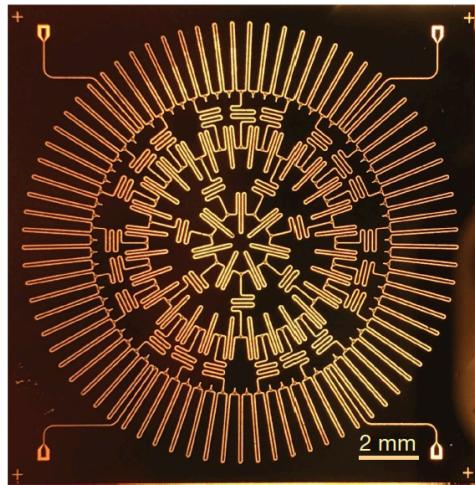


# Hyperbolic band theory



**Joseph Maciejko**  
**University of Alberta**

**Topological Properties of Gauge Theories & their Applications to HEP and CMP**  
**Galileo Galilei Institute for Theoretical Physics**  
**September 16, 2021**



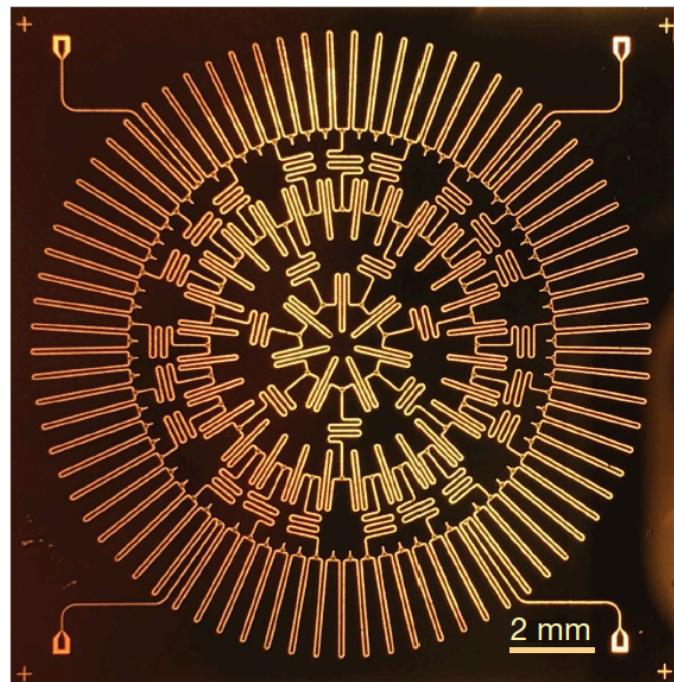
Steven Rayan  
(U. Saskatchewan)

**JM and S. Rayan, Sci. Adv. 7, eabe9170 (2021)**  
**JM and S. Rayan, arXiv:2108.09314**

# Hyperbolic lattices in circuit quantum electrodynamics

Alicia J. Kollár<sup>1,2,3\*</sup>, Mattias Fitzpatrick<sup>1</sup> & Andrew A. Houck<sup>1</sup>

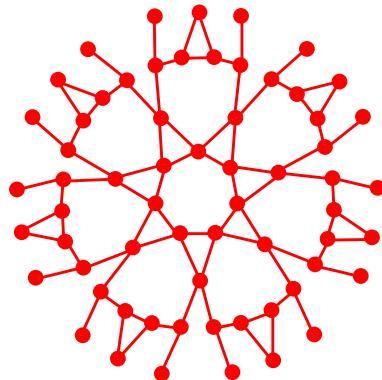
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# Hyperbolic lattices in circuit quantum electrodynamics

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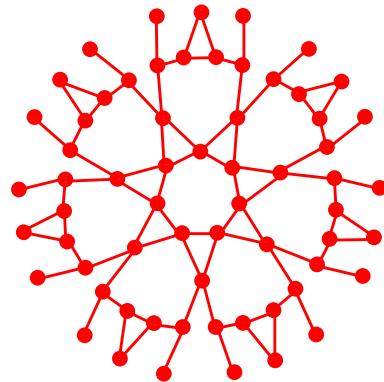


$$H_{\text{TB}} = \omega_0 \sum_i a_i^\dagger a_i - t \sum_{\langle i,j \rangle} (a_i^\dagger a_j + a_j^\dagger a_i)$$

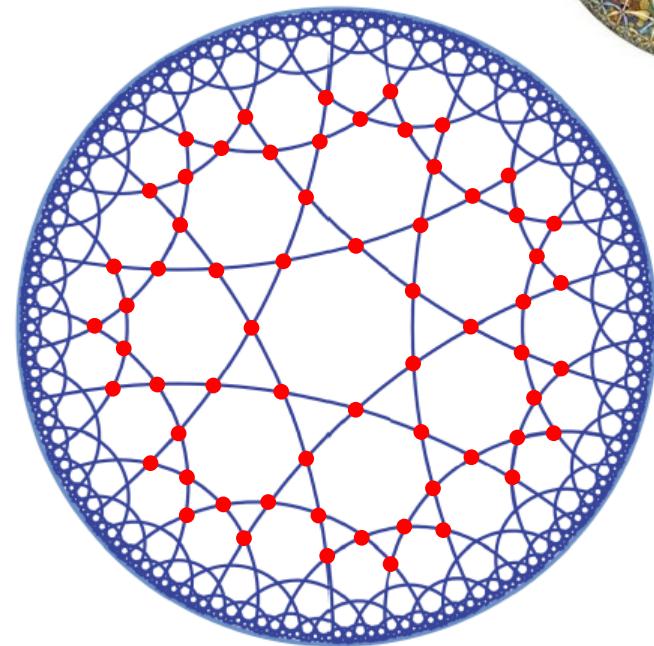
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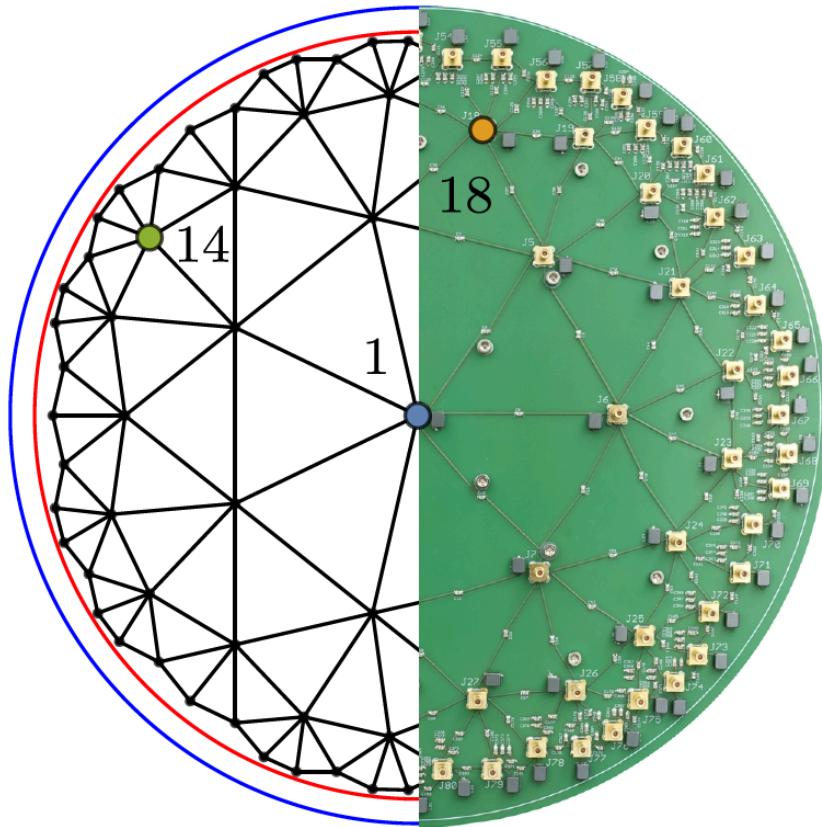
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# Electric-circuit realization of a hyperbolic drum

Patrick M. Lenggenhager <sup>1, 2, 3,\*</sup> Alexander Stegmaier  <sup>4,\*</sup> Lavi K. Upreti  <sup>4</sup> Tobias Hofmann, <sup>4</sup> Tobias Helbig  <sup>4</sup> Achim Vollhardt, <sup>2</sup> Martin Greiter, <sup>4</sup> Ching Hua Lee, <sup>5</sup> Stefan Imhof, <sup>6</sup> Hauke Brand, <sup>6</sup> Tobias Kießling, <sup>6</sup> Igor Boettcher  <sup>7, 8</sup> Titus Neupert  <sup>2, †</sup> Ronny Thomale  <sup>4, †</sup> and Tomáš Bzdušek  <sup>1, 2, †</sup>

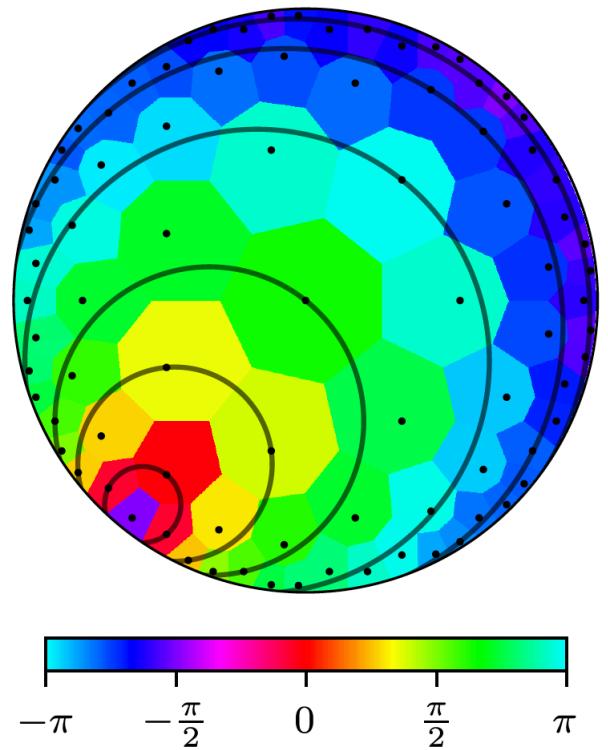
arXiv:2109.01148

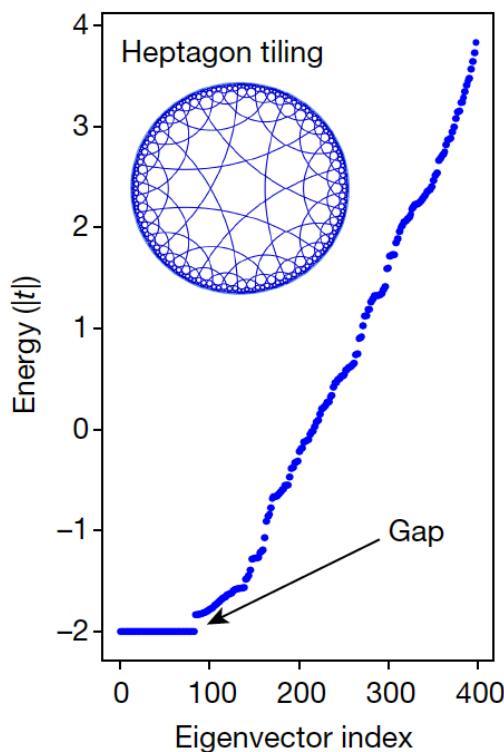
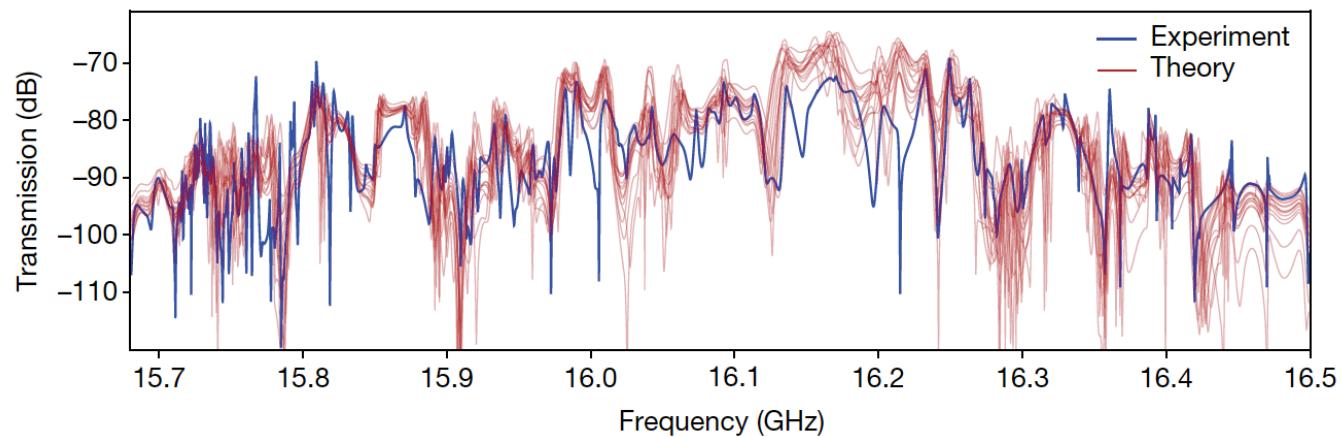
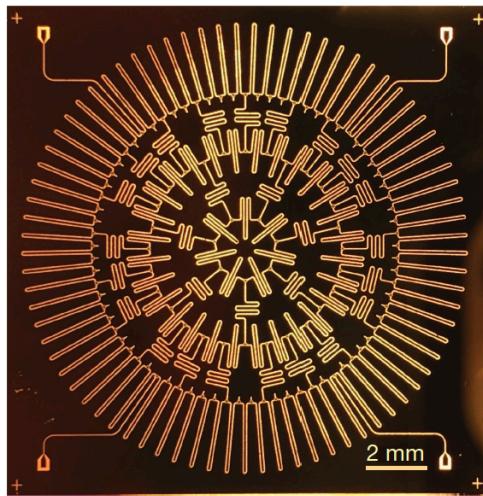


Euclidean drum

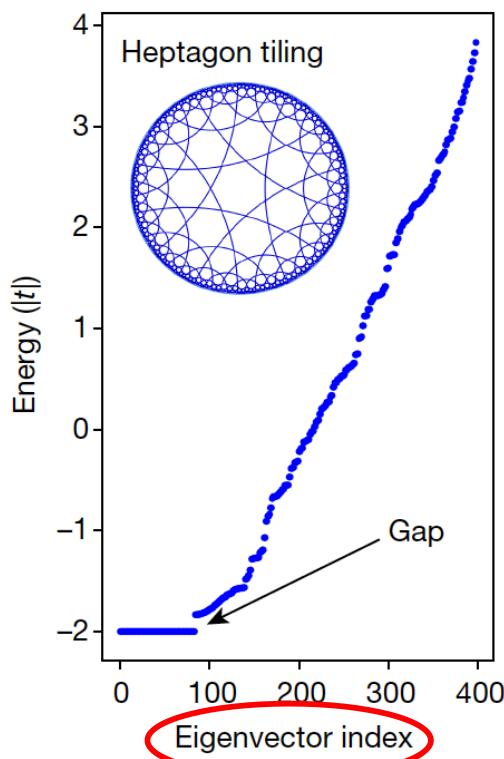
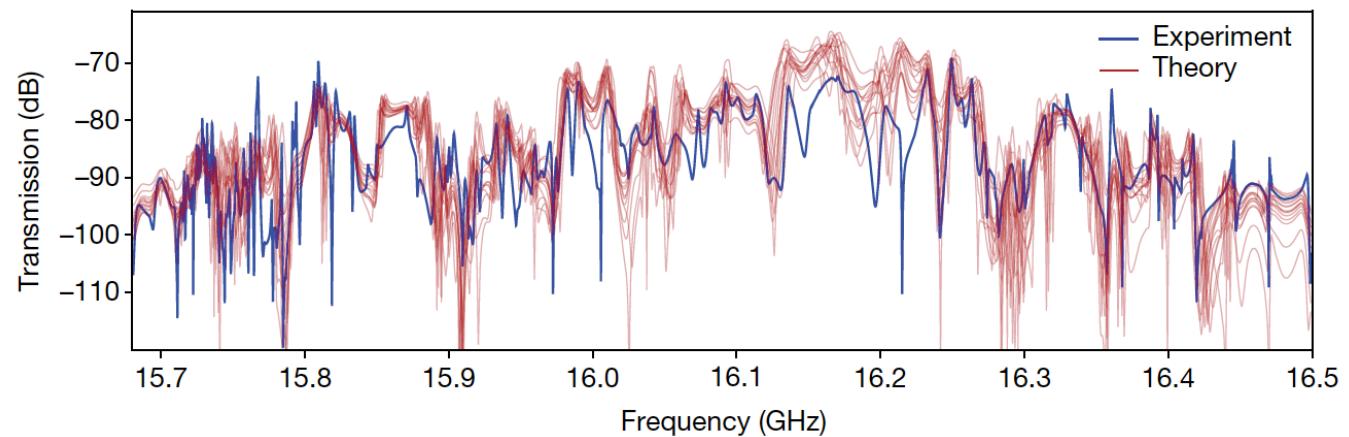
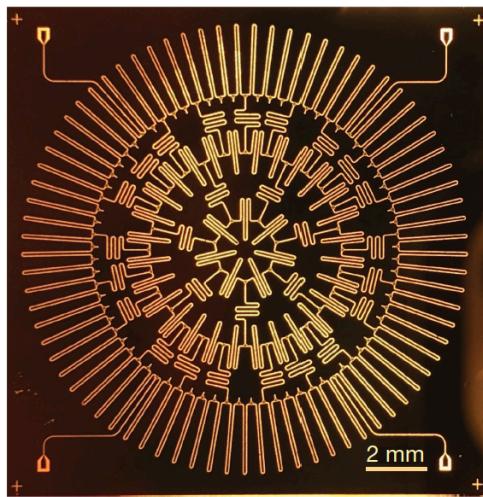
Hyperbolic drum

$t = 2.032 \mu\text{s}$





**Curved-space tight-binding models.** No hyperbolic equivalent of Bloch theory currently exists, and there is no known general procedure for calculating band structures in either the nearly-free-electron or tight-binding limits. Specialized methods are known for the cases of trees<sup>49,50</sup> but fail if there are any closed loops, except in the special case of Cayley graphs of the free products of cyclic groups<sup>32</sup>. The only universal method is numerical diagonalization of the hopping Hamiltonian. This is a brute-force method which yields a list of eigenvectors and eigenvalues, but no classification of eigenstates by a momentum quantum number.

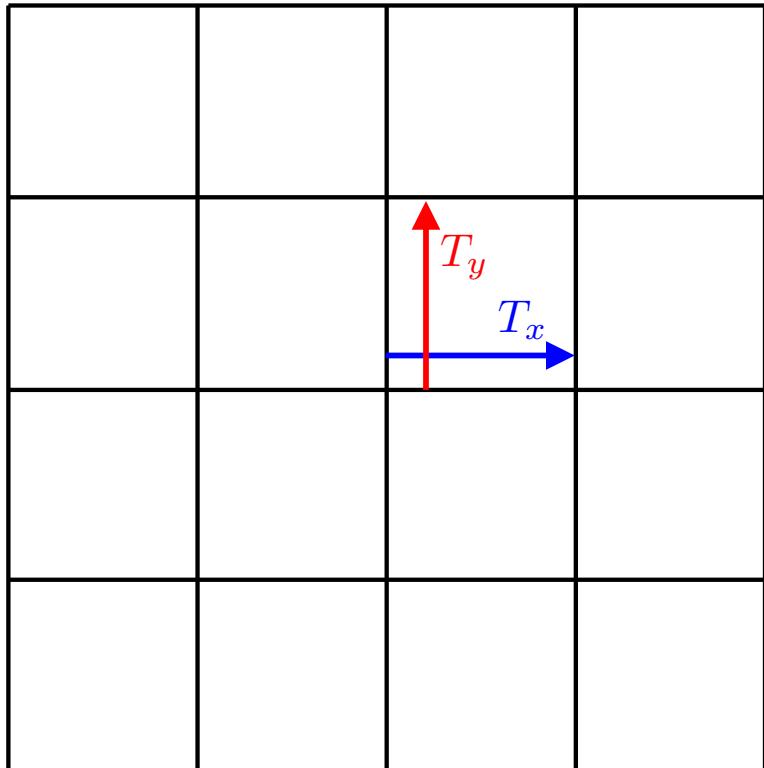


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# Outline

- Euclidean band theory
- Hyperbolic geometry & Fuchsian groups
- The {8,8} lattice
- Periodic boundary conditions & automorphic Bloch theorems
- Conclusion

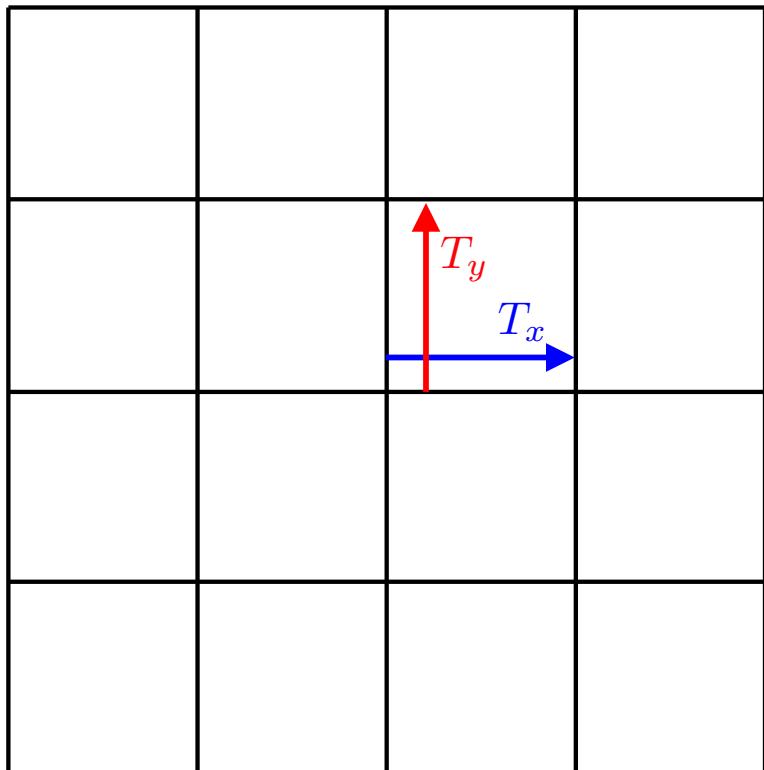
# Euclidean lattices



$$H = \frac{p^2}{2m} + V(x, y)$$

# Euclidean Bloch condition

$$H\psi = E\psi$$

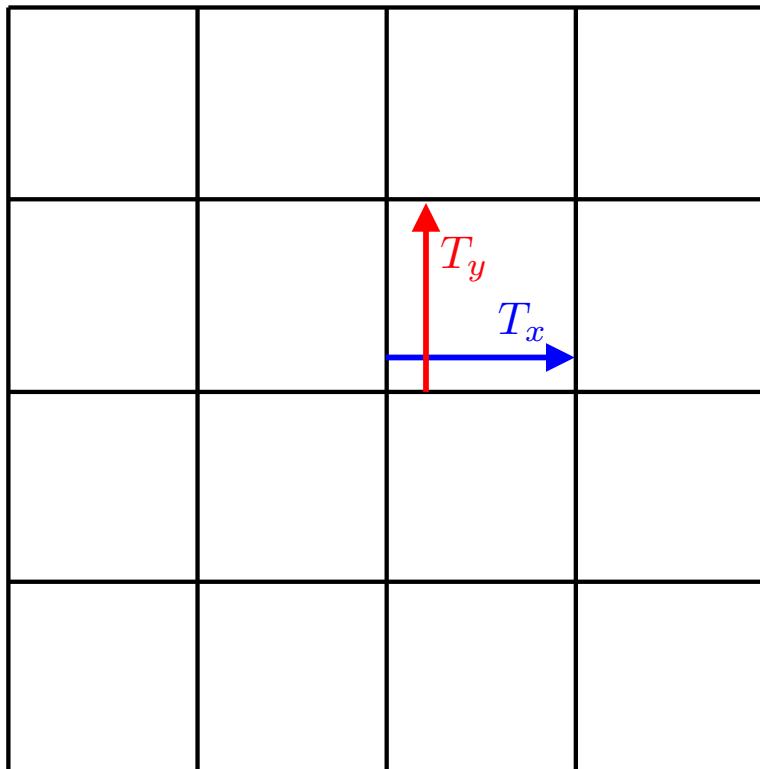


$$\psi(x + 1, y) = e^{ik_x} \psi(x, y)$$

$$\psi(x, y + 1) = e^{ik_y} \psi(x, y)$$

# Brillouin zone torus

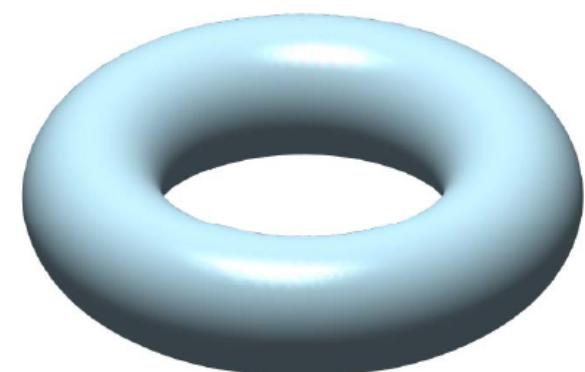
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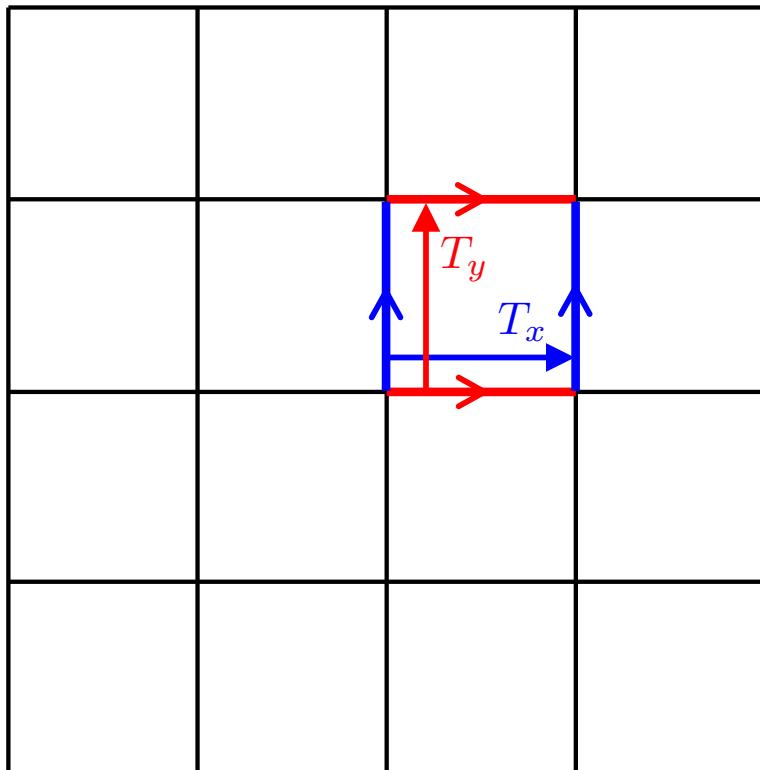
$$\psi(x + 1, y) = e^{ik_x} \psi(x, y)$$
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$$k_x \sim k_x + 2\pi$$

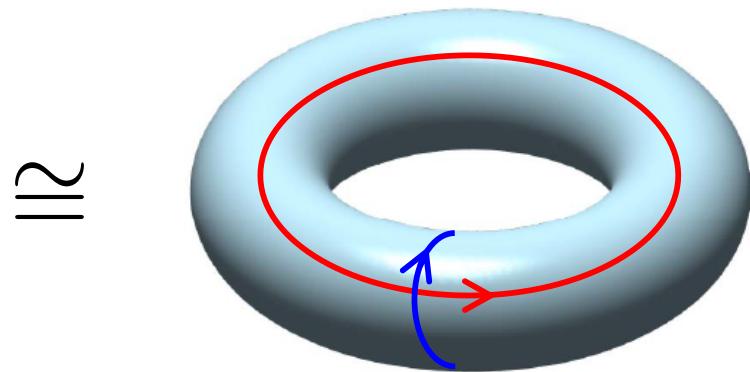
$$k_y \sim k_y + 2\pi$$



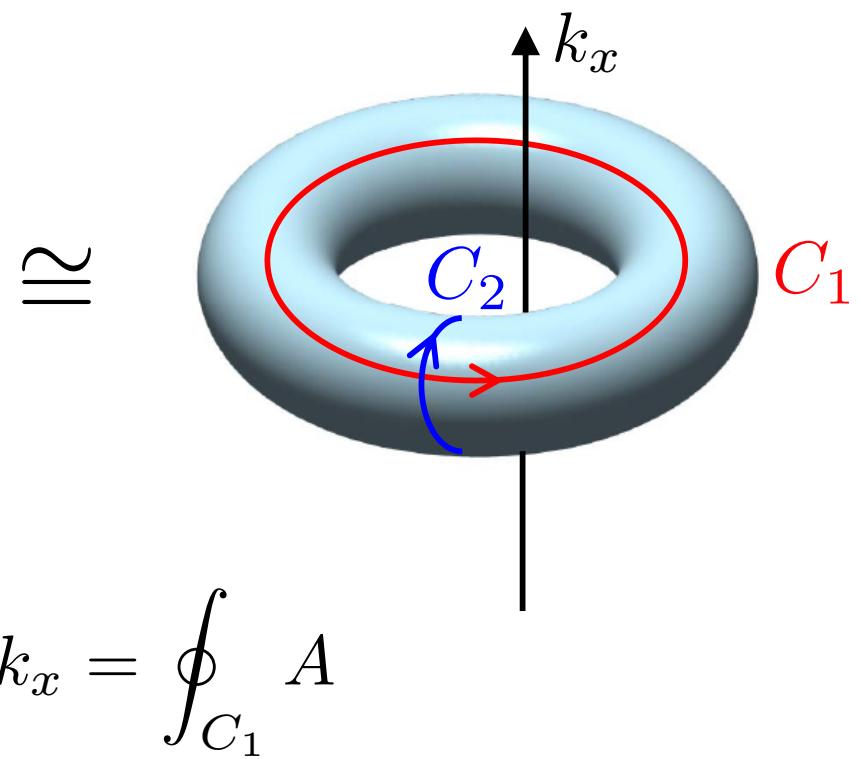
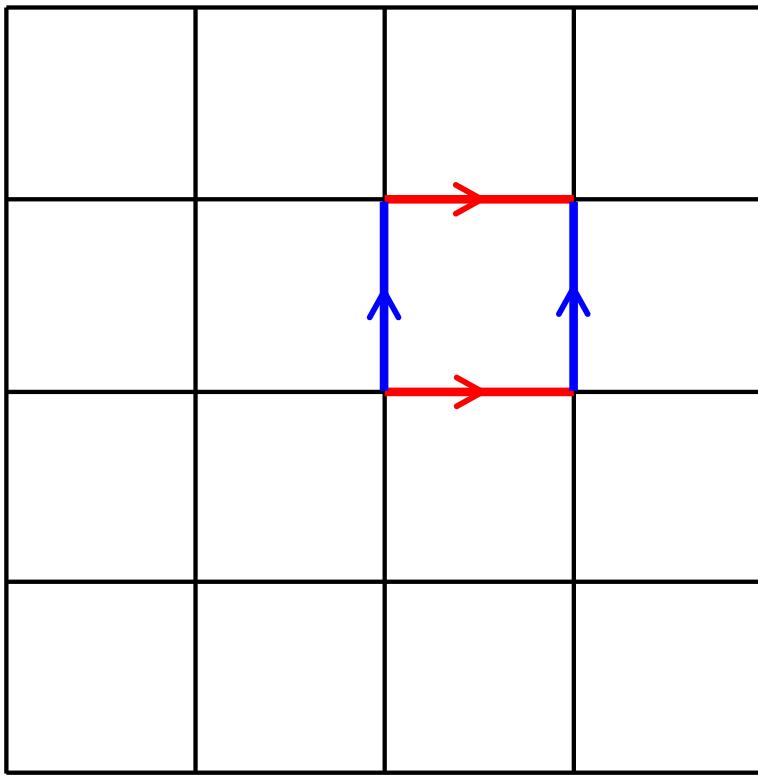
# Compactified unit cell



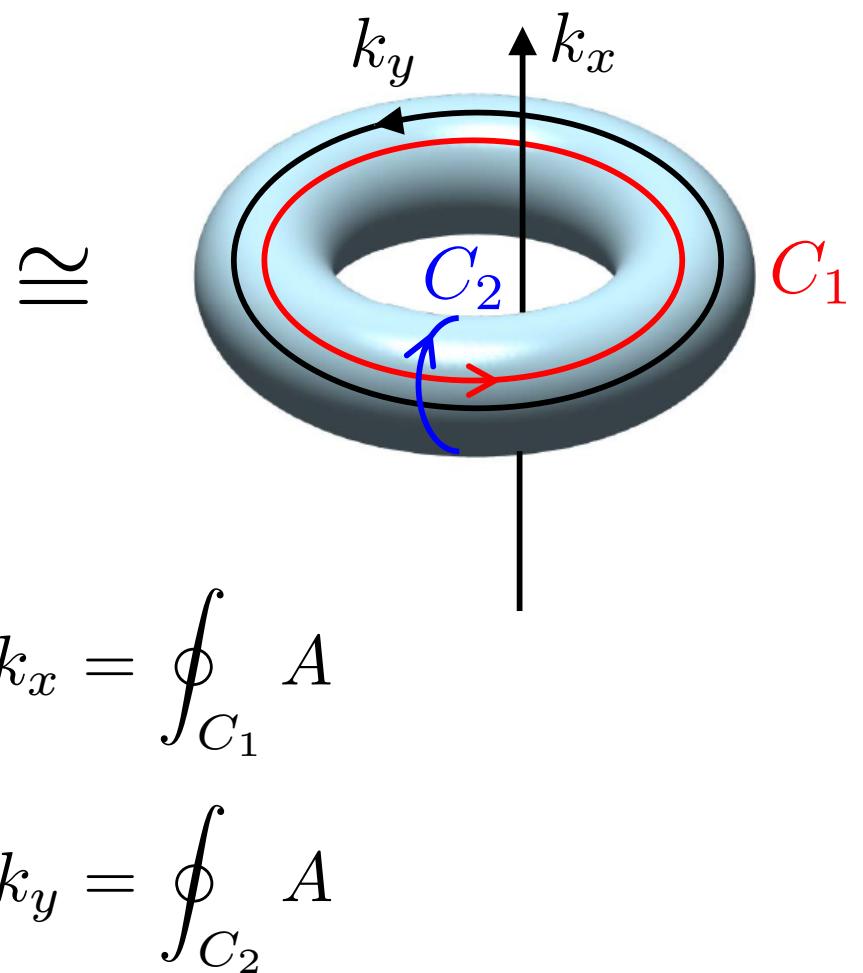
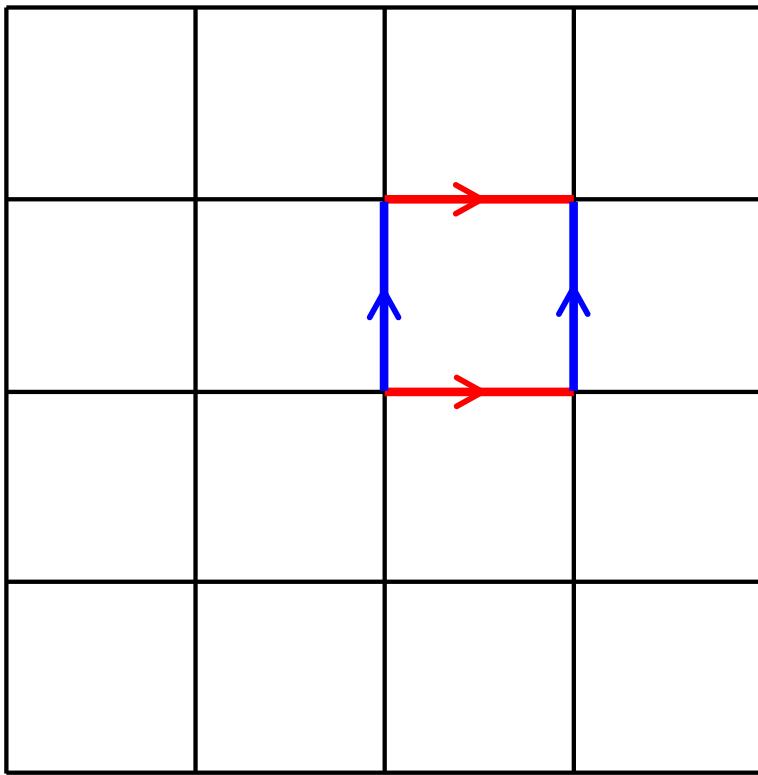
$$\mathbb{R}^2 / \mathbb{Z}^2 \cong T^2$$



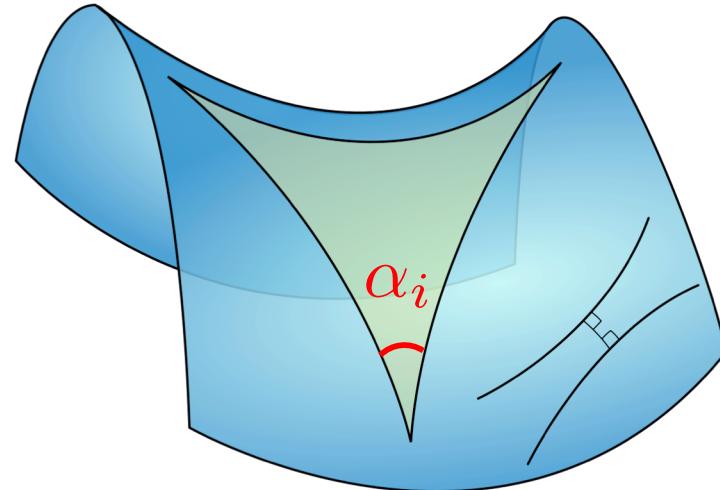
# Aharanov-Bohm fluxes



# Aharanov-Bohm fluxes

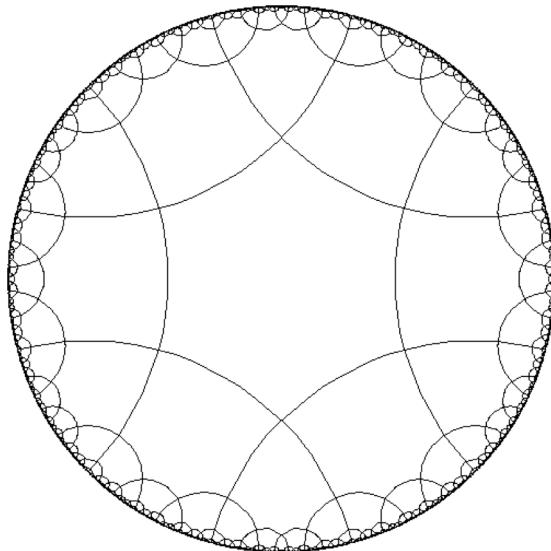


# Hyperbolic lattices

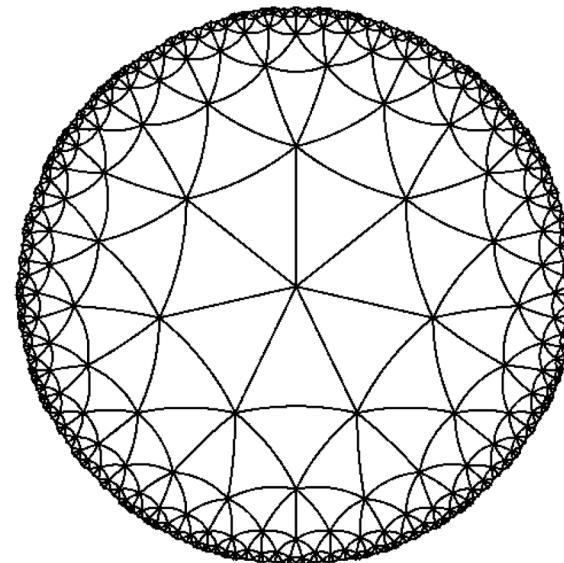


$$\sum_{i=1}^p \alpha_i < (p - 2)\pi$$

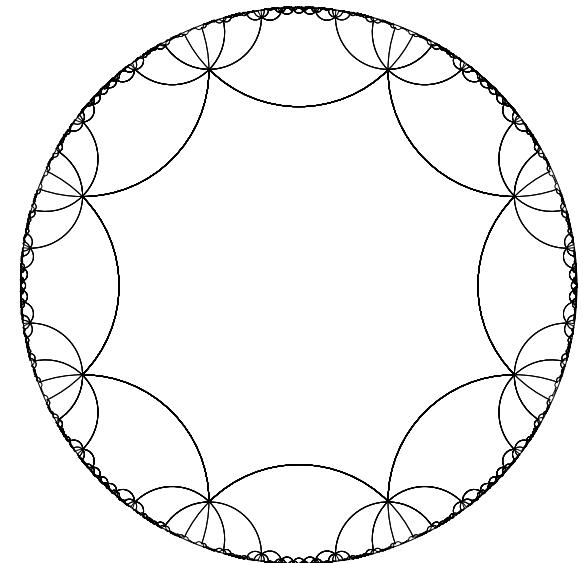
$\{6, 4\}$



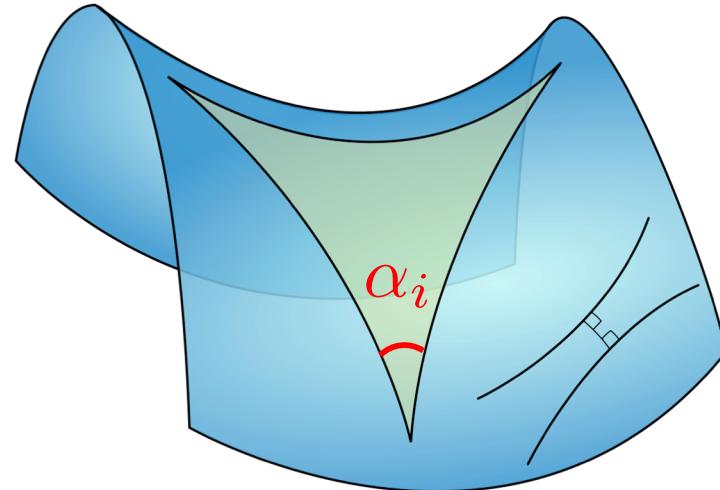
$\{3, 7\}$



$\{8, 8\}$

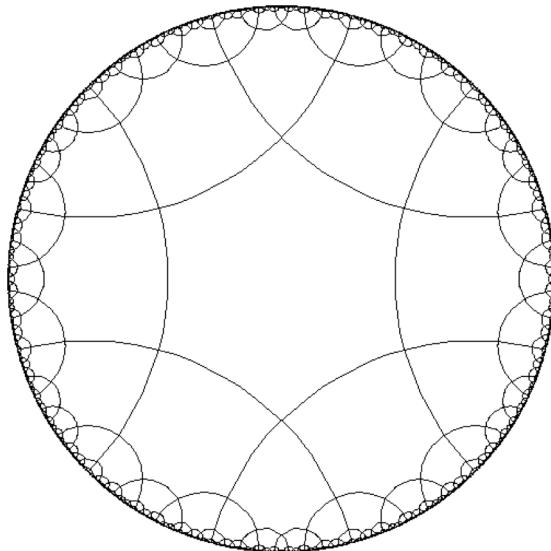


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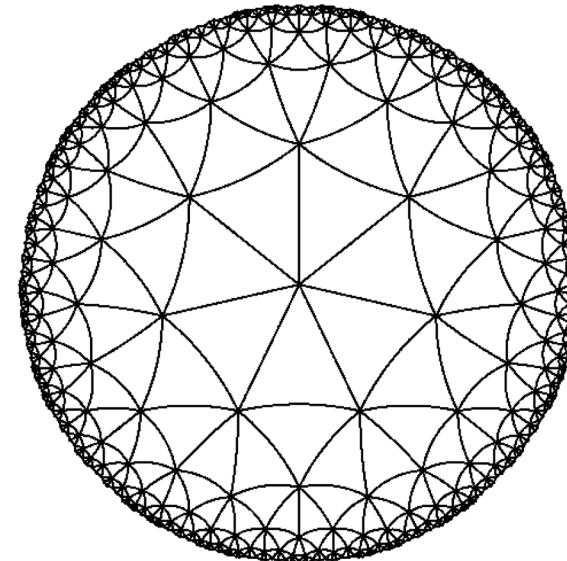


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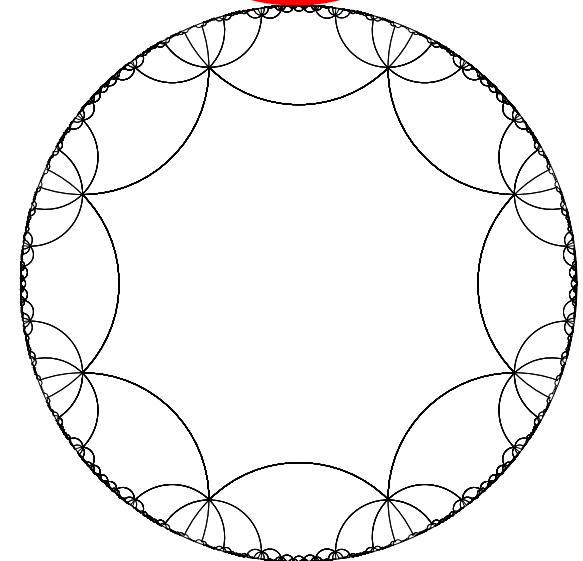
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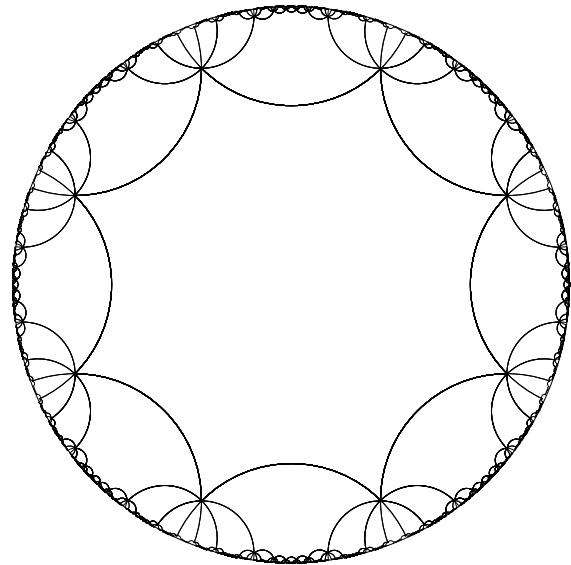


$\{8, 8\}$



# Poincaré disk

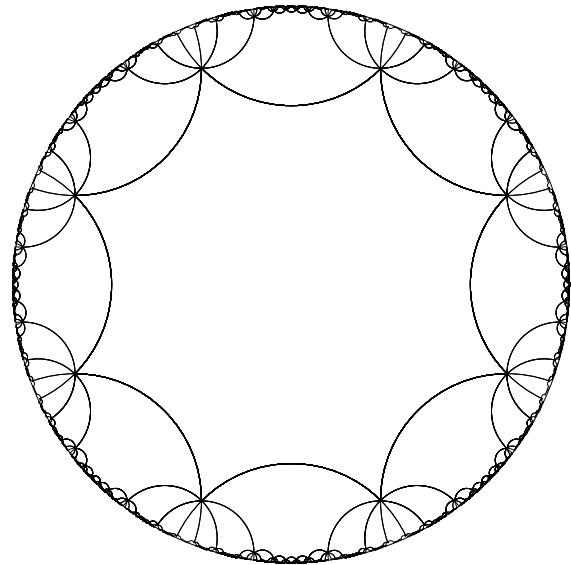
$$\mathbb{H} = \{|z| < 1\}$$



$$ds^2 = \frac{4(dx^2 + dy^2)}{(1 - |z|^2)^2}$$

# Poincaré disk

$$\mathbb{H} = \{|z| < 1\}$$



$\text{PSU}(1,1) \cong \text{PSL}(2,\mathbb{R})$ :

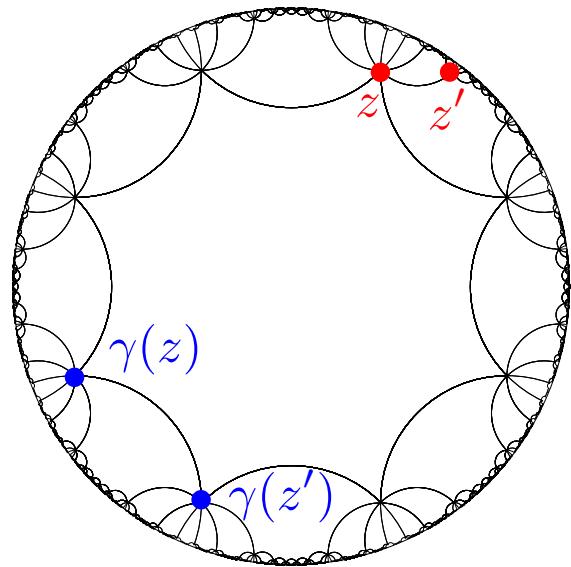
$$\gamma = \begin{pmatrix} \alpha & \beta \\ \beta^* & \alpha^* \end{pmatrix}, \det \gamma = 1$$

$$z \rightarrow \gamma(z) = \frac{\alpha z + \beta}{\beta^* z + \alpha^*}$$

$$ds^2 = \frac{4(dx^2 + dy^2)}{(1 - |z|^2)^2}$$

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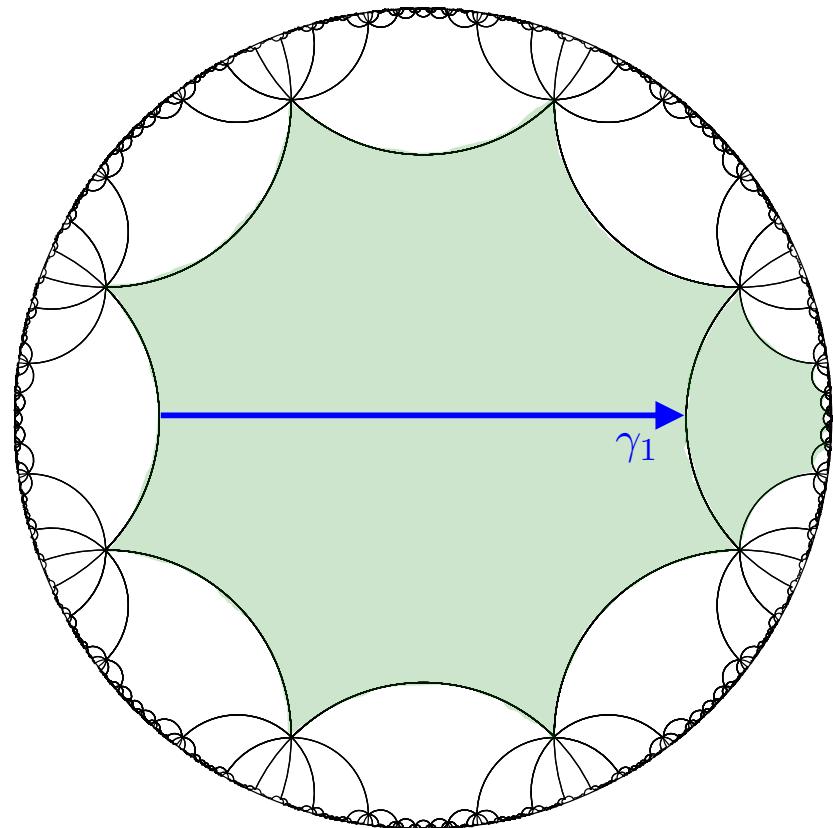
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$$d(z, z') = d(\gamma(z), \gamma(z'))$$

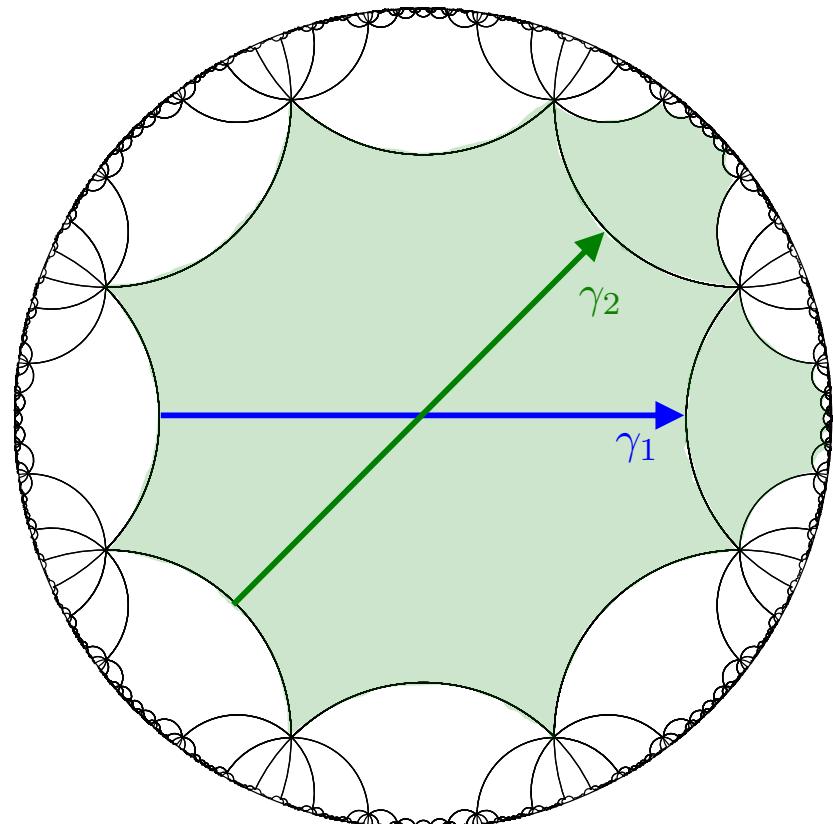
# Fuchsian groups



$$\Gamma \subset PSU(1, 1)$$

$$\gamma_j = \begin{pmatrix} 1 + \sqrt{2} & e^{i(j-1)\pi/4} \sqrt{2 + 2\sqrt{2}} \\ e^{-i(j-1)\pi/4} \sqrt{2 + 2\sqrt{2}} & 1 + \sqrt{2} \end{pmatrix}, \quad j = 1, \dots, 4$$

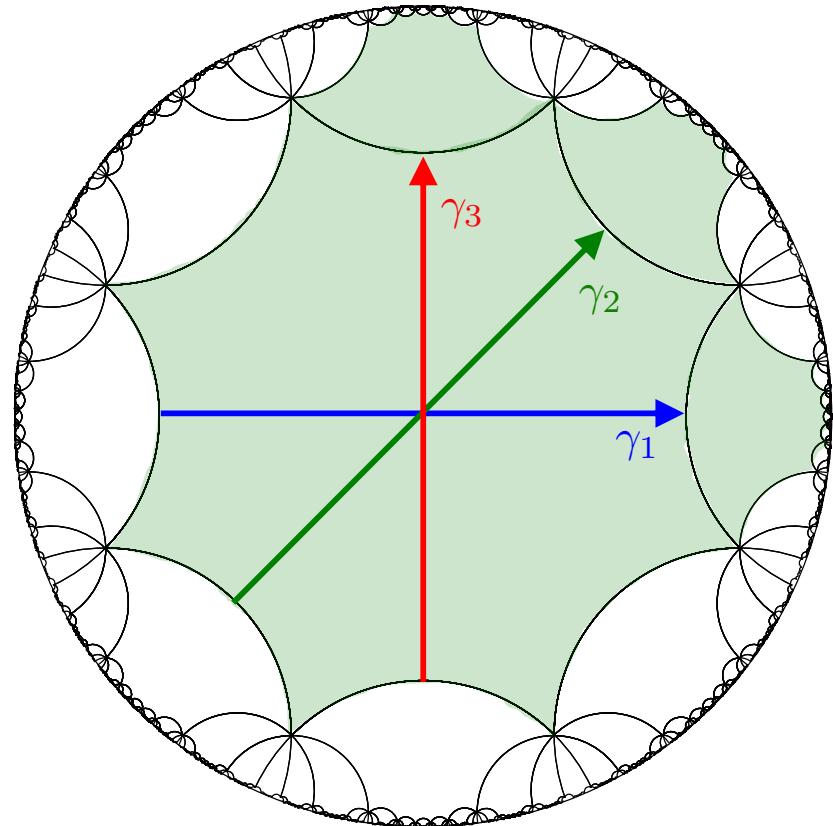
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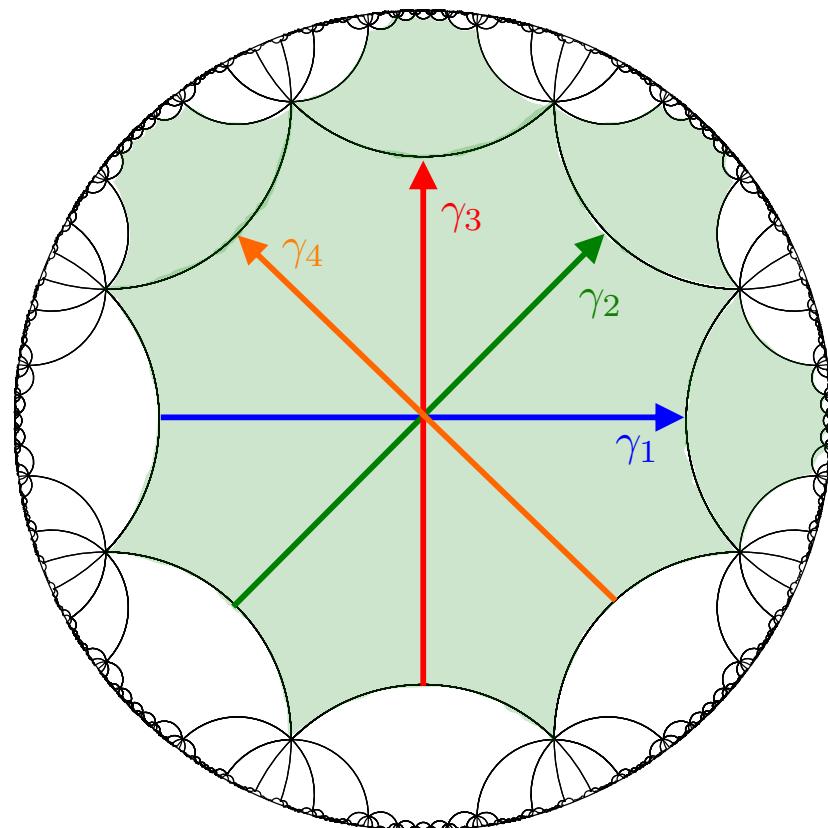
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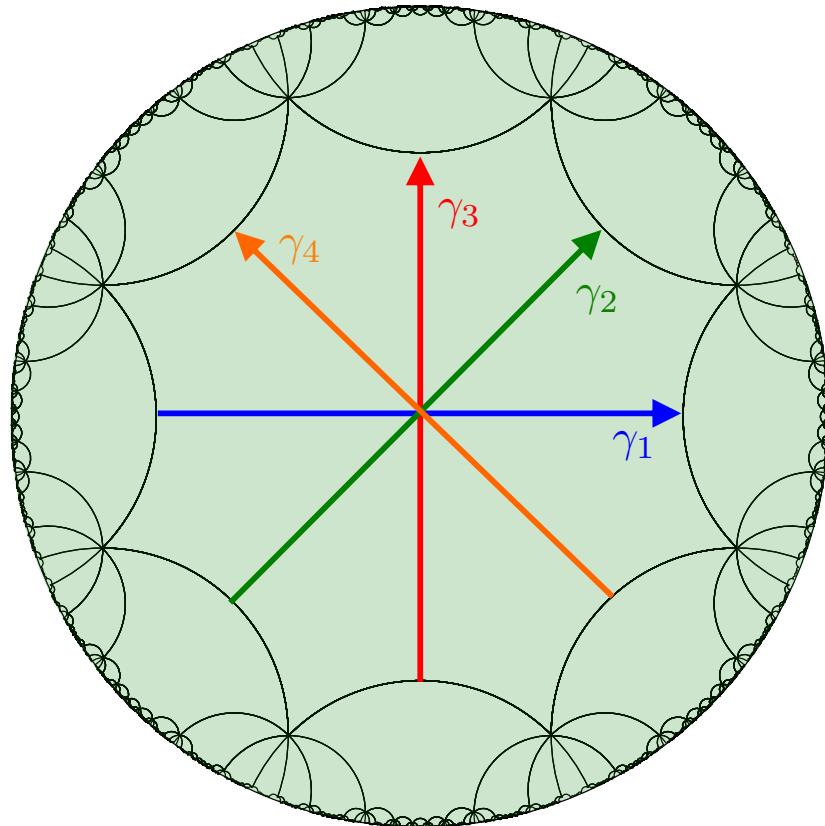
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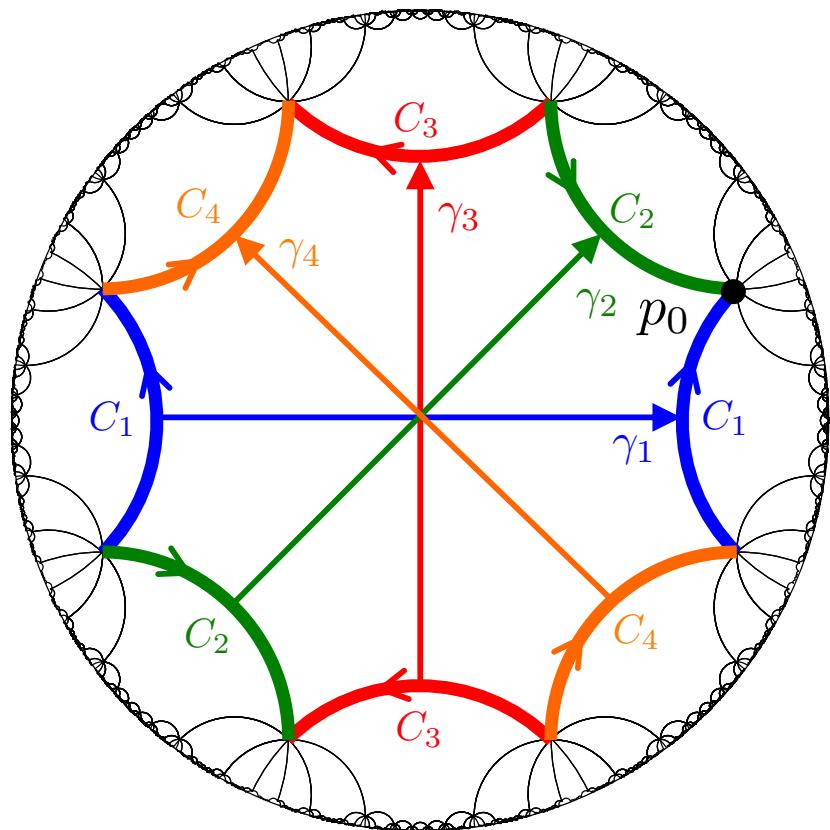
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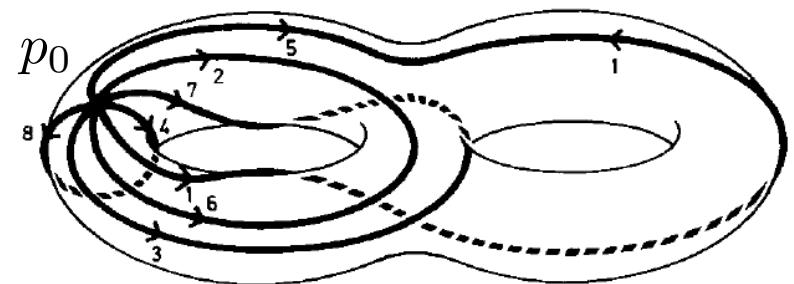
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# Compactified unit cell

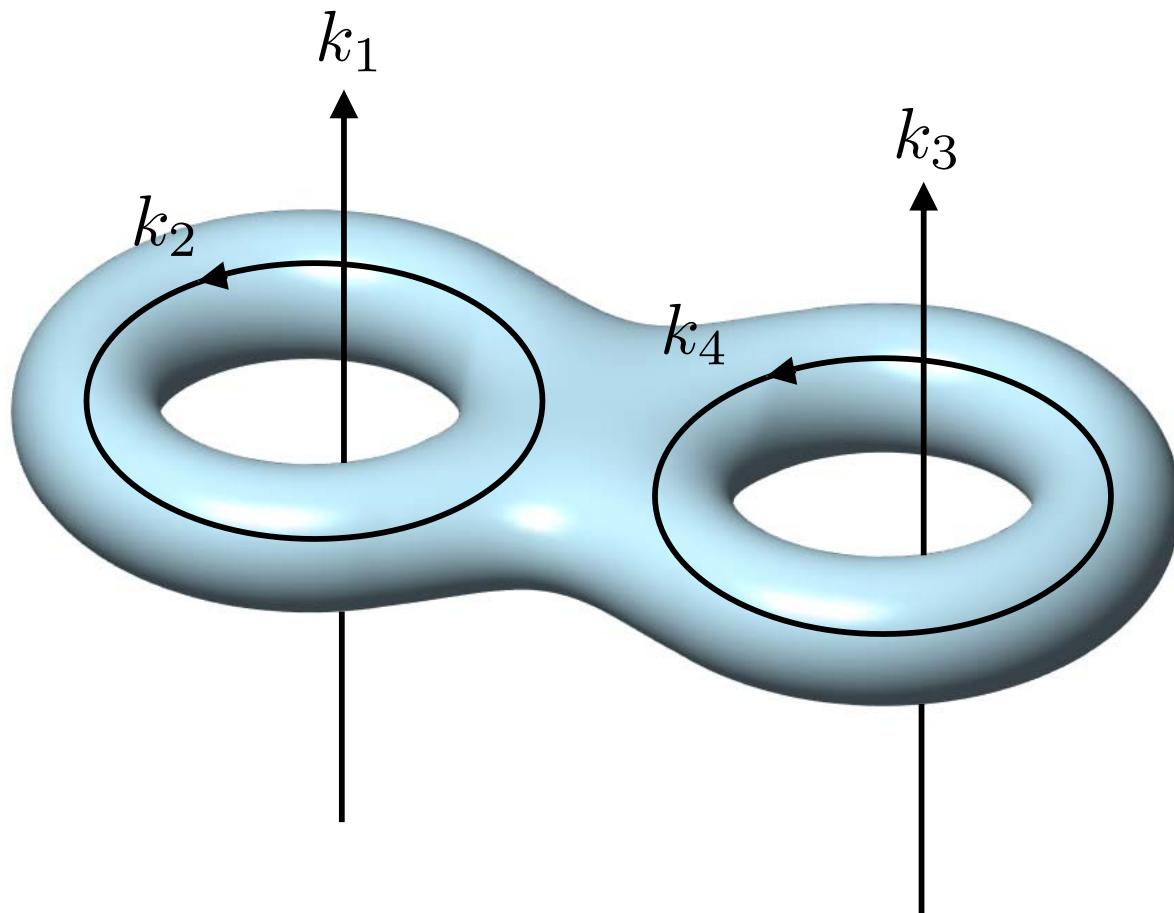


$$\mathbb{H}/\Gamma \cong \Sigma_2$$

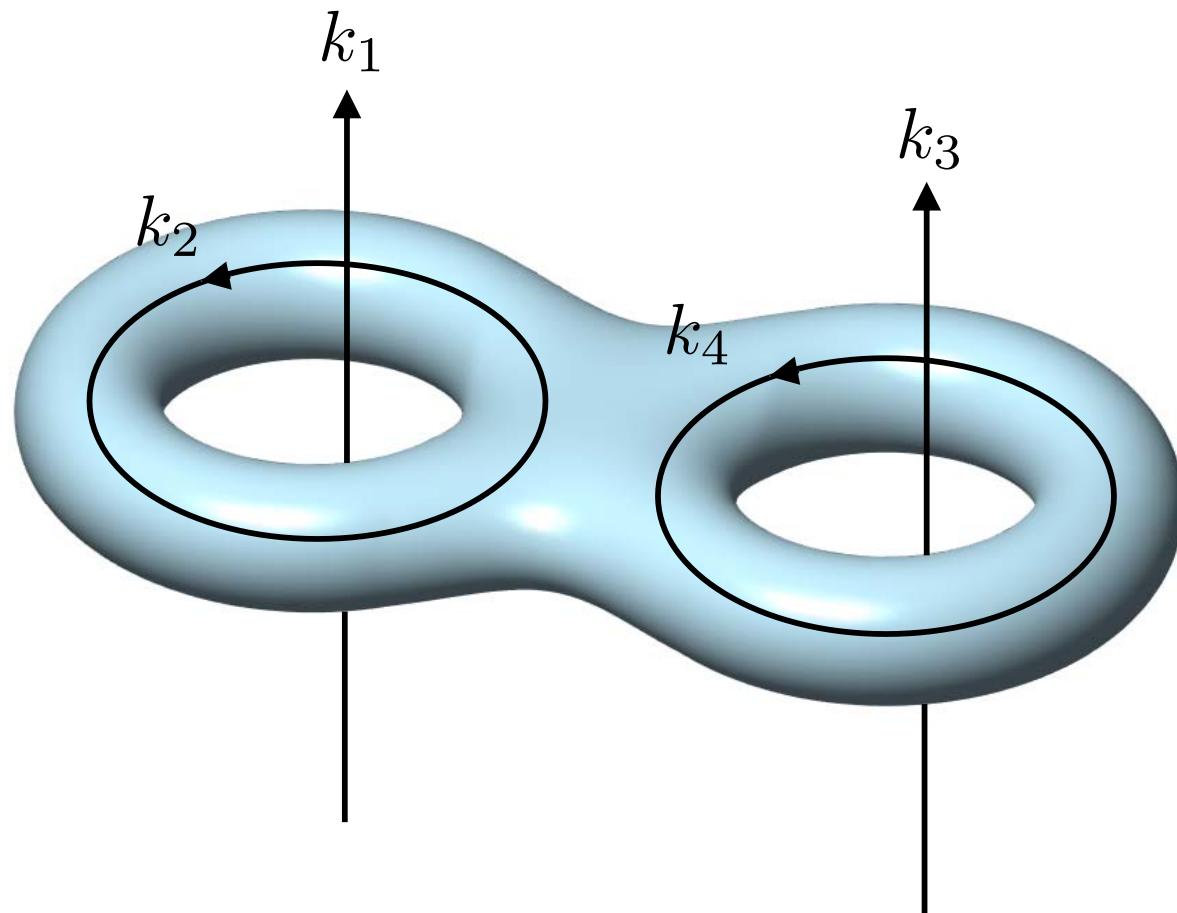


(Figure from Balazs & Voros, Phys. Rep. 143, 109 (1986))

# Aharonov-Bohm fluxes

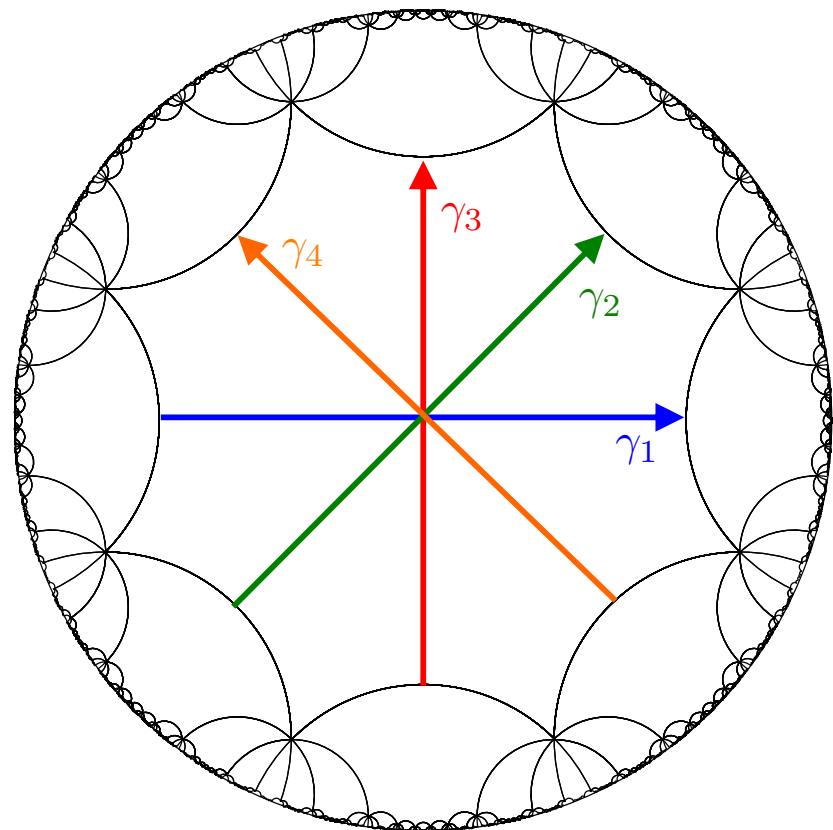


# Hyperbolic crystal momentum



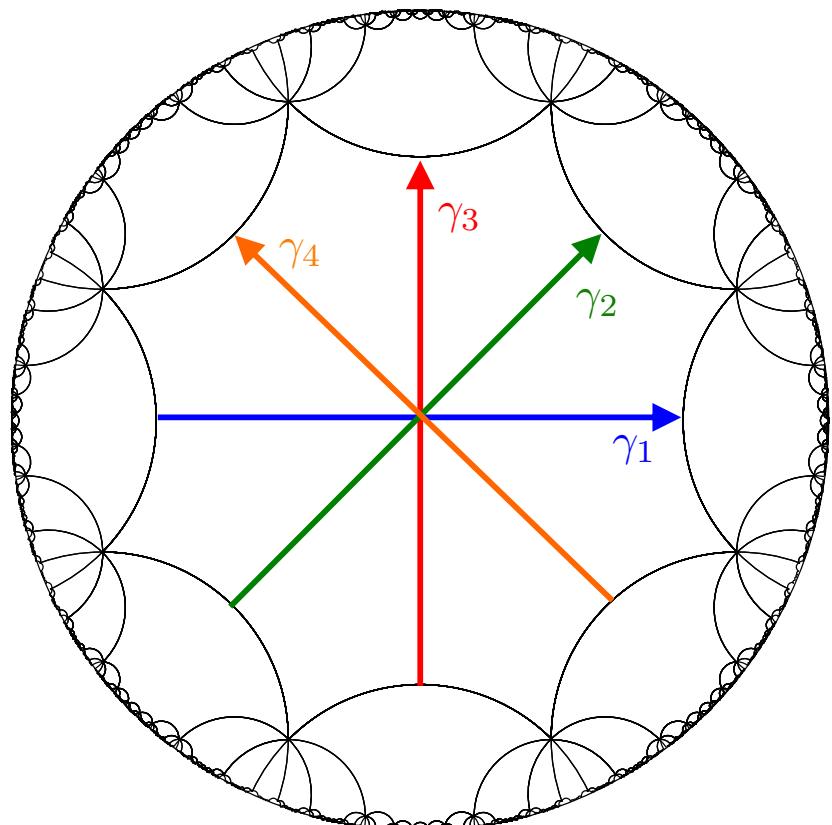
$$\mathbf{k} \equiv (k_1, k_2, \dots, k_{2g-1}, k_{2g}) \in (-\pi, \pi]^{2g} \cong T^{2g} \cong \text{Jac}(\Sigma_g)$$

# The hyperbolic Bloch problem



$$H = -\Delta + V(z)$$

# The hyperbolic Bloch problem



$$H = -\Delta + V(z)$$

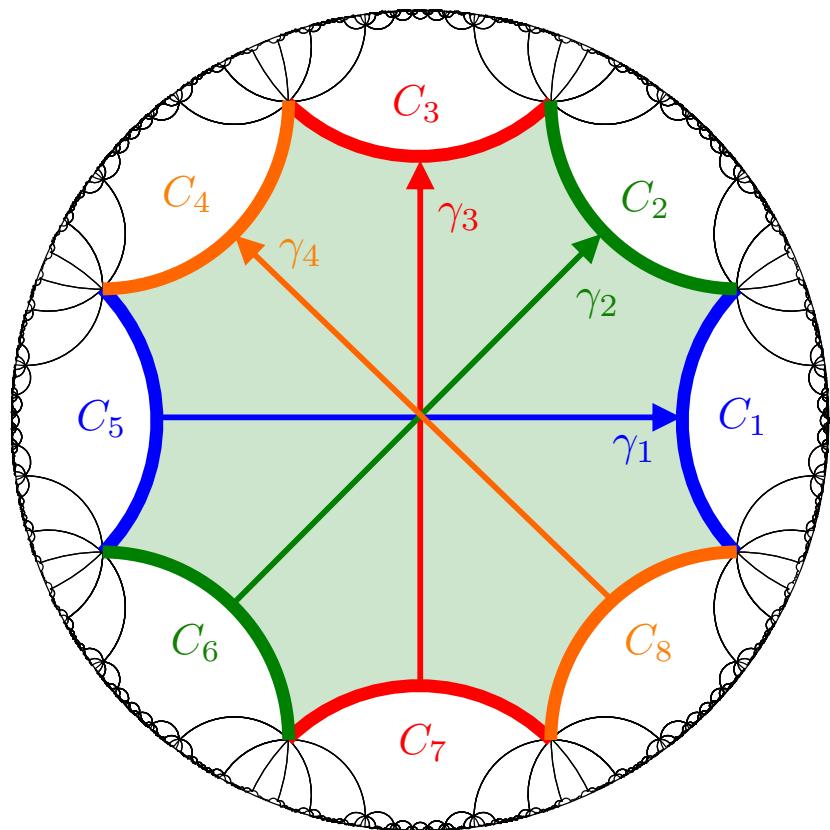
$$\Delta = \frac{1}{4}(1 - |z|^2)^2 \left( \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} \right)$$

Laplace-Beltrami operator

$$V(\gamma(z)) = V(z), \gamma \in \Gamma$$

“periodic” potential

# Automorphic Bloch condition



$$H\psi = E\psi$$

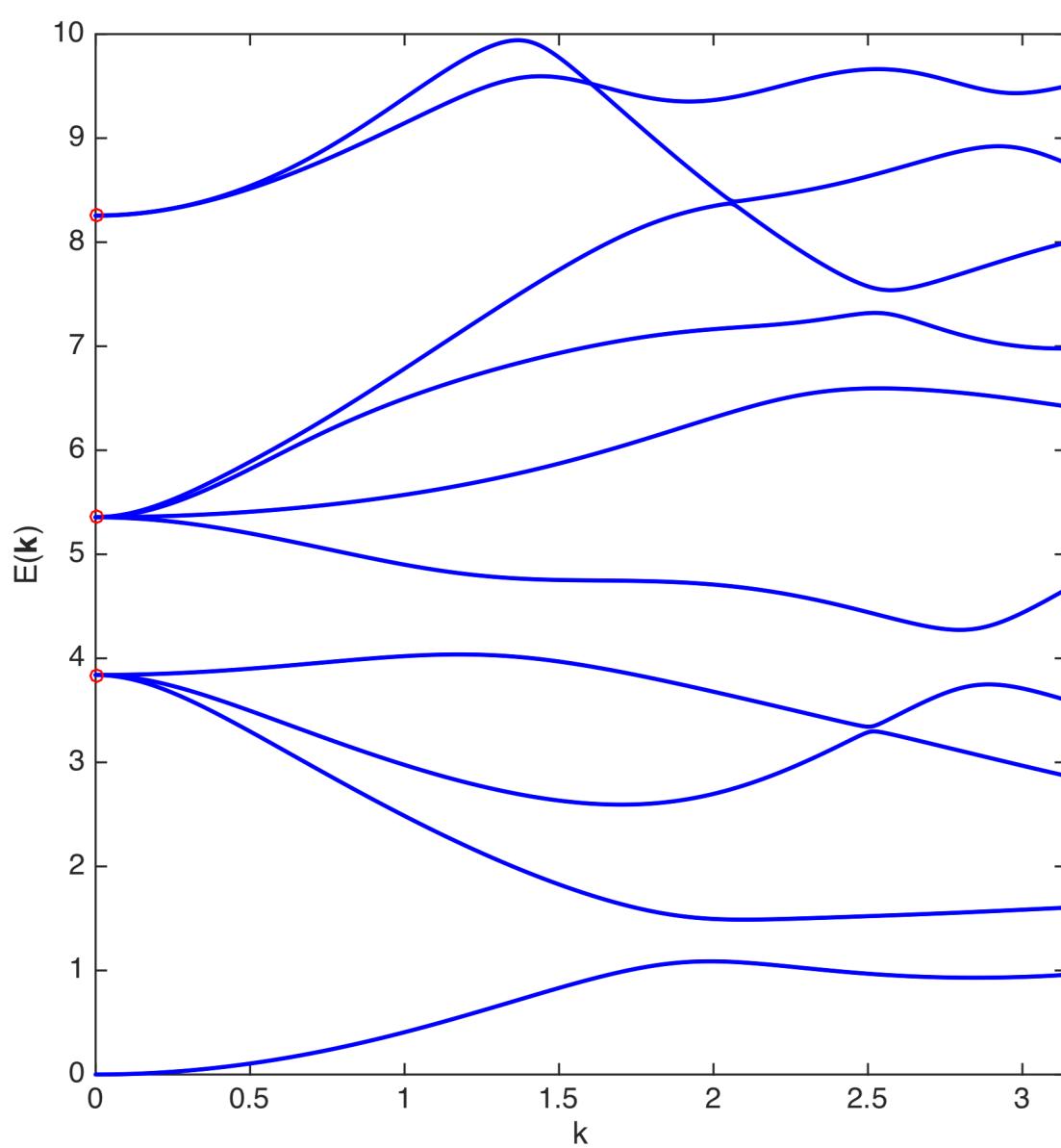
$$\gamma_j(C_{j+4}) = C_j,$$

$$\psi(C_j) = e^{ik_j} \psi(C_{j+4}),$$

$$j = 1, 2, 3, 4$$

$$\psi(\gamma(z)) = \chi(\gamma)\psi(z)$$

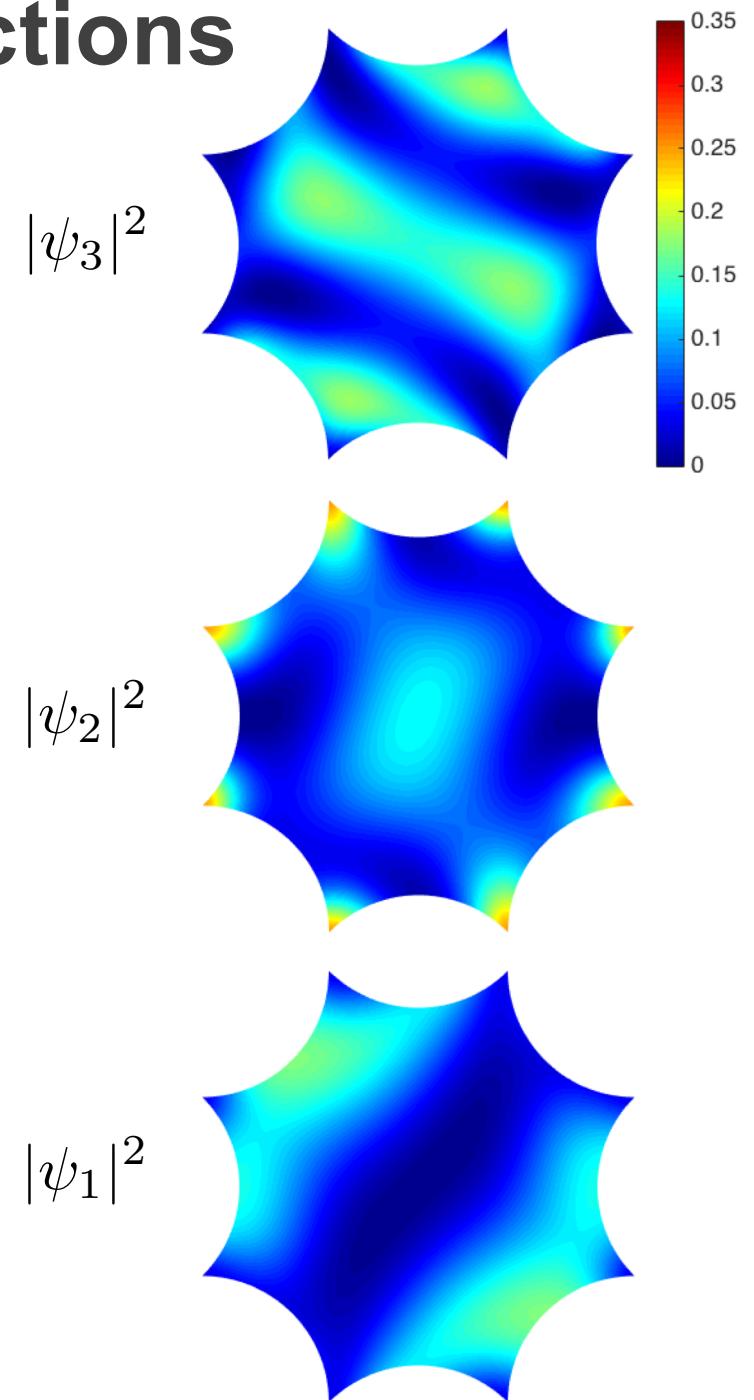
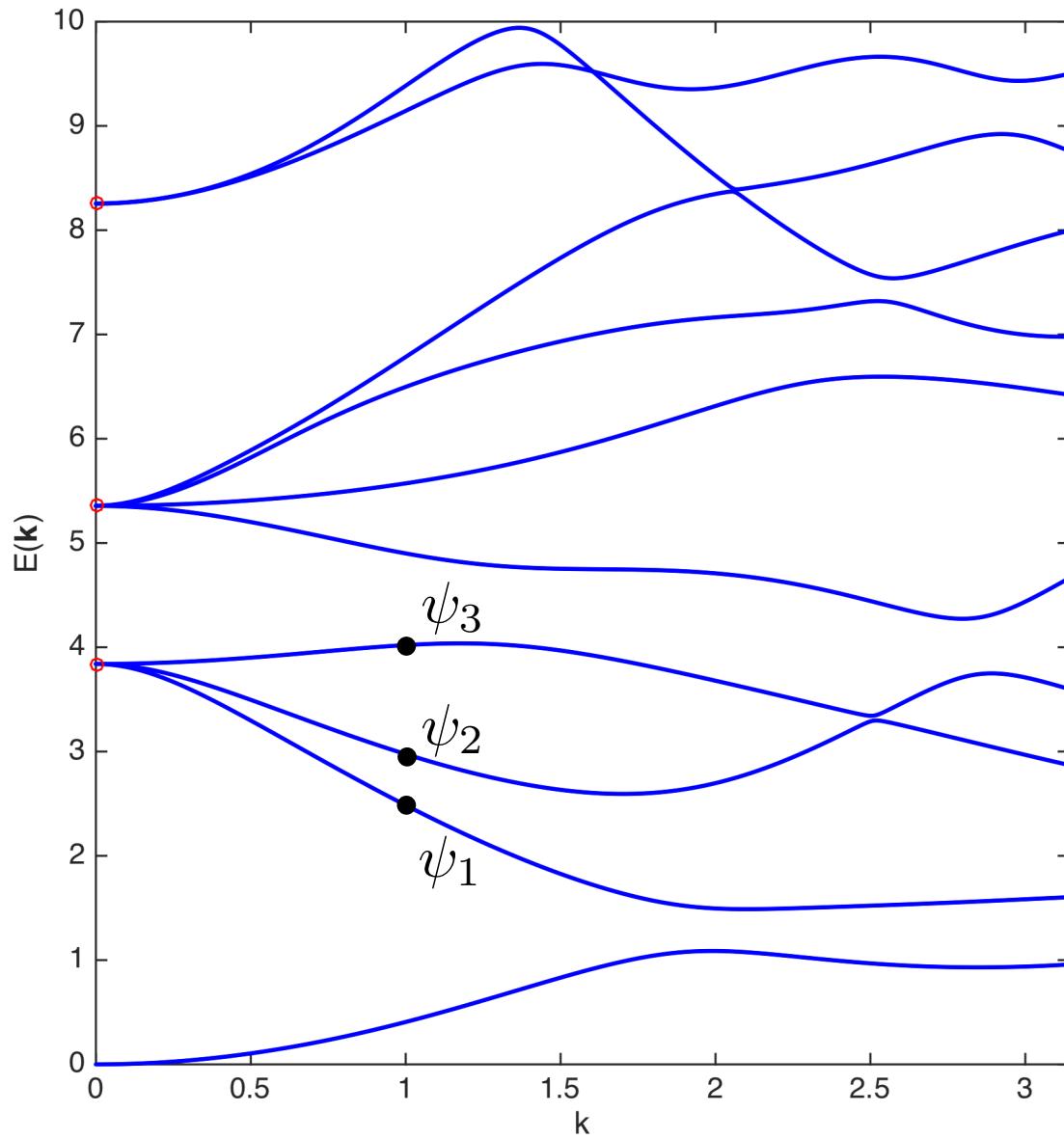
# Hyperbolic bandstructure: empty lattice



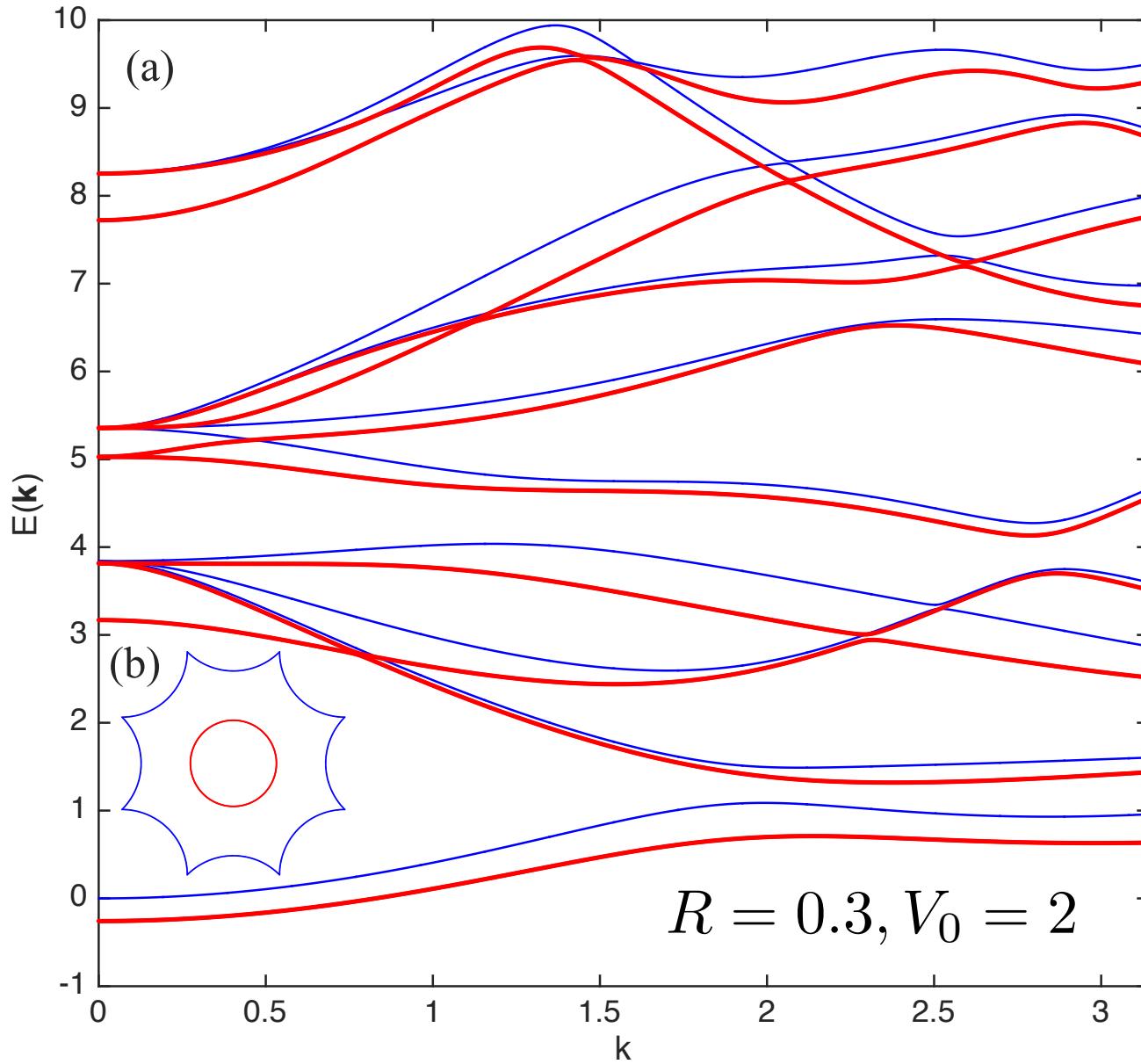
$$H = -\Delta$$

$$\mathbf{k} = (0.8, 0.3, 1.2, 1.7)k$$

# Hyperbolic Bloch wavefunctions



# Particle in an automorphic potential



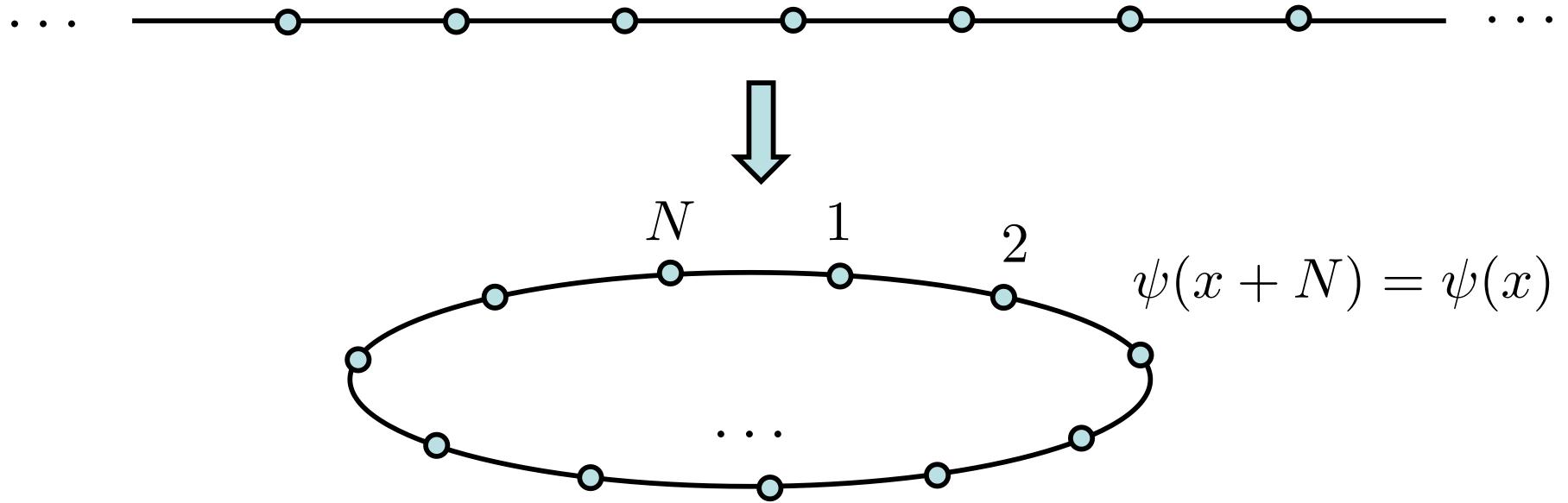
$$V(z) = \sum_{\gamma \in \Gamma} U(\gamma(z)),$$

$$U(z) = \begin{cases} -V_0, & |z| < R, \\ 0, & |z| \geq R. \end{cases}$$

# Two issues

- Do automorphic Bloch states form a complete set (ansatz vs theorem)?
- What about finite lattices (experiment)?

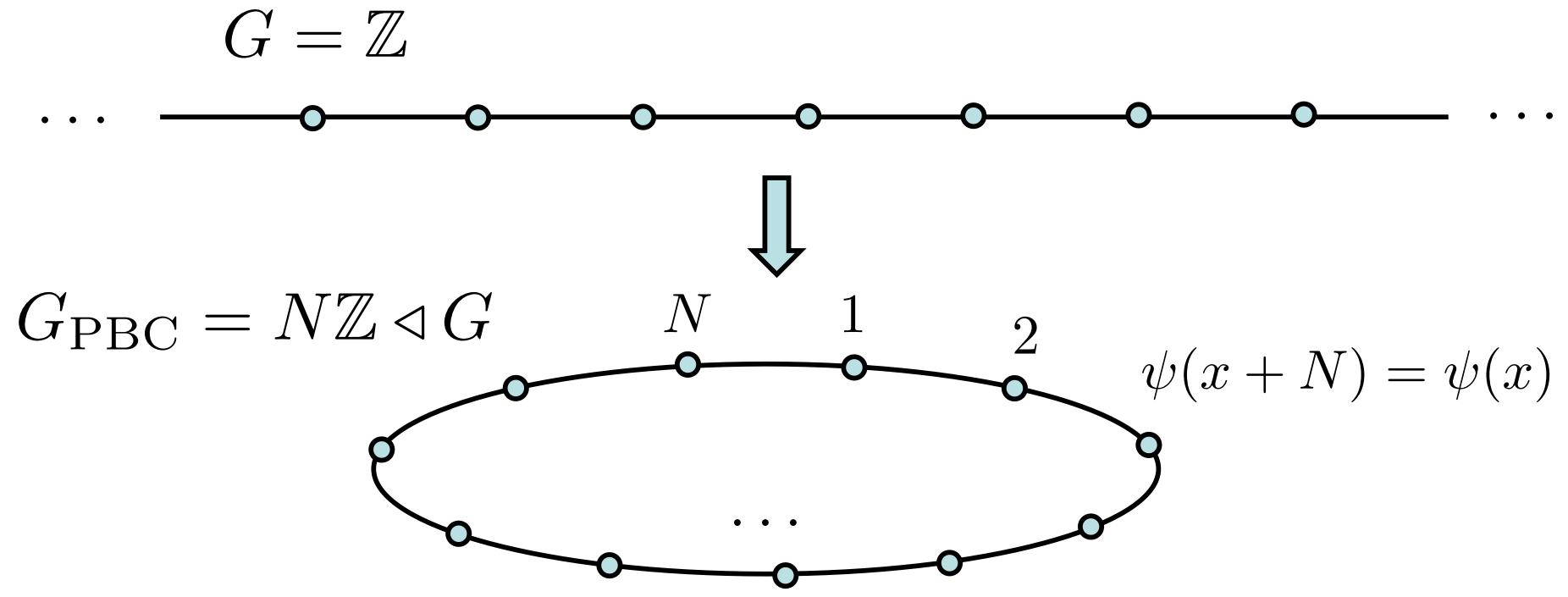
# Euclidean PBC



$$k = \frac{2\pi n}{N}, \quad n = 0, 1, \dots, N - 1$$

$N$  sites =  $N$  Bloch states: complete set

# Euclidean PBC: algebraic viewpoint

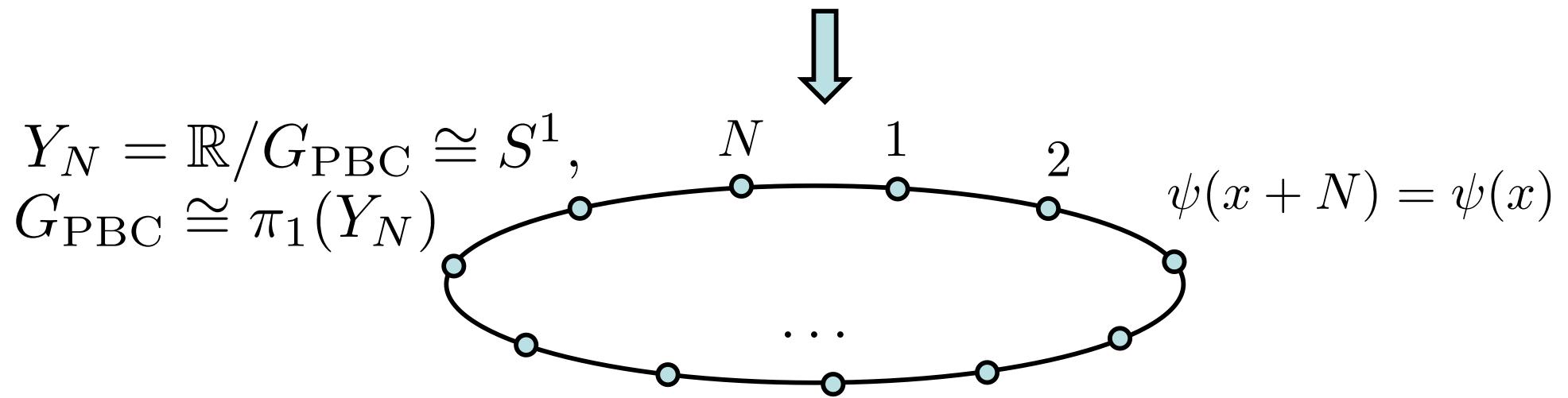


$$G/G_{\text{PBC}} = \mathbb{Z}_N$$

$N$  allowed  $k$  values =  $N$  unitary irreps of  $\mathbb{Z}_N$

# Euclidean PBC: topological viewpoint

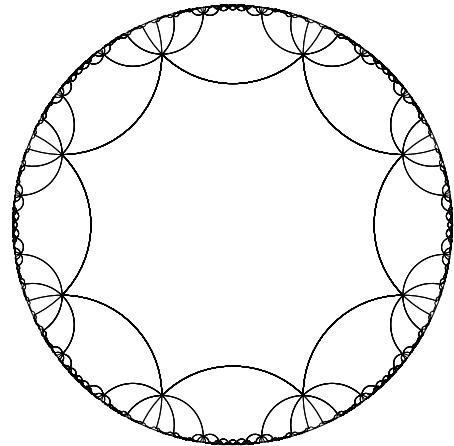
$$X = \mathbb{R}/G \cong S^1, G \cong \pi_1(X)$$



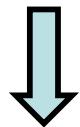
$$X \cong Y_N / \mathbb{Z}_N, \mathbb{Z}_N \cong \pi_1(X) / \pi_1(Y_N)$$

$Y_N$  =  $N$ -sheeted Galois (normal) cover of  $X$ ,  
with group of deck transformations  $Z_N$

# Hyperbolic PBC: algebraic viewpoint



$\Gamma$

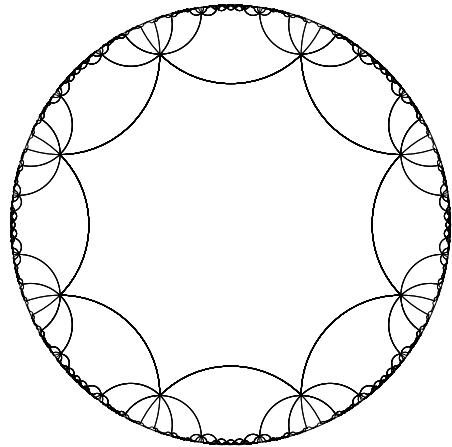


PBC cluster with  $N$  unit cells

$$\psi(\gamma_{\text{PBC}}(z)) = \psi(z)$$

$\Gamma_{\text{PBC}} \triangleleft \Gamma$  : normal subgroup of  
index  $N$

# Hyperbolic PBC: topological viewpoint



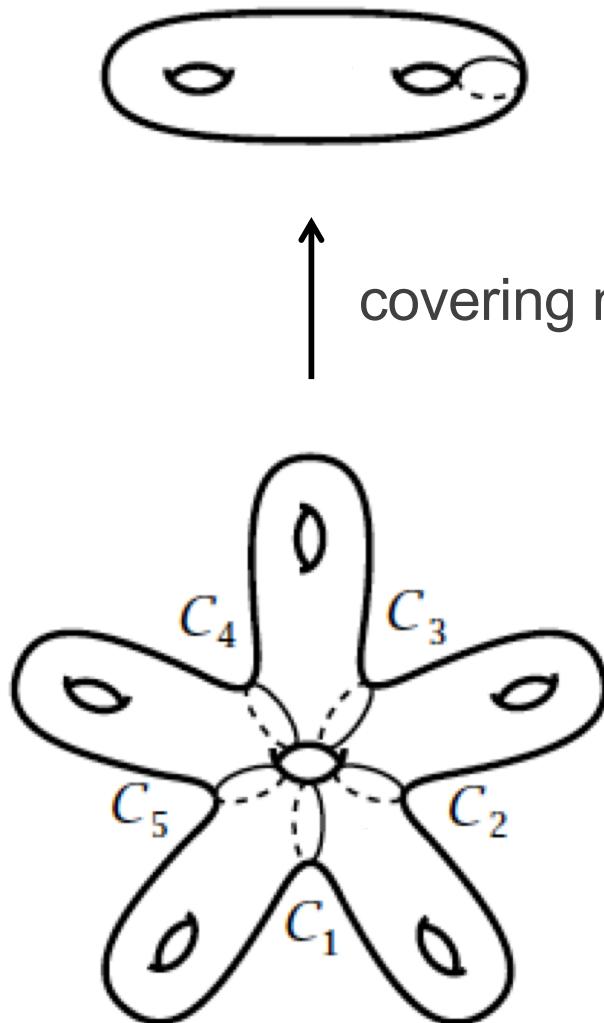
$$X = \mathbb{H}/\Gamma \cong \Sigma_g, \quad \Gamma \cong \pi_1(X)$$



PBC cluster with  $N$  unit cells:  
 $N$ -sheeted Galois cover of  $X$ ,  
with deck group  $\Gamma/\Gamma_{\text{PBC}}$

$$Y_N = \mathbb{H}/\Gamma_{\text{PBC}}, \quad \Gamma_{\text{PBC}} = \pi_1(Y_N)$$

# Hyperbolic PBC: topological viewpoint



(e.g.  $N=5$ )

$$X \cong \Sigma_g$$

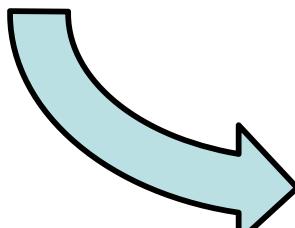
$$Y_N \cong \Sigma_h$$

$$h = N(g - 1) + 1$$

Riemann-Hurwitz formula

# PBC vs the Bloch ansatz

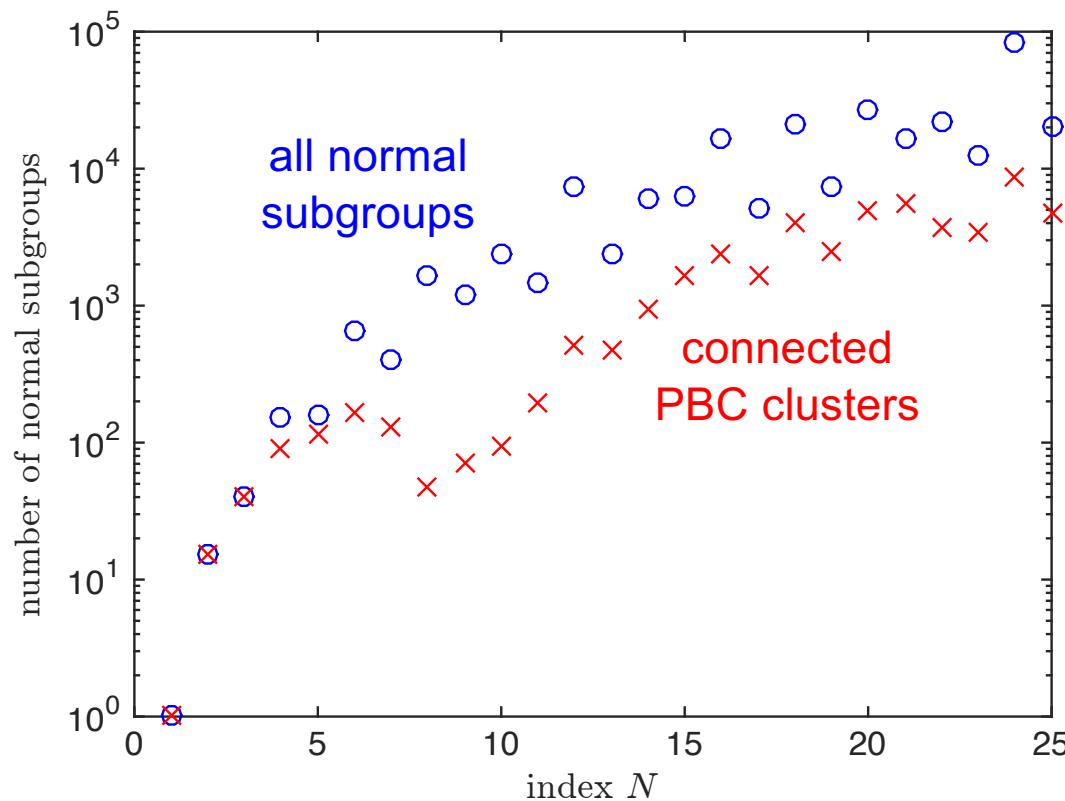
Assume  $\psi(\gamma(z)) = \chi(\gamma)\psi(z)$  :

$$\begin{aligned}\psi(\gamma\gamma_{\text{PBC}}\gamma^{-1}(z)) &= \chi(\gamma\gamma_{\text{PBC}}\gamma^{-1})\psi(z) \\ &= \chi(\gamma)\chi(\gamma_{\text{PBC}})\chi^{-1}(\gamma)\psi(z) \\ &= \psi(z) \text{ since } \chi(\gamma_{\text{PBC}}) = 1\end{aligned}$$


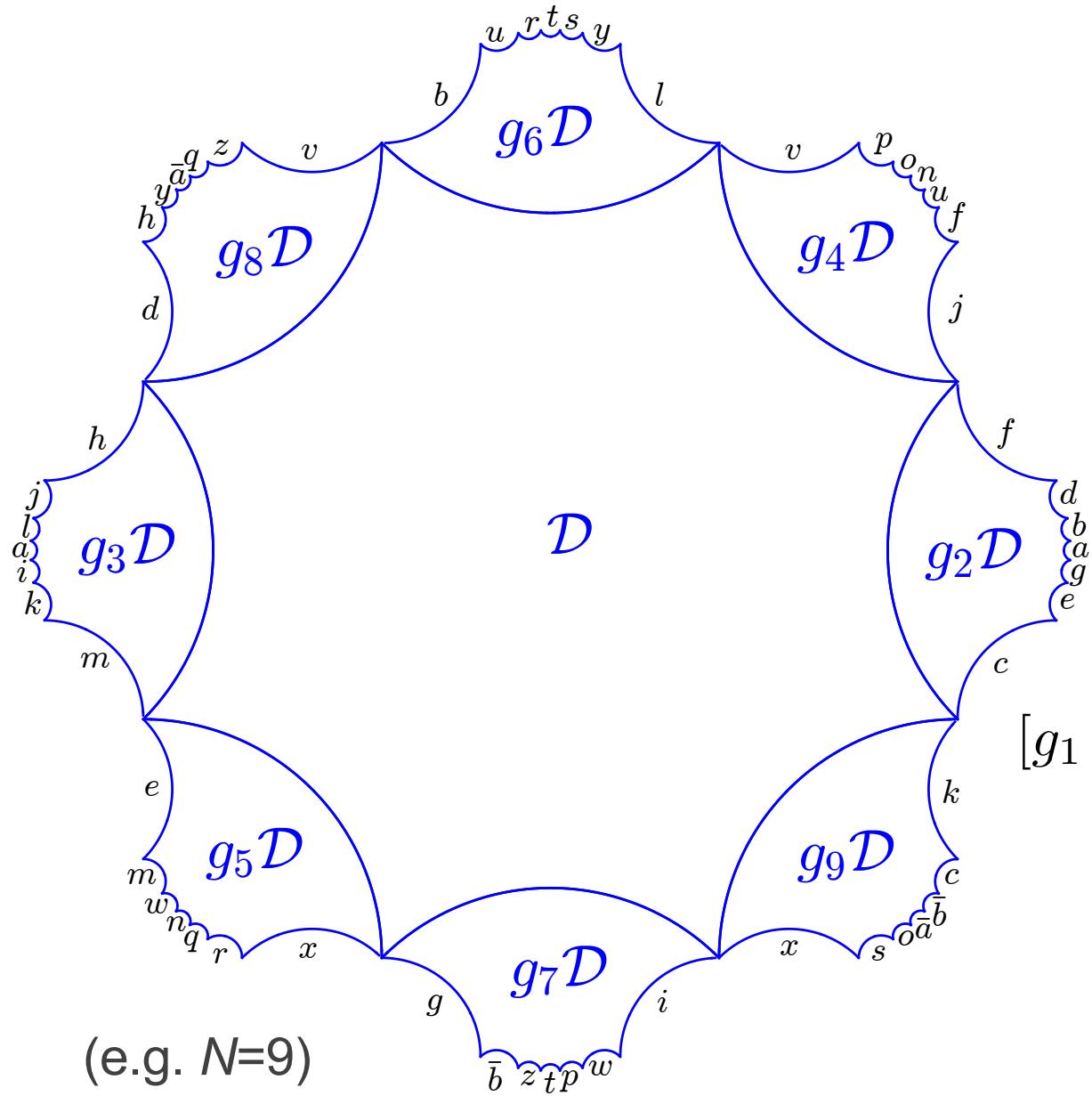
thus  $\gamma\gamma_{\text{PBC}}\gamma^{-1} \in \Gamma_{\text{PBC}}$    $\Gamma_{\text{PBC}} \triangleleft \Gamma$

# Low-index normal subgroups

- For a given  $N$ , many distinct subgroups  $\Gamma_{\text{PBC}} \triangleleft \Gamma$ , although all isomorphic to  $\pi_1(\Sigma_h)$
- Compute all normal subgroups of index up to  $N = 25$  using Firth-Holt algorithm (GAP implementation by F. Rober)

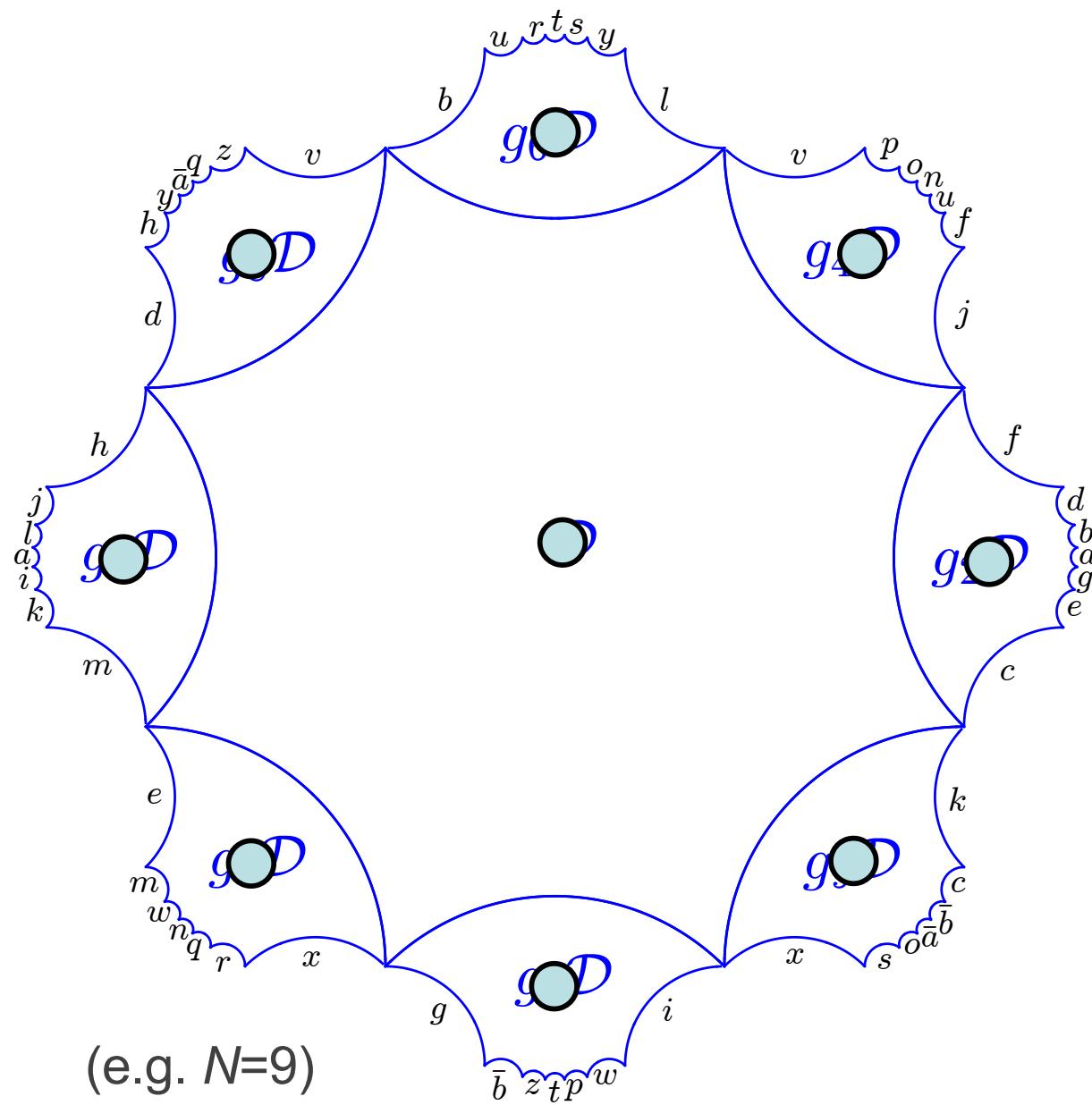


# Connected PBC clusters



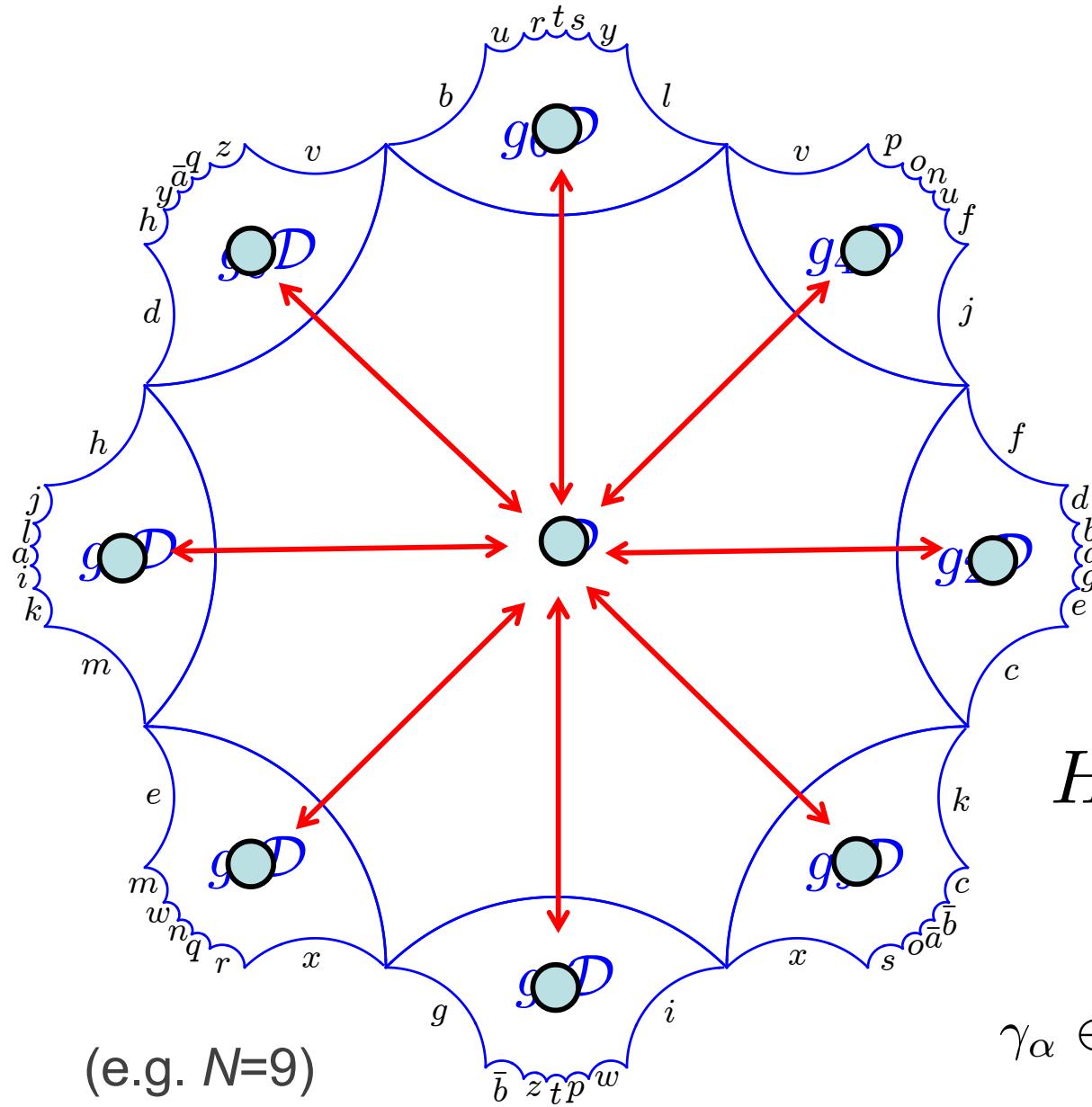
$$\mathcal{C} = \bigsqcup_{i=1}^N g_i \mathcal{D}$$
$$[g_1 = e], [g_2], \dots, [g_N] \in \Gamma / \Gamma_{\text{PBC}}$$

# Hopping matrix



$$H_{ij} = -1 \quad \text{if } i \& j \text{ n.n., 0 otherwise}$$

# Hopping matrix



$$H_{ij} = -1$$

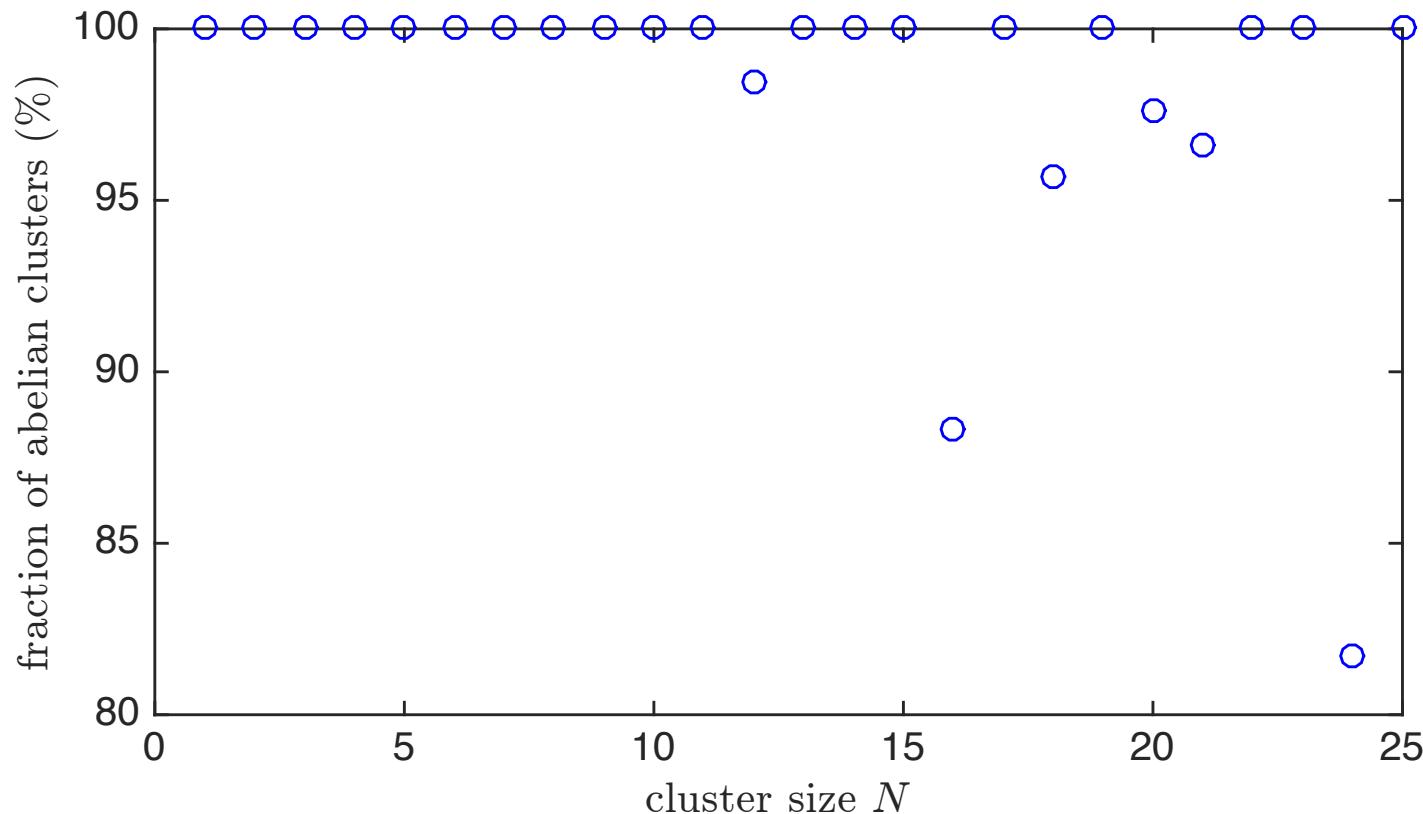
if  $i$  &  $j$  n.n., 0 otherwise

$$H_{ij} = - \sum_{\alpha} \delta_{[g_j], [g_i \gamma_{\alpha}]}$$

$$\gamma_{\alpha} \in \{\gamma_1, \dots, \gamma_4, \gamma_1^{-1}, \dots, \gamma_4^{-1}\}$$

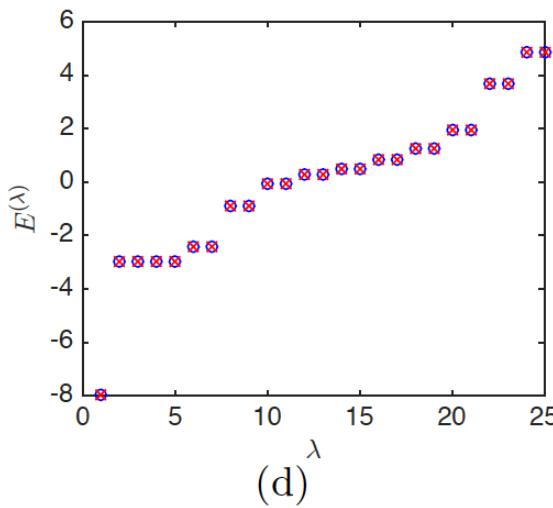
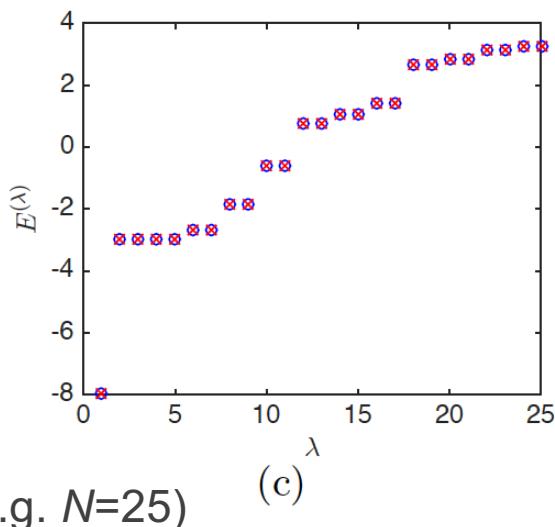
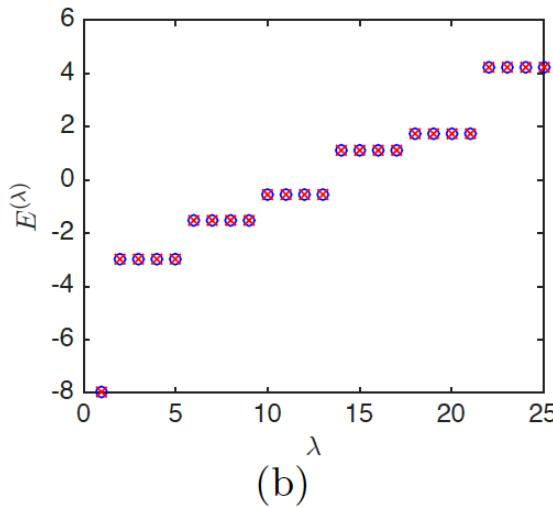
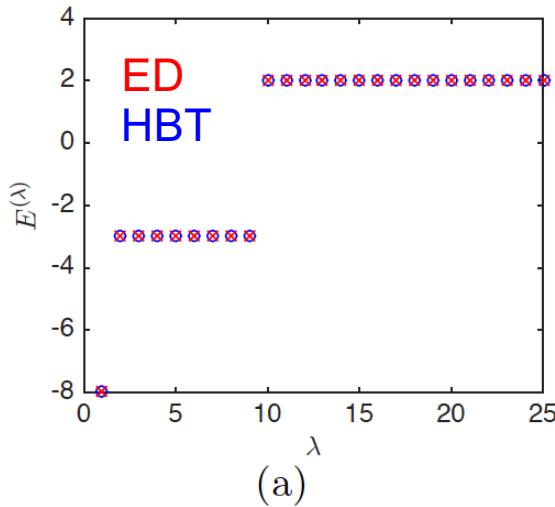
# Abelian clusters

- Eigenstates of the clusters fall into irreps of the residual translation group  $\Gamma/\Gamma_{\text{PBC}} = \text{finite group of order } N$
- For large fraction of clusters, this group is abelian!



# Abelian Bloch theorem

$$\psi^{(\lambda)}(g_k^{-1}(z_i)) = \chi^{(\lambda)}([g_k])\psi^{(\lambda)}(z_i), \quad [g_k] \in \Gamma/\Gamma_{\text{PBC}}$$



$$E^{(\lambda)} = -2 \sum_{j=1}^4 \cos k_j^{(\lambda)},$$

$$\lambda = 1, \dots, N$$

$$k_j^{(\lambda)} \in \frac{2\pi\mathbb{Z}}{N}$$



$\text{Jac}(\Sigma_g) \cong \mathbb{T}^{2g}$  is discretized

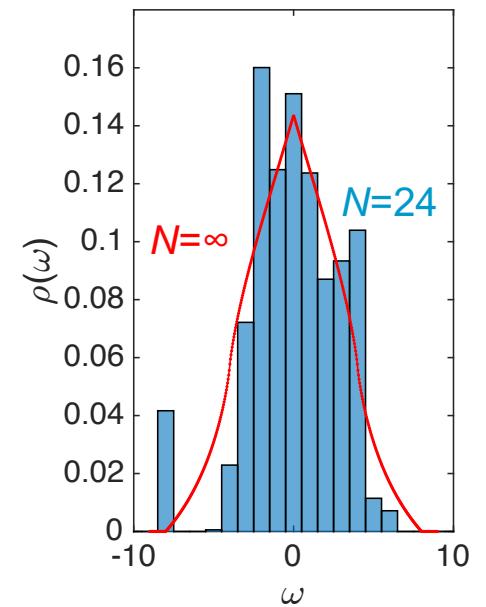
# Maximal abelian cluster

- Finite abelian clusters are subsets of an *infinite* abelian cluster  
 $Y_\infty = \mathbb{H}/\Gamma_{\text{PBC}}$  with  $\Gamma_{\text{PBC}} = [\Gamma, \Gamma]$ :

$$\Gamma/\Gamma_{\text{PBC}} \cong H_1(X, \mathbb{Z}) \cong \mathbb{Z}^{2g}$$

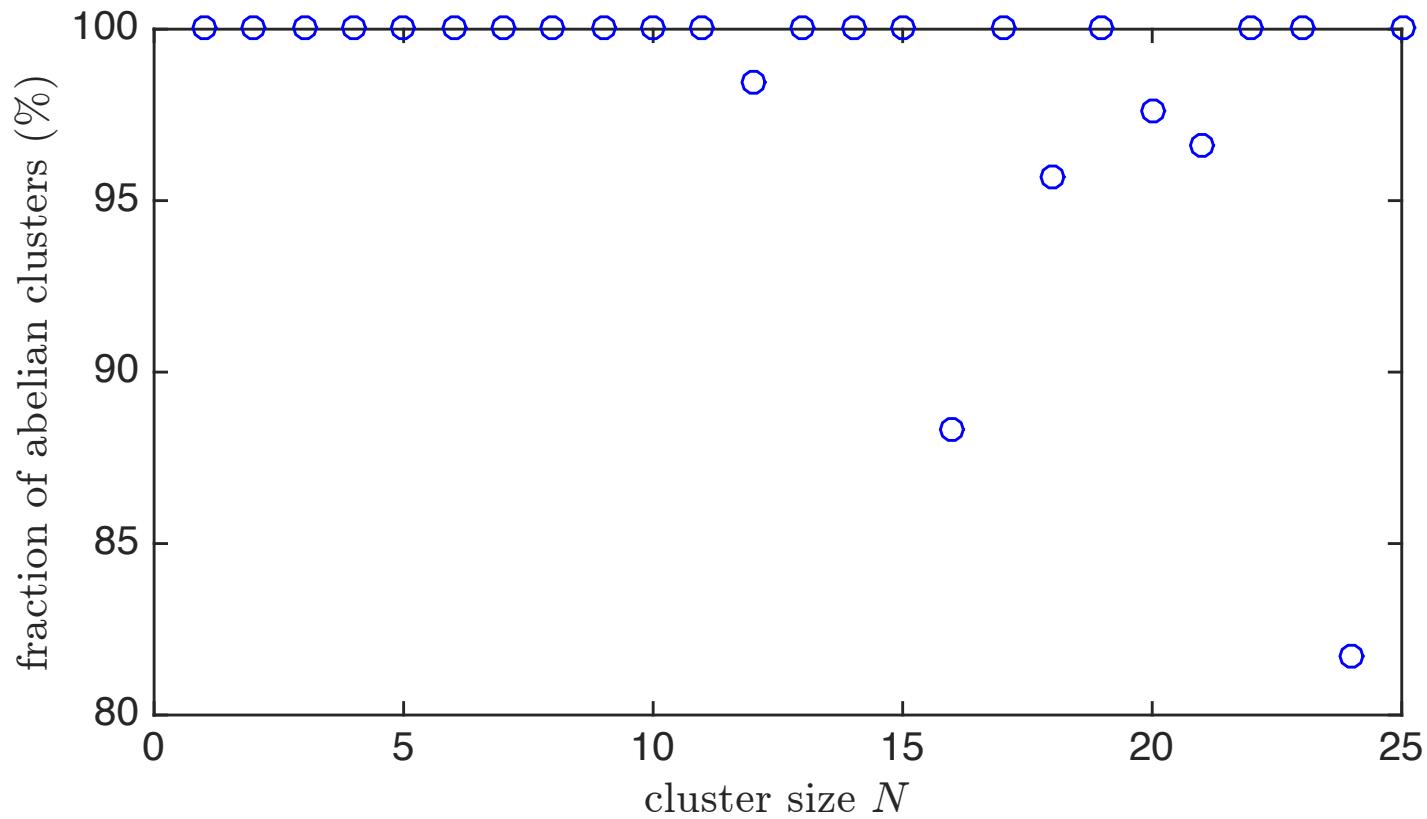


Euclidean lattice in  $2g$  dimensions:  
 $\text{BZ} = \text{Jac}(\Sigma_g) \cong \mathbb{T}^{2g}$



# Nonabelian clusters

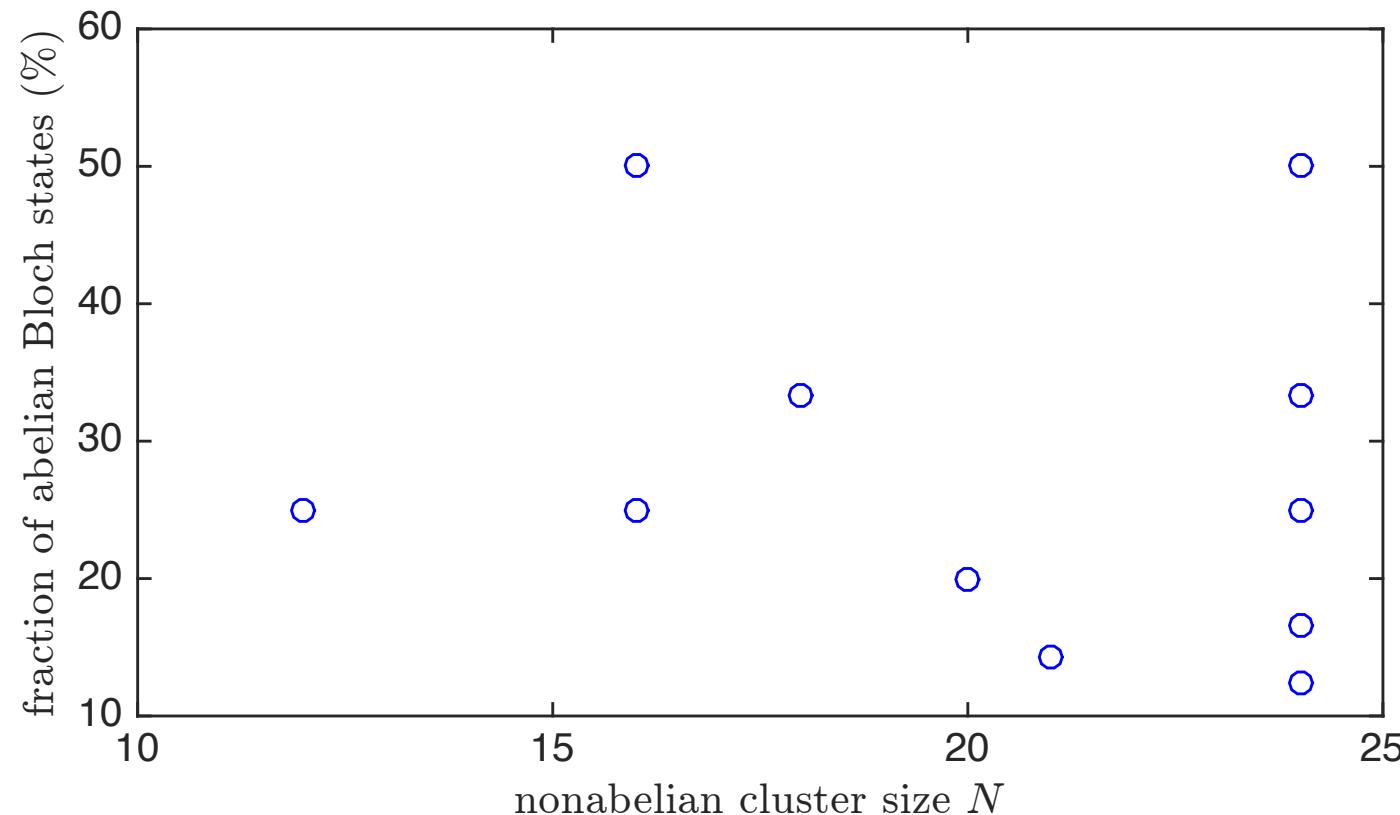
- For  $N < 25$ , nonabelian  $\Gamma/\Gamma_{\text{PBC}}$  found only at  $N = 12, 16, 18, 20, 21, 24$



# Nonabelian Bloch theorem

- Nonabelian  $\Gamma/\Gamma_{\text{PBC}}$  possesses higher-dimensional unitary irreps:

$$\psi_\nu^{(\lambda)}(g_k^{-1}(z_i)) = \sum_{\mu=1}^{r_\lambda} \psi_\mu^{(\lambda)}(z_i) D_{\mu\nu}^{(\lambda)}([g_k]), \quad [g_k] \in \Gamma/\Gamma_{\text{PBC}}$$



# Nonabelian Bloch theorem

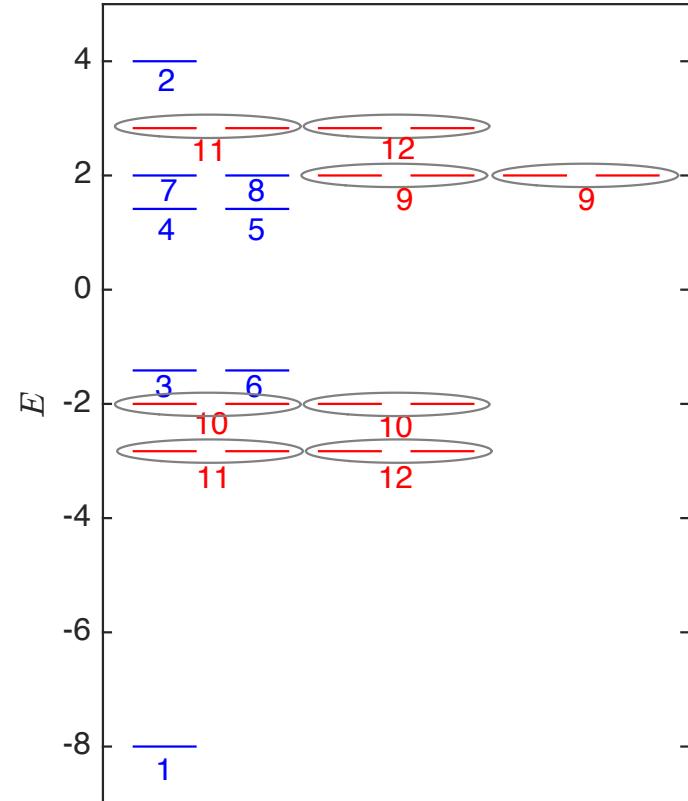
- Translation matrices belong to (reducible) regular representation: irrep of dimension  $r_\lambda$  appears  $r_\lambda$  times
- Example ( $N=24$ ): 8 abelian irreps, 4 nonabelian (2D) irreps

$C$	1	2	3	4	5	6	7	8	9	10	11	12
$n_C$	1	1	1	1	2	2	2	2	3	3	3	3
$D^{(1)}$	1	1	1	1	1	1	1	1	1	1	1	1
$D^{(2)}$	1	1	1	1	1	1	1	1	-1	-1	-1	-1
$D^{(3)}$	1	-1	$a$	$-a$	1	-1	$a$	$-a$	$c$	$-c$	$-1/c$	$1/c$
$D^{(4)}$	1	-1	$a$	$-a$	1	-1	$a$	$-a$	$-c$	$c$	$1/c$	$-1/c$
$D^{(5)}$	1	-1	$-a$	$a$	1	-1	$-a$	$a$	$-1/c$	$1/c$	$c$	$-c$
$D^{(6)}$	1	-1	$-a$	$a$	1	-1	$-a$	$a$	$1/c$	$-1/c$	$-c$	$c$
$D^{(7)}$	1	1	-1	-1	1	1	-1	-1	$a$	$a$	$-a$	$-a$
$D^{(8)}$	1	1	-1	-1	1	1	-1	-1	$-a$	$-a$	$a$	$a$
$D^{(9)}$	2	2	-2	-2	-1	-1	1	1	0	0	0	0
$D^{(10)}$	2	2	2	2	-1	-1	-1	-1	0	0	0	0
$D^{(11)}$	2	-2	$b$	$-b$	-1	1	$-a$	$a$	0	0	0	0
$D^{(12)}$	2	-2	$-b$	$b$	-1	1	$a$	$-a$	0	0	0	0

# Nonabelian Bloch theorem

- Translation matrices belong to (reducible) regular representation: irrep of dimension  $r_\lambda$  appears  $r_\lambda$  times
- Example ( $N=24$ ): 8 abelian irreps, 4 nonabelian (2D) irreps

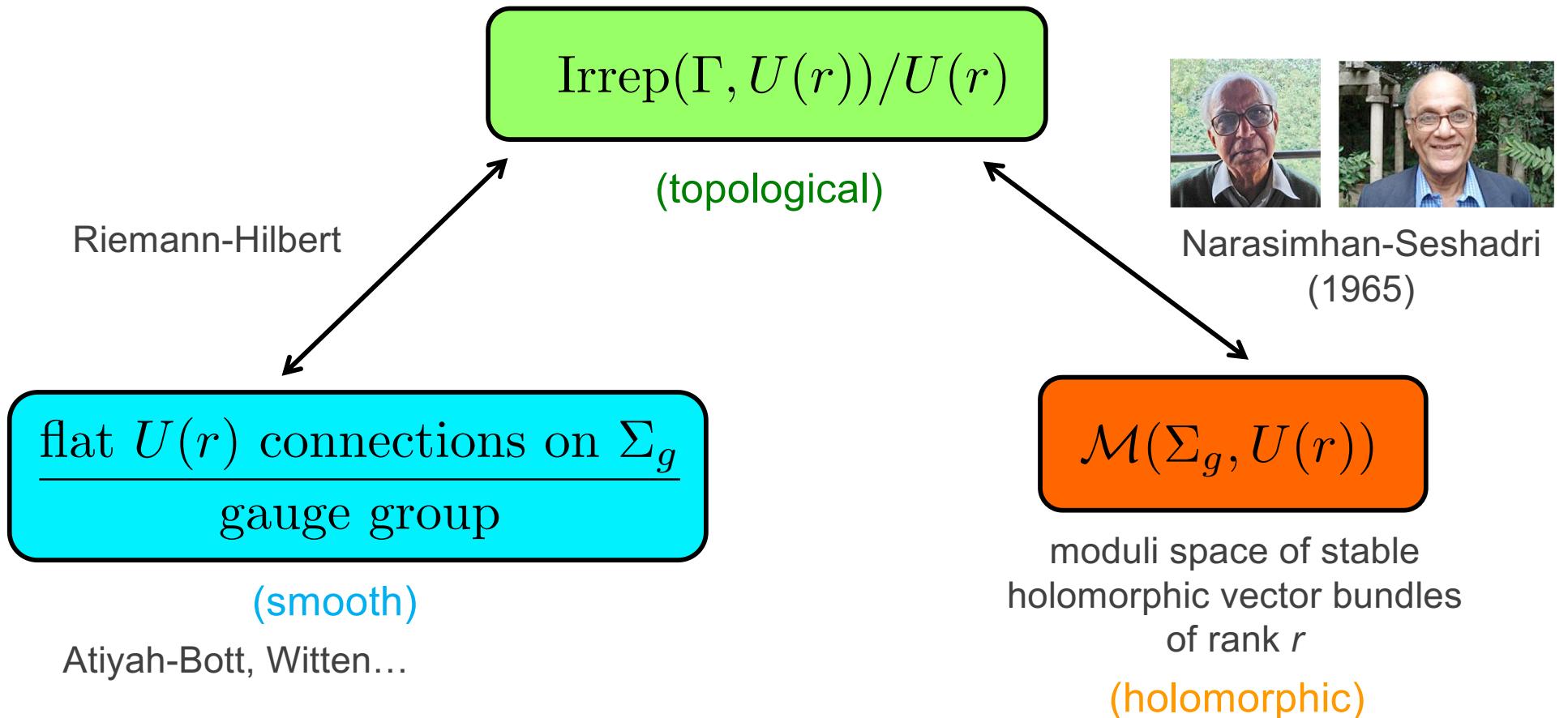
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$D^{(2)}$	1	1	1	1	1	1	1	1	-1	-1	-1	-1
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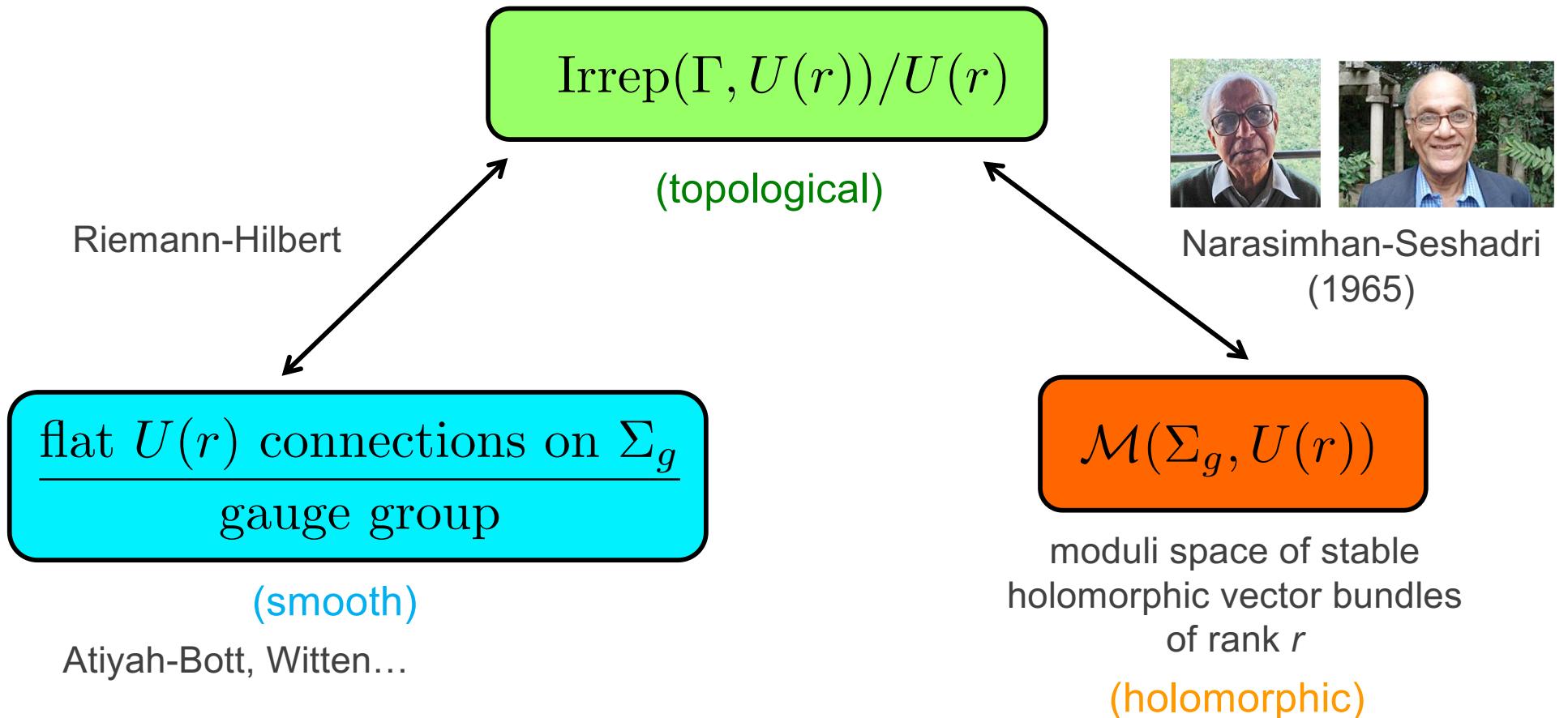
# Nonabelian Bloch theorem

- Translation matrices belong to (reducible) regular representation: irrep of dimension  $r_\lambda$  appears  $r_\lambda$  times
- Choice of  $\Gamma_{\text{PBC}}$  selects  $U(r)$  irreps of  $\Gamma$  whose kernel includes  $\Gamma_{\text{PBC}}$
- For fixed  $r$ , what is the space of all  $U(r)$  irreps of  $\Gamma$ ?

# Nonabelian Hodge theory



# Nonabelian Hodge theory



- Cmplx manifold:  $\dim_{\mathbb{C}} \mathcal{M}(\Sigma_g, U(r)) = r^2(g - 1) + 1$
- For  $g = 2, r = 2$ ,  $\mathcal{M}$  = bundle of  $\mathbb{CP}^3$  fibers over  $\text{Jac}(\Sigma_2)$

# Summary

- $\{4g, 4g\}$  hyperbolic lattices have a (nonabelian) discrete “translation” group  $\Gamma \cong \pi_1(\Sigma_g)$ , the Fuchsian group
- Finite (but arbitrarily large) clusters with PBC correspond to  $\Gamma_{\text{PBC}} \triangleleft \Gamma$
- Hyperbolic “Brillouin zones”:  $\text{Jac}(\Sigma_g) \cong T^{2g} \cong \mathcal{M}(\Sigma_g, U(1))$   
 $\mathcal{M}(\Sigma_g, U(r)), r > 1$
- Automorphic Bloch theorems (abelian/nonabelian)
- Point-group symmetries:  $\text{Aut}(\Sigma_g)$
- Implement PBC clusters in experiment (e.g. topoelectrical circuits)?

**JM and S. Rayan, Sci. Adv. 7, eabe9170 (2021); arXiv:2108.09314**