

Phase Transitions in Particle Physics
Mini Workshop
Galileo Galilei Institute, Firenze – Spring 2022

Interfaces near criticality: results from field theory

Marianna Sorba
SISSA, Trieste



Interfaces

In **particle physics**, the time propagation of a confining flux tube joining a quark-antiquark pair spans an interface.

Interfaces

In **particle physics**, the time propagation of a confining flux tube joining a quark-antiquark pair spans an interface.

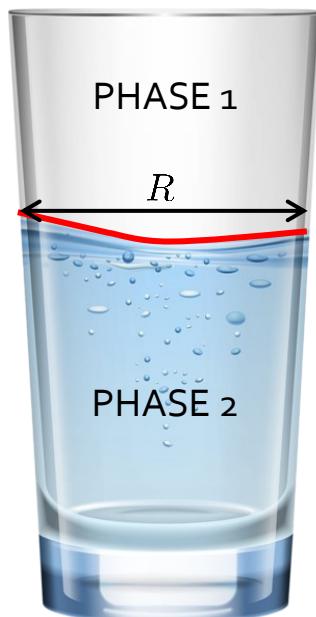
In **statistical systems**, an interface is a surface of separation between two **coexisting phases**:



Interfaces

In **particle physics**, the time propagation of a confining flux tube joining a quark-antiquark pair spans an interface.

In **statistical systems**, an interface is a surface of separation between two **coexisting phases**:

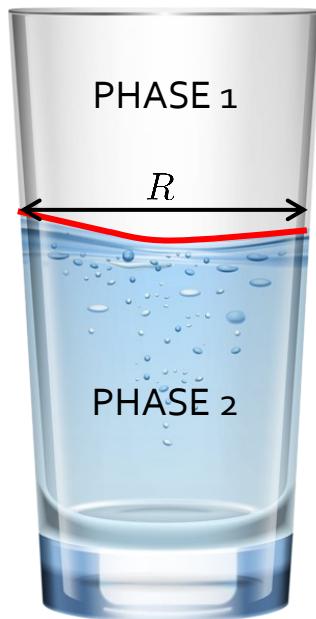


- Discrete internal symmetry
- $T < T_c$
- Suitable boundary conditions on the bulk (homogeneous) system
- Interface linear size $R \gg \xi$

Interfaces

In **particle physics**, the time propagation of a confining flux tube joining a quark-antiquark pair spans an interface.

In **statistical systems**, an interface is a surface of separation between two **coexisting phases**:

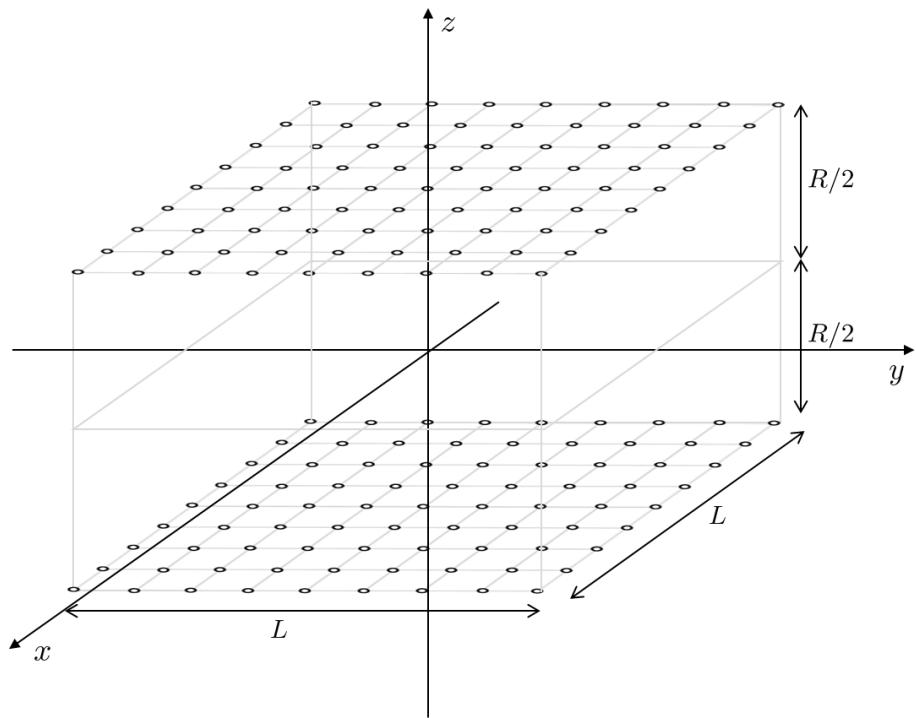


- Discrete internal symmetry
- $T < T_c$
- Suitable boundary conditions on the bulk (homogeneous) system
- Interface linear size $R \gg \xi$

Simplest implementation:
three-dimensional Ising model

Ising model

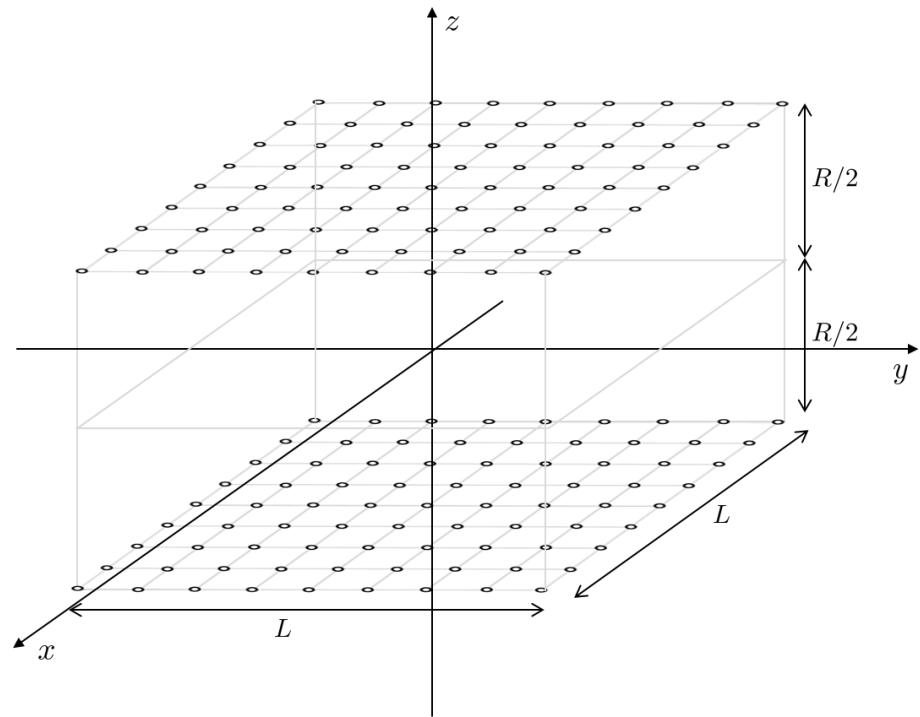
$$\mathcal{H} = -\frac{1}{T} \sum_{\langle i,j \rangle} s_i s_j, \quad s_i = \pm 1 \quad \text{on a cubic lattice} \quad (R \gg \xi, L \rightarrow \infty)$$



Ising model

$$\mathcal{H} = -\frac{1}{T} \sum_{\langle i,j \rangle} s_i s_j, \quad s_i = \pm 1 \quad \text{on a cubic lattice} \quad (R \gg \xi, L \rightarrow \infty)$$

$\begin{cases} T < T_c & \Rightarrow \quad \mathbb{Z}_2 \text{ spontaneously broken: } |\langle s_i \rangle| = M \neq 0 \\ T \rightarrow T_c & \Rightarrow \quad \text{field theory description and universality} \end{cases}$



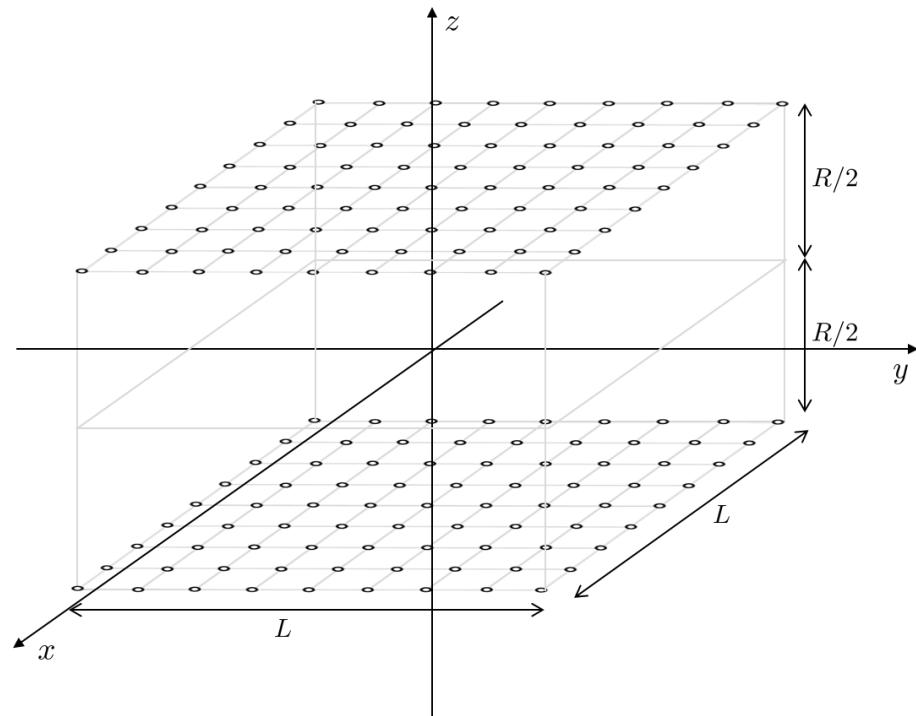
Ising model

$$\mathcal{H} = -\frac{1}{T} \sum_{\langle i,j \rangle} s_i s_j, \quad s_i = \pm 1 \quad \text{on a cubic lattice } (R \gg \xi, L \rightarrow \infty)$$

$$\begin{cases} T < T_c & \Rightarrow \quad \mathbb{Z}_2 \text{ spontaneously broken: } |\langle s_i \rangle| = M \neq 0 \\ T \rightarrow T_c & \Rightarrow \quad \textbf{field theory description and universality} \end{cases}$$

Construct a
non-effective theory:

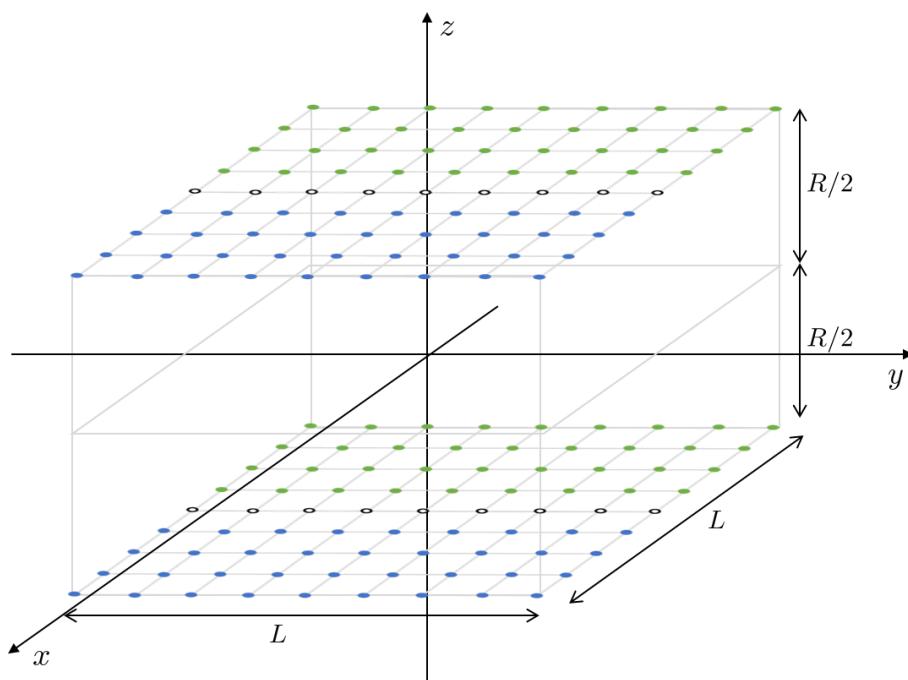
study interfaces
analytically at
fundamental level,
starting from the
bulk field theory.



Standard interface setting

Boundary conditions at $z = \pm R/2$:

$$s_i = \begin{cases} +1 & x < 0 \\ -1 & x > 0 \end{cases}$$

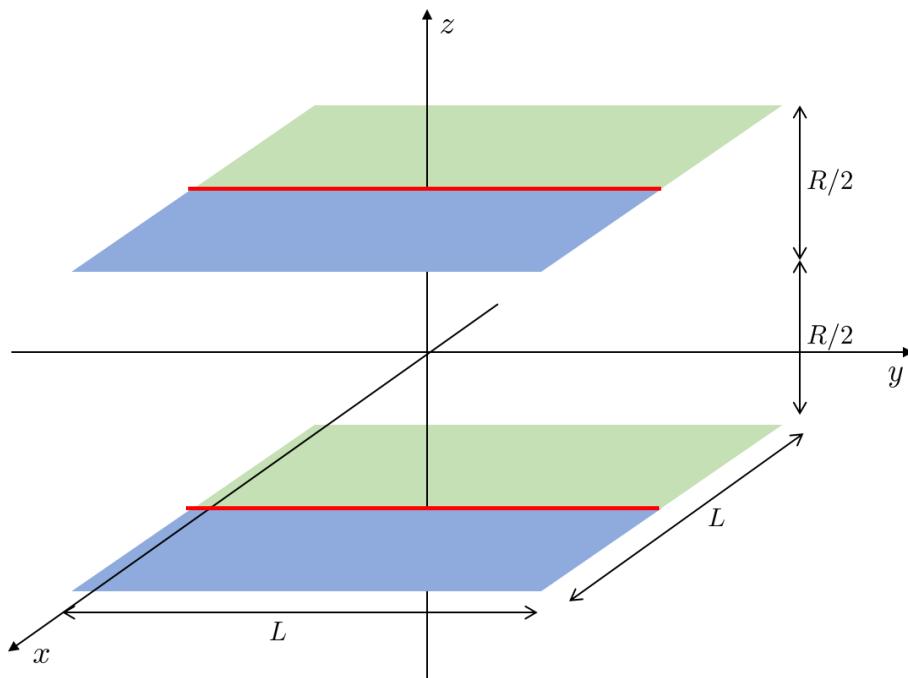


G. Delfino, W. Selke and A. Squarcini,
Nucl. Phys. B 958 (2020) 115139.

Standard interface setting

Boundary conditions at $z = \pm R/2$:

$$s_i = \begin{cases} +1 & x < 0 \\ -1 & x > 0 \end{cases} \quad \rightarrow \quad \begin{cases} \lim_{x \rightarrow -\infty} \langle s(x, y, 0) \rangle_{+-} = +M \\ \lim_{x \rightarrow +\infty} \langle s(x, y, 0) \rangle_{+-} = -M \end{cases}$$



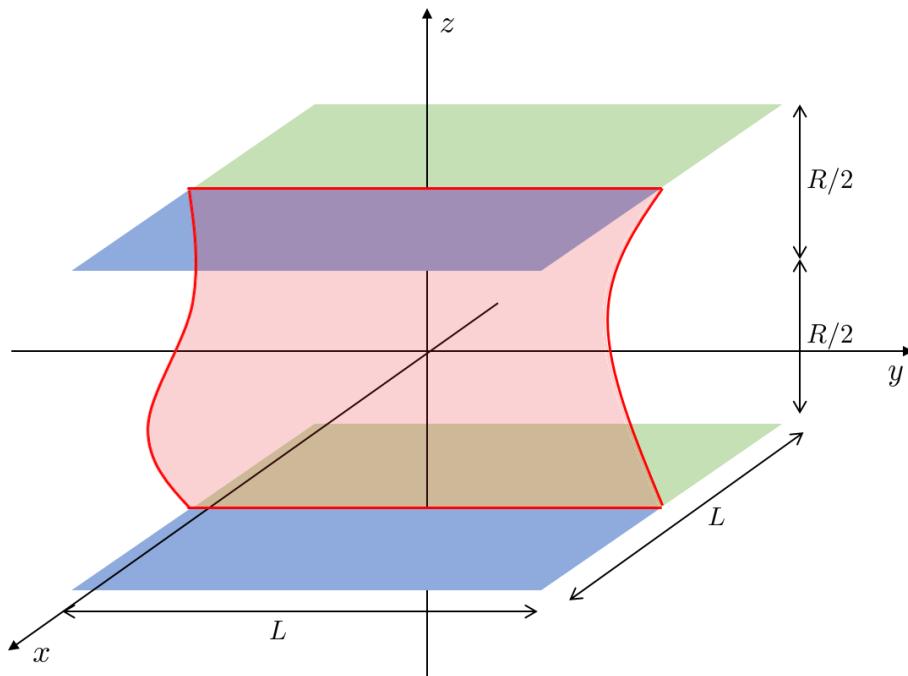
G. Delfino, W. Selke and A. Squarcini,
Nucl. Phys. B 958 (2020) 115139.

Standard interface setting

Boundary conditions at $z = \pm R/2$:

$$s_i = \begin{cases} +1 & x < 0 \\ -1 & x > 0 \end{cases} \quad \rightarrow \quad \begin{cases} \lim_{x \rightarrow -\infty} \langle s(x, y, 0) \rangle_{+-} = +M \\ \lim_{x \rightarrow +\infty} \langle s(x, y, 0) \rangle_{+-} = -M \end{cases}$$

Interface running between the pinning axes $x = 0, z = \pm R/2$.



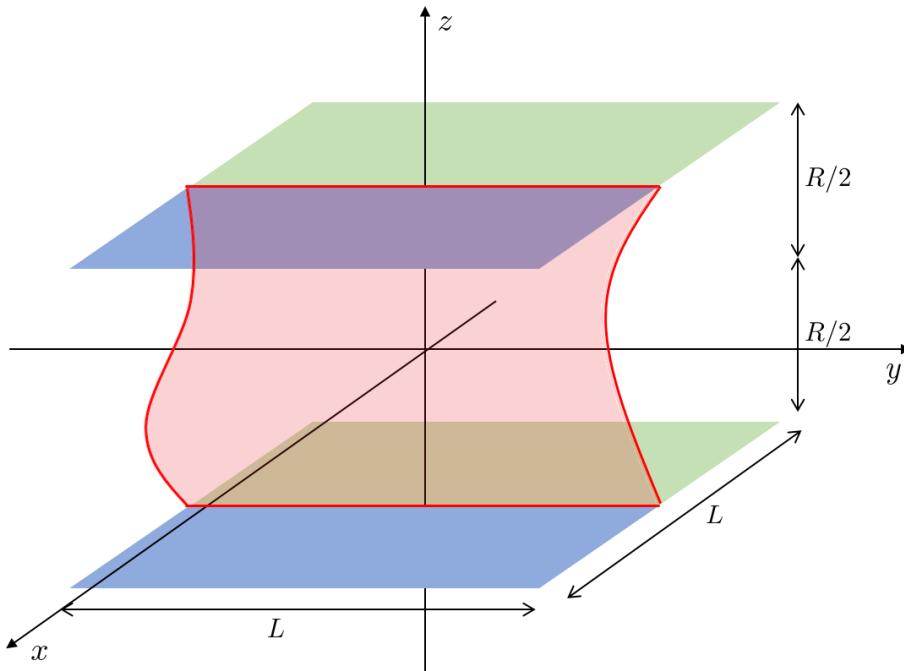
G. Delfino, W. Selke and A. Squarcini,
Nucl. Phys. B 958 (2020) 115139.

Standard interface setting

Boundary conditions at $z = \pm R/2$:

$$s_i = \begin{cases} +1 & x < 0 \\ -1 & x > 0 \end{cases} \quad \rightarrow \quad \begin{cases} \lim_{x \rightarrow -\infty} \langle s(x, y, 0) \rangle_{+-} = +M \\ \lim_{x \rightarrow +\infty} \langle s(x, y, 0) \rangle_{+-} = -M \end{cases}$$

Interface running between the pinning axes $x = 0, z = \pm R/2$.

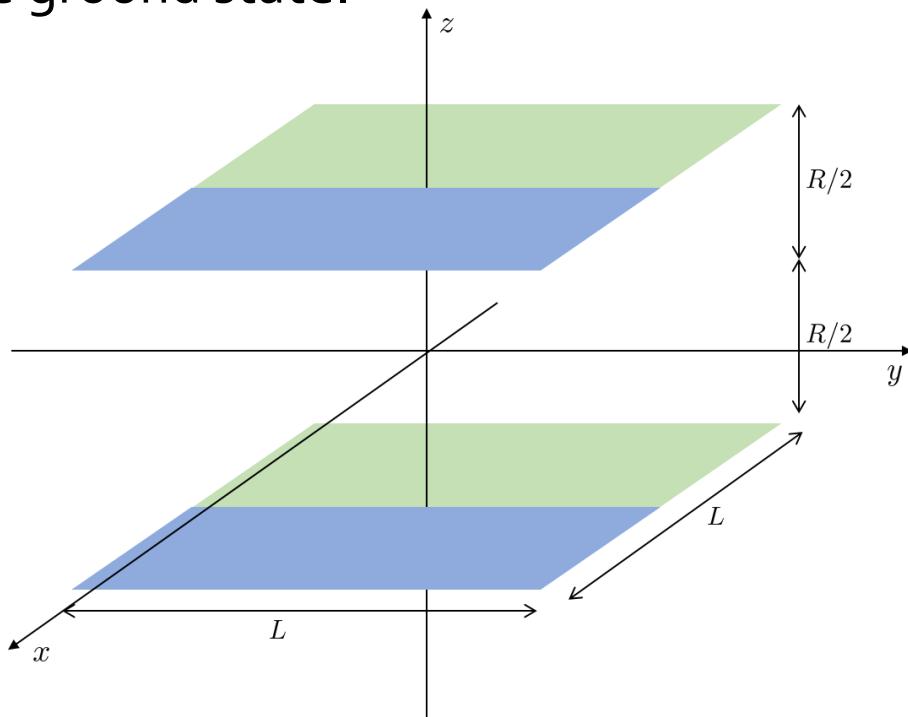


Analytic predictions
about observables
(e.g. $\langle s(r) \rangle_{+-}$)
starting from
first principles ?

G. Delfino, W. Selke and A. Squarcini,
Nucl. Phys. B 958 (2020) 115139.

Particle description

Particles of the bulk field theory in (2+1) dimensions are excitation modes over the ground state.



Particle description

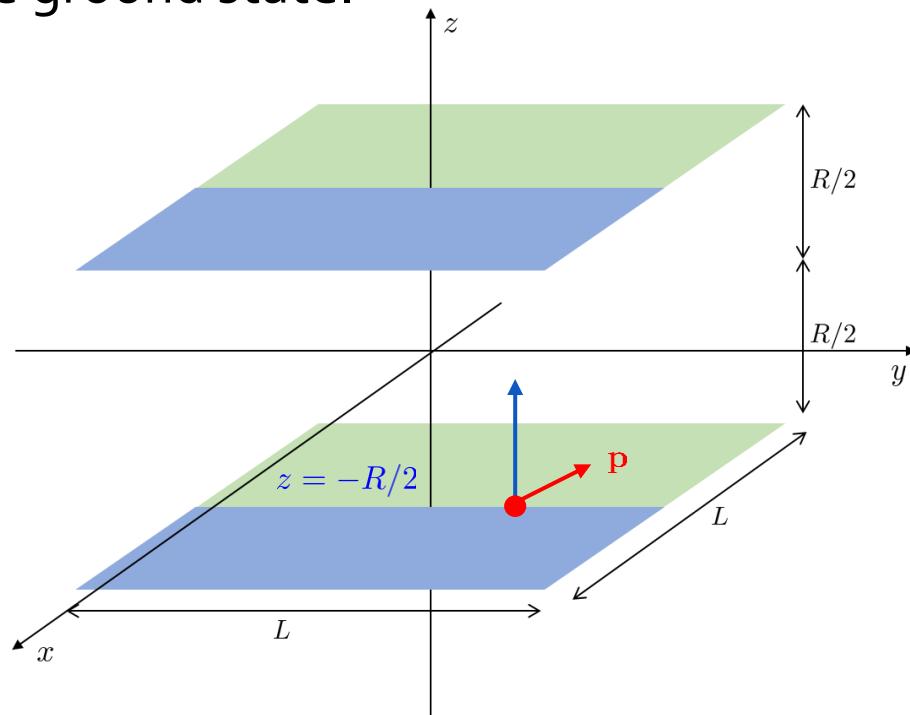
Particles of the bulk field theory in (2+1) dimensions are excitation modes over the ground state.

$$E_{\mathbf{p}} = \sqrt{\mathbf{p}^2 + m^2}$$

$$\mathbf{p} = (p_x, p_y)$$

$$m = 1/\xi$$

lightest particle



Particle description

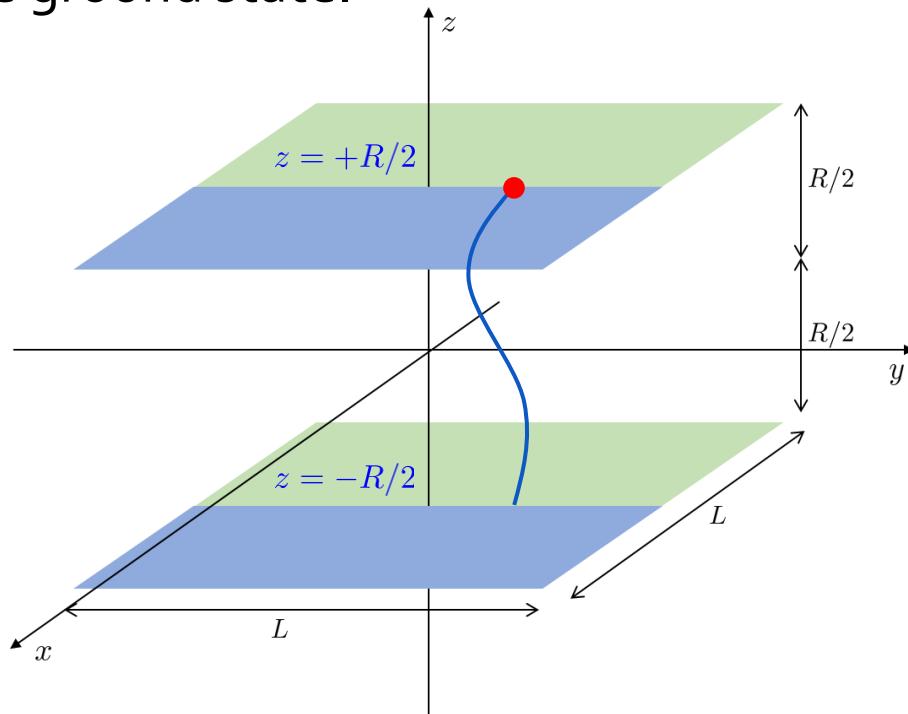
Particles of the bulk field theory in (2+1) dimensions are excitation modes over the ground state.

$$E_{\mathbf{p}} = \sqrt{\mathbf{p}^2 + m^2}$$

$$\mathbf{p} = (p_x, p_y)$$

$$m = 1/\xi$$

lightest particle



Particle description

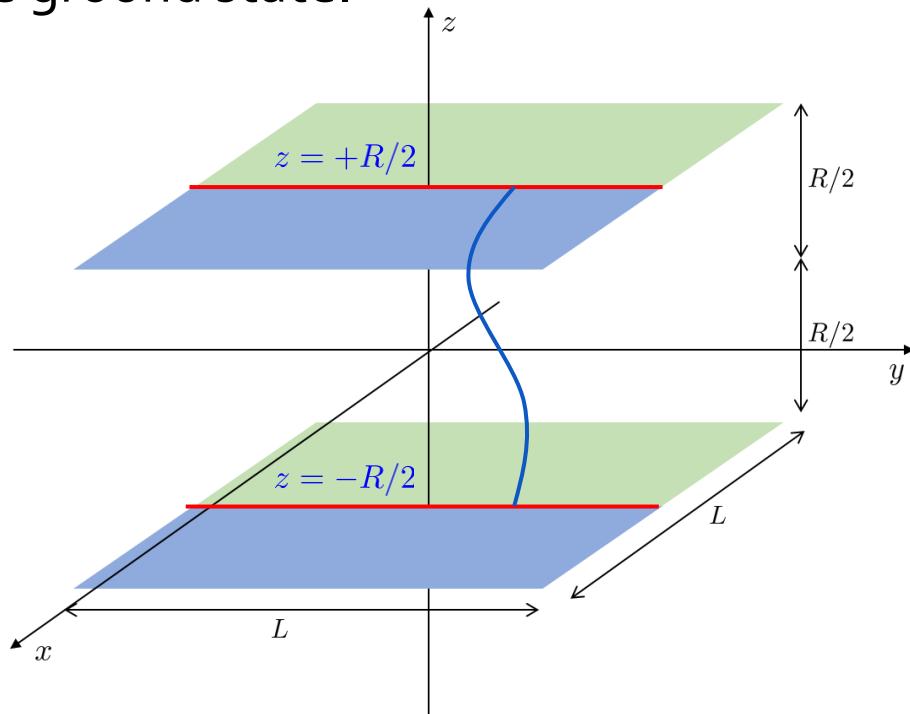
Particles of the bulk field theory in (2+1) dimensions are excitation modes over the ground state.

$$E_{\mathbf{p}} = \sqrt{\mathbf{p}^2 + m^2}$$

$$\mathbf{p} = (p_x, p_y)$$

$$m = 1/\xi$$

lightest particle



Particle description

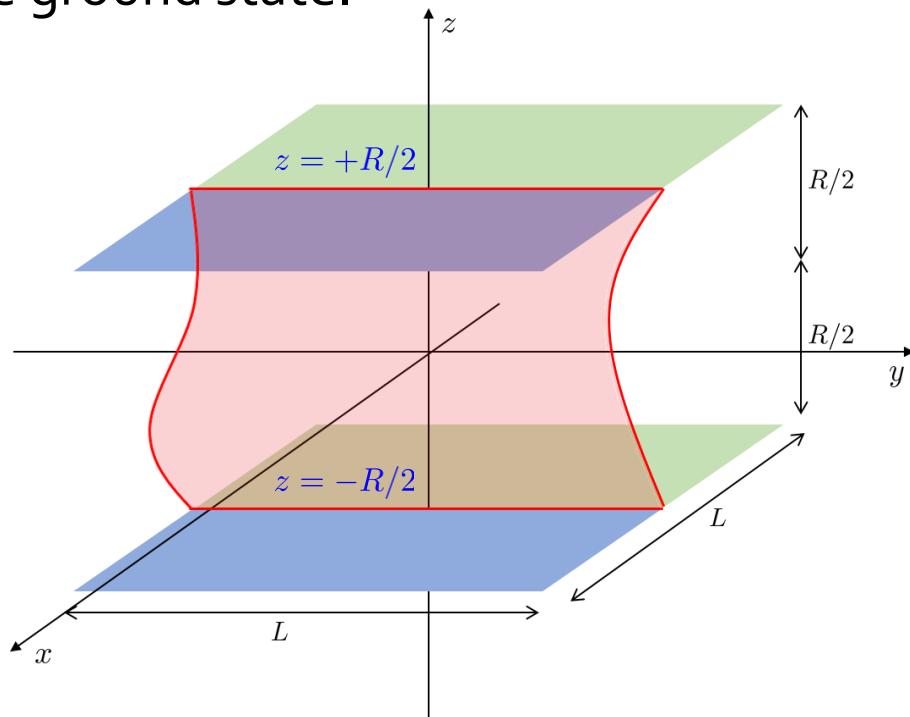
Particles of the bulk field theory in (2+1) dimensions are excitation modes over the ground state.

$$E_{\mathbf{p}} = \sqrt{\mathbf{p}^2 + m^2}$$

$$\mathbf{p} = (p_x, p_y)$$

$$m = 1/\xi$$

lightest particle



The interface is generated by the imaginary time propagation of an excitation (**string**) made of $N \propto L \rightarrow \infty$ particles.

Boundary states

Boundary states $|B(\pm R/2)\rangle$ specify that particles are emitted/absorbed at $z = \pm R/2$.

Expansion over the basis of **asymptotic particle states**:

$\{|\mathbf{p}_1, \dots, \mathbf{p}_N\rangle\}$ with eigenvalues $\sum_{i=1}^N E_{\mathbf{p}_i}, \sum_{i=1}^N \mathbf{p}_i$

normalization $\langle \mathbf{p} | \mathbf{q} \rangle = (2\pi)^2 E_{\mathbf{p}} \delta(\mathbf{p} - \mathbf{q})$

Boundary states

Boundary states $|B(\pm R/2)\rangle$ specify that particles are emitted/absorbed at $z = \pm R/2$.

Expansion over the basis of **asymptotic particle states**:

$\{|\mathbf{p}_1, \dots, \mathbf{p}_N\rangle\}$ with eigenvalues $\sum_{i=1}^N E_{\mathbf{p}_i}$, $\sum_{i=1}^N \mathbf{p}_i$

normalization $\langle \mathbf{p} | \mathbf{q} \rangle = (2\pi)^2 E_{\mathbf{p}} \delta(\mathbf{p} - \mathbf{q})$

$$|B(\pm R/2)\rangle = e^{\pm \frac{R}{2} H} |B(0)\rangle$$

$$= \frac{1}{\sqrt{N!}} \int \prod_{i=1}^N \frac{d\mathbf{p}_i}{(2\pi)^2 E_{\mathbf{p}_i}} f(\mathbf{p}_1, \dots, \mathbf{p}_N) e^{\pm \frac{R}{2} \sum_{i=1}^N E_{\mathbf{p}_i}} \delta \left(\sum_{i=1}^N p_{y,i} \right) |\mathbf{p}_1, \dots, \mathbf{p}_N\rangle + \dots$$

Boundary states

Boundary states $|B(\pm R/2)\rangle$ specify that particles are emitted/absorbed at $z = \pm R/2$.

Expansion over the basis of **asymptotic particle states**:

$\{|\mathbf{p}_1, \dots, \mathbf{p}_N\rangle\}$ with eigenvalues $\sum_{i=1}^N E_{\mathbf{p}_i}, \sum_{i=1}^N \mathbf{p}_i$

normalization $\langle \mathbf{p} | \mathbf{q} \rangle = (2\pi)^2 E_{\mathbf{p}} \delta(\mathbf{p} - \mathbf{q})$

$$|B(\pm R/2)\rangle = e^{\pm \frac{R}{2} H} |B(0)\rangle$$

$$= \frac{1}{\sqrt{N!}} \int \prod_{i=1}^N \frac{d\mathbf{p}_i}{(2\pi)^2 E_{\mathbf{p}_i}} f(\mathbf{p}_1, \dots, \mathbf{p}_N) e^{\pm \frac{R}{2} \sum_{i=1}^N E_{\mathbf{p}_i}} \delta \left(\sum_{i=1}^N p_{y,i} \right) |\mathbf{p}_1, \dots, \mathbf{p}_N\rangle + \dots$$

emission/absorption
amplitude

translation
invariance in
 y -direction

Partition function

$$Z_{+-} = \langle B(R/2) | B(-R/2) \rangle$$

$$= \frac{L}{2\pi} \int \prod_{i=1}^N \frac{d\mathbf{p}_i}{(2\pi)^2 E_{\mathbf{p}_i}} |f(\mathbf{p}_1, \dots, \mathbf{p}_N)|^2 \delta \left(\sum_{i=1}^N p_{y,i} \right) e^{-R \sum_{i=1}^N E_{\mathbf{p}_i}}$$

Partition function

$$Z_{+-} = \langle B(R/2) | B(-R/2) \rangle$$

$$= \frac{L}{2\pi} \int \prod_{i=1}^N \frac{d\mathbf{p}_i}{(2\pi)^2 E_{\mathbf{p}_i}} |f(\mathbf{p}_1, \dots, \mathbf{p}_N)|^2 \delta \left(\sum_{i=1}^N p_{y,i} \right) e^{-R \sum_{i=1}^N E_{\mathbf{p}_i}}$$

low-energy limit
for $R \gg \xi$

Then $f(\mathbf{p}_1, \dots, \mathbf{p}_N) \rightarrow f(0, \dots, 0) \equiv f_0$

Partition function

$$Z_{+-} = \langle B(R/2) | B(-R/2) \rangle$$

$$= \frac{L}{2\pi} \int \prod_{i=1}^N \frac{d\mathbf{p}_i}{(2\pi)^2 E_{\mathbf{p}_i}} |f(\mathbf{p}_1, \dots, \mathbf{p}_N)|^2 \delta \left(\sum_{i=1}^N p_{y,i} \right) e^{-R \sum_{i=1}^N E_{\mathbf{p}_i}}$$

low-energy limit
for $R \gg \xi$

Then $f(\mathbf{p}_1, \dots, \mathbf{p}_N) \rightarrow f(0, \dots, 0) \equiv f_0$

$$Z_{+-} \sim \frac{L |f_0|^2 e^{-RNm}}{(2\pi)^{2(N+1)}} \left(\frac{2\pi}{R} \right)^N \sqrt{\frac{2\pi R}{Nm}}$$

Partition function

$$Z_{+-} = \langle B(R/2) | B(-R/2) \rangle$$

$$= \frac{L}{2\pi} \int \prod_{i=1}^N \frac{d\mathbf{p}_i}{(2\pi)^2 E_{\mathbf{p}_i}} |f(\mathbf{p}_1, \dots, \mathbf{p}_N)|^2 \delta \left(\sum_{i=1}^N p_{y,i} \right) e^{-R \sum_{i=1}^N E_{\mathbf{p}_i}}$$

low-energy limit
for $R \gg \xi$

Then $f(\mathbf{p}_1, \dots, \mathbf{p}_N) \rightarrow f(0, \dots, 0) \equiv f_0$

$$Z_{+-} \sim \frac{L |f_0|^2 e^{-RNm}}{(2\pi)^{2(N+1)}} \left(\frac{2\pi}{R} \right)^N \sqrt{\frac{2\pi R}{Nm}}$$

Interfacial tension = interfacial free energy per unit area

$$\sigma = - \lim_{R \rightarrow \infty} \frac{1}{LR} \ln Z_{+-} = \frac{\kappa}{\xi^2} \quad \text{with } \kappa = \frac{N\xi}{L}$$

Partition function

$$Z_{+-} = \langle B(R/2) | B(-R/2) \rangle$$

$$= \frac{L}{2\pi} \int \prod_{i=1}^N \frac{d\mathbf{p}_i}{(2\pi)^2 E_{\mathbf{p}_i}} |f(\mathbf{p}_1, \dots, \mathbf{p}_N)|^2 \delta \left(\sum_{i=1}^N p_{y,i} \right) e^{-R \sum_{i=1}^N E_{\mathbf{p}_i}}$$

low-energy limit
for $R \gg \xi$

Then $f(\mathbf{p}_1, \dots, \mathbf{p}_N) \rightarrow f(0, \dots, 0) \equiv f_0$

$$Z_{+-} \sim \frac{L |f_0|^2 e^{-RNm}}{(2\pi)^{2(N+1)}} \left(\frac{2\pi}{R} \right)^N \sqrt{\frac{2\pi R}{Nm}}$$

Interfacial tension = interfacial free energy per unit area

$$\sigma = - \lim_{R \rightarrow \infty} \frac{1}{LR} \ln Z_{+-} = \frac{\kappa}{\xi^2} \quad \text{with} \quad \kappa = \frac{N\xi}{L} \quad (\text{universal})$$

Monte Carlo estimate $\kappa = 0.1084(11)$

M. Caselle, M. Hasenbusch and
M. Panero, JHEP 09 (2007) 117.

Partition function

$$Z_{+-} = \langle B(R/2) | B(-R/2) \rangle$$

$$= \frac{L}{2\pi} \int \prod_{i=1}^N \frac{d\mathbf{p}_i}{(2\pi)^2 E_{\mathbf{p}_i}} |f(\mathbf{p}_1, \dots, \mathbf{p}_N)|^2 \delta \left(\sum_{i=1}^N p_{y,i} \right) e^{-R \sum_{i=1}^N E_{\mathbf{p}_i}}$$

low-energy limit
for $R \gg \xi$

Then $f(\mathbf{p}_1, \dots, \mathbf{p}_N) \rightarrow f(0, \dots, 0) \equiv f_0$

$$Z_{+-} \sim \frac{L |f_0|^2 e^{-RNm}}{(2\pi)^{2(N+1)}} \left(\frac{2\pi}{R} \right)^N \sqrt{\frac{2\pi R}{Nm}}$$

Interfacial tension = interfacial free energy per unit area

$$\sigma = - \lim_{R \rightarrow \infty} \frac{1}{LR} \ln Z_{+-} = \frac{\kappa}{\xi^2} \quad \text{with} \quad \kappa = \frac{N\xi}{L} \quad (\text{universal})$$

Monte Carlo estimate $\kappa = 0.1084(11)$

M. Caselle, M. Hasenbusch and
M. Panero, JHEP 09 (2007) 117.

Partition function

$$Z_{+-} = \langle B(R/2) | B(-R/2) \rangle$$

$$= \frac{L}{2\pi} \int \prod_{i=1}^N \frac{d\mathbf{p}_i}{(2\pi)^2 E_{\mathbf{p}_i}} |f(\mathbf{p}_1, \dots, \mathbf{p}_N)|^2 \delta \left(\sum_{i=1}^N p_{y,i} \right) e^{-R \sum_{i=1}^N E_{\mathbf{p}_i}}$$

low-energy limit
for $R \gg \xi$

Then $f(\mathbf{p}_1, \dots, \mathbf{p}_N) \rightarrow f(0, \dots, 0) \equiv f_0$

$$Z_{+-} \sim \frac{L |f_0|^2 e^{-RNm}}{(2\pi)^{2(N+1)}} \left(\frac{2\pi}{R} \right)^N \sqrt{\frac{2\pi R}{Nm}}$$

widely separated
particles

Interfacial tension = interfacial free energy per unit area

$$\sigma = - \lim_{R \rightarrow \infty} \frac{1}{LR} \ln Z_{+-} = \frac{\kappa}{\xi^2} \quad \text{with} \quad \kappa = \frac{N\xi}{L} \quad (\text{universal})$$

Monte Carlo estimate $\kappa = 0.1084(11)$

M. Caselle, M. Hasenbusch and
M. Panero, JHEP 09 (2007) 117.

Order parameter profile

$$G_s(x) \equiv \langle s(x, y, 0) \rangle_{+-} = \frac{1}{Z_{+-}} \langle B(R/2) | s(x, y, 0) | B(-R/2) \rangle$$

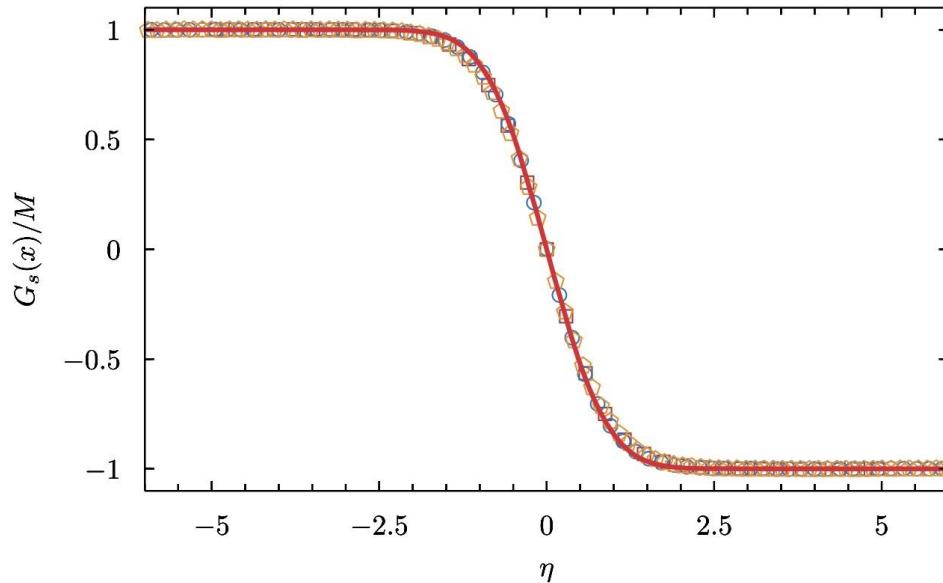
Order parameter profile

$$G_s(x) \sim -M \operatorname{erf}(\eta) \quad \eta = \sqrt{\frac{2}{R\xi}} x$$

Order parameter profile

$$G_s(x) \sim -M \operatorname{erf}(\eta) \quad \eta = \sqrt{\frac{2}{R\xi}} x$$

Confirmed by Monte Carlo simulations for different values of R and $T \lesssim T_c$, with $L \gg R$:

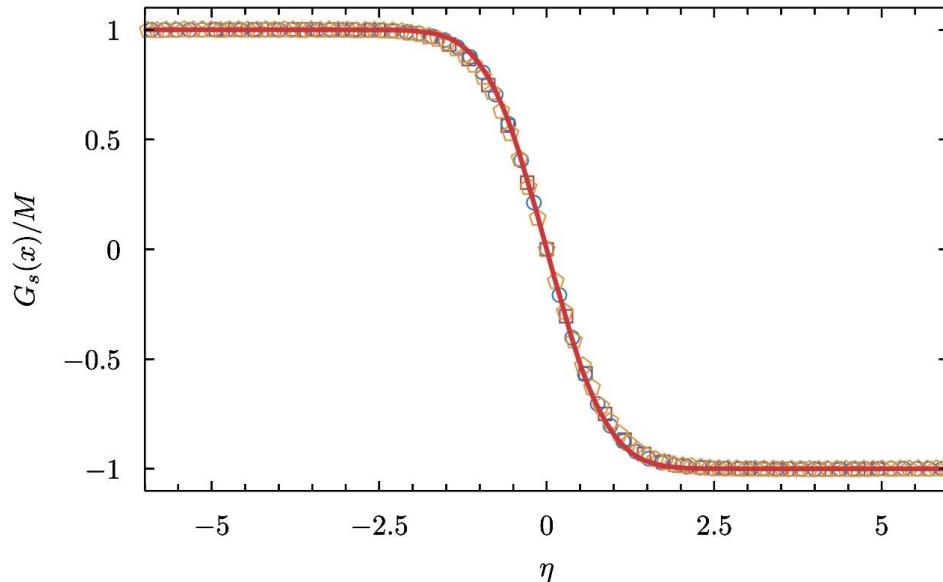


G. Delfino, W. Selke and A. Squarcini,
Nucl. Phys. B 958 (2020) 115139.

Order parameter profile

$$G_s(x) \sim -M \operatorname{erf}(\eta) \quad \eta = \sqrt{\frac{2}{R\xi}} x$$

Confirmed by Monte Carlo simulations for different values of R and $T \lesssim T_c$, with $L \gg R$: (No adjustable parameters)



$$\xi \simeq \xi_0 (T_c - T)^{-\nu}$$

$$\xi_0 \simeq 0.668$$

$$\nu = 0.6310(15)$$

$$T_c \simeq 4.51153$$

M. Caselle, M. Hasenbusch,
J. Phys. A 30 (1997) 4963.

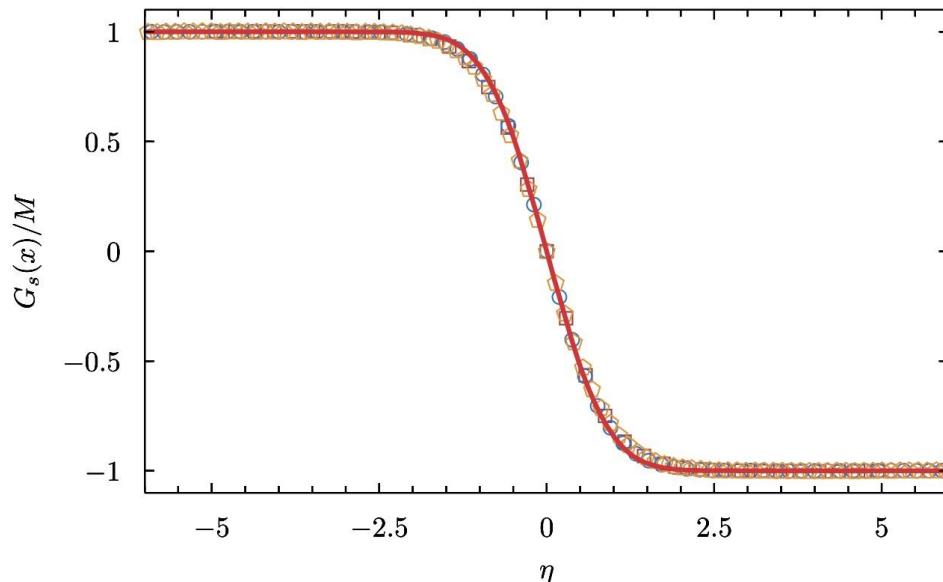
G. Delfino, W. Selke and A. Squarcini,
Nucl. Phys. B 958 (2020) 115139.

Order parameter profile

$$G_s(x) \sim -M \operatorname{erf}(\eta)$$

$$\eta = \sqrt{\frac{2}{R\xi}} x$$

Confirmed by Monte Carlo simulations for different values of R and $T \lesssim T_c$, with $L \gg R$: (No adjustable parameters)



$$\xi \simeq \xi_0 (T_c - T)^{-\nu}$$

$$\xi_0 \simeq 0.668$$

$$\nu = 0.6310(15)$$

$$T_c \simeq 4.51153$$

M. Caselle, M. Hasenbusch,
J. Phys. A 30 (1997) 4963.

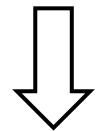
First analytical determination and numerical confirmation of the profile.

G. Delfino, W. Selke and A. Squarcini,
Nucl. Phys. B 958 (2020) 115139.

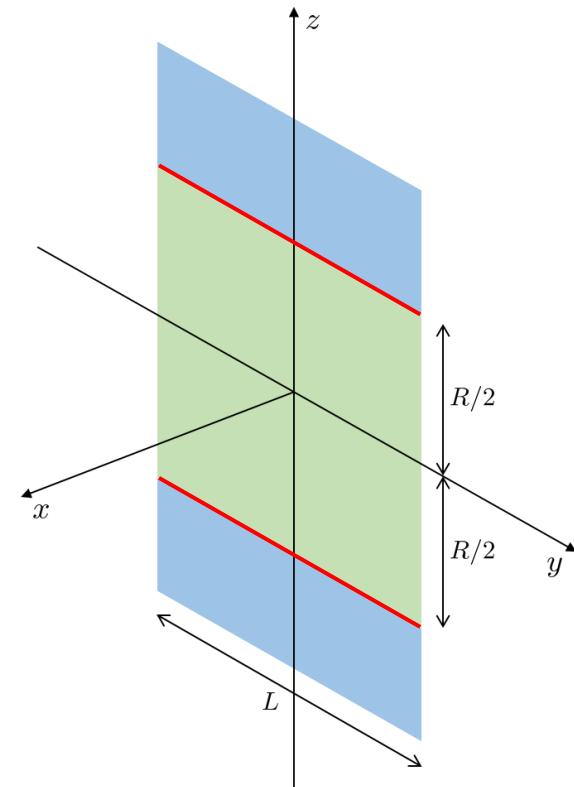
Interface in presence of a wall

Boundary conditions on the wall $x = 0$:

$$s_i = \begin{cases} +1 & |z| < R/2 \\ -1 & |z| > R/2 \end{cases}$$



$$\lim_{x \rightarrow +\infty} \langle s(x, y, 0) \rangle_{+-} = \begin{cases} -M & R < \infty \\ +M & R = \infty \end{cases}$$

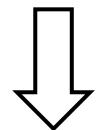


G. Delfino, M. Sorba and A. Squarcini,
Nucl. Phys. B 967 (2021) 115396.

Interface in presence of a wall

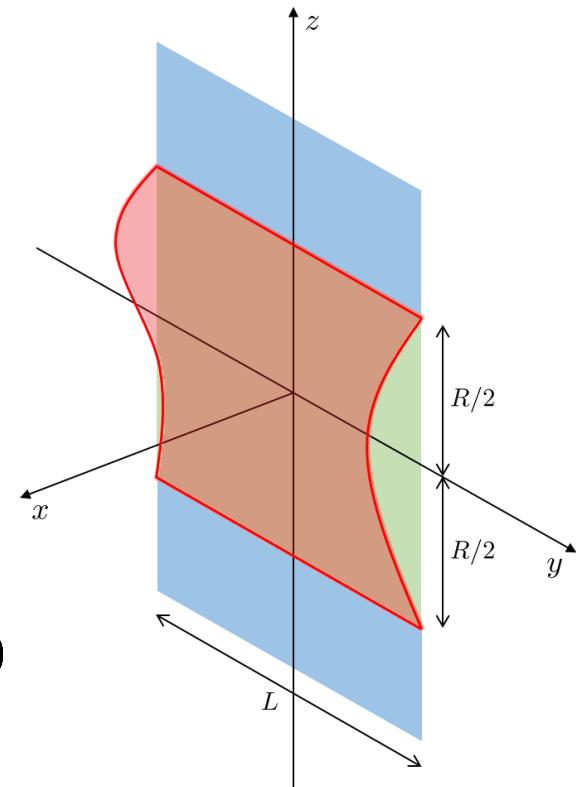
Boundary conditions on the wall $x = 0$:

$$s_i = \begin{cases} +1 & |z| < R/2 \\ -1 & |z| > R/2 \end{cases}$$



$$\lim_{x \rightarrow +\infty} \langle s(x, y, 0) \rangle_{+-} = \begin{cases} -M & R < \infty \\ +M & R = \infty \end{cases}$$

Interface constrained in the region $x > 0$
with an average distance from the wall
growing with R .

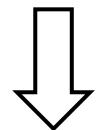


G. Delfino, M. Sorba and A. Squarcini,
Nucl. Phys. B 967 (2021) 115396.

Interface in presence of a wall

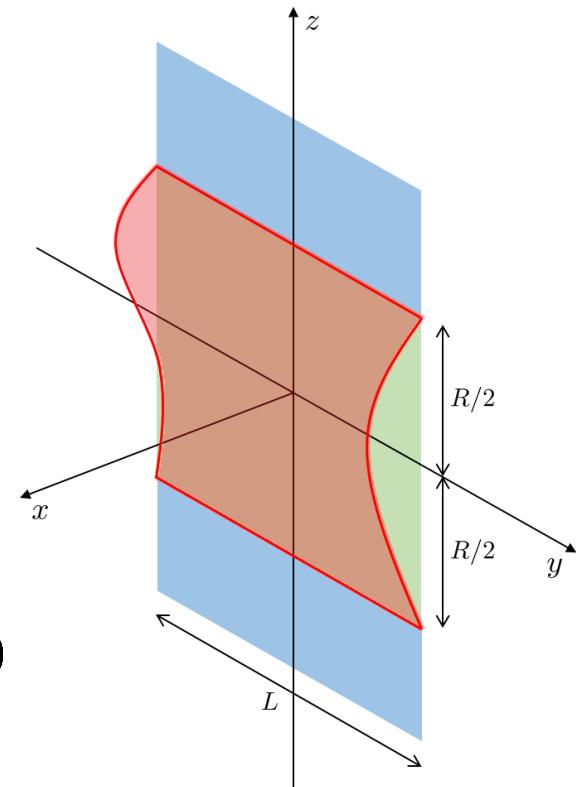
Boundary conditions on the wall $x = 0$:

$$s_i = \begin{cases} +1 & |z| < R/2 \\ -1 & |z| > R/2 \end{cases}$$



$$\lim_{x \rightarrow +\infty} \langle s(x, y, 0) \rangle_{+-} = \begin{cases} -M & R < \infty \\ +M & R = \infty \end{cases}$$

Interface constrained in the region $x > 0$
with an average distance from the wall
growing with R .



Field theoretical description ?

G. Delfino, M. Sorba and A. Squarcini,
Nucl. Phys. B 967 (2021) 115396.

Partition function

- Particle description
 - Boundary states
 - Large R limit (small momenta limit)
- 
- same as before

Partition function

- Particle description
- Boundary states
- Large R limit (small momenta limit)

same as before

$$f(\mathbf{p}_1, \dots, \mathbf{p}_N) \rightarrow f_0 \prod_{i=1}^N p_{x,i}$$

None of the particles can stay along the **impenetrable wall** $x = 0$ (i.e. $f = 0$ when $\exists p_{x,i} = 0$).

Partition function

- Particle description
- Boundary states
- Large R limit (small momenta limit)

same as before

$$f(\mathbf{p}_1, \dots, \mathbf{p}_N) \rightarrow f_0 \prod_{i=1}^N p_{x,i}$$

None of the particles can stay along the **impenetrable wall** $x = 0$ (i.e. $f = 0$ when $\exists p_{x,i} = 0$).

$$Z_{+-} \sim \frac{L |f_0|^2 e^{-R N m}}{(2\pi)^{2(N+1)}} \left(\frac{2\pi}{R^2} \right)^N \sqrt{\frac{2\pi R}{Nm}}$$

Partition function

- Particle description
- Boundary states
- Large R limit (small momenta limit)

same as before

$$f(\mathbf{p}_1, \dots, \mathbf{p}_N) \rightarrow f_0 \prod_{i=1}^N p_{x,i}$$

None of the particles can stay along the **impenetrable wall** $x = 0$ (i.e. $f = 0$ when $\exists p_{x,i} = 0$).

$$Z_{+-} \sim \frac{L |f_0|^2 e^{-R N m}}{(2\pi)^{2(N+1)}} \left(\frac{2\pi}{R^2} \right)^N \sqrt{\frac{2\pi R}{Nm}}$$

Interfacial tension:

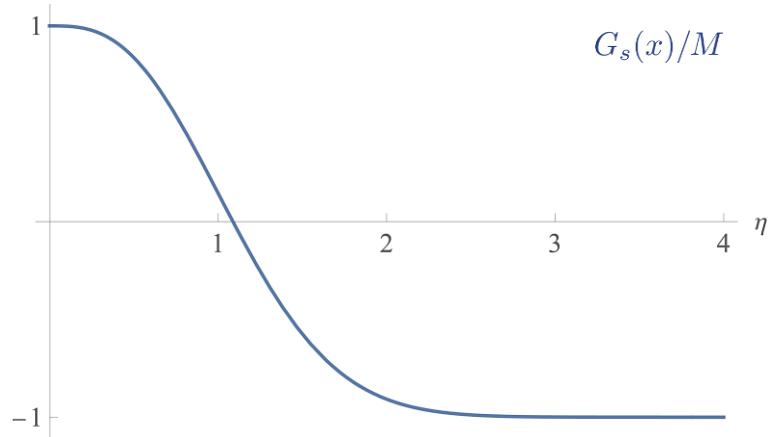
$$\sigma = - \lim_{R \rightarrow \infty} \frac{1}{LR} \ln Z_{+-} = \frac{\kappa}{\xi^2}$$

same as in absence of the wall

Order parameter profile

$$G_s(x) \sim M + 2M \left\{ \frac{2}{\sqrt{\pi}} \eta e^{-\eta^2} - \text{erf}(\eta) \right\}$$

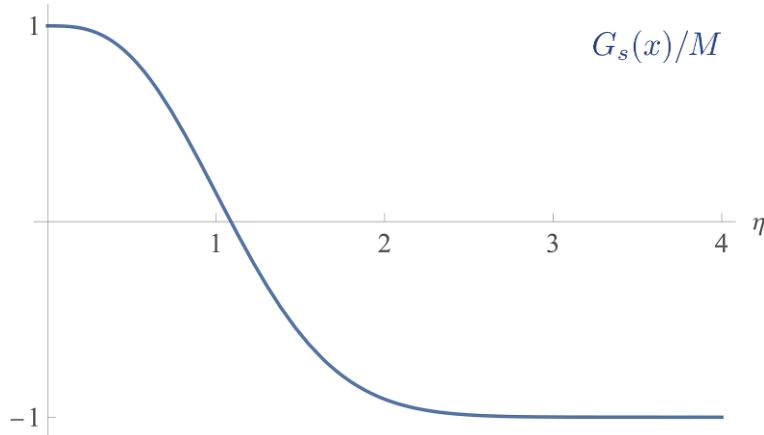
$$\eta = \sqrt{\frac{2}{R\xi}} x$$



Order parameter profile

$$G_s(x) \sim M + 2M \left\{ \frac{2}{\sqrt{\pi}} \eta e^{-\eta^2} - \text{erf}(\eta) \right\}$$

$$\eta = \sqrt{\frac{2}{R\xi}} x$$

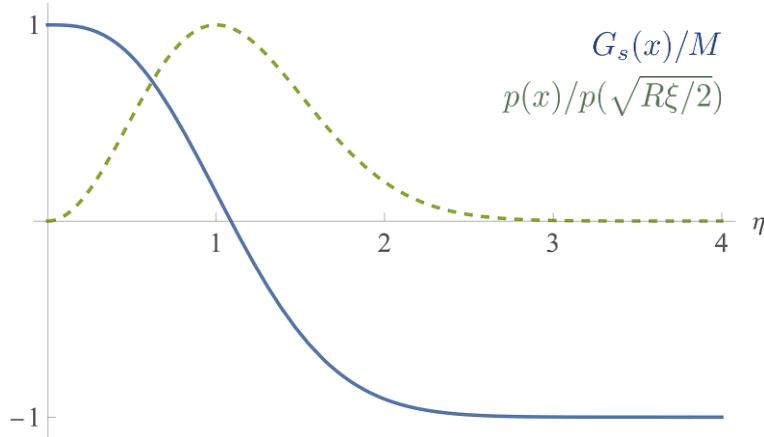


At leading order in the large R expansion, the interface is a **sharp separation** between two pure phases $\pm M$.

Order parameter profile

$$G_s(x) \sim M + 2M \left\{ \frac{2}{\sqrt{\pi}} \eta e^{-\eta^2} - \text{erf}(\eta) \right\}$$

$$\eta = \sqrt{\frac{2}{R\xi}} x$$



At leading order in the large R expansion, the interface is a **sharp separation** between two pure phases $\pm M$.

Probability density of intersection at a point $x = u$:

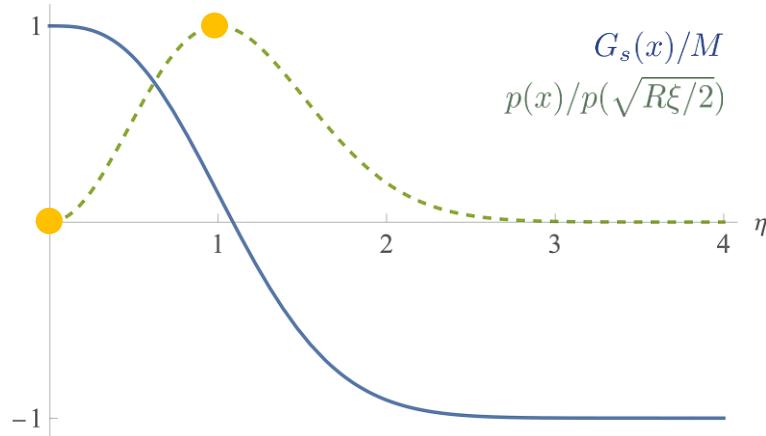
$$\overline{s(x)} = \int_0^{+\infty} du p(u) s(x|u), \quad s(x|u) = M \theta(u - x) - M \theta(x - u)$$

$$p(x) = 4 \sqrt{\frac{2}{\pi R\xi}} \eta^2 e^{-\eta^2}$$

Order parameter profile

$$G_s(x) \sim M + 2M \left\{ \frac{2}{\sqrt{\pi}} \eta e^{-\eta^2} - \text{erf}(\eta) \right\}$$

$$\eta = \sqrt{\frac{2}{R\xi}} x$$



At leading order in the large R expansion, the interface is a **sharp separation** between two pure phases $\pm M$.

Probability density of intersection at a point $x = u$:

$$\overline{s(x)} = \int_0^{+\infty} du p(u) s(x|u), \quad s(x|u) = M \theta(u - x) - M \theta(x - u)$$

$$p(x) = 4 \sqrt{\frac{2}{\pi R \xi}} \eta^2 e^{-\eta^2}$$

- Correct normalization
- Impenetrability of the wall
- Average distance $x \propto \sqrt{R}$

Binding transition

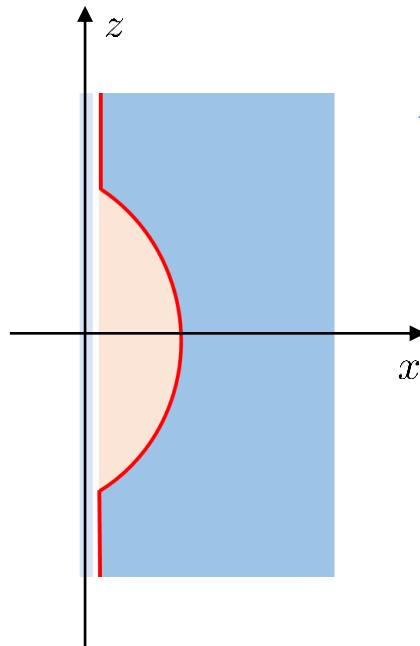
The wall at $x = 0$ contributes to the theory with a boundary Hamiltonian $h \int dy dz \Phi_B(0, y, z)$.

It is possible to tune h so that the wall-interface interaction becomes **very attractive**:

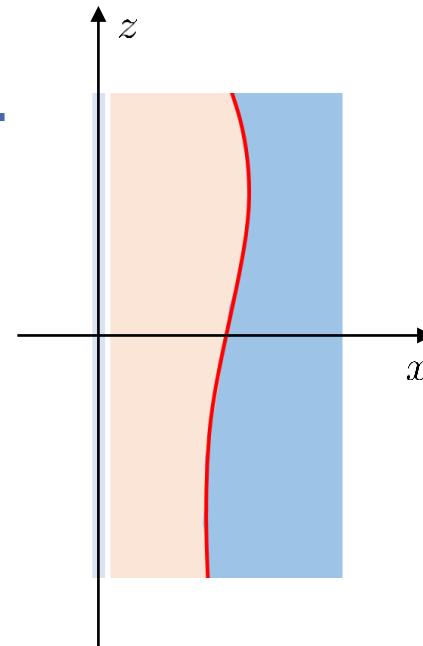
Binding transition

The wall at $x = 0$ contributes to the theory with a boundary Hamiltonian $h \int dy dz \Phi_B(0, y, z)$.

It is possible to tune h so that the wall-interface interaction becomes **very attractive**:



Binding regime

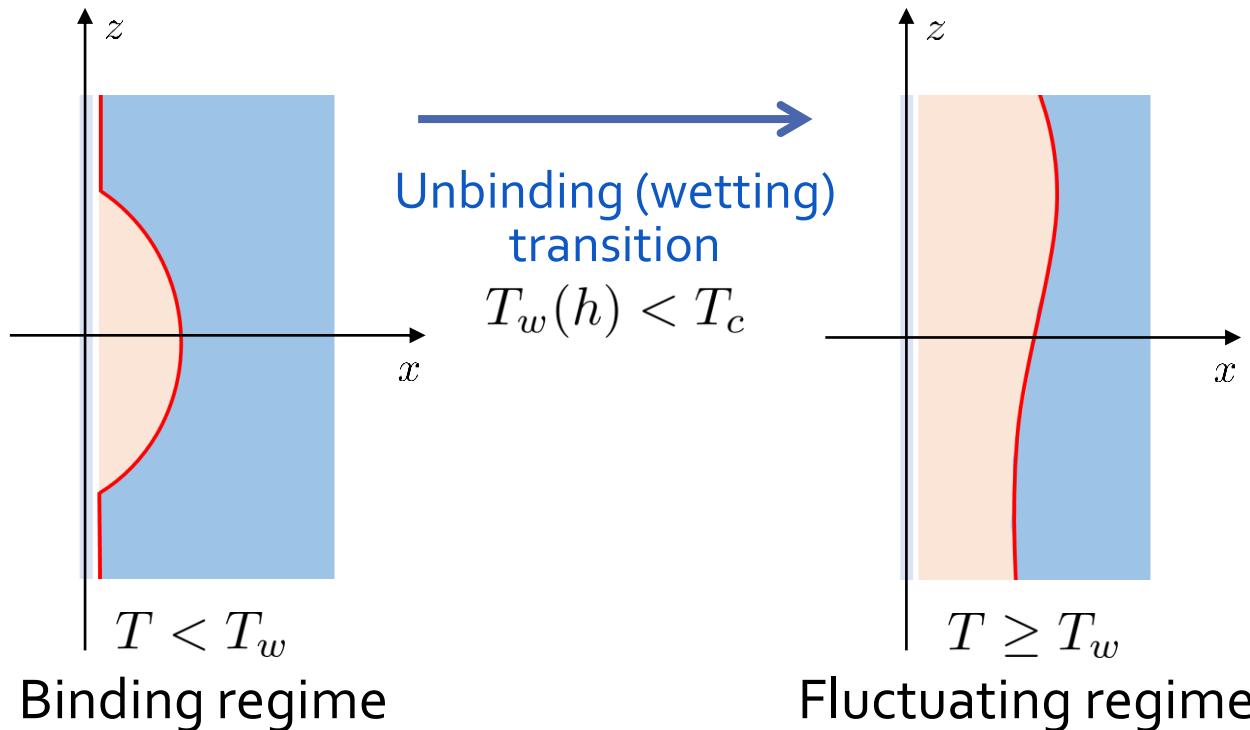


Fluctuating regime

Binding transition

The wall at $x = 0$ contributes to the theory with a boundary Hamiltonian $h \int dy dz \Phi_B(0, y, z)$.

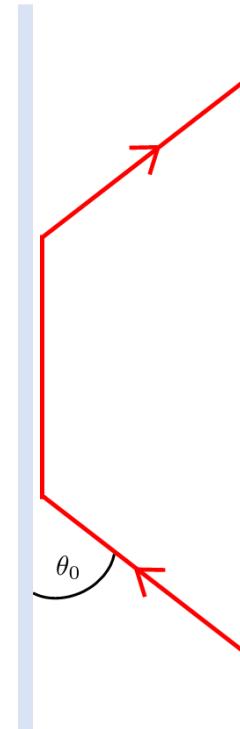
It is possible to tune h so that the wall-interface interaction becomes **very attractive**:



Scattering theory

Interaction of a particle with the wall:

- Rapidity parametrization $\begin{cases} E = m \cosh \beta \\ |\mathbf{p}| = m \sinh \beta \end{cases}$
- Stable bound state for $\beta = i\theta_0$, $\theta_0 \in (0, \pi)$ with energy $E = m \cos \theta_0 < m$



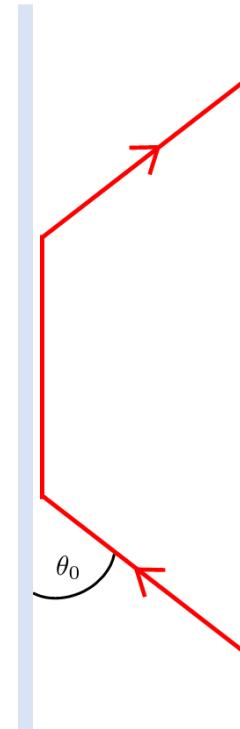
Scattering theory

Interaction of a particle with the wall:

- Rapidity parametrization $\begin{cases} E = m \cosh \beta \\ |\mathbf{p}| = m \sinh \beta \end{cases}$
- Stable bound state for $\beta = i\theta_0$, $\theta_0 \in (0, \pi)$ with energy $E = m \cos \theta_0 < m$

Total energy of the system per unit length in the bound regime:

$$\tilde{e} = e + \frac{N}{L} m \cos \theta_0 = e + \sigma \cos \theta_0$$



Scattering theory

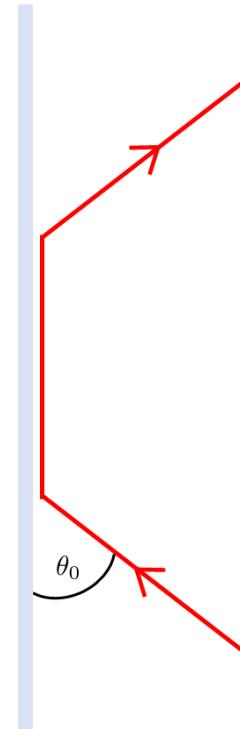
Interaction of a particle with the wall:

- Rapidity parametrization $\begin{cases} E = m \cosh \beta \\ |\mathbf{p}| = m \sinh \beta \end{cases}$
- Stable bound state for $\beta = i\theta_0$, $\theta_0 \in (0, \pi)$ with energy $E = m \cos \theta_0 < m$

Total energy of the system per unit length in the bound regime:

$$\tilde{e} = e + \frac{N}{L} m \cos \theta_0 = e + \sigma \cos \theta_0$$

\Rightarrow unbinding (wetting) at $\boxed{\theta_0(T_w) = 0}$



Scattering theory

Interaction of a particle with the wall:

- Rapidity parametrization $\begin{cases} E = m \cosh \beta \\ |\mathbf{p}| = m \sinh \beta \end{cases}$
- Stable bound state for $\beta = i\theta_0$, $\theta_0 \in (0, \pi)$ with energy $E = m \cos \theta_0 < m$

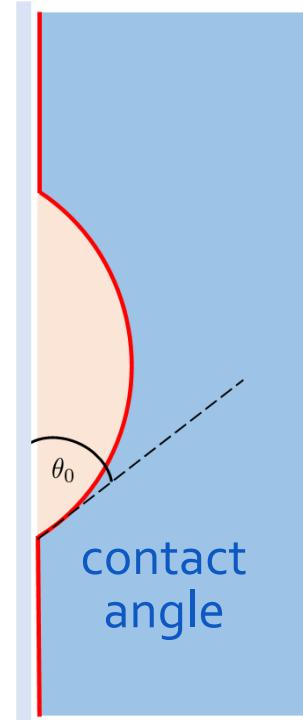
Total energy of the system per unit length in the bound regime:

$$\tilde{e} = e + \frac{N}{L} m \cos \theta_0 = e + \sigma \cos \theta_0$$

Equilibrium condition
for a liquid drop on a surface

T. Young (1805).

\Rightarrow unbinding (wetting) at $\boxed{\theta_0(T_w) = 0}$



Scattering theory

Interaction of a particle with the wall:

- Rapidity parametrization $\begin{cases} E = m \cosh \beta \\ |\mathbf{p}| = m \sinh \beta \end{cases}$
- Stable bound state for $\beta = i\theta_0$, $\theta_0 \in (0, \pi)$ with energy $E = m \cos \theta_0 < m$

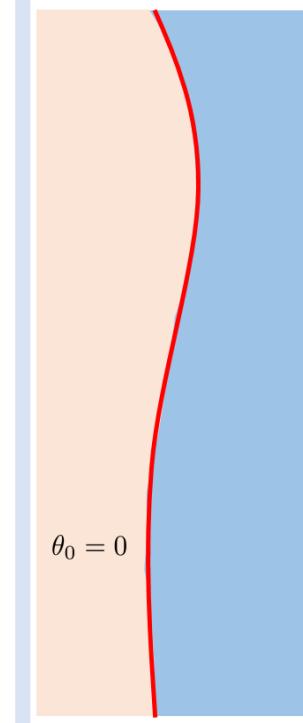
Total energy of the system per unit length
in the bound regime:

$$\tilde{e} = e + \frac{N}{L} m \cos \theta_0 = e + \sigma \cos \theta_0$$

Equilibrium condition
for a liquid drop on a surface

T. Young (1805).

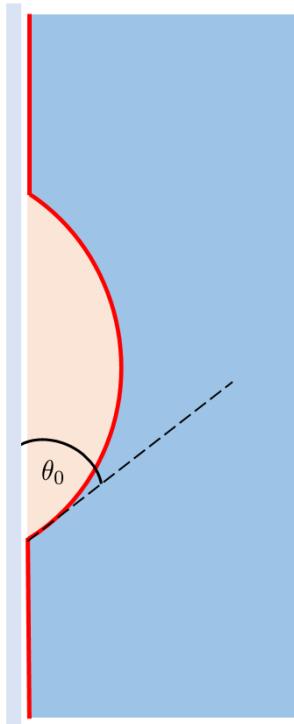
\Rightarrow unbinding (wetting) at $\boxed{\theta_0(T_w) = 0}$



Exponent α_S

Close to the wetting transition point $T \rightarrow T_w^-$
the contact angle vanishes according to:

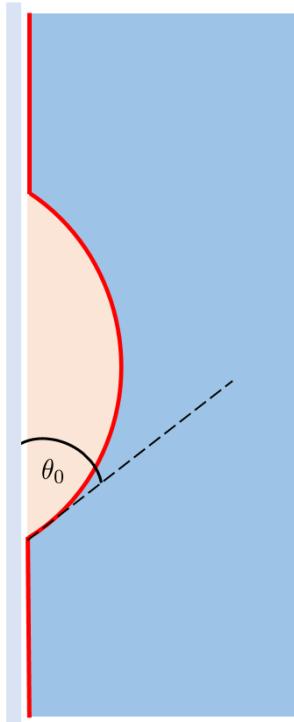
$$\theta_0^2 \propto (T_w - T)^{2-\alpha_S}$$



Exponent α_S

Close to the wetting transition point $T \rightarrow T_w^-$
the contact angle vanishes according to:

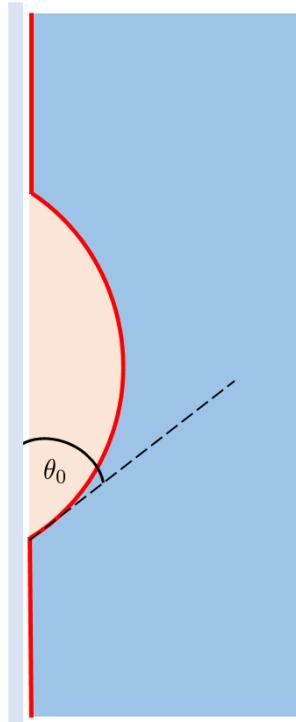
$$\theta_0^2 \propto (T_w - T)^{2-\alpha_S}$$



The pole of the scattering amplitude slides through a branch point at $E = m$, while θ_0 changes sign:

$$\theta_0 \propto (T_w - T)$$

Exponent α_S



Close to the wetting transition point $T \rightarrow T_w^-$
the contact angle vanishes according to:

$$\theta_0^2 \propto (T_w - T)^{2-\alpha_S}$$

The pole of the scattering amplitude slides
through a branch point at $E = m$, while θ_0
changes sign:

$$\theta_0 \propto (T_w - T) \quad \Rightarrow \boxed{\alpha_S = 0}$$

in agreement with Monte Carlo simulations.

K. Binder, D.P. Landau and D.M. Kroll,
Phys. Rev. Lett. 56 (1986) 2272.

Summary

- ❖ Formulation of a field theory of interfaces near criticality, giving analytic universal predictions with no need of free parameters.
- ❖ Interfacial tension, magnetization profile, passage probability density from the definition of the partition function.
- ❖ Implementing in momentum space the presence of an impenetrable wall.
- ❖ Binding transition and its key features using the particle description of the bulk field theory.

**Thank you
for the attention !**

msorba@sissa.it