

QFT in AdS

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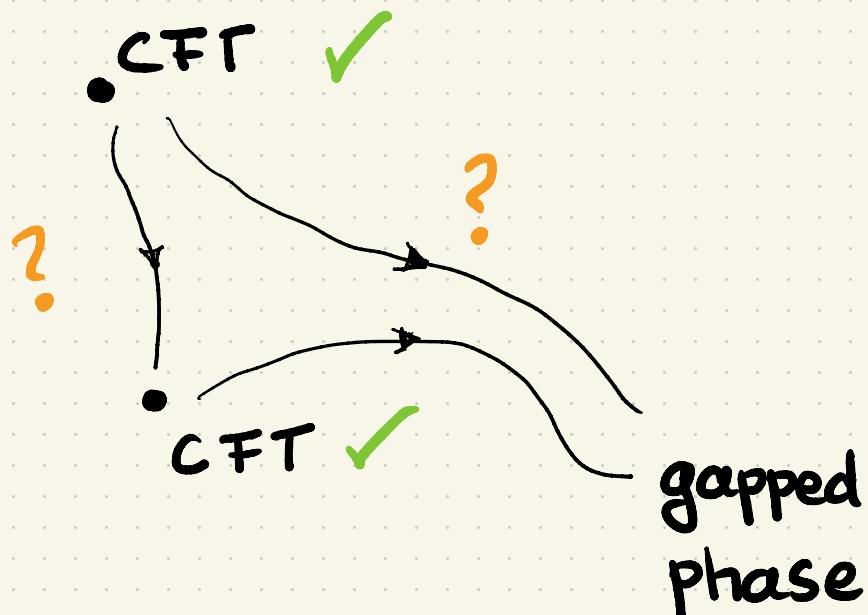
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based on work with :

A. Antunes, E. Lauria, M. Milam (to appear)

A. Antunes, M. Costa, J. Penedones, A. Salgarkar (appeared)

Motivation



Q: can we bootstrap an RG flow?

Q': are there any conformal structures along an RG flow?

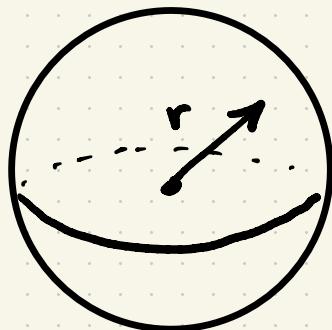
AdS_{d+1}

$$\text{Ric} \sim \frac{1}{R^2}$$

$$ds^2 = dr^2 + \sinh^2(r/R) R^2 d\Omega_d^2$$

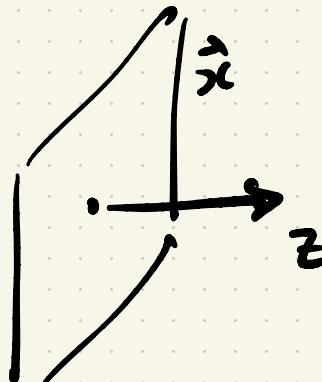
$$r \geq 0, \quad \hat{n} \quad \text{s.t.} \quad |\hat{n}|^2 = 1$$

↗
cliffo



$$ds^2 = \frac{R^2}{z^2} (d\vec{x}^2 + dz^2)$$

$$z > 0, \quad \vec{x} \in \mathbb{R}^d$$



isometries: $\text{so}(d+1, 1)$ conf. tr. on \vec{x} !

AdS_{d+1}

consider, for a QFT

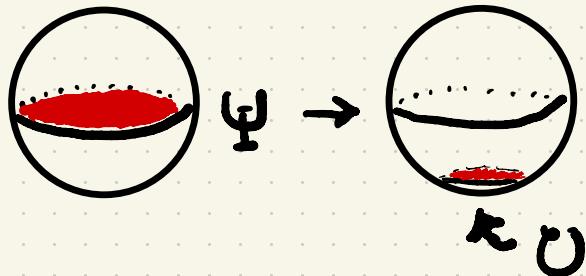
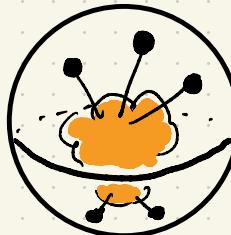
in AdS_{d+1} , the correlator:

$$\lim_{r_i \rightarrow \infty} \left(\prod_i e^{m_i r_i} \right) \langle O_1(r_1, \vec{n}_1) \dots O_k(r_k, \vec{n}_k) \rangle$$

$$=: \langle \hat{O}_1(\vec{n}_1) \dots \hat{O}_k(\vec{n}_k) \rangle$$

"boundary correlation function"

- conf. invt.
 - state-op. correspondence
- \Rightarrow conf. block dec.
- \Rightarrow bootstrap!



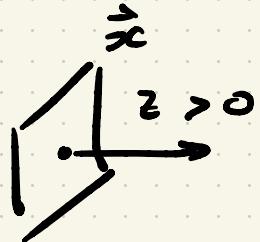
Expectations

- ① BCFTs
- ② RG flows

Expectations

BCFTs

$$ds^2 = \left(\frac{R}{z}\right)^2 (dz^2 + d\vec{x}^2) \xleftrightarrow{\text{Weyl}} ds^2 = dz^2 + d\vec{x}^2$$



$$\langle O(\vec{x}, z) \rangle_{\text{flat}} = \frac{a_p}{z^{\Delta_O}}$$

$$\langle O_1(\vec{x}_1, z_1) O_2(\vec{x}_2, z_2) \rangle_{\text{flat}} =$$

$$f(\xi) / z_1^{\Delta_1} z_2^{\Delta_2}$$

$$\xi = \frac{(\vec{x}_1 - \vec{x}_2)^2 + (z_1 - z_2)^2}{4 z_1 z_2}$$

Expectations

BCFTs

$$ds^2 = \left(\frac{R}{z}\right)^2 (dz^2 + d\vec{x}^2) \xrightarrow{\text{Weyl}} ds^2 = dz^2 + d\vec{x}^2$$

$$\langle O(\vec{x}, z) \rangle_{\text{AdS}} = \frac{a_O}{R^{\Delta_O}}$$

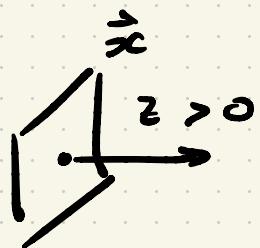
just constant

$$\leftarrow \quad \langle O(\vec{x}, z) \rangle_{\text{flat}} = \frac{a_O}{z^{\Delta_O}}$$

$$\langle O_1(\vec{x}_1, z_1) O_2(\vec{x}_2, z_2) \rangle_{\text{AdS}} = f(\xi) / R^{\Delta_1 + \Delta_2}$$

$$2\xi + 1 = \cosh(d_{1,2}/R)$$

just geodesic distance



$$\langle O_1(\vec{x}_1, z_1) O_2(\vec{x}_2, z_2) \rangle_{\text{flat}} =$$

$$f(\xi) / z_1^{\Delta_1} z_2^{\Delta_2}$$

$$\xi = \frac{(\vec{x}_1 - \vec{x}_2)^2 + (z_1 - z_2)^2}{4z_1 z_2}$$

→ at fixed R , no new conf. W.I.!

Boundary OPE

In BCFT:

$$\langle \hat{O}(\vec{x}, z) \dots \rangle_{\text{flat}} = \sum_{\hat{\phi}} z^{\Delta_{\hat{\phi}}} b_0 \langle \hat{\phi}(\vec{x}) \dots \rangle_{\text{flat}}$$

so in AdS:

$$\langle \hat{O}(\vec{x}, z) \dots \rangle_{\text{AdS}} = \sum_{\hat{\phi}} z^{\Delta_{\hat{\phi}}} R^{\Delta_{\hat{\phi}}} b_0 \langle \hat{\phi}(\vec{x}) \dots \rangle_{\text{AdS}}$$

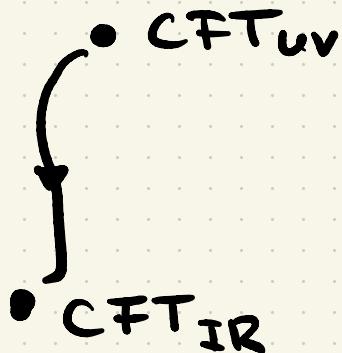
}

$\Delta_{\hat{\phi}} > 0$, so falloff as $z \rightarrow 0$

BCFT "UV" \rightarrow IR in AdS

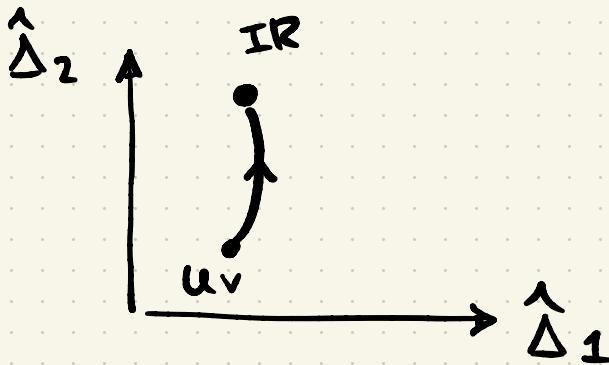
RG flows (naive)

In AdS,



is normally covariant,

so bdy. correlators should be conformal :



RG flows (naive)

$$ds^2 = dr^2 + \sinh^2(r/R) R^2 d\Omega^2$$

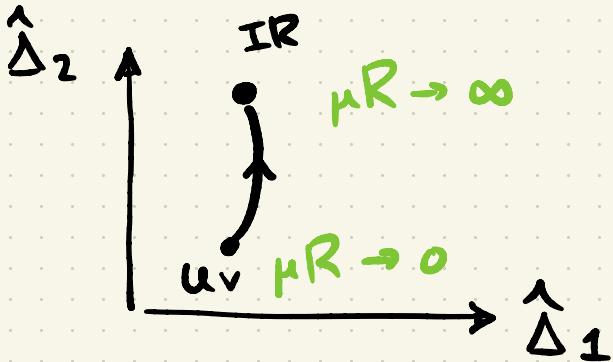
$$\text{Ric} \sim 1/R^2$$

In AdS,

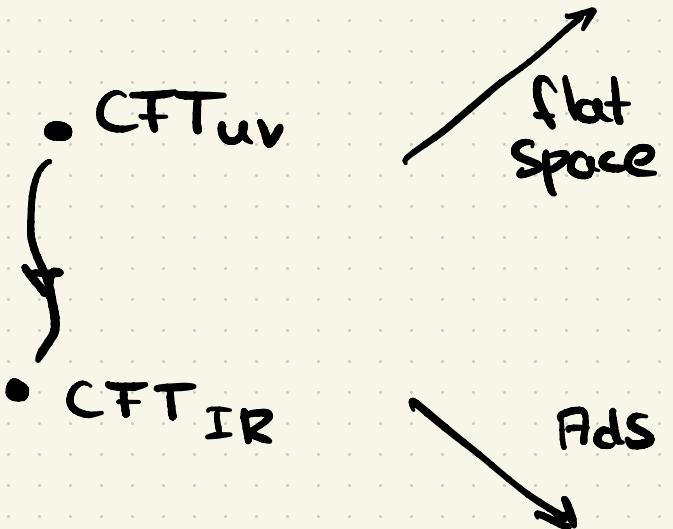
- CFT_{uv} $R \rightarrow 0$
 - CFT_{IR} $R \rightarrow \infty$
- $\left. \begin{matrix} \bullet \\ \bullet \end{matrix} \right\} \mu R$

is normally covariant,

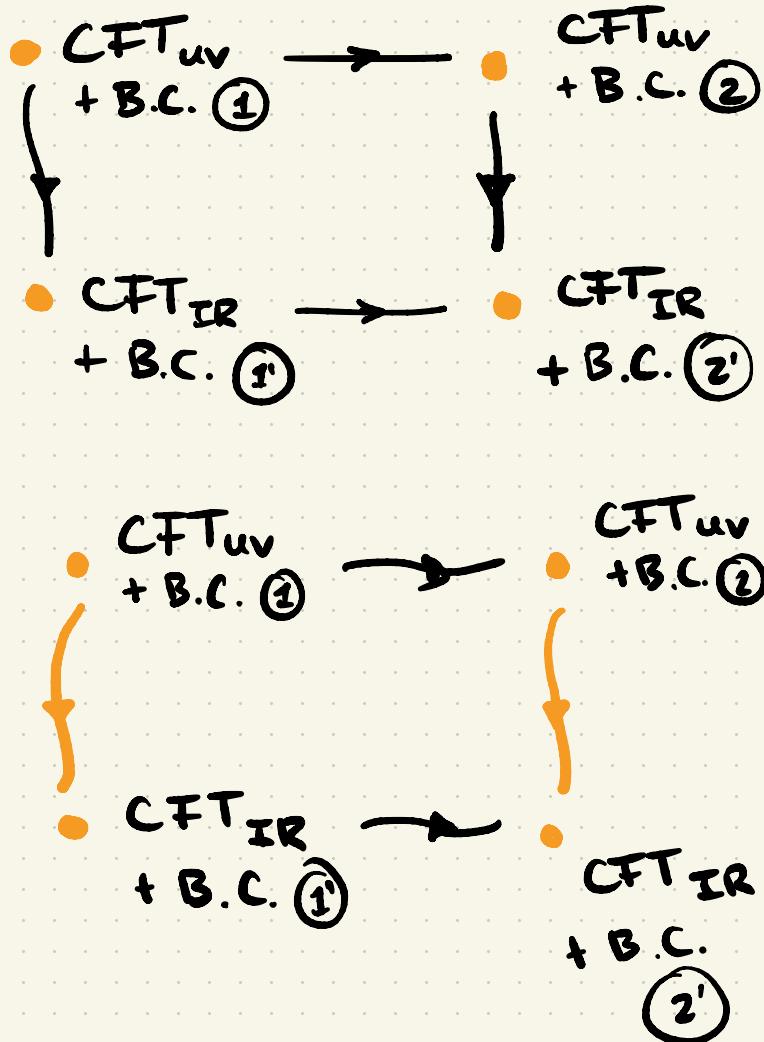
so bdy. correlators should be conformal :



RG flows (less naive)



Claim : bdy flows are
not covariant \Rightarrow bdy
follows bulk.

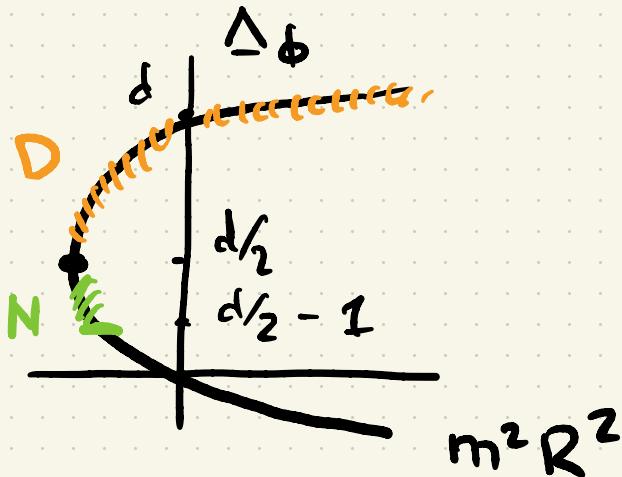


Boundary RG flows - example

free scalar in AdS_{d+1}

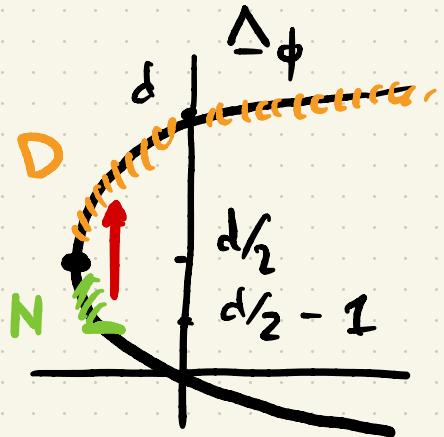
$$S = \frac{1}{2} \int d^d \bar{x} dz \sqrt{g} (\partial_\mu \Phi \partial^\mu \Phi + m^2 \Phi^2)$$

bdy. correlators : G.F.F. ϕ



$$\left(\langle \phi_1 \dots \phi_n \rangle_{\text{connected}}^{n \geq 2} = 0 \right)$$

$$\Delta_\phi (\Delta - d) = m^2 R^2$$



$$m^2 R^2$$

$\uparrow : \text{RG flow by } \int d^d \vec{x} \phi^2 \text{ (at bdy!)}$

$$< \phi(\vec{x}) \phi(0) \exp(-g \int d^d \vec{x} \phi^2) >$$

$$\propto \left(\int d^d \vec{k} e^{i \vec{k} \cdot \vec{x}} \frac{\# k^{2\Delta-d}}{1 + \# g k^{2\Delta-d}} \right)$$

$$\rightarrow \begin{cases} \frac{1}{x^{2\Delta\phi^-}} & \text{as } x \rightarrow 0 \\ \frac{1}{x^{2\Delta\phi^+}} & \text{as } x \rightarrow \infty \end{cases}$$

is not conformal!
(i.e. not covariant)

Bulk flows - boundary terms

BCFT + relevant deformation
in AdS

$$\langle \dots \exp \left(-\lambda \int_0^d \tilde{x}^i dz \sqrt{g} O(\tilde{x}, z) \right) \rangle$$

$$= \langle \dots \lambda \int_0^d \tilde{x}^i \int dz \left(\frac{z}{R} \right)^{-(d+1)} z^{\hat{\Delta}_O} R^{-\Delta_O} b_O \hat{O}(\tilde{x}) \rangle$$

\Rightarrow IR divergence if $\hat{\Delta}_O \leq d$. Covariant counterterm:

$$\left. ds^2 \right|_{z=\epsilon} = \frac{d\tilde{x}^2}{(z/R)^2} \sim \sqrt{h} = \left(\frac{z}{R} \right)^{-d}$$

$$S_{ct} = -\frac{\lambda}{\hat{\Delta}_O - d} \int_0^d \tilde{x}^i \sqrt{h} O(\tilde{x}, z) \Big|_{z=\epsilon}$$

bulk operator

Recap

Callan-Wilczek 1990: .. constant negative curvature provides a very convenient IR regulator..

RG flows in AdS can have IR divergences, but
 \exists covariant RG flows in AdS, where bdy follows the bulk.

(q: what happens for $\hat{\Delta}\hat{g} = d$?)

Q1 : perturbative checks?

Q2 : can we bootstrap a flow?

Perturbative checks

- WF in AdS
- Minimal models in AdS

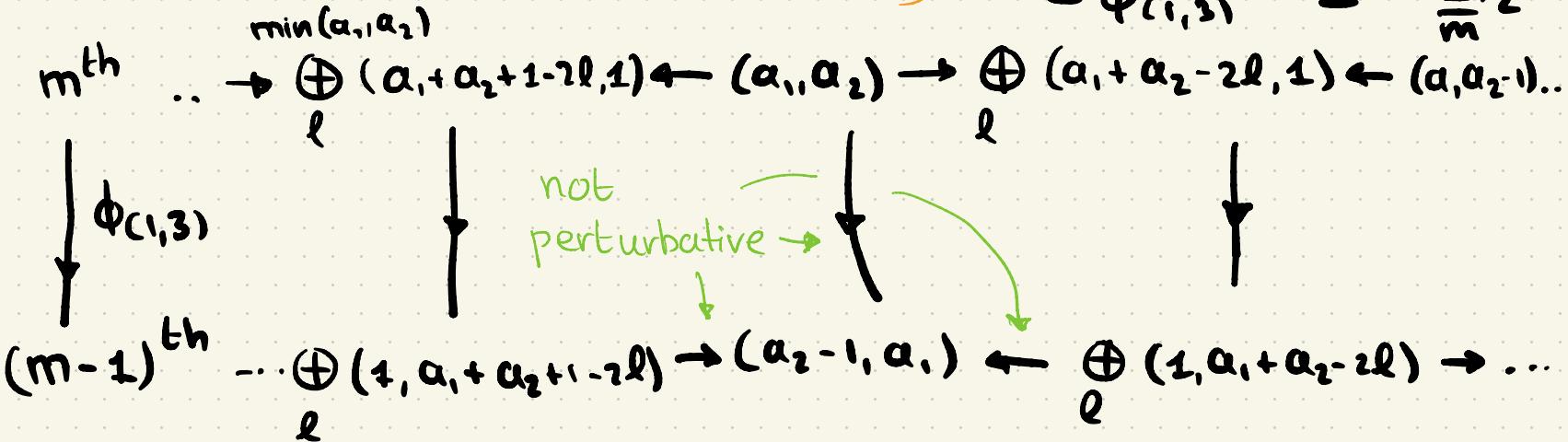
Subtle : $\log(d_{1-2}/R)$, $\log(\mu R)$

so RG improvement is of limited use.

Perturbative checks - minimal models

- flat space BCFT

Fredenhagen, Gaberdiel, Schmidt-Colinet '09



- AdS (α_1, α_2) Lauria, Milam, BrR '22

$$(\alpha_2, \alpha_1)$$

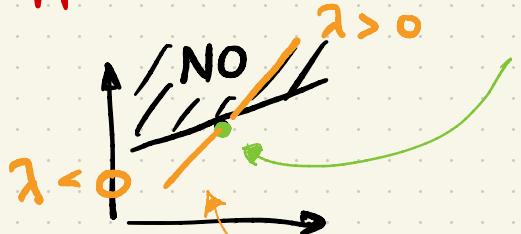
- bdy spectrum: $\Psi_{r,s}^m \rightarrow \Psi_{s,r}^{m-1}$ ✓
- 1-pt. functions ?

Numerical results

- ① EFT - type bounds
- ② RG flow constraints

EFT - type bounds

idea:



QFT in AdS

+
irrelevant bulk
interaction

$$\lambda \int d^{d+1}x \sqrt{g} \propto 0$$

implies $\lambda < 0$ necessarily!

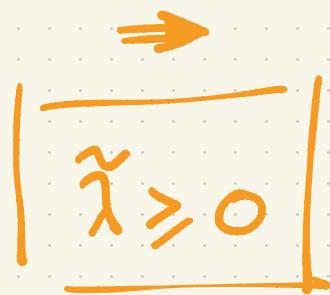
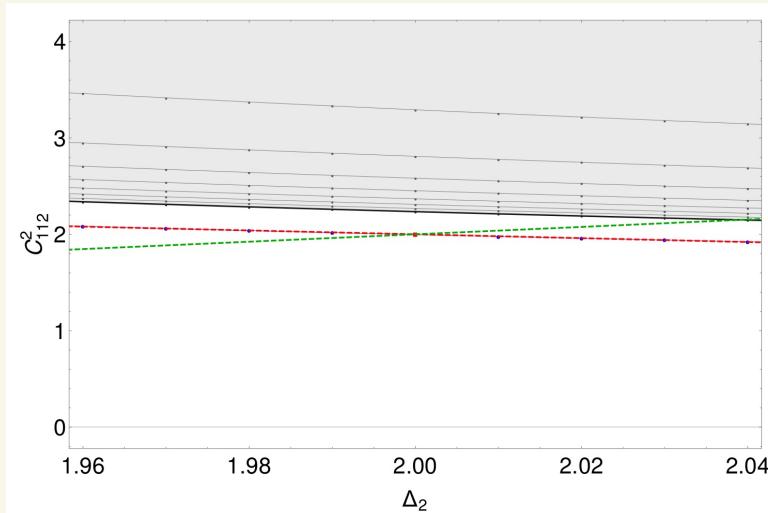
EFT - type bounds

$$\textcircled{1} \quad S = \int d^{d+1}x \sqrt{g} \left(\frac{1}{2} |\partial \Phi|^2 - \tilde{\lambda} |\Phi|^4 \right)$$

pert
thy : $(\Delta_\Phi, \Delta_{\Phi^2}, C_{112}^2) = (1, 2 - \frac{\tilde{\lambda}}{6\pi}, 2 - \tilde{\lambda} \frac{23}{36\pi}) + \mathcal{O}(\tilde{\lambda}^2)$

numerics:

$(\langle \Phi \Phi \Phi \Phi \rangle)$
in 1d



EFT-type bounds

$$\textcircled{2} \quad S = S_{\substack{\text{min.} \\ \text{model}}} + \lambda \int T \bar{T} d^2x$$

Pert. thy: $(\hat{\Delta}_{\phi}, \hat{\Delta}_D, \hat{\Delta}_{D^2}) = (\hat{\Delta}_{\phi} + \frac{\lambda\pi}{2}\hat{\Delta}_{\phi}(\hat{\Delta}_{\phi}-1), 2+\lambda\pi, 4+6\lambda\pi) + \mathcal{O}(\lambda^2)$

numerics:

$$\langle \phi \phi \phi \phi \rangle$$

$$\phi \times \phi = 1 + D + \text{gap}$$

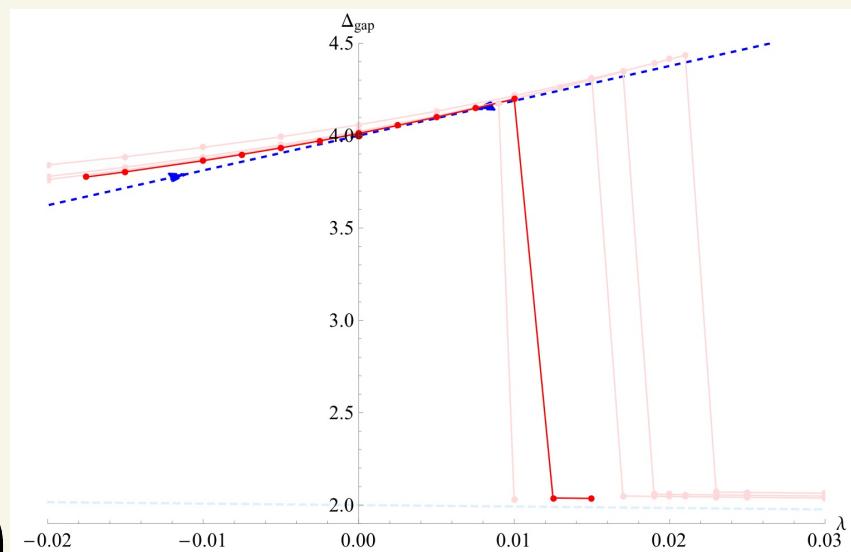
$$\hat{\Delta}_{\phi} = \frac{1}{2} - \frac{\lambda\pi}{8}$$

$$\hat{\Delta}_D = 2 + \lambda\pi$$

max Δ_{gap}

$(\phi = \psi_{(1,3)} \text{ in Ising})$

Antunes, Lauria,
Milam, BrR
(to appear)



⇒

$\lambda \leq 0$
 for 2d
 Ising

EFT - type bounds

numerics.

$$\langle \phi \phi \phi \phi \rangle$$

$$\phi \times \phi = 1 + D + \text{gap}$$

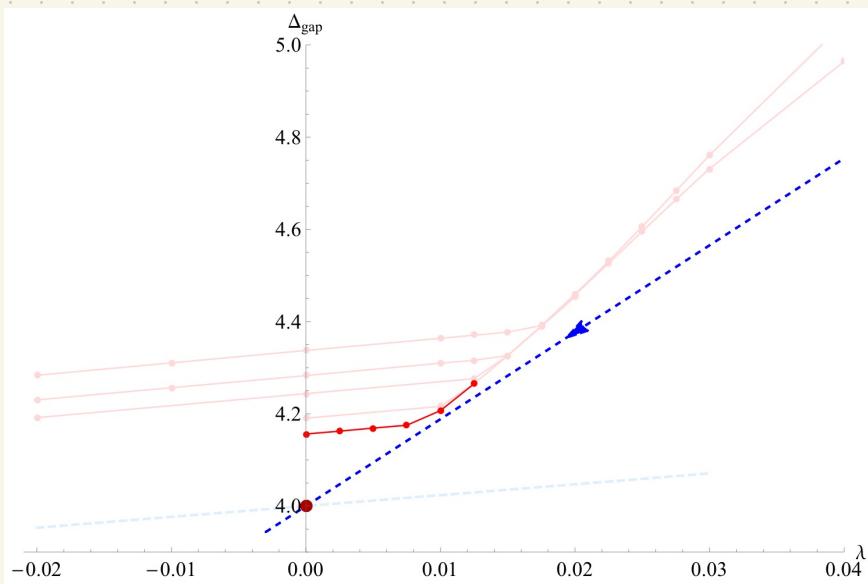
$$\hat{\Delta}_\phi = \frac{3}{2} + \frac{3\lambda}{8}$$

$$\hat{\Delta}_D = 2 + \lambda \pi$$

max Δ_{gap}

$$(\phi = \Psi_{3,1}$$

in tri-crit.
Ising)



no bound
for tri-crit.
Ising
(from this
corr.)

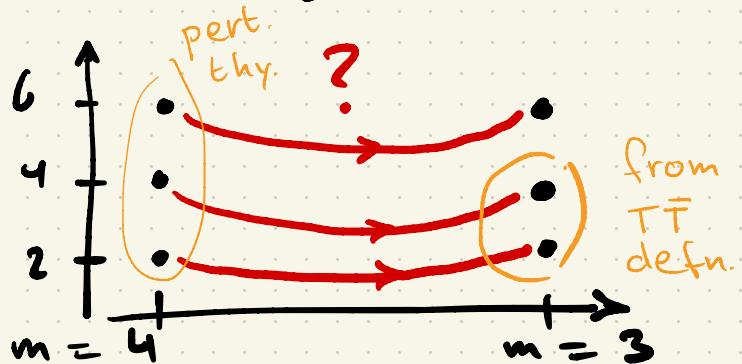
RG flow constraints

minimal model flow

$m=4$ $m=3$
tri-crit. Ising
Ising

$$(2,1) \Big| 11, \Psi_{3,1} \Big| \Delta_{3,1} = \frac{3}{2} \quad (1,2) \Big| 11, \Psi_{1,3} \Big| \Delta_{1,3} = \frac{1}{2}$$

singlets



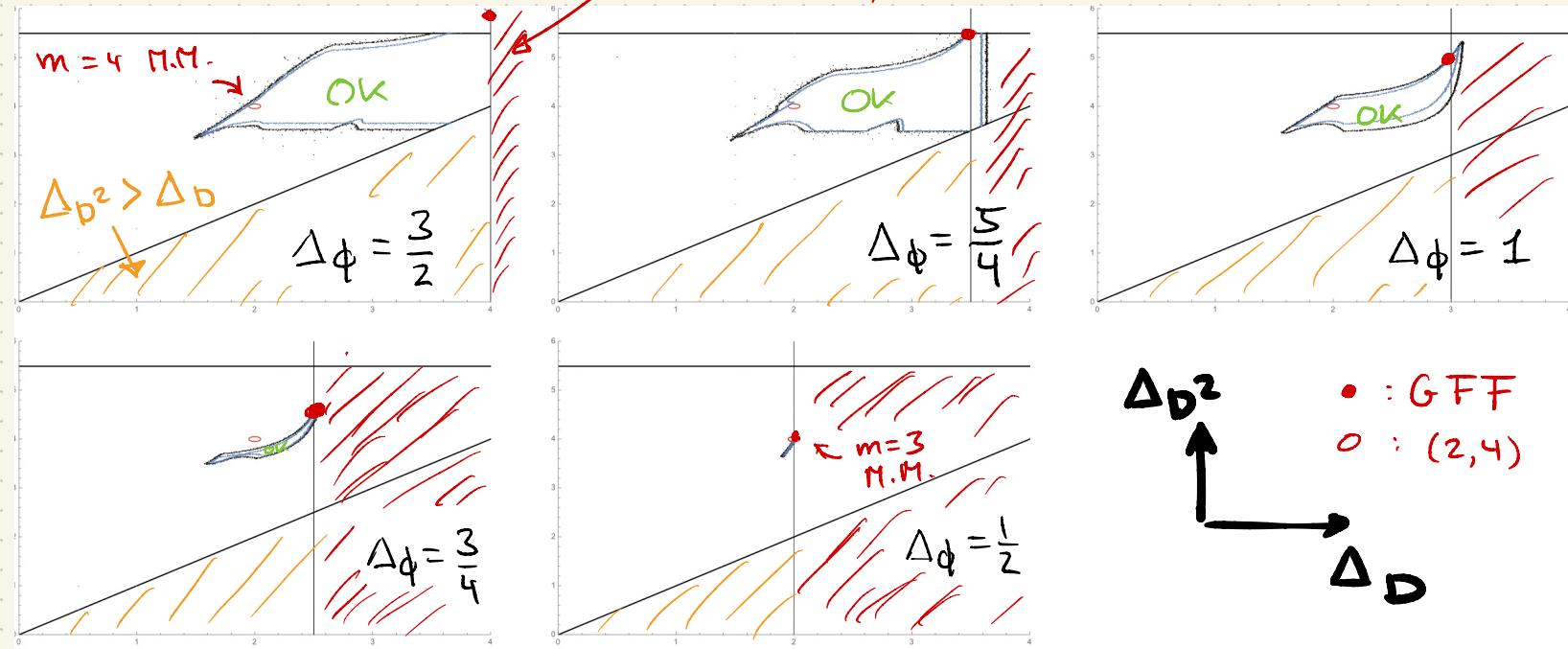
RG flow constraints - numerics

$\langle \phi\phi\phi\phi \rangle, \langle \phi\phi DDD \rangle, \langle DDDDD \rangle$

$\phi \times \phi \sim 1 + D + D^2 + \text{operators with } \Delta > 5.5$

$D \times D \sim \text{"}$

$$\Delta_D < \Delta_{D,GFF} = 2\Delta_\phi + 1$$



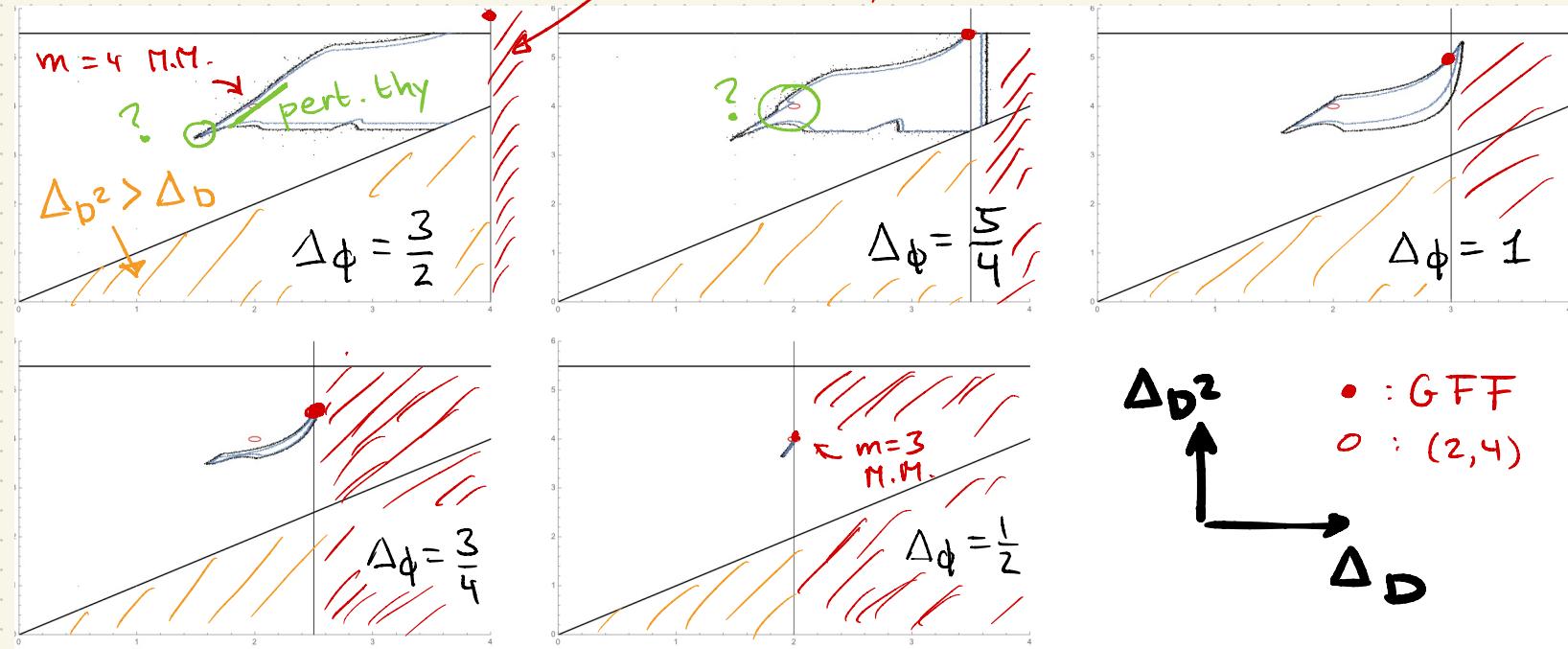
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$$\Delta_D < \Delta_{D,GFF} = 2\Delta_\phi$$



Conclusions

QFT in AdS offers a promising way to constrain RG flows.

- Systematize EFT-type bounds?
Caron-Huot, Mazac, Rastelli, Simmons-Duffin '21
- Physics of extremal allowed points?
- Connect with other analyses?

Carmi, Di Pietro, Komatsu '18 (large N)
Hogervorst, Meineri, Penedones, Salehi Vaziri '21 (HT)

- flat-space limit and S-matrix bootstrap

Thank you !