# Local-equilibrium transport of oscillating neutrinos



LA-UR-24-26818

### Overview

Flavor mixing in neutron star mergers & core-collapse supernovae is estimated to have significant effects, but the quantum kinetic equation (QKE) is computationally intractable.



This talk: A coarse-grained transport theory based on local mixing equilibrium.

## I. Foundations

Studies have questioned the accuracy of quantum kinetics by comparing many-body and mean-field evolution.

$$H = \sum_{\mathbf{p}} \omega_{\mathbf{p}} \mathbf{B} \cdot \hat{\mathbf{J}}_{\mathbf{p}} + \mu \sum_{\mathbf{p}, \mathbf{q}} \left(1 - \cos \theta_{\mathbf{p}\mathbf{q}}\right) \hat{\mathbf{J}}_{\mathbf{q}} \cdot \hat{\mathbf{J}}_{\mathbf{p}}$$
Pairing interaction

where each J is a flavor isospin operator.

Studies have questioned the accuracy of quantum kinetics by comparing many-body and mean-field evolution.

$$H = \sum_{\mathbf{p}} \omega_{\mathbf{p}} \mathbf{B} \cdot \hat{\mathbf{J}}_{\mathbf{p}} + \mu \sum_{\mathbf{p}, \mathbf{q}} (1 - \cos \theta_{\mathbf{pq}}) \hat{\mathbf{J}}_{\mathbf{q}} \cdot \hat{\mathbf{J}}_{\mathbf{p}}$$
Pairing interaction

where each **J** is a flavor isospin operator.

The many-body Hamiltonian used in this literature *requantizes the forward-scattering Hamiltonian*—which fundamentally alters the physics.

Johns, 2305.04916 (Neutrino many-body correlations)

### Calculations have recently been done using the correct many-body Hamiltonian:

#### Neutrino many-body flavor evolution: the full Hamiltonian

Vincenzo Cirigliano,<sup>1</sup> Srimoyee Sen,<sup>2</sup> and Yukari Yamauchi<sup>1</sup>

<sup>1</sup>Institute for Nuclear Theory, University of Washington, Seattle, WA 98195, USA <sup>2</sup>Department of Physics and Astronomy, Iowa State University, Ames, IA, 50011 (Dated: April 26, 2024)

We study neutrino flavor evolution in the quantum many-body approach using the full neutrinoneutrino Hamiltonian, including the usually neglected terms that mediate non-forward scattering processes. Working in the occupation number representation with plane waves as single-particle states, we explore the time evolution of simple initial states with up to N = 10 neutrinos. We discuss the time evolution of the Loschmidt echo, one body flavor and kinetic observables, and the one-body entanglement entropy. For the small systems considered, we observe 'thermalization' of both flavor and momentum degrees of freedom on comparable time scales, with results converging towards expectation values computed within a microcanonical ensemble. We also observe that the inclusion of non-forward processes generates a faster flavor evolution compared to the one induced by the truncated (forward) Hamiltonian.

### Calculations have recently been done using the correct many-body Hamiltonian:

#### Neutrino many-body flavor evolution: the full Hamiltonian

Vincenzo Cirigliano,<sup>1</sup> Srimoyee Sen,<sup>2</sup> and Yukari Yamauchi<sup>1</sup>

<sup>1</sup>Institute for Nuclear Theory, University of Washington, Seattle, WA 98195, USA <sup>2</sup>Department of Physics and Astronomy, Iowa State University, Ames, IA, 50011 (Dated: April 26, 2024)

We study neutrino flavor evolution in the quantum many-body approach using the full neutrinoneutrino Hamiltonian, including the usually neglected terms that mediate non-forward scattering processes. Working in the occupation number representation with plane waves as single-particle states, we explore the time evolution of simple initial states with up to N = 10 neutrinos. We discuss the time evolution of the Loschmidt echo, one body flavor and kinetic observables, and the one-body entanglement entropy. For the small systems considered, we observe 'thermalization' of both flavor and momentum degrees of freedom on comparable time scales, with results converging towards expectation values computed within a microcanonical ensemble. We also observe that the inclusion of non-forward processes generates a faster flavor evolution compared to the one induced by the truncated (forward) Hamiltonian.

Coherent forward scattering ( $G_F$  vs.  $G_F^2$ ) is a kinetic phenomenon. Most neutrino many-body models aren't actually expected to be described by kinetics.

Johns, 2305.04916 (*Neutrino many-body correlations*) Also see Shalgar & Tamborra, 2304.13050

### We devised the **once-in-a-lifetime encounter (OILE) model** to illustrate the emergence of coherent flavor evolution in many-body systems with brief interaction times.

Kost, **Johns**, & Duan, 2402.05022 (*Once-in-a-lifetime encounter models for neutrino media*)

Compare with Friedland & Lunardini 2003



# II. Paradigms

1. Direct numerical simulation

### 1. Direct numerical simulation

### 2. Asymptotic-state subgrid modeling

### 1. Direct numerical simulation

### 2. Asymptotic-state subgrid modeling

3. Miscidynamics (*i.e.*, transport near local mixing equilibrium)

$$i\left(\partial_t + \hat{\mathbf{p}} \cdot \partial_{\mathbf{r}}\right)\rho = [H, \rho] + iC$$

Vanishes in Vanishes in *mixing* eq *collisional* eq



Johns, 2306.14982 (Thermodynamics of oscillating neutrinos)

Asymptotic-state subgrid models are not self-consistent: If fast instabilities are instantaneous, then angular crossings can't form in the first place.

Asymptotic-state subgrid models are not self-consistent: If fast instabilities are instantaneous, then angular crossings can't form in the first place.

In fact, the assumption of **rapid relaxation to local mixing equilibrium** is pathway #1 to the miscidynamic equation:

Asymptotic-state subgrid models are not self-consistent: If fast instabilities are instantaneous, then angular crossings can't form in the first place.

In fact, the assumption of **rapid relaxation to local mixing equilibrium** is pathway #1 to the miscidynamic equation:

$$[H_{\boldsymbol{p}}, \rho_{\boldsymbol{p}}] \approx \gamma \left( \rho_{\boldsymbol{p}}^{\mathrm{eq}} - \rho_{\boldsymbol{p}} \right)$$

with  $\gamma 
ightarrow \infty$ 

Nagakura, **Johns**, & Zaizen, 2312.16285 (*BGK subgrid model for neutrino quantum kinetics*)

Johns, 2401.15247 (Subgrid modeling of neutrino oscillations in astrophysics)

Asymptotic-state subgrid models are not self-consistent: If fast instabilities are instantaneous, then angular crossings can't form in the first place.

In fact, the assumption of **rapid relaxation to local mixing equilibrium** is pathway #1 to the miscidynamic equation:

$$[H_{\boldsymbol{p}}, \rho_{\boldsymbol{p}}] \approx \gamma \left( \rho_{\boldsymbol{p}}^{\mathrm{eq}} - \rho_{\boldsymbol{p}} \right)$$
  
with  $\gamma \to \infty$ 

Nagakura, **Johns**, & Zaizen, 2312.16285 (*BGK subgrid model for neutrino quantum kinetics*)  $i \left(\partial_t + \hat{\boldsymbol{p}} \cdot \partial_{\boldsymbol{x}}\right) \rho_{\boldsymbol{p}}^{\text{eq}} = i C_{\boldsymbol{p}}^{\text{eq}}$ with  $\left[H_{\boldsymbol{p}}^{\text{eq}}, \rho_{\boldsymbol{p}}^{\text{eq}}\right] = 0$ 

Johns, 2401.15247 (Subgrid modeling of neutrino oscillations in astrophysics)

Asymptotic-state subgrid models are not self-consistent: If fast instabilities are instantaneous, then angular crossings can't form in the first place.

In fact, the assumption of **rapid relaxation to local mixing equilibrium** is pathway #1 to the miscidynamic equation:

$$[H_{\boldsymbol{p}}, \rho_{\boldsymbol{p}}] \approx \gamma \left( \rho_{\boldsymbol{p}}^{\mathrm{eq}} - \rho_{\boldsymbol{p}} \right)$$
  
with  $\gamma \to \infty$ 

Nagakura, **Johns**, & Zaizen, 2312.16285 (*BGK subgrid model for neutrino quantum kinetics*)

$$\begin{split} i\left(\partial_t + \hat{\boldsymbol{p}} \cdot \partial_{\boldsymbol{x}}\right) \rho_{\boldsymbol{p}}^{\mathrm{eq}} &= iC_{\boldsymbol{p}}^{\mathrm{eq}} \\ \text{with } \left[H_{\boldsymbol{p}}^{\mathrm{eq}}, \rho_{\boldsymbol{p}}^{\mathrm{eq}}\right] = 0 \end{split}$$

Essential qualities of miscidynamics:

- Astrophysical driving
- Self-consistent equilibrium

Johns, 2401.15247 (Subgrid modeling of neutrino oscillations in astrophysics)

$$\rho_{\boldsymbol{p}} = \langle \rho_{\boldsymbol{p}} \rangle + \delta \rho_{\boldsymbol{p}}$$

$$\begin{split} \rho_{\boldsymbol{p}} &= \langle \rho_{\boldsymbol{p}} \rangle + \delta \rho_{\boldsymbol{p}} \\ & \text{No correlations} \\ \langle [H_{\boldsymbol{p}}, \rho_{\boldsymbol{p}}] \rangle &= [\langle H_{\boldsymbol{p}} \rangle, \langle \rho_{\boldsymbol{p}} \rangle] + \langle [\delta H_{\boldsymbol{p}}, \delta \rho_{\boldsymbol{p}}] \rangle \\ & \text{This must vanish,} \\ & \text{otherwise coarse-grained} \\ & \text{averages vary on the} \\ & \text{oscillation scale.} \\ & \downarrow (\cdot)^{\text{eq}} \equiv \langle \cdot \rangle \\ & i \left( \partial_t + \hat{\boldsymbol{p}} \cdot \partial_{\boldsymbol{x}} \right) \rho_{\boldsymbol{p}}^{\text{eq}} = i C_{\boldsymbol{p}}^{\text{eq}} \\ & \text{with} \left[ H_{\boldsymbol{p}}^{\text{eq}}, \rho_{\boldsymbol{p}}^{\text{eq}} \right] = 0 \end{split}$$

Why would there be rapid relaxation or uncorrelated fluctuations? We'll explore ergodicity and thermodynamics as possible explanations.

$$\begin{split} i\left(\partial_t + \hat{\boldsymbol{p}} \cdot \partial_{\boldsymbol{x}}\right) \rho_{\boldsymbol{p}}^{\mathrm{eq}} &= iC_{\boldsymbol{p}}^{\mathrm{eq}} \\ \text{with } \left[H_{\boldsymbol{p}}^{\mathrm{eq}}, \rho_{\boldsymbol{p}}^{\mathrm{eq}}\right] = 0 \end{split}$$

 $\langle \rangle$ 

### Miscidynamics generalizes adiabatic quantum evolution.



Miscidynamic adiabaticity

### Miscidynamics generalizes adiabatic quantum evolution.

MSW adiabaticity In particular, we can have adiabatic *incoherent* evolution. Control parameters include flavor polarizations |P| as well as couplings  $\lambda$  and  $\mu$ .



Raffelt-Smirnov adiabaticity Raffelt & Smirnov 2007



Miscidynamic adiabaticity

### Miscidynamics generalizes adiabatic quantum evolution.



In particular, we can have adiabatic incoherent evolution. Control parameters include flavor polarizations |P| as well as couplings  $\lambda$  and  $\mu$ .

Kost, Johns, & Duan, in preparation

μt

3000

2000

Stage 3

Group 1 miscidynamic solution

Group 2 miscidynamic solution

4000

5000

### **Current status of miscidynamics**

- The adiabatic theory (what I'm presenting) is largely complete, *i.e.*, simulation-ready.
- The nonadiabatic theory (with finite and/or correlated deviations from mixing eq) is still being developed.
- The nonadiabatic theory is definitely needed for some models, but it's unclear whether it's needed for astrophysics. There may be adiabatic adjustment. *The adiabatic hypothesis*

# III. Flavor ergodicity

### **Ergodic hypothesis** Neutrinos uniformly explore all accessible parts of flavor space.

Johns, 2306.14982 (Thermodynamics of oscillating neutrinos)

### **Ergodic hypothesis** Neutrinos uniformly explore all accessible parts of flavor space. Johns, 2306.14982 (Thermodynamics of oscillating neutrinos)

"In a sense, ergodicity is universal, and the central question is to define the subspace over which it exists."

Lichtenberg & Lieberman

### **Ergodic hypothesis** Neutrinos uniformly explore all accessible parts of flavor space. Johns, 2306.14982 (Thermodynamics of oscillating neutrinos)

"In a sense, ergodicity is universal, and the central question is to define the subspace over which it exists." Lichtenberg & Lieberman



Urquilla & Richers, 2401.01936

# If ergodicity is assumed, the iff relationship between **fast instabilities & angular crossings** becomes easy to prove.

Johns, 2402.08896 (*Ergodicity demystifies fast neutrino flavor instability*) Proofs without ergodicity: Morinaga 2022 and Fiorillo & Raffelt 2024

# If ergodicity is assumed, the iff relationship between **fast instabilities & angular crossings** becomes easy to prove.

Johns, 2402.08896 (*Ergodicity demystifies fast neutrino flavor instability*) Proofs without ergodicity: Morinaga 2022 and Fiorillo & Raffelt 2024

> The insufficiency of angular crossings for instability of the **fast flavor pendulum** is also illuminated by ergodicity: the infinite number of invariants restricts the ergodic subspace.

Johns, Nagakura, Fuller, & Burrows, 1910.05682 (Neutrino oscillations in supernovae: Angular moments and fast instabilities)

# IV. Neutrino quantum thermodynamics

Guided by ergodicity, equate coarse-grained averages with **ensemble expectation values**:

$$\langle \rho \rangle \cong \rho^{\mathrm{eq}}$$

Guided by ergodicity, equate coarse-grained averages with **ensemble expectation values**:

$$\langle \rho \rangle \cong \rho^{\rm eq}$$

 $\rho^{eq}$  maximizes entropy *S* subject to constraints. Assuming there are no fine-grained spatial correlations,

$$S = -V \int \frac{d^3 \boldsymbol{p}}{(2\pi)^3} \operatorname{Tr}\left[\rho_{\boldsymbol{p}} \log \rho_{\boldsymbol{p}} + (1 - \rho_{\boldsymbol{p}}) \log(1 - \rho_{\boldsymbol{p}})\right]$$

Guided by ergodicity, equate coarse-grained averages with **ensemble expectation values**:

$$\langle \rho \rangle \cong \rho^{\rm eq}$$

 $\rho^{eq}$  maximizes entropy *S* subject to constraints. Assuming there are no fine-grained spatial correlations,

$$S = -V \int \frac{d^3 \boldsymbol{p}}{(2\pi)^3} \operatorname{Tr}\left[\rho_{\boldsymbol{p}} \log \rho_{\boldsymbol{p}} + (1 - \rho_{\boldsymbol{p}}) \log(1 - \rho_{\boldsymbol{p}})\right]$$

Fix total energy & neutrino number at each *p* 

$$\rho_{\boldsymbol{p}}^{\mathrm{eq}} = \frac{1}{e^{\beta(H_{\boldsymbol{p}}^{\mathrm{eq}} - \mu_{\boldsymbol{p}})} + 1}$$

Pathway #3 to the miscidynamic equation goes through neutrino quantum thermodynamics:

 $\left(\left[H_{\boldsymbol{p}},\rho_{\boldsymbol{p}}\right]\right)^{\mathrm{eq}}\cong\left[H_{\boldsymbol{p}}^{\mathrm{eq}},\rho_{\boldsymbol{p}}^{\mathrm{eq}}\right]$ 

Negligible thermal fluctuations in the thermodynamic limit

Pathway #3 to the miscidynamic equation goes through neutrino quantum thermodynamics:

 $\left(\left[H_{\boldsymbol{p}},\rho_{\boldsymbol{p}}\right]\right)^{\mathrm{eq}}\cong\left[H_{\boldsymbol{p}}^{\mathrm{eq}},\rho_{\boldsymbol{p}}^{\mathrm{eq}}\right]$ 

Negligible thermal fluctuations in the thermodynamic limit

### Adiabaticity ( $\Delta S = 0$ ) makes *S*-maximization unnecessary:

- 1. Evolve |*P*| under collisions & advection.
- 2. Obtain new  $P^{eq}$  using |P| from step 1.

## V. Collisional flavor instabilities

# The asymmetry between neutrino & antineutrino interaction rates can cause collisional flavor instability.

Johns, 2104.11369 (Collisional flavor instabilities of supernova neutrinos)

# The asymmetry between neutrino & antineutrino interaction rates can cause collisional flavor instability.

Johns, 2104.11369 (Collisional flavor instabilities of supernova neutrinos)



Akaho, Liu, Nagakura, Zaizen, & Yamada 2024

### Estimating CFI effects is uniquely challenging because of kilometer-scale growth.

Johns & Nagakura, 2206.09225 [collision-related modeling issues] (*Self-consistency in models of neutrino scattering and fast flavor conversion*) Xiong, Wu, Martinez-Pinedo, Fischer, George, Lin, & Johns, 2210.08254 [a first attempt at estimates] (*Evolution of collisional neutrino flavor instabilities in spherically symmetric supernova models*)

### Estimating CFI effects is uniquely challenging because of kilometer-scale growth.

Johns & Nagakura, 2206.09225 [collision-related modeling issues] (*Self-consistency in models of neutrino scattering and fast flavor conversion*) Xiong, Wu, Martinez-Pinedo, Fischer, George, Lin, & Johns, 2210.08254 [a first attempt at estimates] (*Evolution of collisional neutrino flavor instabilities in spherically symmetric supernova models*)

### Analytic insights from idealized models:

### Slow flavor pendulum Gravitational field due to mass splitting

Hannestad, Raffelt, Sigl, & Wong 2006 [SNe] Johns & Fuller, 1709.00518 [early universe] (*Strange mechanics of the neutrino flavor pendulum*)

### Fast flavor pendulum Gravitational field due to spectral asymmetry

Johns, Nagakura, Fuller, & Burrows, 1910.05682 (*Neutrino oscillations in supernovae: Angular moments and fast instabilities*) Padilla-Gay, Tamborra, & Raffelt 2022 Fiorillo & Raffelt 2023

### Collisional flavor pendulum Adiabatic spin reversal due to collisional asymmetry

Johns & Rodriguez, 2312.10340 (*Collisional flavor pendula and neutrino quantum thermodynamics*) Miscidynamic adiabaticity

# Summary

### I. Foundations

- ➢ No evidence that QKE is inadequate.
- ➢ New type of many-body model: OILE.

### **II.** Paradigms

- Asymptotic-state models aren't self-consistent.
- Miscidynamics = local-equilibrium transport.

### **III. Flavor ergodicity**

A hypothetical but potentially useful concept (*e.g.*, understanding instability).

### IV. Neutrino quantum thermodynamics

Sheds new light on neutrino oscillations.

### V. Collisional flavor instabilities

Challenging even to estimate their effects.