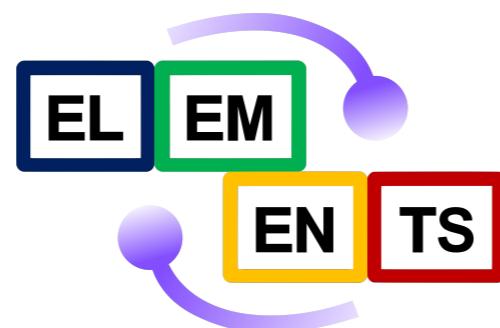


The QCD chiral phase transition for different numbers of quark flavours

Owe Philipsen

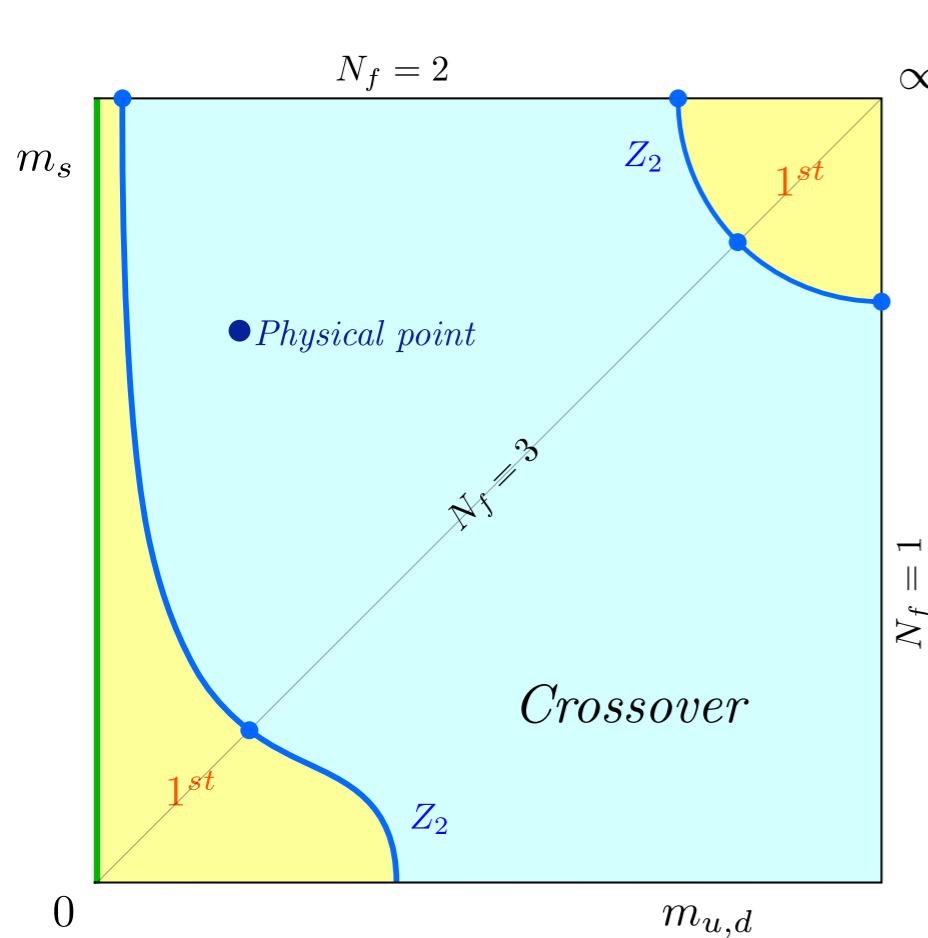
with Francesca Cuteri and Alessandro Sciarra

- Chiral phase transition in massless limit constrains the QCD phase diagram
- 40 years of common wisdom and inconclusive lattice results
- Resolution by a new approach + data [Cuteri, O.P., Sciarra JHEP 21]

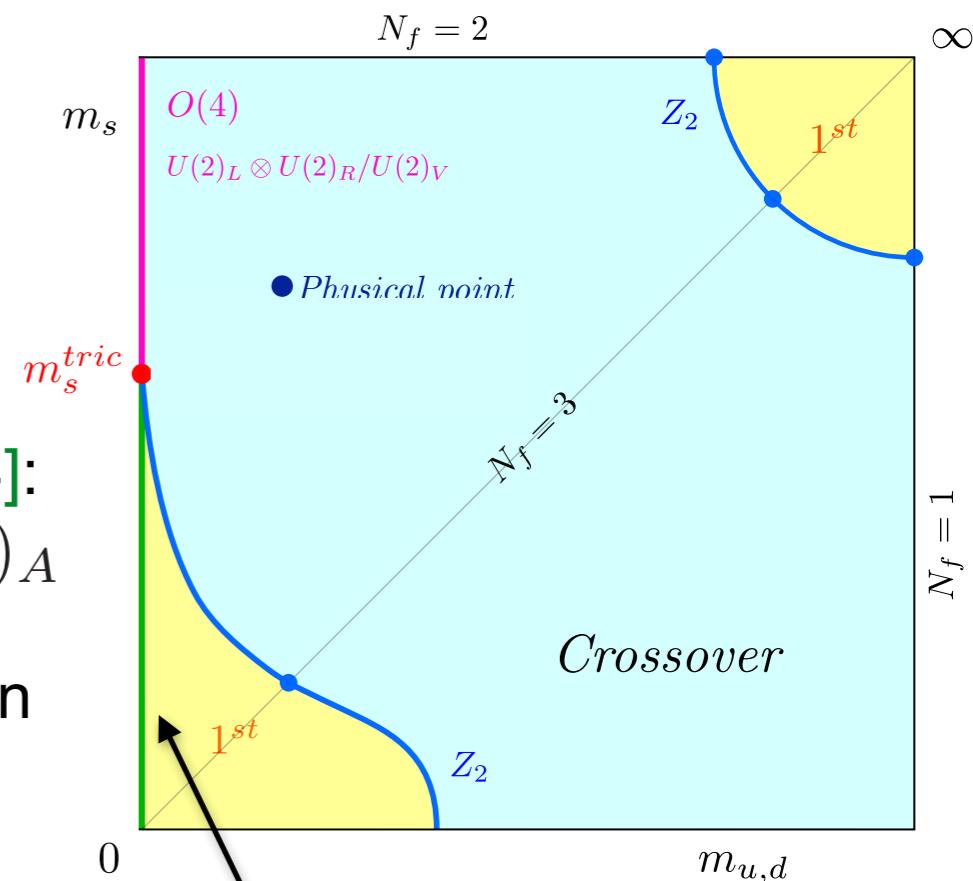


The nature of the QCD thermal transition

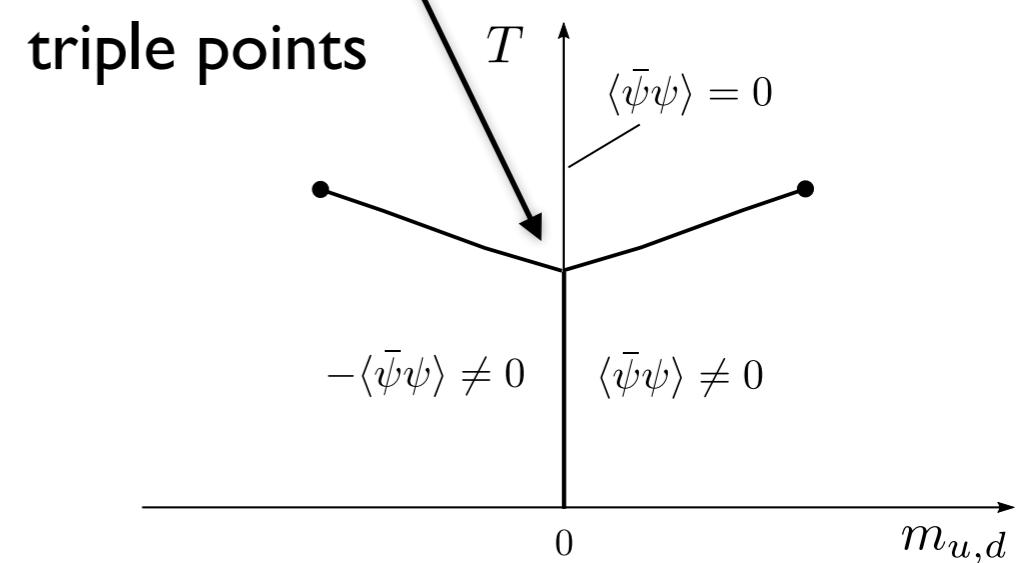
deconfinement p.t.:
breaking of global $Z(3)$ symmetry



[Pisarski, Wilczek, PRD 84]:
 $N_f = 2$ depends on $U(1)_A$
restored
 $N_f \geq 3$ 1st order

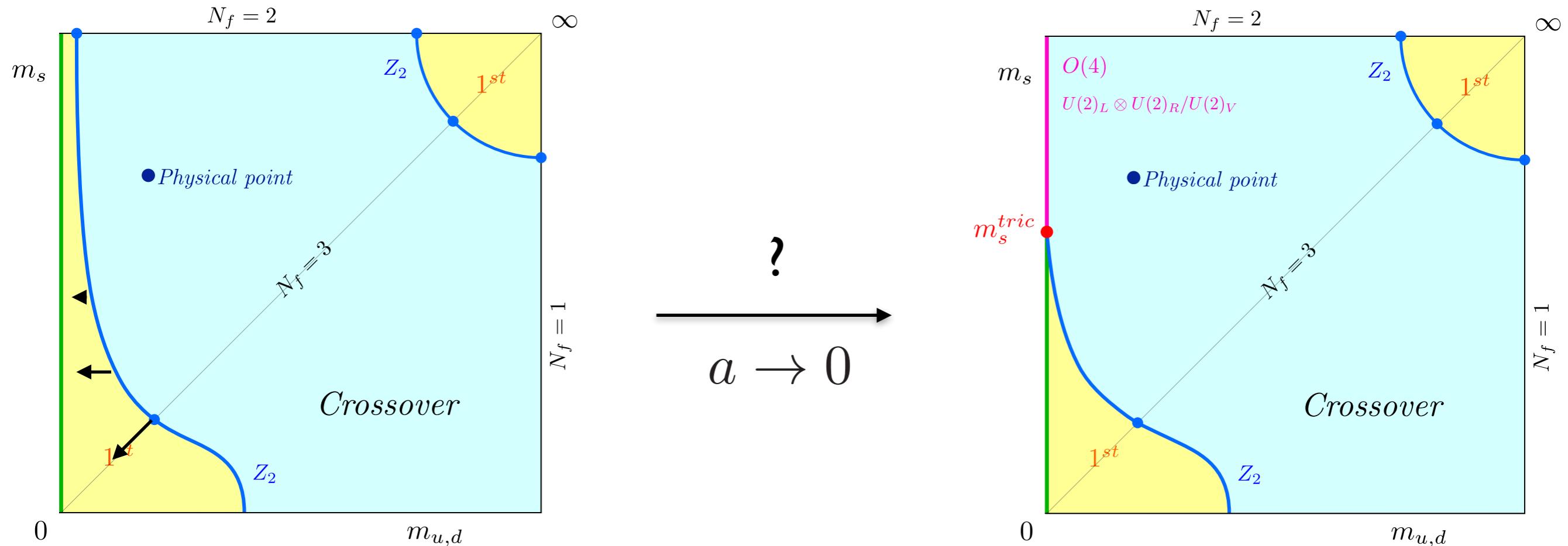


chiral p.t.
restoration of global symmetry in flavour space
 $SU(2)_L \times SU(2)_R \times U(1)_A$
anomalous



The nature of the QCD chiral transition

...is elusive, massless limit **not simulable!**

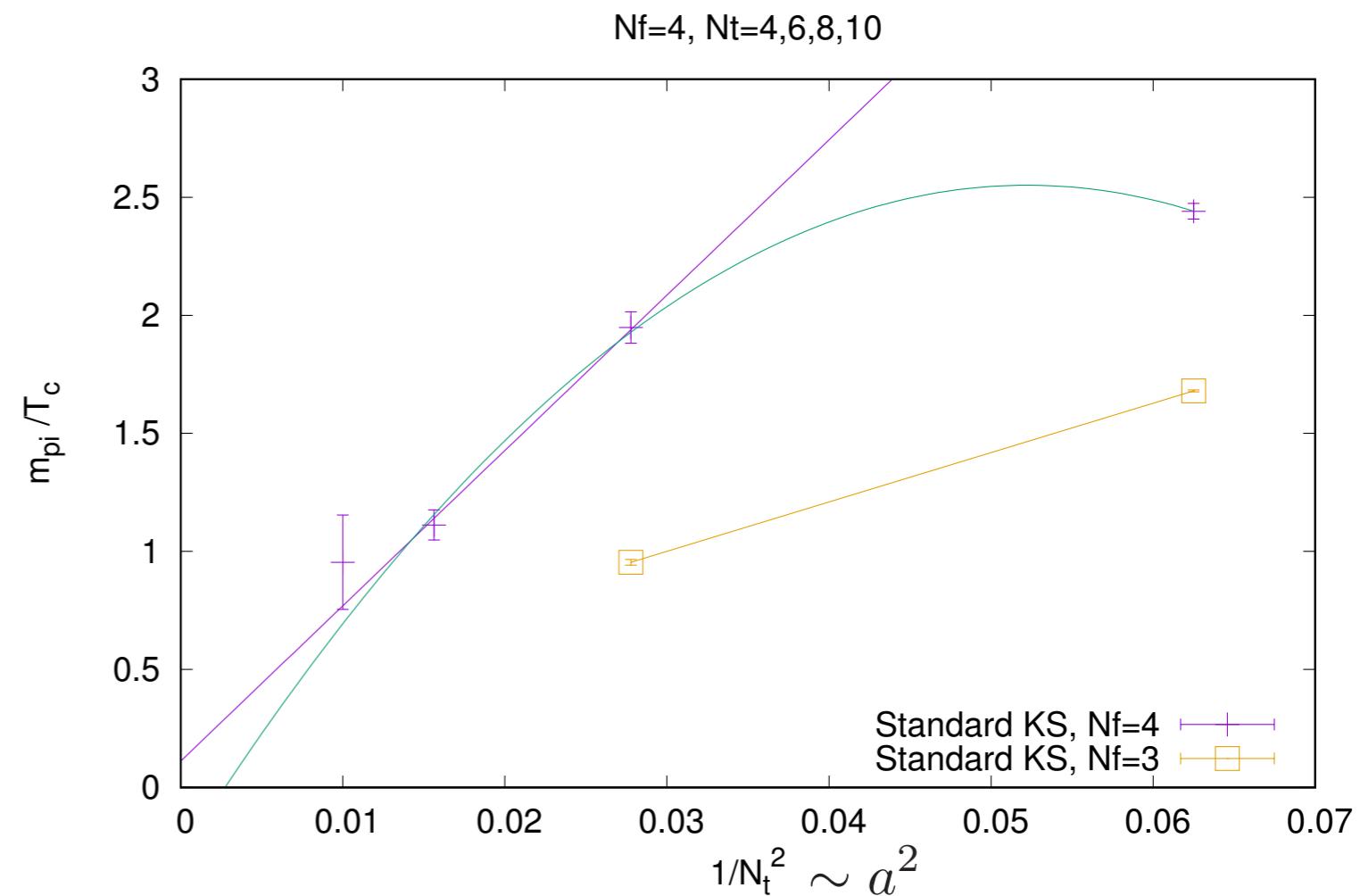
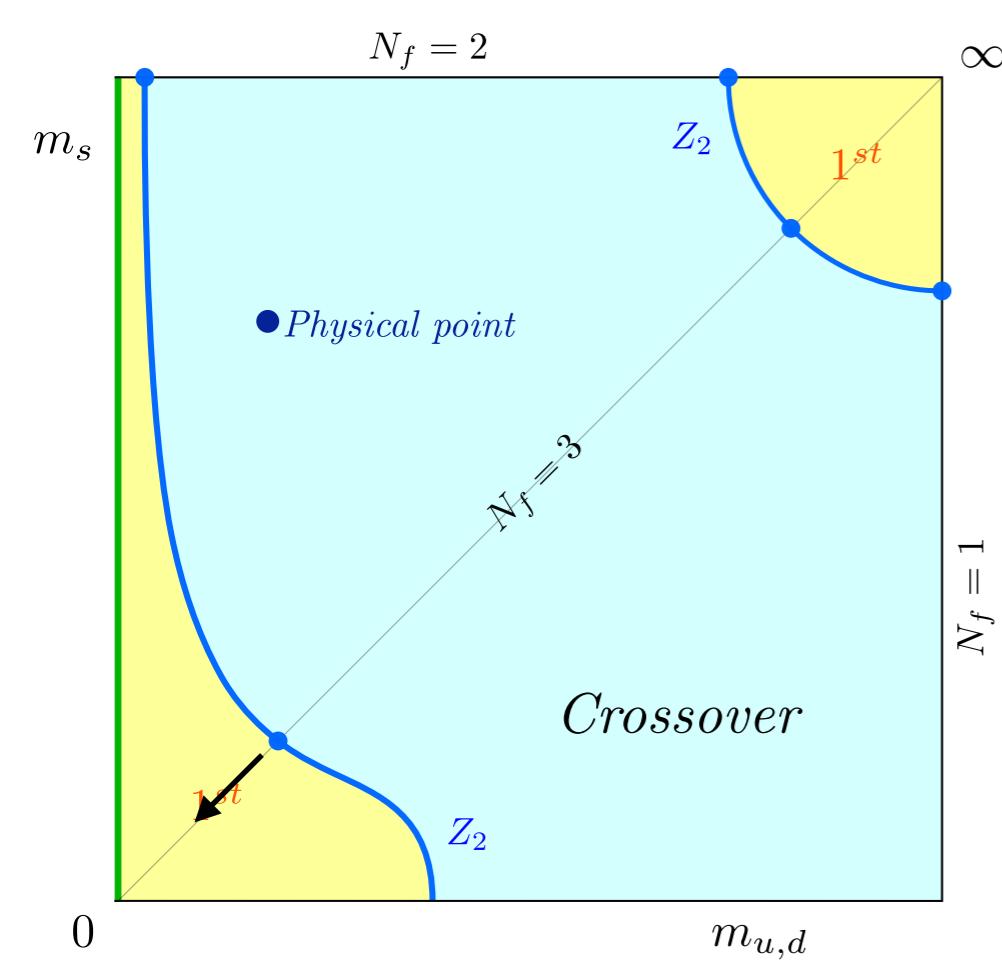


- Coarse lattices or unimproved actions: 1st order for $N_f = 2, 3$
- 1st order region shrinks rapidly as $a \rightarrow 0$
- Improved staggered actions: no 1st order region so far, even for $N_f = 3$ $m_{PS} > 45\text{MeV}$

Details and references: [O.P., Symmetry 13, 2021]

The nature of the QCD chiral transition, $N_f=3,4$

...has enormously large cut-off effects!



Unimproved staggered:

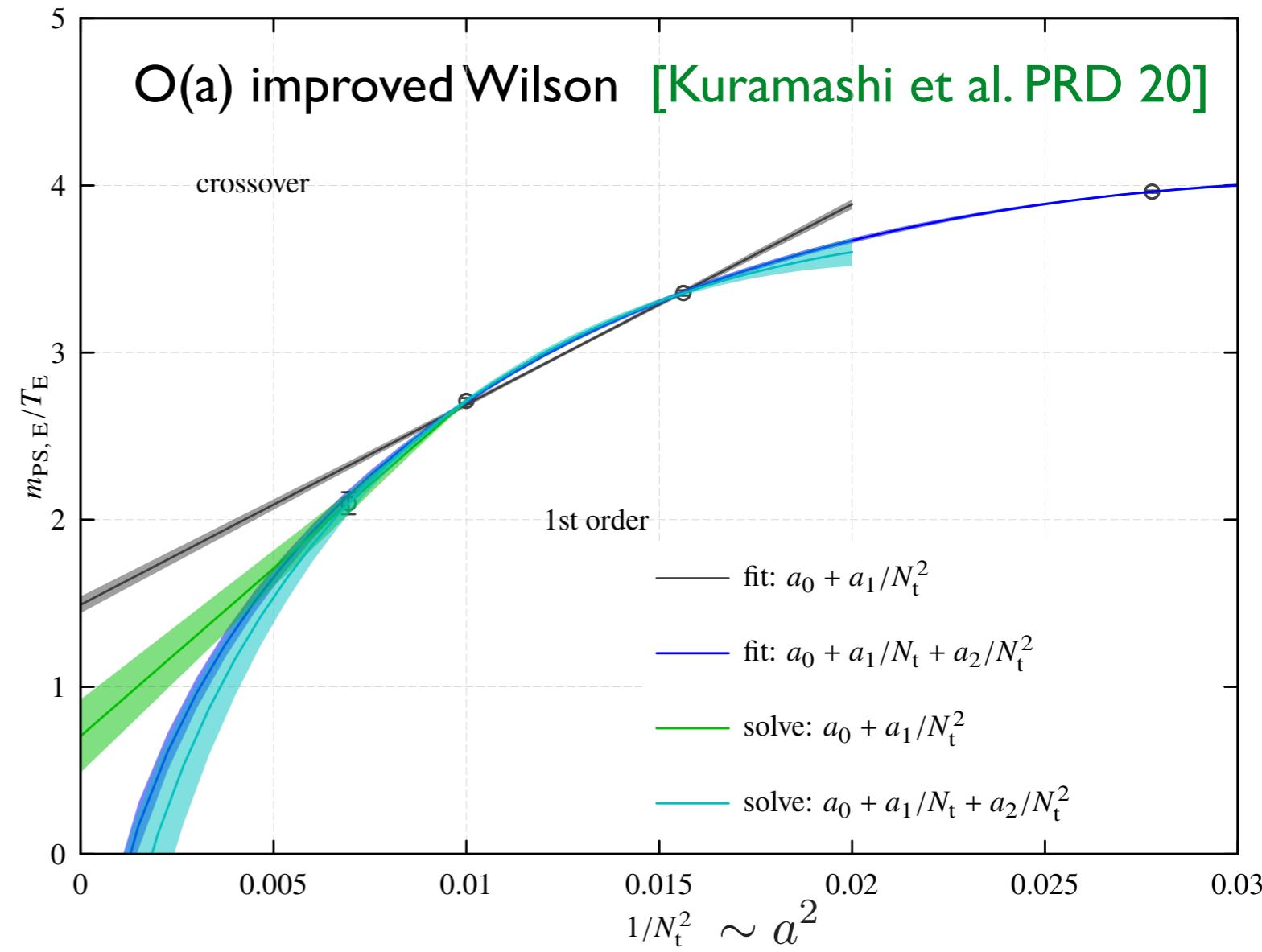
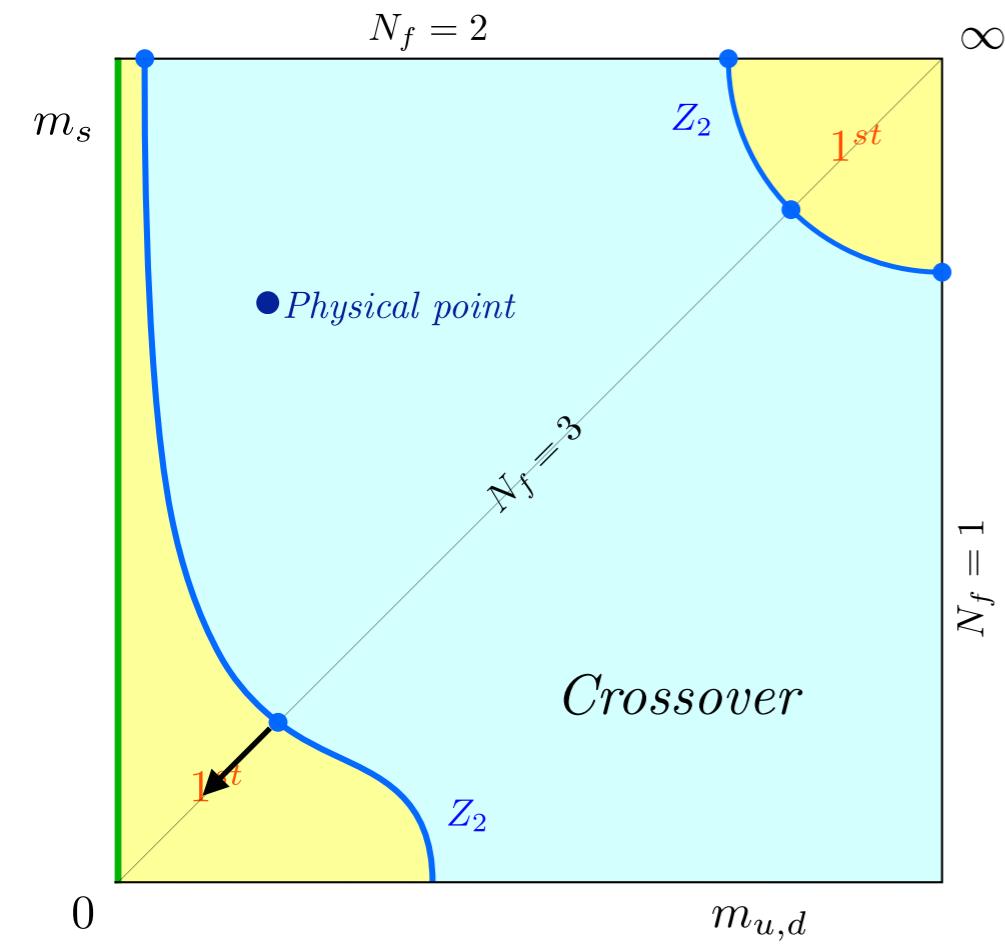
1st order region shrinks for $a \rightarrow 0$, both for $N_f = 3, 4$

[de Forcrand, D'Elia, PoS LAT 17]

No 1st order seen for improved staggered actions, even $N_f = 3$ [HotQCD PRD 17, PRD 22]

The nature of the QCD chiral transition, Nf=3

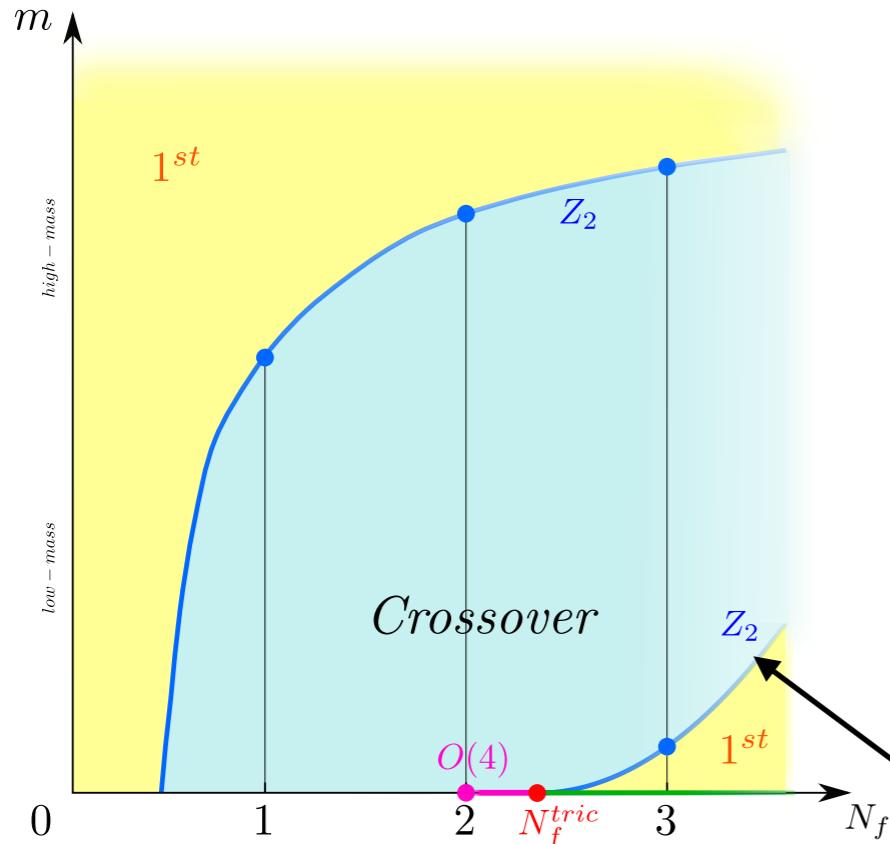
...has enormously large cut-off effects!



O(a)-improved Wilson:
1st order region shrinks for $a \rightarrow 0$

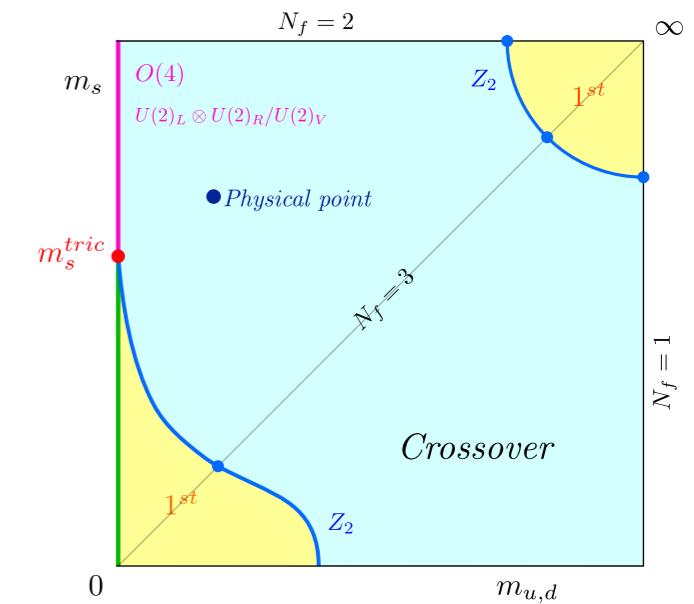
$$m_\pi^c \leq 110 \text{ MeV} \quad N_\tau = 4, 6, 8, 10, 12$$

Different view point: mass degenerate quarks



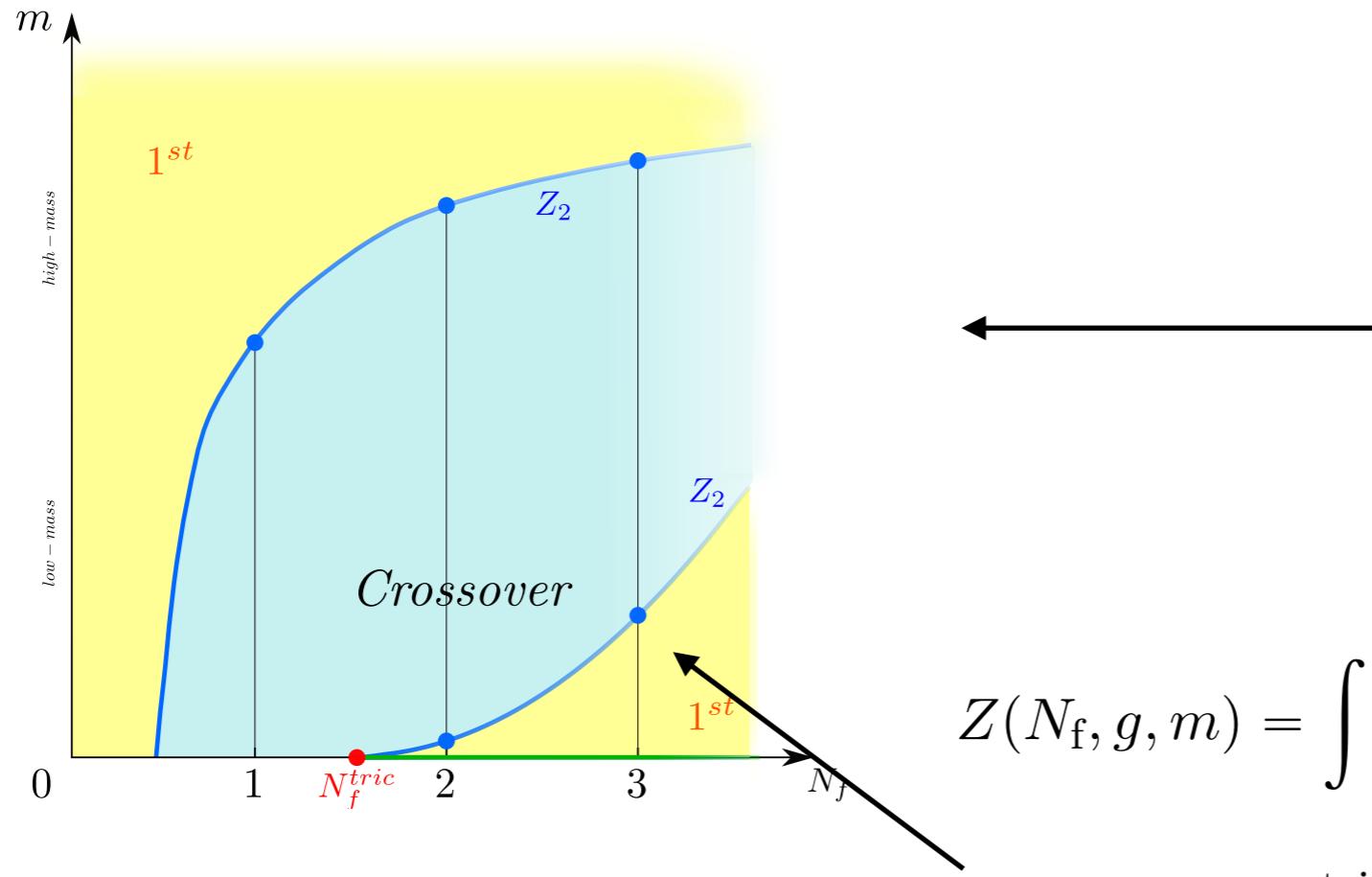
$$Z(N_f, g, m) = \int \mathcal{D}A_\mu (\det M[A_\mu, m])^{N_f} e^{-S_{\text{YM}}[A_\mu]}$$

$$N_f^c(am) = N_f^{\text{tric}} + \mathcal{B}_1 \cdot (am)^{2/5} + \mathcal{O}((am)^{4/5})$$



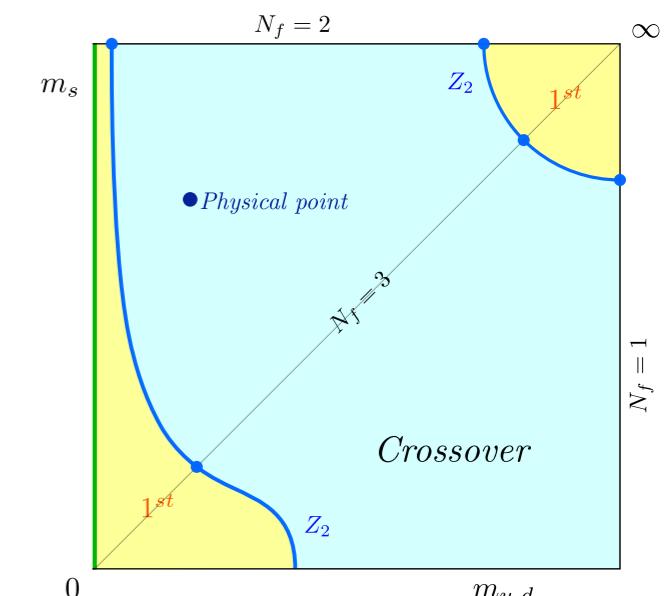
- Consider analytic continuation to continuous N_f
- Tricritical point **guaranteed** to exist if there is 1st order at any N_f
- Known exponents for critical line entering tric. point!
- Continuation to $a \neq 0$: $Z(2)$ surface ends in tricritical line

Different view point: mass degenerate quarks



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Methodology to determine order of transition

Finite size scaling of generalised cumulants

$$B_n = \frac{\langle (\bar{\psi}\psi - \langle\bar{\psi}\psi\rangle)^n \rangle}{\langle (\bar{\psi}\psi - \langle\bar{\psi}\psi\rangle)^2 \rangle^{n/2}}$$

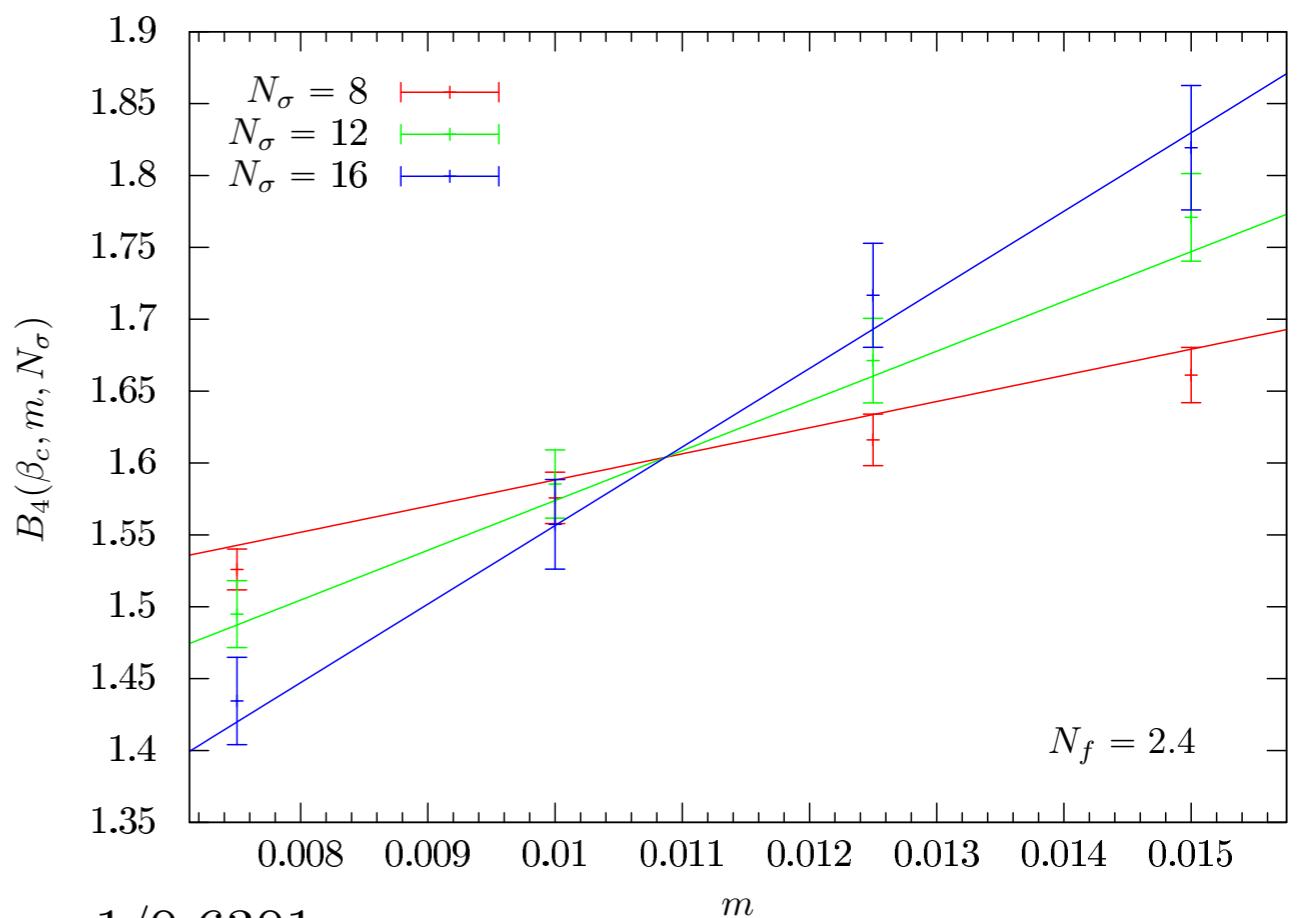
Standard staggered fermions, bare parameters: β, am, N_f, N_τ

(Pseudo-critical) phase boundary: $B_3 = 0$ 3d manifold

Second-order 3D Ising:

2d chiral critical surface separates 1st order from crossover

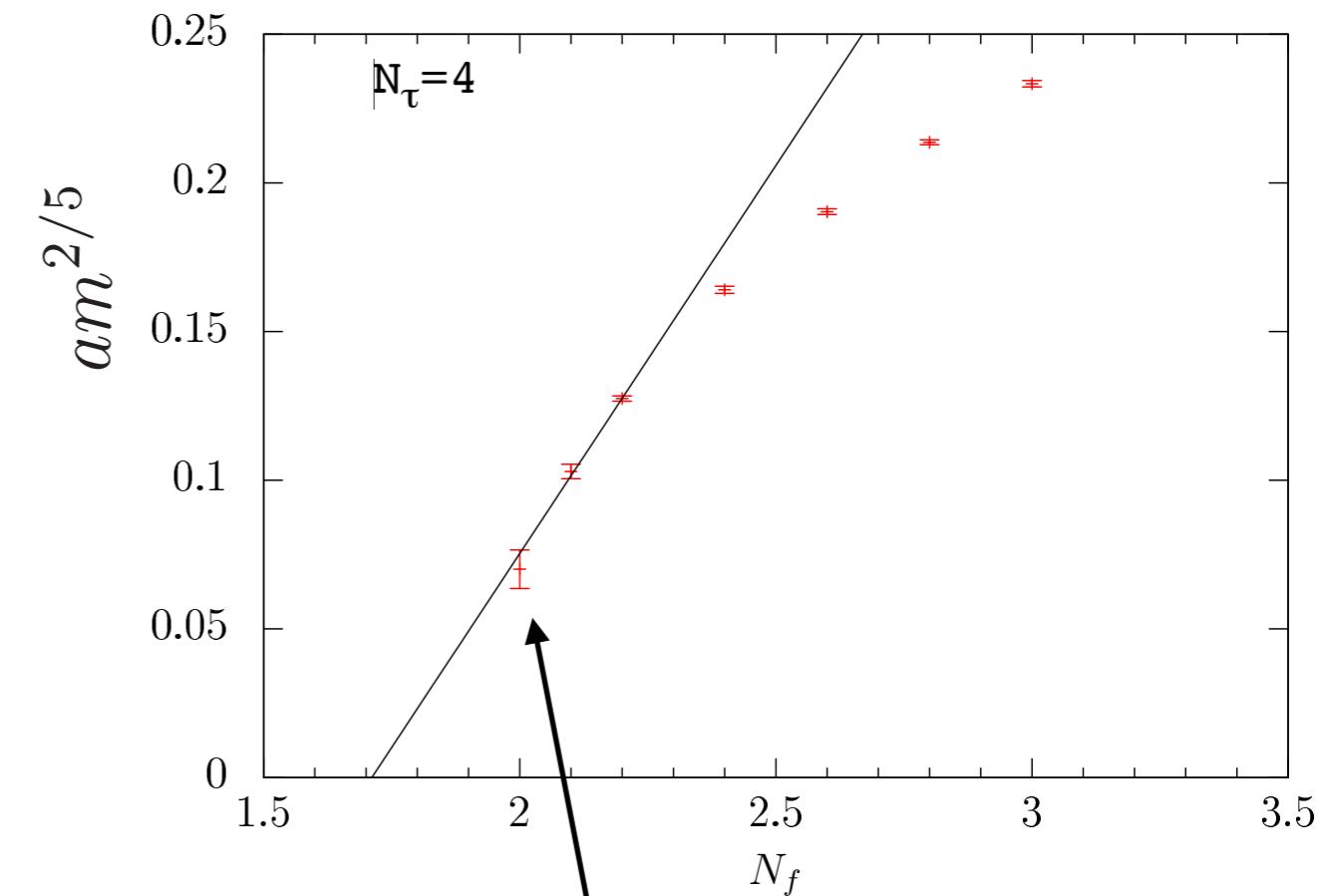
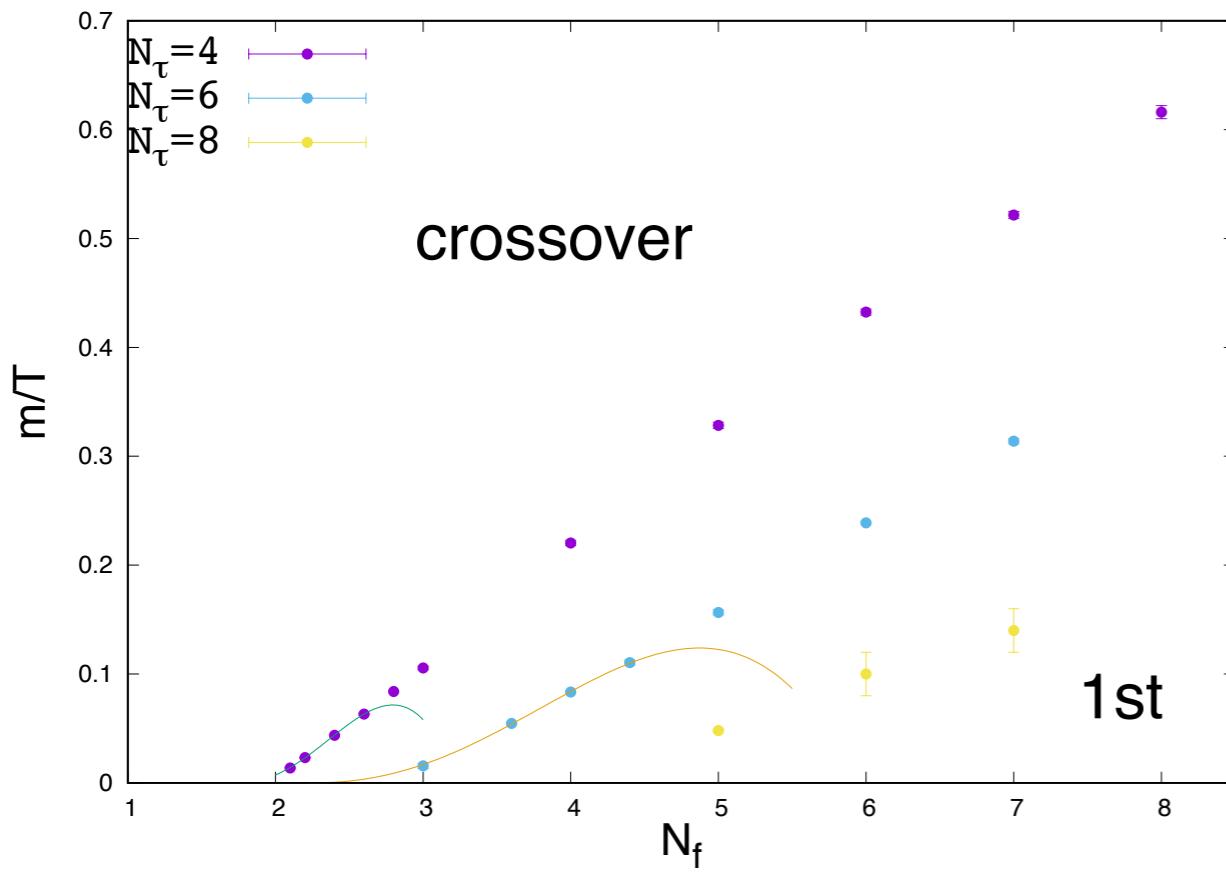
$$B_4(\beta_c, am, N_\sigma) \approx 1.604 + c(am - am_c) N_\sigma^{1/0.6301}$$



Bare parameter space of unimproved staggered LQCD

[Cuteri, O.P., Sciarra 21]

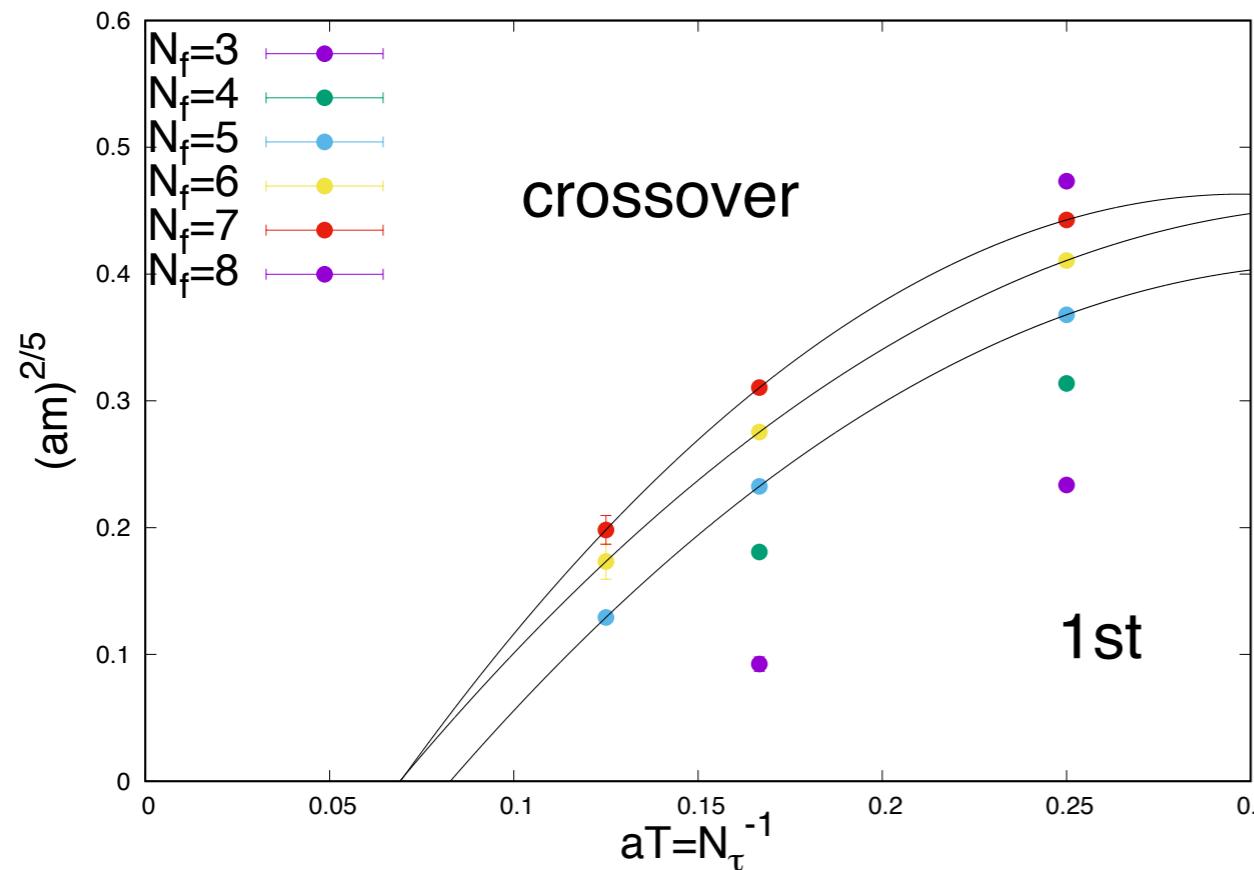
$\sim 120\text{ M}$ Monte Carlo trajectories with light fermions,
aspect ratios 3,4,5



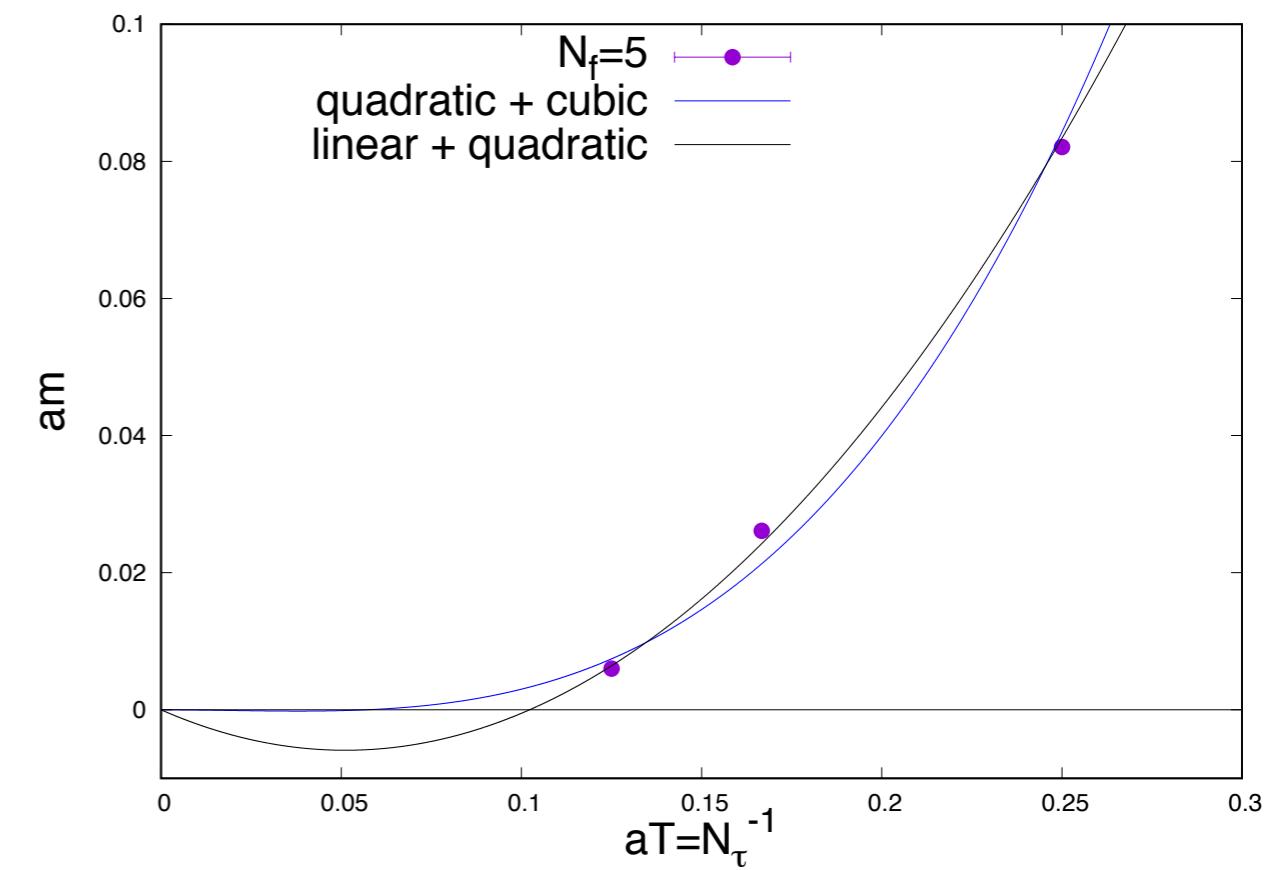
- Tricritical scaling observed also in different variable pairings
- Extrapolated by tric. scaling from finite imaginary μ [Bonati et al. PRD 14]
- Old question: $m_c/T = 0$ or $\neq 0$? Answered for $N_f = 2$
- New question: will N_f^{tric} slide beyond $N_f = 3$?

Bare parameter space of unimproved staggered LQCD

[Cuteri, O.P., Sciarra 21]



1st order scenario does not fit!

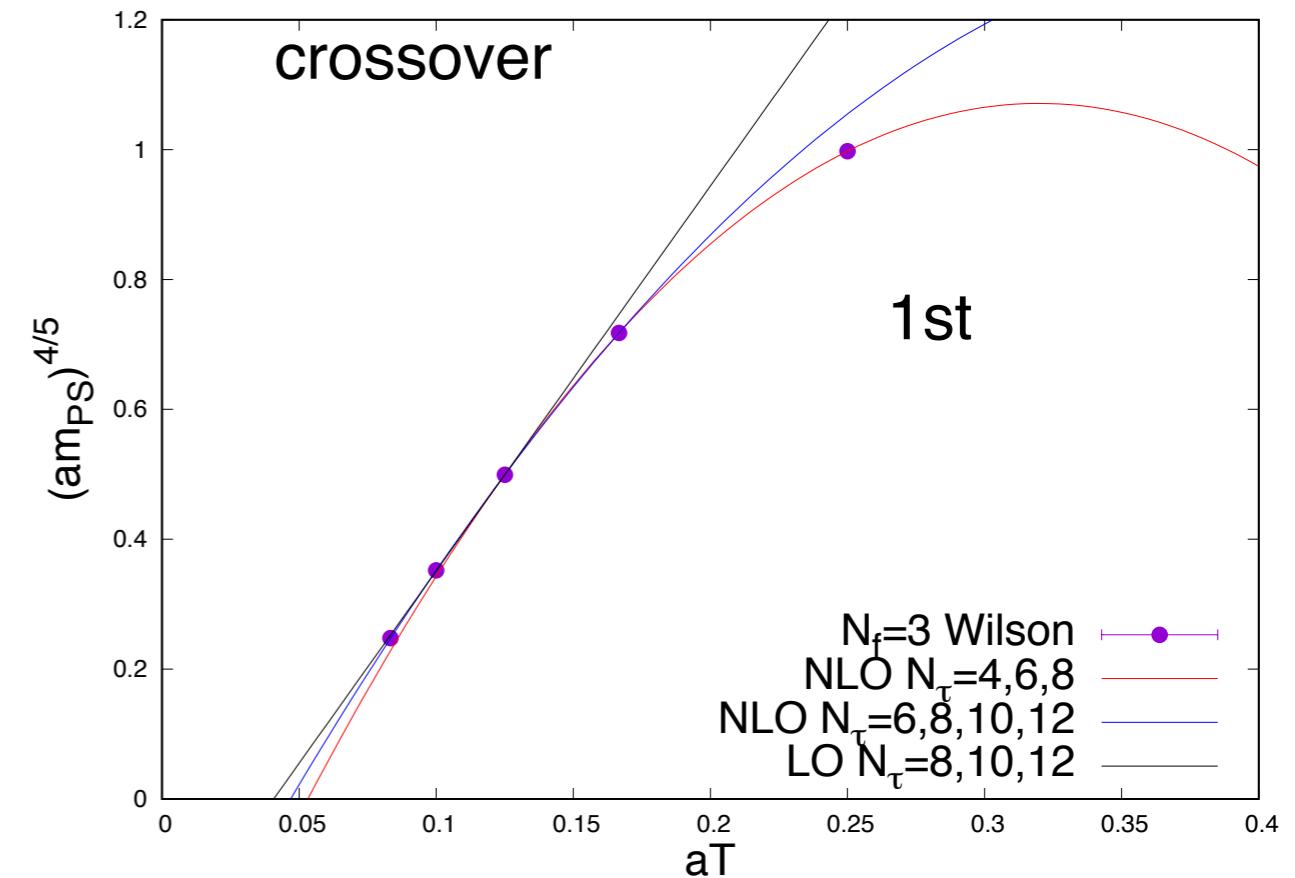
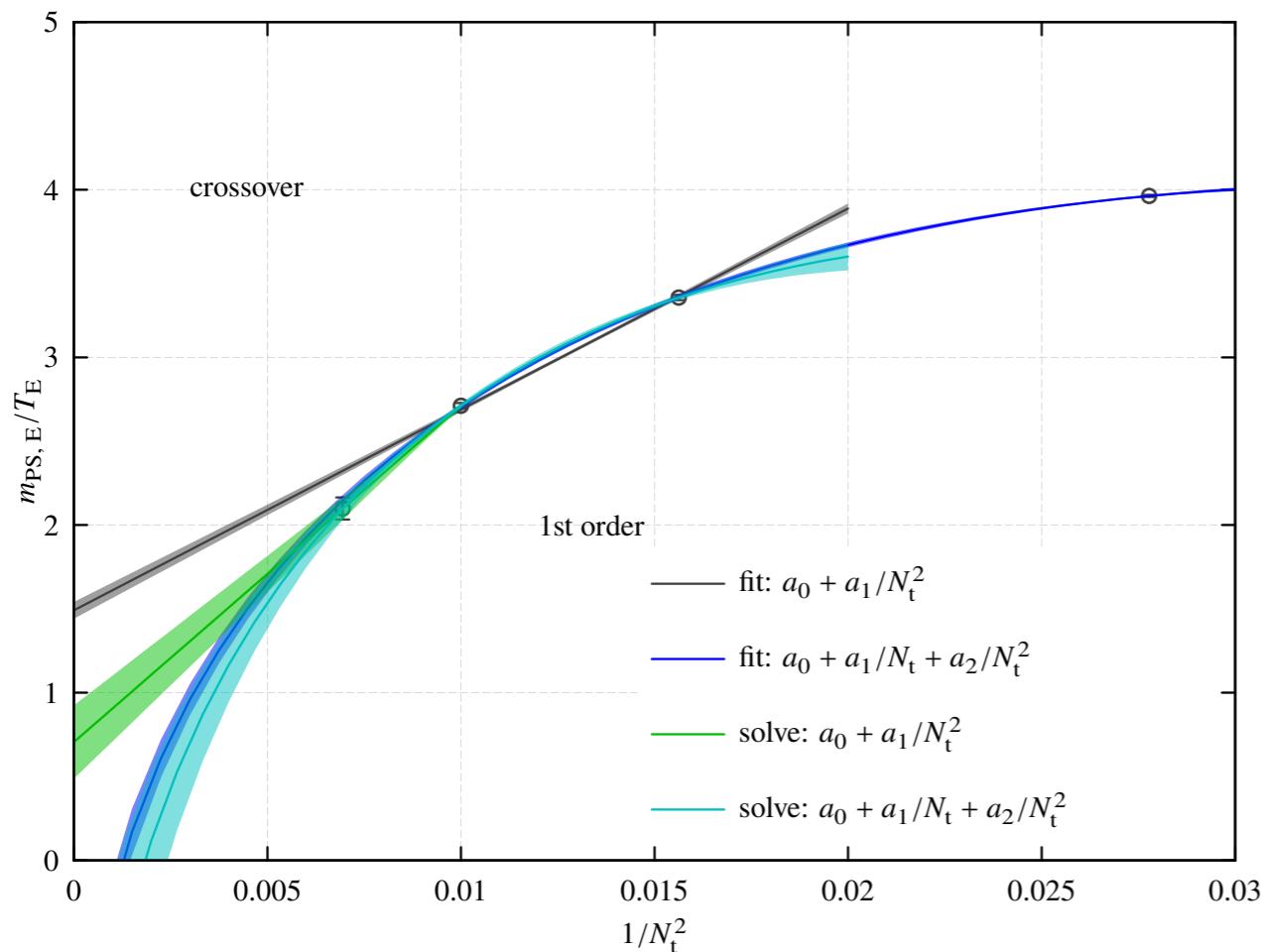


- Tricritical scaling observed also in plane of mass vs. lattice spacing
 - Allows extrapolation to lattice chiral limit, tricritical points $N_{\tau}^{\text{tric}}(N_f)$
 - 1st order scenario: $m_c(a) = m_c(0) + c_1(aT) + c_2(aT)^2 + \dots$
- Incompatible with data! $\chi^2_{\text{dof}} > 10$

Nf=3 O(a)-improved Wilson fermions

[Kuramashi et al. PRD 20] $m_\pi^c \leq 110$ MeV $N_\tau = 4, 6, 8, 10, 12$

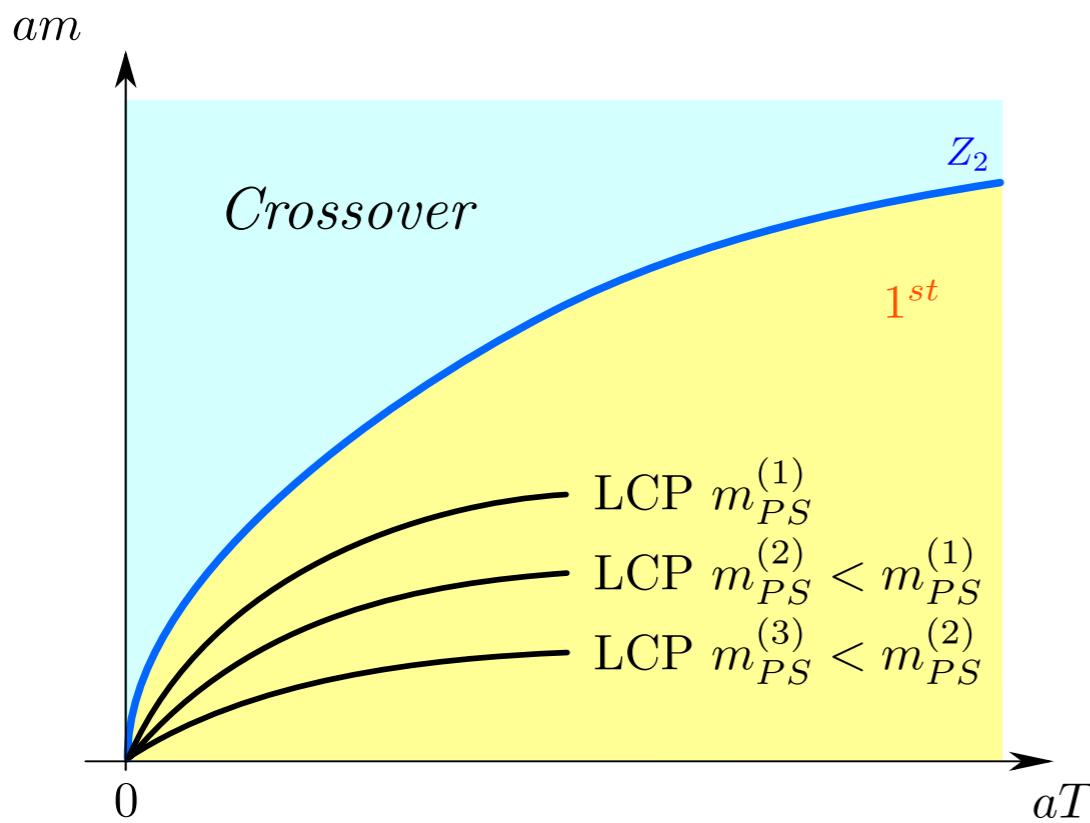
Re-analysis using: $am_{PS}^2 \propto am_q$



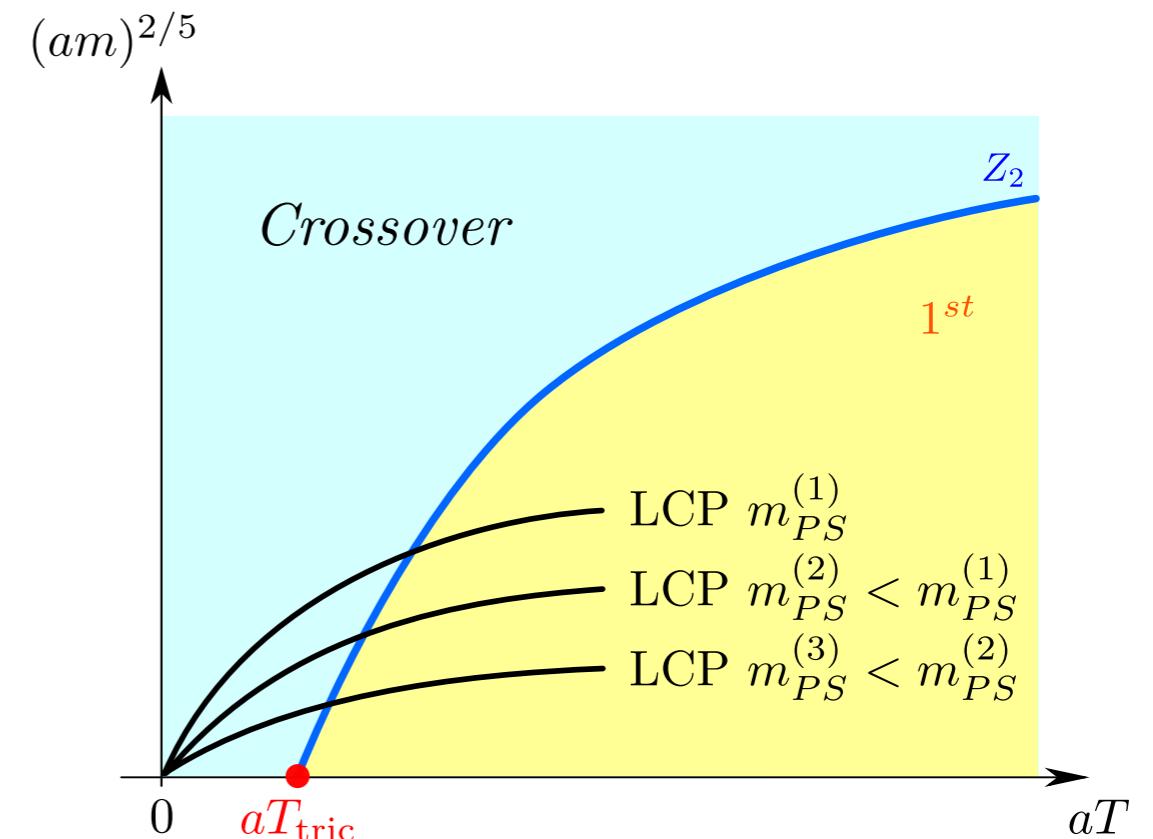
Tricritical scaling, Nf=3 consistent with unimproved staggered!

Implications for the continuum

- Finite $N_\tau^{\text{tric}}(N_f)$ implies that 1st order transition is not connected to continuum
- Approaching continuum first, then chiral limit:
Continuum chiral phase transition second-order!

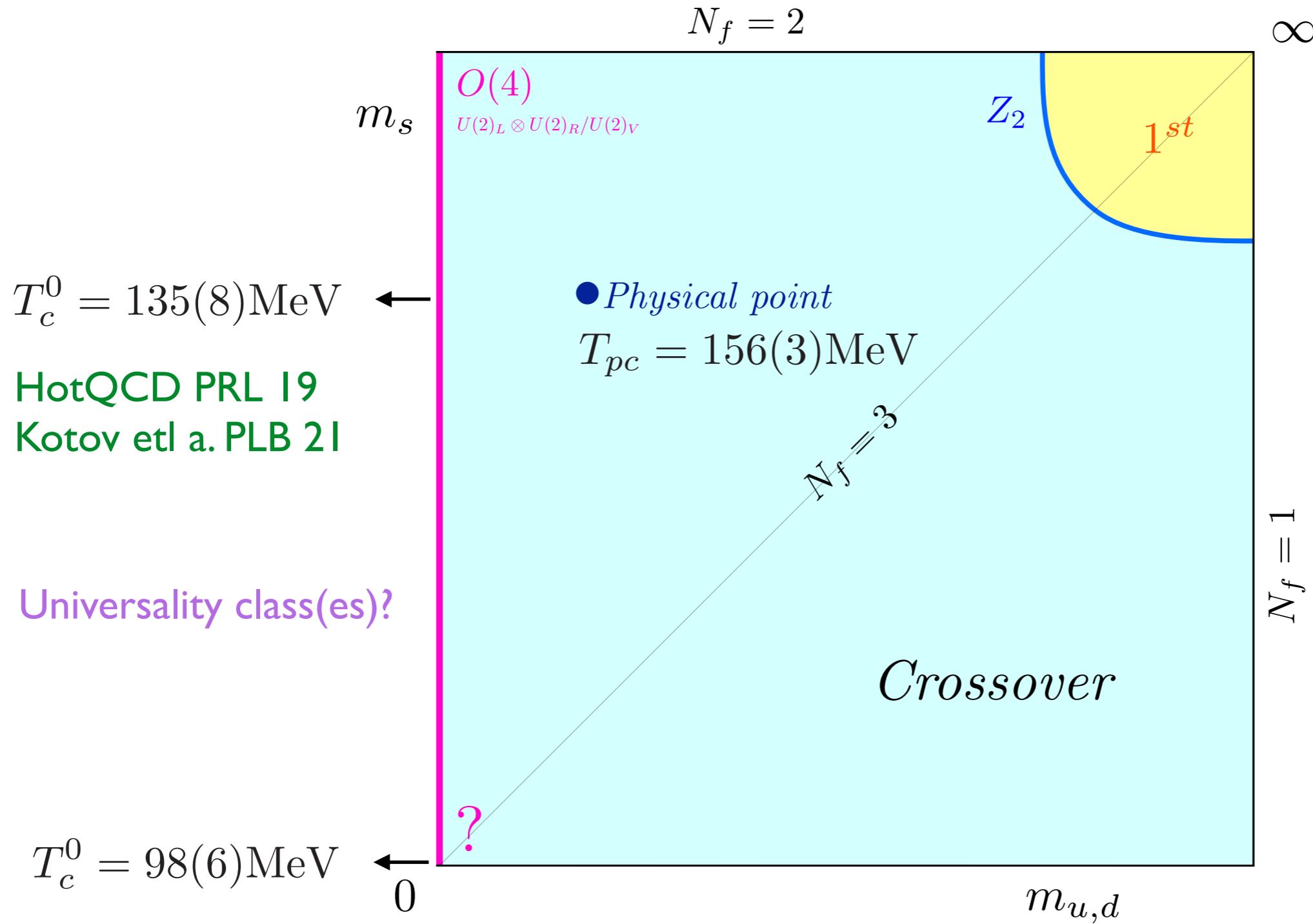


1st order scenario

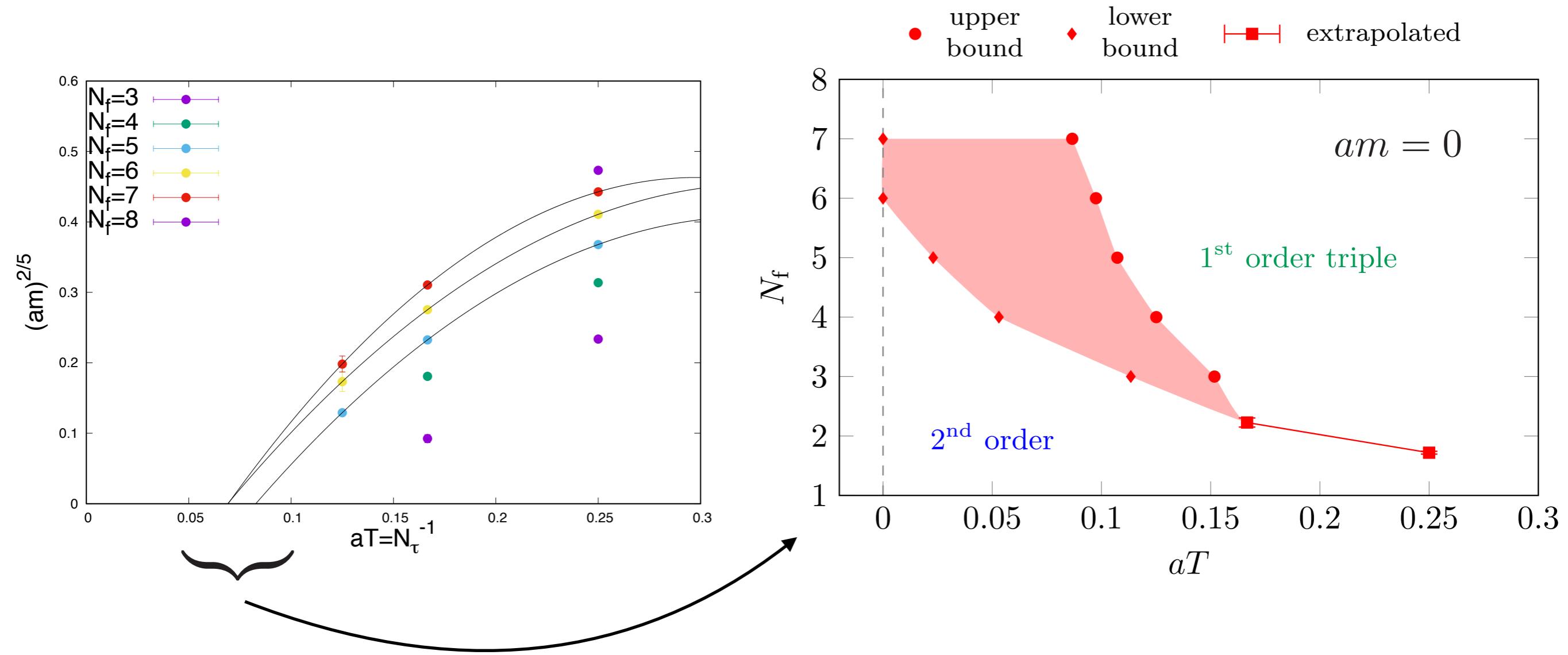


2nd order scenario

The Columbia plot in the continuum



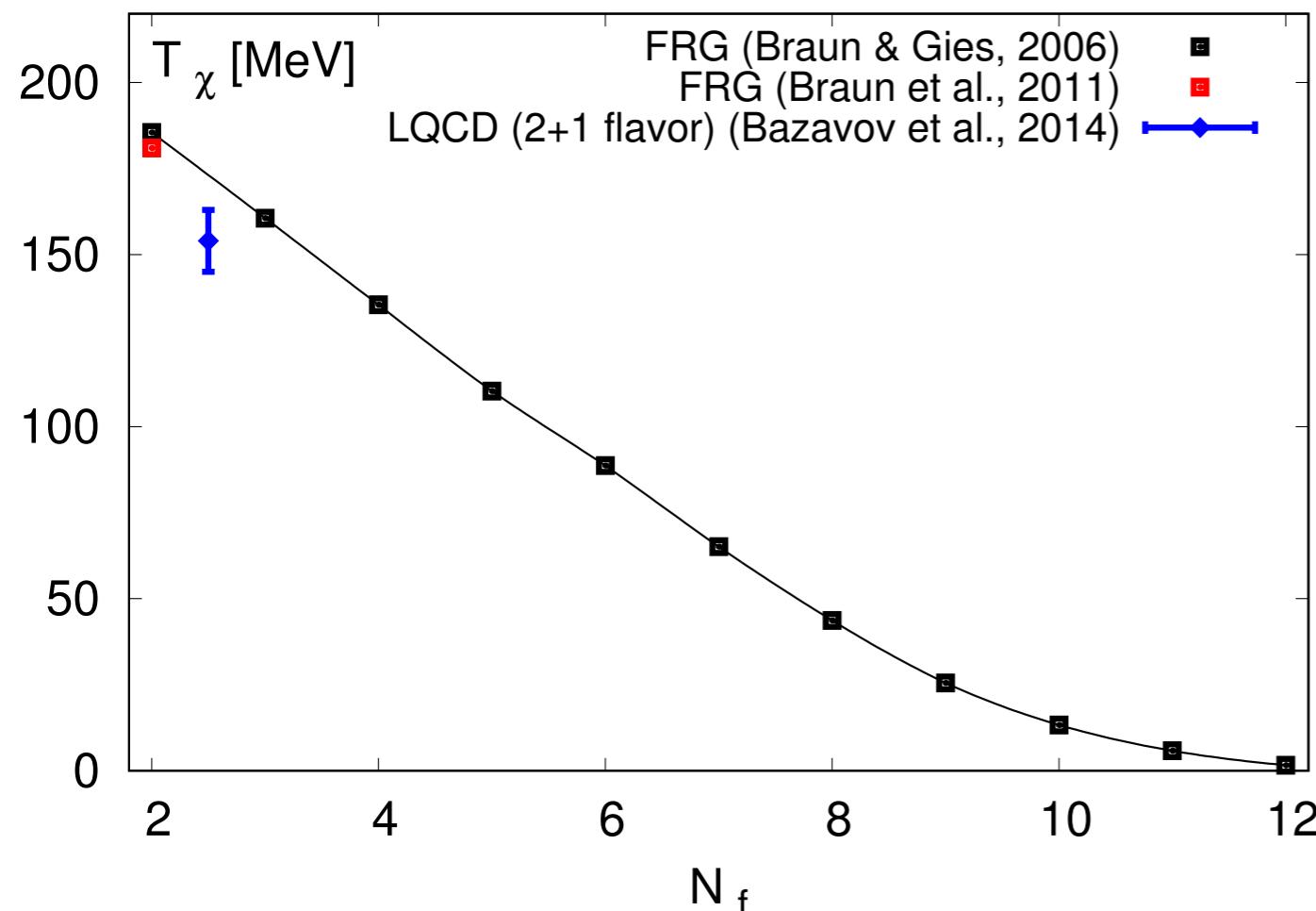
Staggered: tricritical points as function of N_f



- $N_\tau^{\text{tric}}(N_f)$ increasing function
- Tricritical line in the plane of the lattice chiral limit, separates 1st from 2nd

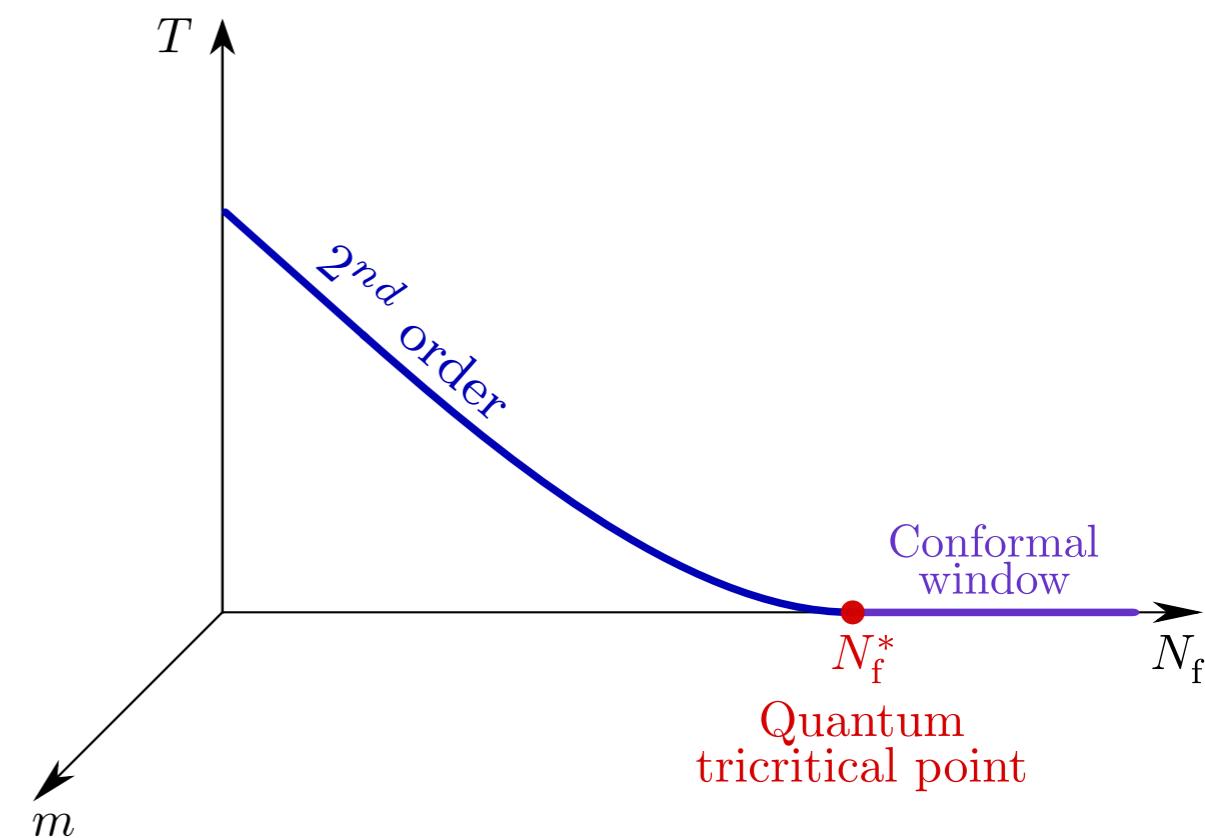
The chiral phase transition for different N_f

Temperature dependence:



For lattice, see [Miura, Lombardo, NPB 13]

Order of the transition:



[Cuteri, O.P., Sciarra 21]

The chiral phase transition in the massless limit is likely second-order for all N_f

What about Pisarski, Wilczek?

- Investigated 3d sigma model,
i.e. ϕ^4 - Ginzburg-Landau-Wilson theory for chiral condensate
- Results based on epsilon expansion about $\epsilon = 1$
- All conclusions confirmed by [Butti, Pelissetto, Vicari, JHEP 03]
(High order perturbative expansion in fixed d)
- Support also from simulation of 3d sigma model [Gausterer, Sanielovici, PLB 88]

Suggested resolution:

ϕ^6 term should be included in 3d, renormalisable

FRG: 3d ϕ^6 has infrared fixed points and 2nd order transitions [Litim, Tetradis, NPB 96]

[Fejos, arxiv:2201.07909] 3d ϕ^6 with t'Hooft term: 2nd order transition for restored anomaly!

Conclusions

- Zero density, unimproved staggered, $N_f=2-6$, $O(a)$ -improved Wilson $N_f=3$:
1st order chiral transition region not connected to continuum limit
- Chiral transition second order up to conformal window
- Check: Wilson $N_f=4$
- Domain wall and overlap fermions..... + FRG!

Long history, key results

- [Pisarski, Wilczek PRD 84]

Universality, 3D sigma models: 1st order for

$$N_f \geq 3$$

$N_f = 2$ depends on $U(1)_A$

- Shrinking of 1st-order region towards continuum, will something remain?:

Standard staggered

$N_f = 2$ [Bonati et al. PRD 14, Cuteri et al. PoS LAT 18]

$N_f = 3$ [de Forcrand, O.P. PoS LAT 07]

$N_f = 4$ [de Forcrand, D'Elia PoS LAT 16]

$\mathcal{O}(a)$ -improved Wilson

$N_f = 3$ [Jin et al. PRD 15,17; Kuramashi et al. PRD 20]

$N_f = 4$ [Ohno et al. PoS LAT 18]

- No 1st order transition seen at all:

HISQ, $N_f = 3, m_{PS} \geq 50$ MeV

[HotQCD. PRD 17, 21]