Parity Doubling in Cold and Dense QCD

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Outline

- I. Introduction to chiral mixing via WZW
- II. Parity doubling of V&AV mesons, spectral functions, dilepton production rates
- III. In-medium widths of vector mesons
 - i. rho-N resonances
 - ii. Kaon and anti-kaon in dense matter

Why chiral mixing?

- Q. Do we see any signal of chiral symmetry restoration in dilepton measurement?
- Light vector mesons change their properties in hot/dense matter --- χ -sym. restoration?
- □Strategy: vector and axial-vector states
- Axial-vector mesons can show up in vector spectrum in a medium!

<VV> \leftarrow chiral mixing \rightarrow <AA>

My fingers crossed, FAIR/SIS/NICA/J-PARC/RHIC-BES!



[Domokos, Harvey ('07)] Holographic approach at finite μ B $S_{\text{4dim}} = \int d^4x \left| \frac{1}{2} \left(\partial_\mu \pi \right)^2 - \frac{1}{2} m_\pi^2 \pi^2 - \frac{1}{4} \left(\rho_{\mu\nu} \right)^2 - \frac{1}{4} \left(a_{\mu\nu} \right)^2 \right|^2$ $+\frac{1}{2}m_{\rho}^{2}\rho_{\nu}^{2}+\frac{1}{2}m_{a}^{2}a_{\mu}^{2}+C\epsilon^{ijk}\left(\rho_{i}\partial_{j}a_{k}+a_{i}\partial_{j}\rho_{k}\right)$ $p_0^2 - |\vec{p}|^2 = \frac{1}{2} \left[m_\rho^2 + m_{a_1}^2 \pm \sqrt{(m_{a_1}^2 - m_\rho^2)^2 + 16C^2 |\vec{p}|^2} \right]$

 ρ meson

 a_1 meson



Spectral function: Not BW 10⁰ longitudinal transverse longitudinal transv: C=0.5 GeV average 10⁻¹ transv: C=1 GeV ImG_V [GeV²] 2.5 p₀ [GeV] 2 10⁻² 1.5 1 10⁻³ 0.5 0 0.5 1.5 2 0 p [GeV] 10-4 0.5 1.5 2 0 s^{1/2} [GeV]

C = 1 GeV, 3-momentum p = 0.5 GeV
1 bump of transv. rho, 1 bump of transv. a1

[Harada, CS ('09)] Chiral mixing induced from WZW

Wess-Zumino-Witten term [Kaiser, Meissner ('90)]

$$\mathcal{L}_{\omega\rho a_{1}} = g_{\omega\rho a_{1}} \epsilon^{\mu\nu\lambda\sigma} \omega_{\mu} \left[\partial_{\nu} V_{\lambda} \cdot A_{\sigma} + \partial_{\nu} A_{\lambda} \cdot V_{\sigma} \right]$$
$$\omega_{0} \rangle = g_{\omega NN} \cdot n_{B} / m_{\omega}^{2} \qquad C = g_{\omega\rho a_{1}} \cdot g_{\omega NN} \cdot \frac{n_{B}}{m_{\omega}^{2}}$$

 $\Box Mixing strength: C = 0.1 GeV at \rho_{0}$

- AdS/QCD \rightarrow C = 1 GeV at $\rho_0 \rightarrow$ vector cond.!?
- Why so large? --- higher-lying states in large Nc cf. VMD in SS $C_{hQCD} \sim C_{\omega\rho a_1} + \sum C_{\omega^n \rho a_1}$

Weak mixing ... No impact?

A missing piece: χ sym. restoration

 $<AA> \rightarrow <VV>$

CS, 2019

Chiral restoration vs. mixing

Dispersion relations for small 3-momenta

$$p_0^2 \simeq m_{a_1,\rho}^2 + \left(1 \pm \frac{4C^2}{m_{a_1}^2 - m_{\rho}^2}\right) \bar{p}^2$$

The mixing effect will be enhanced as δ m decreases!

- \geq In-medium δ m
- ≻In-medium mixing C
- ← Quark-nucleon hybrid model[NS: Marczenko et al. (19,20)]

Mass difference vs. mixing : T=50 MeV



Vector-current correlator

$$G_V^L = \left(\frac{g_{\rho}}{m_{\rho}}\right)^2 \frac{-s}{D_V}, \quad G_V^T = \left(\frac{g_{\rho}}{m_{\rho}}\right)^2 \frac{-sD_A + 4C^2\bar{p}^2}{D_V D_A - 4C^2\bar{p}^2},$$

 $D_{V,A} = s - m_{\rho,a_1}^2 + im_{\rho,a_1}\Gamma_{\rho,a_1}(s),$

□m and Γ: *in-medium* masses and widths
□Strategy of an illustrative computation:

- Modify only mass and width of axial-vector states.
- Set G_A equal to G_V at CSR, according to

 $\label{eq:constraint} \Gamma_a1 = \ \Gamma(a1 {\rightarrow} \rho \ \pi) + \delta \ \Gamma(f_pi) {\rightarrow} \ \Gamma_\rho$

Spectral function at T = 50 MeV



(top) chiral restoration (bottom) no restoration--- longitudinal --- transverse --- average







In-medium broadening

$\Gamma_v, vac \rightarrow \Gamma_v, med(\rho)$





ϕ meson in nuclear matter

 \Box No ϕ N resonances, but the kaon cloud.

Generation Kaplan, Kelson (86)

$$m_{K}^{*} = \left[m_{K}^{2} - a_{K}\rho_{S} + (b_{K}\rho)^{2}\right]^{1/2} + b_{K}\rho,$$

$$* \left[a_{K}^{2} - a_{K}\rho_{S} + (b_{K}\rho)^{2}\right]^{1/2} + b_{K}\rho,$$

$$m_{\bar{K}}^* = \left[m_{K}^2 - a_{\bar{K}}\rho_{S} + (b_{K}\rho)^2\right]^{1/2} - b_{K}\rho,$$

$$b_K = 3/(8f_\pi^2)$$
 $a_K = a_{\bar{K}} = \Sigma_{KN}/f_\pi^2$

Li, Lee, Brown (97): kaon production in Ni+Ni at 1 & 1.8 A GeV $a_K \approx 0.22 \text{ GeV}^2 \text{ fm}^3$ and $a_{\bar{K}} \approx 0.45 \text{ GeV}^2 \text{ fm}^3$





[CS, in preparation]

Spectral function of ϕ meson

□Chiral mixing ≈ C x 3-momenta

□Spectral function at chiral crossover with p = 0.1, 0.5, 1.0 GeV



□Binning of DL data in momenta > 0.5 GeV



[CS, in preparation] Int.over p > 0.5 GeV Dilepton rates at T=50 MeV



Summary

Chiral sym. restoration in cold dense matter

- Chiral mixing induced by WZW, exists at any *ρ* !
- Clear structural change in spectra/dilepton rates
- Big discovery potential at FAIR/NICA/RHIC-BES

Coarse-grained approach

➢ Mixing strength from lattice Q₂CD, FRG, AdS/QFT

Parity doubling of nucleons in neutron stars

- Parity doublet model with nucleons and deltas
- Hadronic picture w/ CSR compatible with all modern constraints! [Marczenko et al. , ApjL (22)]

Backup

Low-energy theorem at T \neq 0

$$G_V^{\mu\nu}(T) = (1 - \epsilon) G_V^{\mu\nu}(0) + \epsilon G_A^{\mu\nu}(0)$$
$$G_A^{\mu\nu}(T) = (1 - \epsilon) G_A^{\mu\nu}(0) + \epsilon G_V^{\mu\nu}(0)$$

[Dey, Eletsky and loffe (90)]

[finite *ρ* : Krippa (98)]



 $\epsilon = \frac{T^2}{6F^2}$



ε →1/2: chiral restoration? NO! Higher T: reducing *π* ρ a1 int.: 2 → 1 bump Chiral mixing ≈ 0.06 mpi at Tc [Harada, CS, Weise (08)]

[Sakai and Sugimoto (2005)] VDM from holography

Infinite tower of vector mesons

$$F(q^2) = \sum_{n=0}^{\infty} \frac{g_{\rho_n} \cdot g_{\rho_n \pi \pi}}{m_{\rho_n}^2 - q^2} \stackrel{q^2 \to 0}{\to} 1$$

Approximately saturated by the lowest 4

n	0	1	2
PDG	776	1465	1720
SS	776	1607	2435

F(0) = 1.31 - 0.35 + 0.05 - 0.01 = 1.00

[Harada et al. (2006)] From large Nc to Nc=3

□Vector mesons integrated out except the lowest → effective Lagrangian of π and ρ
 □Power counting, loop corrections ≈ 1/Nc corr.

✓ a ≈ 2 cf. SS a ≈ 1

✓ KSRF II O.K.

 $\frac{m_{\rho}^2}{g_{\rho\pi\pi}^2 F_{\pi}^2} \bigg|_{ss} \simeq 3.0, \qquad \left(\frac{m_{\rho}^2}{g_{\rho\pi\pi}^2 F_{\pi}^2}\bigg|_{exp} \simeq 2.0\right)$ 1/Nc corrections

[Friman, Pirner (97)]

p meson self-energy



$$\begin{split} \Sigma_{\rho}^{R}(\omega, q) &= \frac{4}{3} \frac{f_{N^{*}N\rho}^{2}}{m_{\rho}^{2}} F(q^{2}) q^{2} \rho_{B} \frac{(\varepsilon_{q}^{N^{*}} - m_{N})}{\omega^{2} - (\varepsilon_{q}^{N^{*}} - m_{N})^{2}} \\ &+ \frac{2}{5} \frac{f_{\Delta N\rho}^{2}}{m_{\rho}^{2}} F(q^{2}) q^{2} \rho_{B} \frac{(\varepsilon_{q}^{\Delta} - m_{N})}{\omega^{2} - (\varepsilon_{q}^{\Delta} - m_{N})^{2}}, \end{split}$$

where

$$\varepsilon_q^{N^\star} = \sqrt{q^2 + m_{N^\star}} - \frac{i}{2} \Gamma_{N^\star}, \ \varepsilon_q^{\varDelta} = \sqrt{q^2 + m_{\varDelta}} - \frac{i}{2} \Gamma_{\varDelta},$$

[Hohler, Rapp ('14,'16)]

From low T to high T



Weinberg SRs [Weinberg ('67); Kapusta, Shuryak ('94)]

□Vector SF & ansatz for a1 mass and width

✓ Reduction of a1 mass, width broadening

✓ Role of higher-lying states: ρ' , a1', ...