

Hexagons, QSL and the Bottom bridge

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Summary

- Hexagons are a powerful tool for computing structure constants.
- As it happens for the anomalous dimensions finite size effects (Wrapping) are difficult to control
- We will mainly focus on a particular contribution to the hexagons (Bottom bridge)

Overview

- $N=4$ Super Yang-Mills (SYM)
- A Bit of Quantum Spectral Curve (QSC)
- Hexagons, Wrapping & some conjectures
- Results (5-loops from QSC)
- Conclusions

Integrability in $N=4$ SU(N) SYM

- $d=4$ interacting CFT. Composed of:
 $\{A^\mu, \bar{\Phi}^I, \Psi_\alpha^a, \bar{\Psi}_{\dot{\alpha}}{}^a\}$ all transforming in the Adjoint $SU(N)$
1 6 4 4
- Theory is dual to String theory on $AdS_5 \times S^5$
- Planar limit ($N \gg 0$) the theory is integrable
- Supersymmetry + conformal invariance \Rightarrow $PSU(2,2|4)$

γ and C_{ijk}

- Given conformal invariance, we are interested into two quantities anomalous dimensions & structure const.

- $O(x) \xrightarrow{x \rightarrow \lambda x} \lambda^{-\Delta} O(\lambda x)$ scaling dimension ↑

if op is not protected this gets corrections!

- OPE expansion

$$O_i(x) O_j(y) \approx \sum_k C_{ijk} g(\Delta, x-y) O_k^{(x-y)}$$

structure constants $\langle O_i O_j O_k \rangle = C_{ijk}/...$

- Higher points correlation functions can be reduced with OPE

Spectrum

- In the planar limit we need to focus on single trace operators only: $O(x) = \text{Tr}[\dots\dots]$

↳ field+derivatives
of $N=4$
spin-chains

- Single trace operators, can be seen as with periodic boundary conditions

$$\text{Tr}[Z X Z Z Z Z X] \rightarrow \uparrow \downarrow \uparrow \uparrow \uparrow \uparrow \downarrow$$

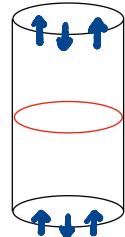
$$\begin{array}{ccc} D & \longrightarrow & H \\ \gamma_n & \longrightarrow & E_n \end{array}$$

non-protected
can have γ_n

protected, Δ is
correct ($\frac{1}{2}B\delta^2 - \text{Tr}[Z^4]$)

Spectrum of $N=5$

- Problem in trying to move up in perturbative order (away from large chains) are finite size corrections.

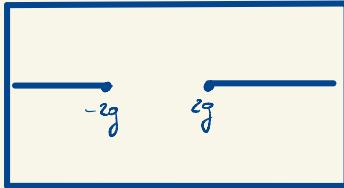
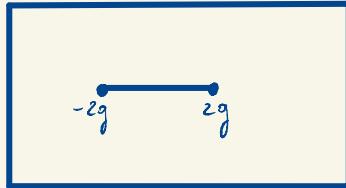


- TBA first solution for obtaining γ_n at finite coupling (complicated set of nested integral eq.)
- TBA \rightarrow Y-System \rightarrow T (hirata) \rightarrow QSC

A Bit of QSC

- A non-perturbative approach to computing Δ .
- 8 basic Q-functions

$$(P_1, P_2, P_3, P_4 \mid Q_1, Q_2, Q_3, Q_4)$$



- Global information about the state is encoded into asymptotics ($t \rightarrow \infty$)
- Obtaining these functions is a R-H problem + finite difference eq.
- Can obtain T-matrices from combinations of Q-functions

- QQ-Relations :

$$Q_i = -P^\alpha Q_{\alpha i}^+, \quad Q^i = P_\alpha Q^{\alpha i+}$$

$$Q_{\alpha i}^+ - Q_{\alpha i}^- = P_\alpha Q_i \quad P_\alpha P^\alpha = 0 = Q_i Q^i \quad Q^{\alpha i} = -(Q_{\alpha i}^+)^*$$

- Two extra i -periodic functions

$M_{ab}, w_{ij} \leftarrow$ used for analytical continuation

P- μ System

$P_a, P^\alpha, \tilde{P}_a, \tilde{P}^\alpha$

M_{ab}, \tilde{M}_{ab}

$Q_{\alpha i}$

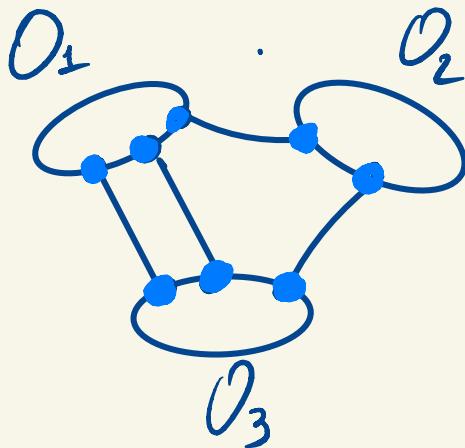
Q-w System

$Q_i, Q^i, \tilde{Q}_i, \tilde{Q}^i$

w_{ij}, \tilde{w}_{ij}

Structure Constants

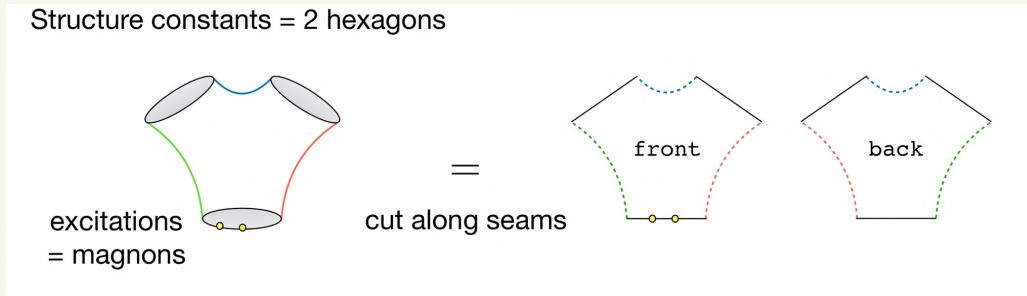
- As we have seen structure constants are associated with 3-point function:



- We will focus on the simplest case where only O_3 is non-protected

Hexagons

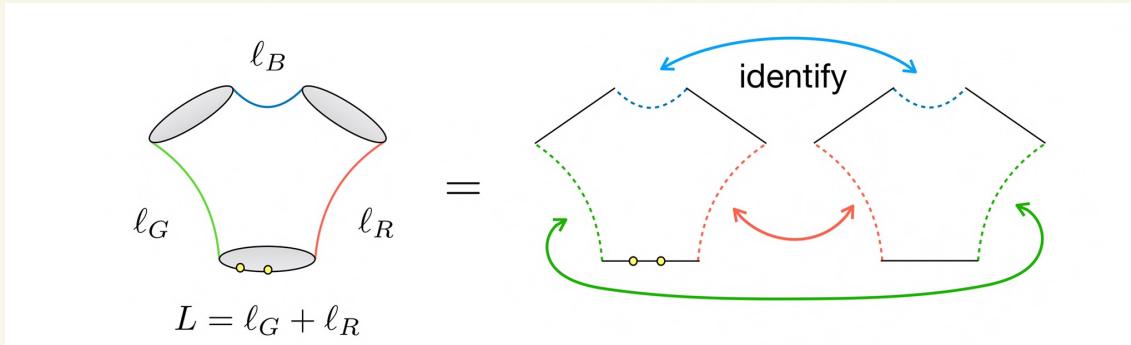
- Building blocks for correlation functions (single trace Op)



- Hexagon form factor conjecturally known to all loops, fixed by symmetry and form-factors axioms

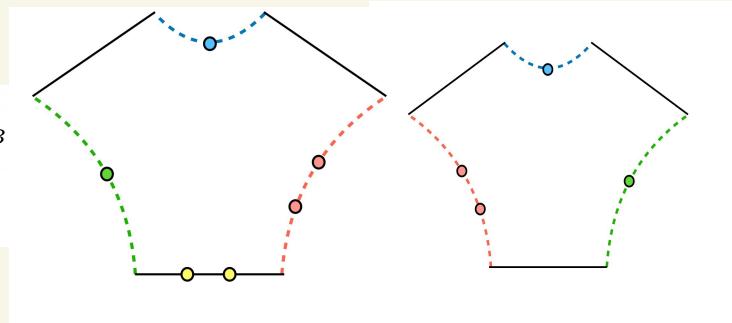
A hexagon with boundary points u_1 and u_2 at the bottom. The hexagon is composed of two pentagons: a green "front" pentagon and a red "back" pentagon. The boundary between them is a dashed arc. An equals sign follows this diagram, and to its right is the expression $= h(u_1, u_2)$.

• How does it work?

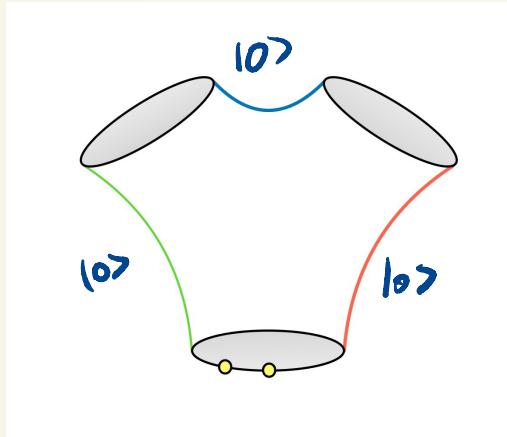


To obtain C_{ijk} , we need to sum over the 3-mirror channels

$$\sum_{N_\text{green}, N_\text{red}, N_\text{blue}} e^{-\ell_G E_G - \ell_R E_R - \ell_B E_B}$$



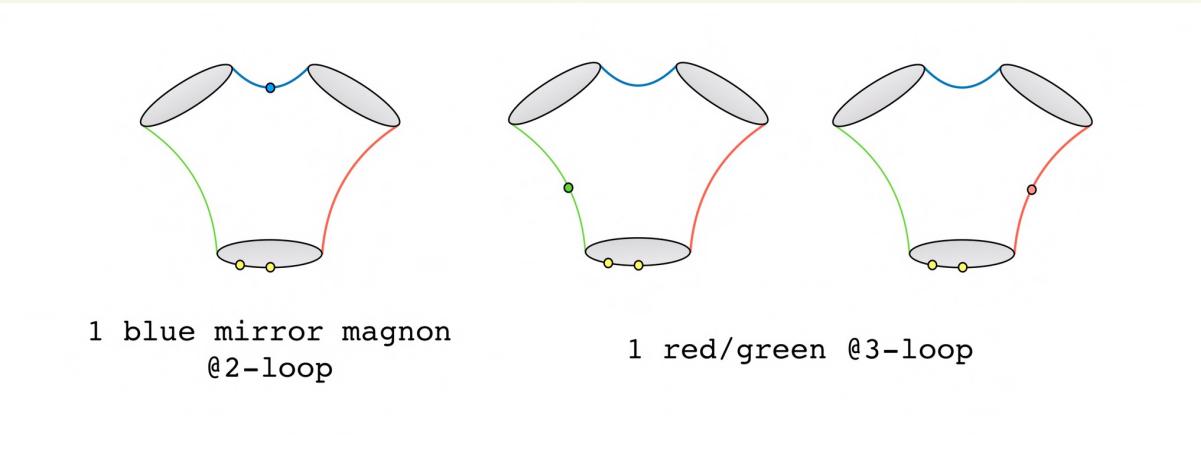
The asymptotic contribution, when we have 3-vacua



$$\approx \text{Gaudin}^{-1} \times \begin{matrix} \text{Sum over partitions} \\ \text{of Bethe roots} \end{matrix}$$

- for large enough bridges ($l_q, l_r, l_b \gg 1$) we can suppress the other contribution.

The other contributions look like (we specialize to the "lowest" operator)

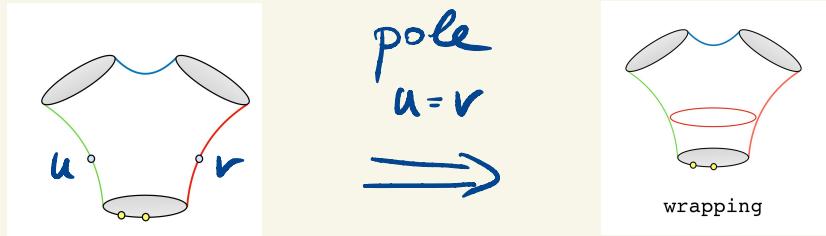


- Still under control, for example the first term

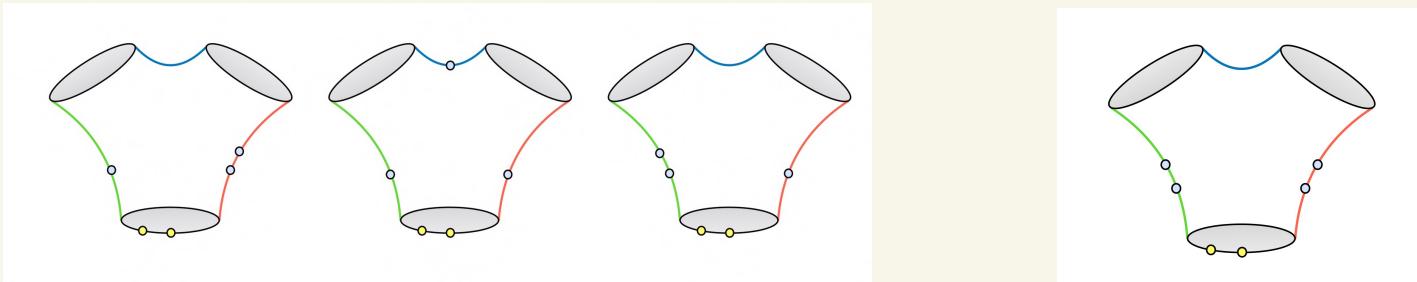
Asymptotic contribution $\times \sum_{\alpha=1}^{\infty} \int_{\partial D} \frac{T_\alpha(\bar{v}^\delta) M_\alpha(\bar{v}^\delta)}{(y^{[\alpha]} y^{[-\alpha]})^c \prod_j h_\alpha(\bar{v}^\delta, u_j)} \Rightarrow$

$36 g_3$ for Konishi

- At 5-loops we encounter the first divergence



- We get way more terms at the next orders.



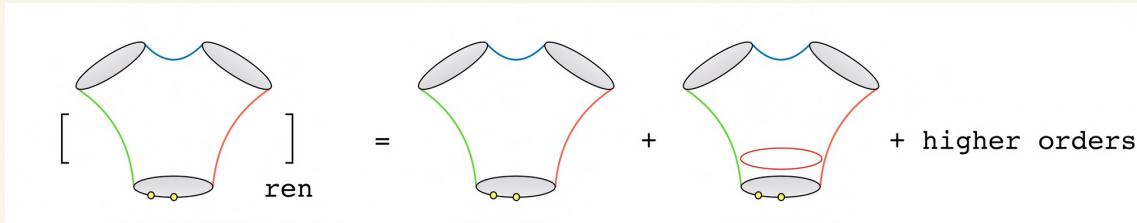
- Let us look at the bare integrand:

$$\frac{e^{-\ell_G E_G} W_G(u_G) e^{-\ell_B E_B} W_B(u_B) e^{-\ell_R E_R} W_R(u_R)}{\Delta(u_G, u_R)}$$

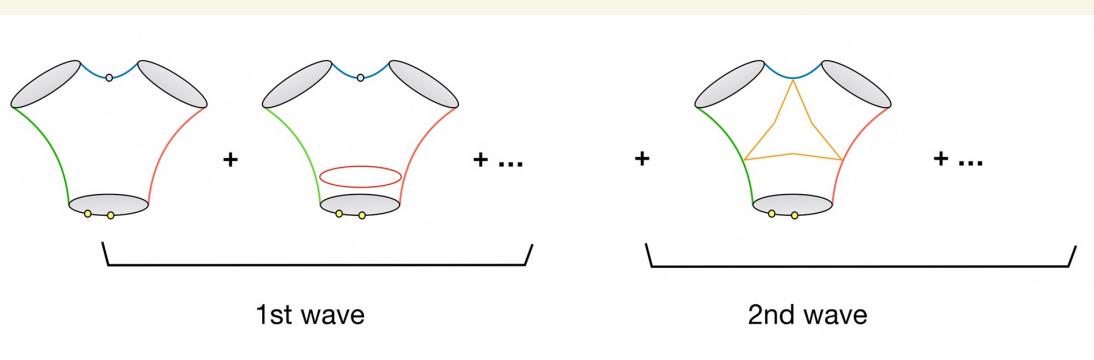
- divergence comes from $\Delta(u_G, u_R)$, we could shift the integration contour and add a contact term (proper derivation using form factors)

| | | |
|-------------------|-----|---|
| [] connected | $=$ | Bulk: 2-fold integral with $i0$ shifted contours |
| | $+$ | Contact term: localized 1-fold integral |

- In this process the contact term modifies the asymptotic contribution



- By doing the same procedure @ 3-loops discovered a new contribution



Renormalized Integrand

We can now write down a renormalized version of our bare integrand

$$C_{ijk} = N \sum \int \frac{W_G W_B W_R}{\Delta_{GR}^{+io}}$$

$$W_{G,R}(u) = e^{-\ell_{G,R} E} \mathbb{T}_{+,-}(u^\gamma)$$

$$W_B(u) = e^{-\ell_B E} \mathbb{T}_-(u^{-\gamma}) + e^{-\ell_D E} \mathbb{T}_+(u^{+\gamma})$$

Dressed Transf. f.

$$e^{-LE} \mathbb{T}_+(u^\gamma) \mathbb{T}_-(u^\gamma) = \frac{Y(u^\gamma)}{1 + Y(u^\gamma)}$$

$$\mathbb{T}_+(u^\gamma)/\mathbb{T}_-(u^\gamma) = e^{i \int \partial_v \ln \Delta(v^\gamma, u^\gamma) \log(1+Y(v^\gamma))}$$

QSC-Integrand

- Using a relation between the different W_i we can write down:

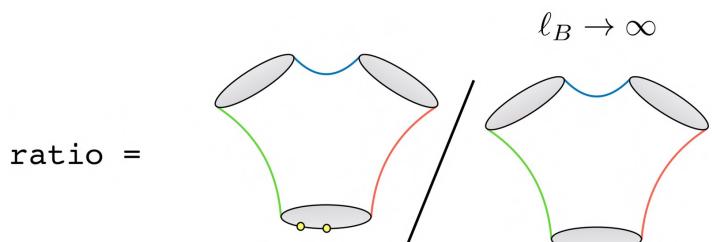
$$W_B \approx -\frac{Q_i^{(+\alpha)} Q^{(-\alpha)}}{y_i^{(+\alpha)} y^{(-\alpha)}} \mu_\alpha(v^r).$$

- $Q_i^{(+\alpha)} Q^{(-\alpha)}$ can be extracted perturbatively from the QSC.

Match with perturbative data

C_{ijk} for our "lowest" operator (Konishi) with two protected ones ($\langle \text{Tr}[z^i] \text{Tr}[D^2 z^j] \text{Tr}[z^k] \rangle$) has been computed using perturbation theory (Feynman diagrams) up to 5-loops.

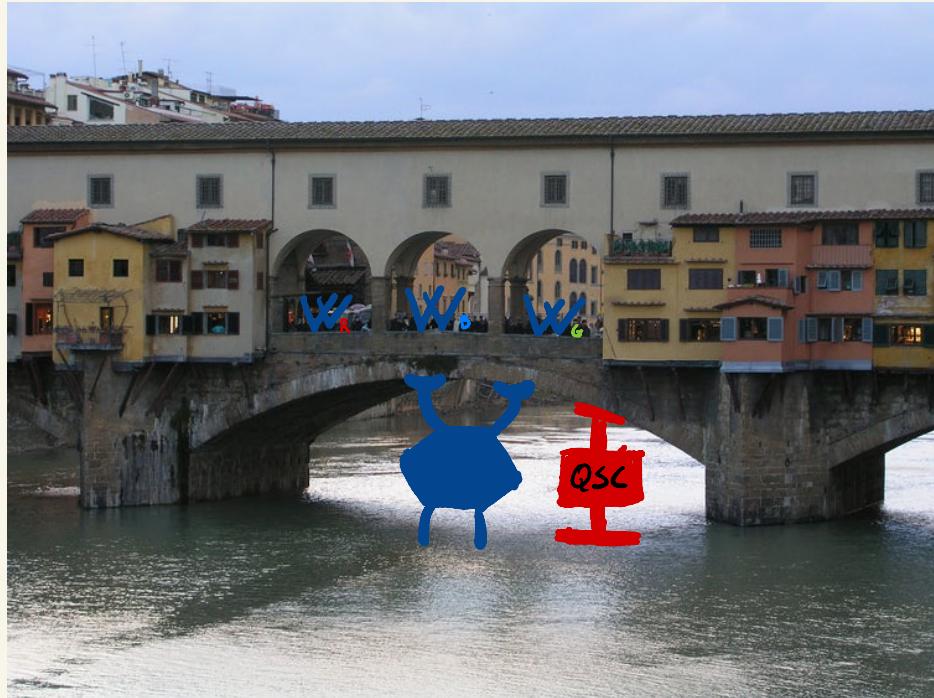
- We can extract the bottom contribution by:



Matches with
⇒ W_B !

Conclusions

- Interpret "triple-wrapping" as new fundamental object in the Bottom bridge
- Constructed corrected integrand to account for new effect using QSC building blocks
(Non trivial 5-loop check with perturbative data)
- We also have several strong coupling checks that support our conjectures. Some ideas also can be applied to adjacent bridges.



Thank you!