspondence among:

![](_page_27_Figure_5.jpeg)

Some open questions:

- Extension to non-supersymmetric cases?
- Closer relations to other conjectures?
- Closer connections to Nielsen-Olesen strings or cosmic strings in literature?

Thank you!

# BACKUP SLIDES

# THE STRING BACKREACTION AS RG FLOW

The flow of the tension is

$$\mathcal{T}_{\rm str} = M_{\rm P}^2 e\,\ell(r) = \frac{e\,n\,M_{\rm P}^2}{2s_0 - \frac{e}{\pi}\log\frac{r}{r_0}}$$

which breaks down at a distance

$$r_{\rm IR} = r_0 \exp\left[\frac{n\pi M_P^2}{\mathcal{T}_{\rm str}(r_0)}\right]$$

In the limit  $r \to 0$ :  $\mathcal{T}_{str} \to 0$ 

![](_page_30_Picture_6.jpeg)

We regard the profile of the tension as RG-flow of the tension

$$\mathcal{T}_{\rm str}^{\rm eff}(\Lambda) = \frac{\mathcal{T}_{\rm str}^0}{1 + \frac{\mathcal{T}_{\rm str}^0}{2\pi M_{\rm P}^2} \log(\Lambda r_0)}$$

and the EFT breaks at the strong coupling scale

$$\Lambda_{\rm strong} = \Lambda_0 \exp\left[-\frac{n\pi \, M_P^2}{\mathcal{T}_{\rm str}(\Lambda_0)}\right]$$

On the other hand, the limit  $r \rightarrow 0$  corresponds to weak coupling.

![](_page_30_Picture_12.jpeg)

![](_page_30_Picture_13.jpeg)

### EXPLORING THE MODULI SPACE VIA STRINGS

The string tension has the following symmetry

$$\mathcal{T}_{\mathbf{e}}(\Lambda e^{2\pi\sigma}, \bar{s}) = \mathcal{T}_{\mathbf{e}}(\Lambda, \bar{s} + \mathbf{e}\sigma)$$

that is, a change in the cutoff can be seen as a change in the field configuration

![](_page_31_Figure_4.jpeg)

$$\bar{s}_0 \to \bar{s}(\sigma) = \bar{s}_0 + \sigma \mathbf{e} \quad \Leftrightarrow \quad \Lambda \to \Lambda e^{2\pi\sigma}$$