

Non-modal effects in holographic models

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Based on:

2406.06685 with J. Carballo

2503.05871 with J. Carballo & C. Pantelidou

2505.04696 with P. Arnaudo & J. Carballo

THE
ROYAL
SOCIETY



Foundations and Applications
of Relativistic Hydrodynamics
GGI, 16 May 2025

Plan

- Introduction
- Perturbations of Poiseuille flow
- Perturbations of black holes in AdS
- (QNM orthogonality relations)
- Summary

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- **Introduction**
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This talk is about some basic questions
regarding perturbations of black holes

Through AdS/CFT: AdS black holes provide real-time, non-perturbative
description of a strongly coupled plasma

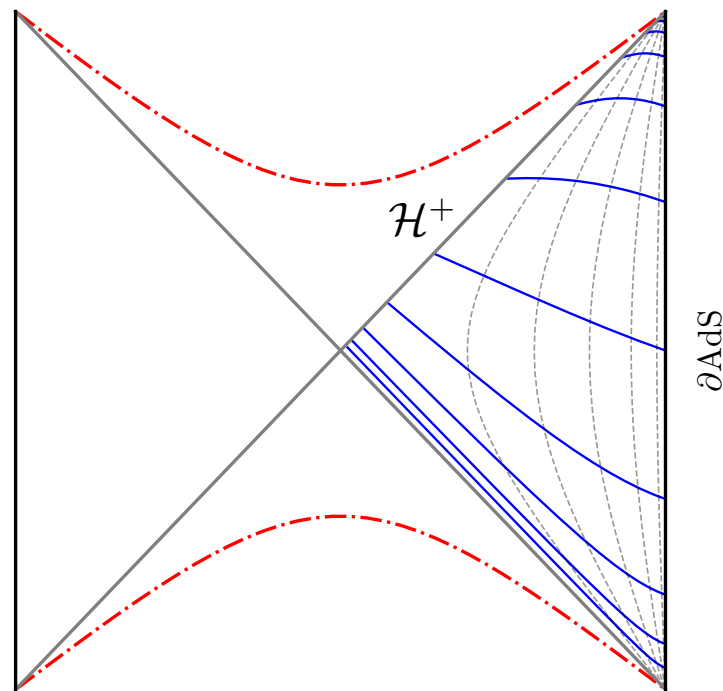
Focus on linear perturbations around thermal equilibrium

Often characterised by eigenfunctions of the Hamiltonian $\mathcal{H} = i\partial_t$
 $(\sim e^{-i\omega t})$

$$\omega(k) = c_s k + \dots \quad \text{sound}$$

$$\omega(k) = -iDk^2 + \dots \quad \text{shear}$$

$$\omega(k) = -i\Gamma + \dots \quad \text{non-hydrodynamic}$$

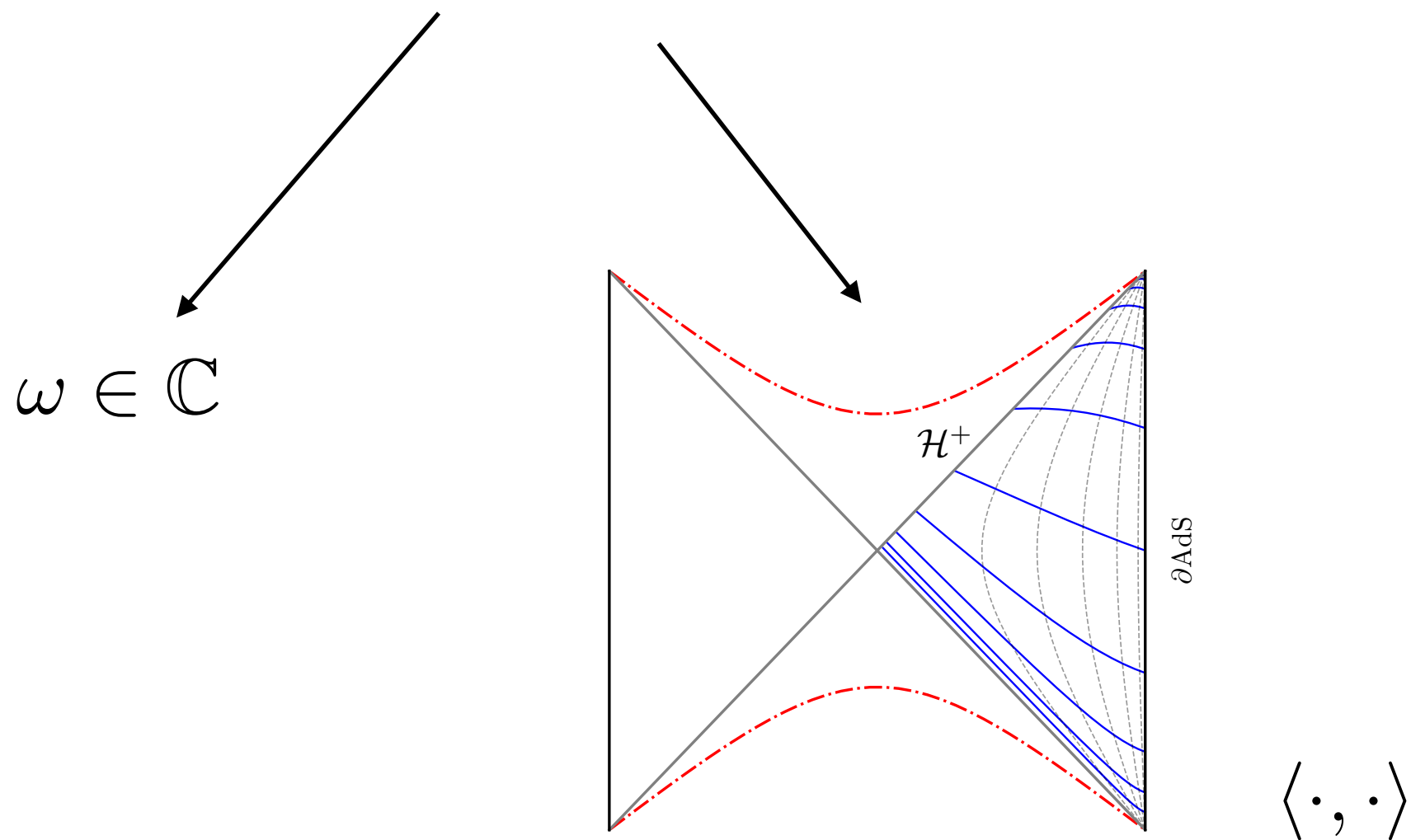


These are **quasinormal modes** (QNMs) of the black hole

- Ingoing at the horizon (**regularity**)
- Normalisable at the AdS boundary

For black holes $\mathcal{H} = i\partial_t$ is not Hermitian

This arises due to dissipation through the horizon



$$\langle a, \mathcal{H}b \rangle = \langle \mathcal{H}a, b \rangle + \text{fluxes through } \mathcal{H}^+$$

So, for black holes $\mathcal{H} \neq \mathcal{H}^\dagger$

Moreover the Hamiltonian is **non-normal**: $[\mathcal{H}, \mathcal{H}^\dagger] \neq 0$

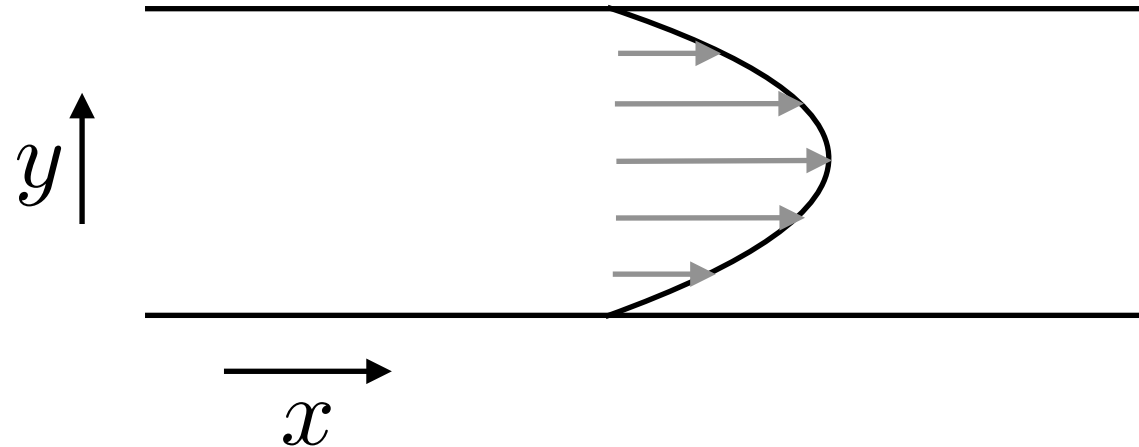
This is important because there exists
a complete, orthonormal basis made of (spectral theorem)
eigenfunctions iff $[\mathcal{H}, \mathcal{H}^\dagger] = 0$

- QNMs are generically not orthogonal to each other
- Leads to transient effects that a spectral analysis misses

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Plane Poiseuille flow



$$\vec{u} = (1 - y^2, 0, 0)$$

Linear perturbations

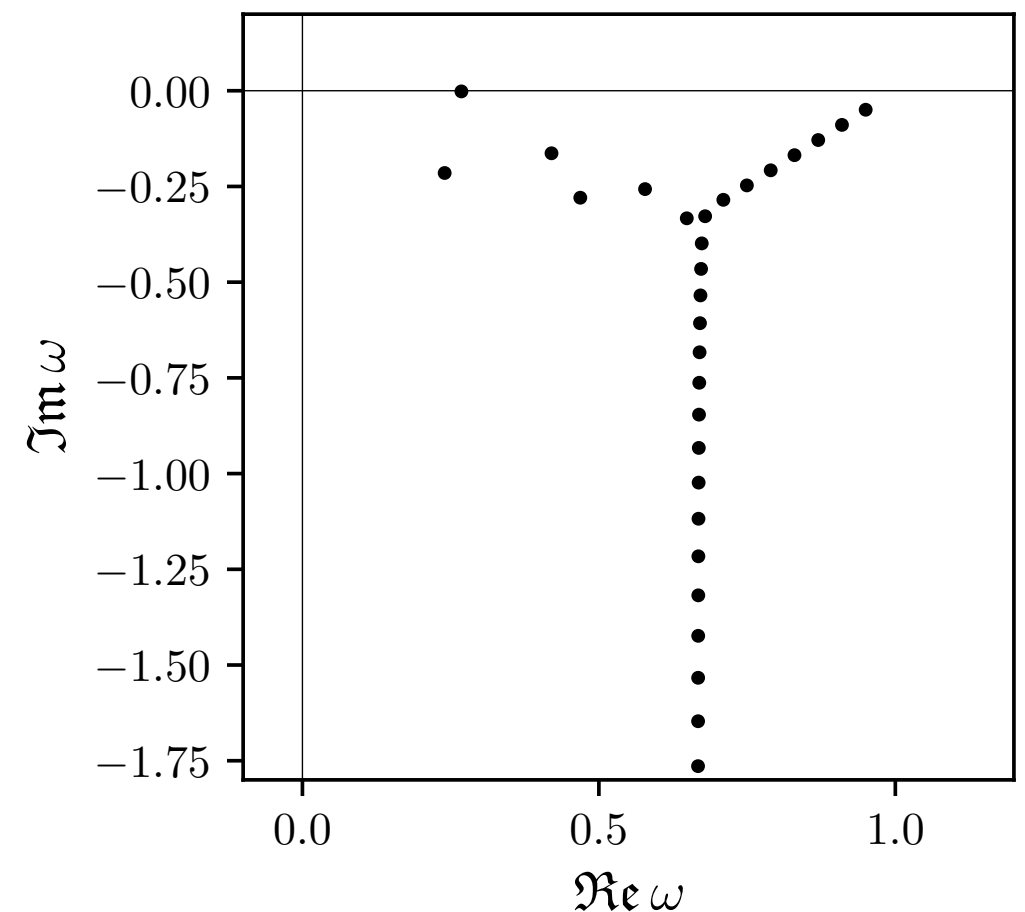
$$\vec{u} = (1 - y^2 + \partial_y \Phi, -\partial_x \Phi, 0)$$

Eigenvalue analysis

$$\Phi = \phi(y)e^{-i\omega t + i\alpha x}$$

$$\alpha O_{OS} \phi(y) = \omega \phi(y)$$

Spectrum at Re = 5000



Critical Re = 5772.22 [Orszag] (1971)

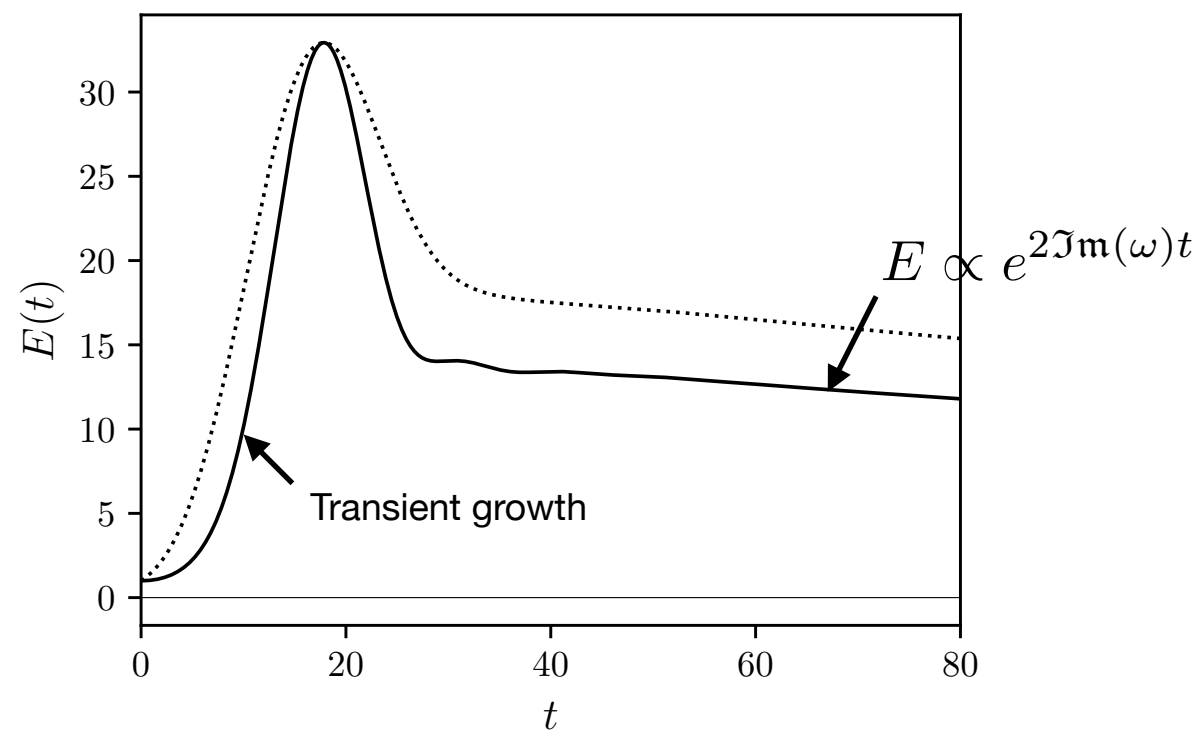
Evolution of linear perturbations

Kinetic energy $E = \int_{-1}^1 (|\partial_y \phi|^2 + \alpha^2 |\phi|^2) dy$

For a single mode:

$$E \propto e^{2\Im(\omega)t} \quad \begin{cases} \text{Re} < \text{Re}_c & \text{decay} \\ \text{Re} > \text{Re}_c & \text{growth} \end{cases}$$

For a sum of modes at $\text{Re} < \text{Re}_c$:



[Reddy, Schmid, Henningson] (1993)

Why?

\exists natural inner-product associated to E

$$\langle a, b \rangle = \int_{-1}^1 (\partial_y a^* \partial_y b + \alpha^2 a^* b) dy$$

s.t. $E[\phi] = \langle \phi, \phi \rangle$

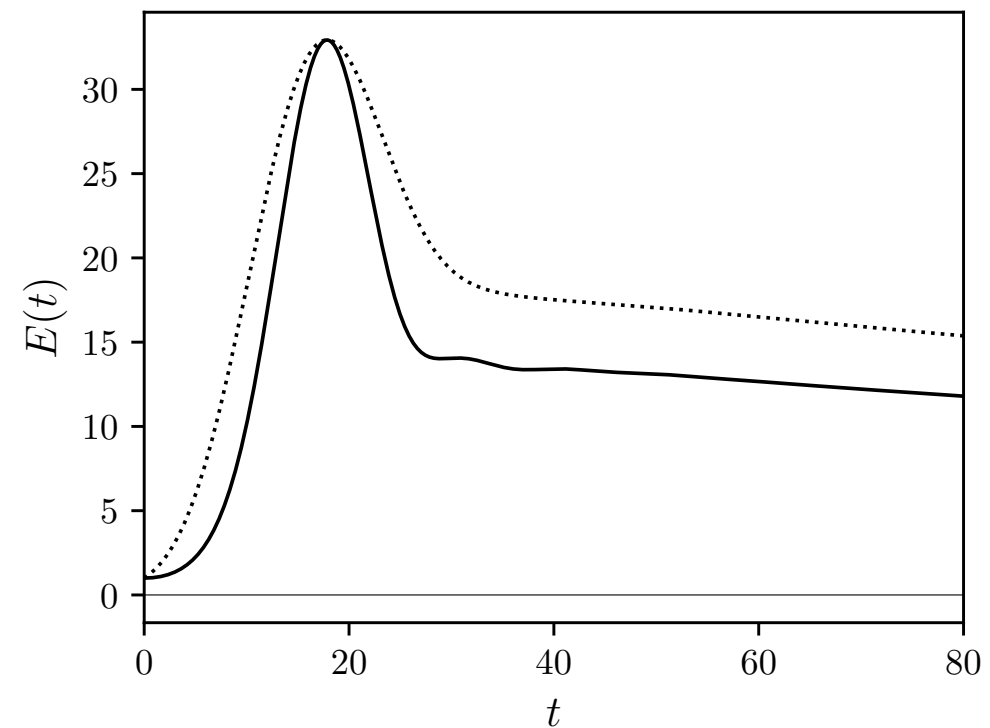
& the Hamiltonian O_{OS} is non-normal wrt to $\langle \cdot, \cdot \rangle$

In particular,

$$\langle \phi_n, \phi_m \rangle \neq 0$$

So that

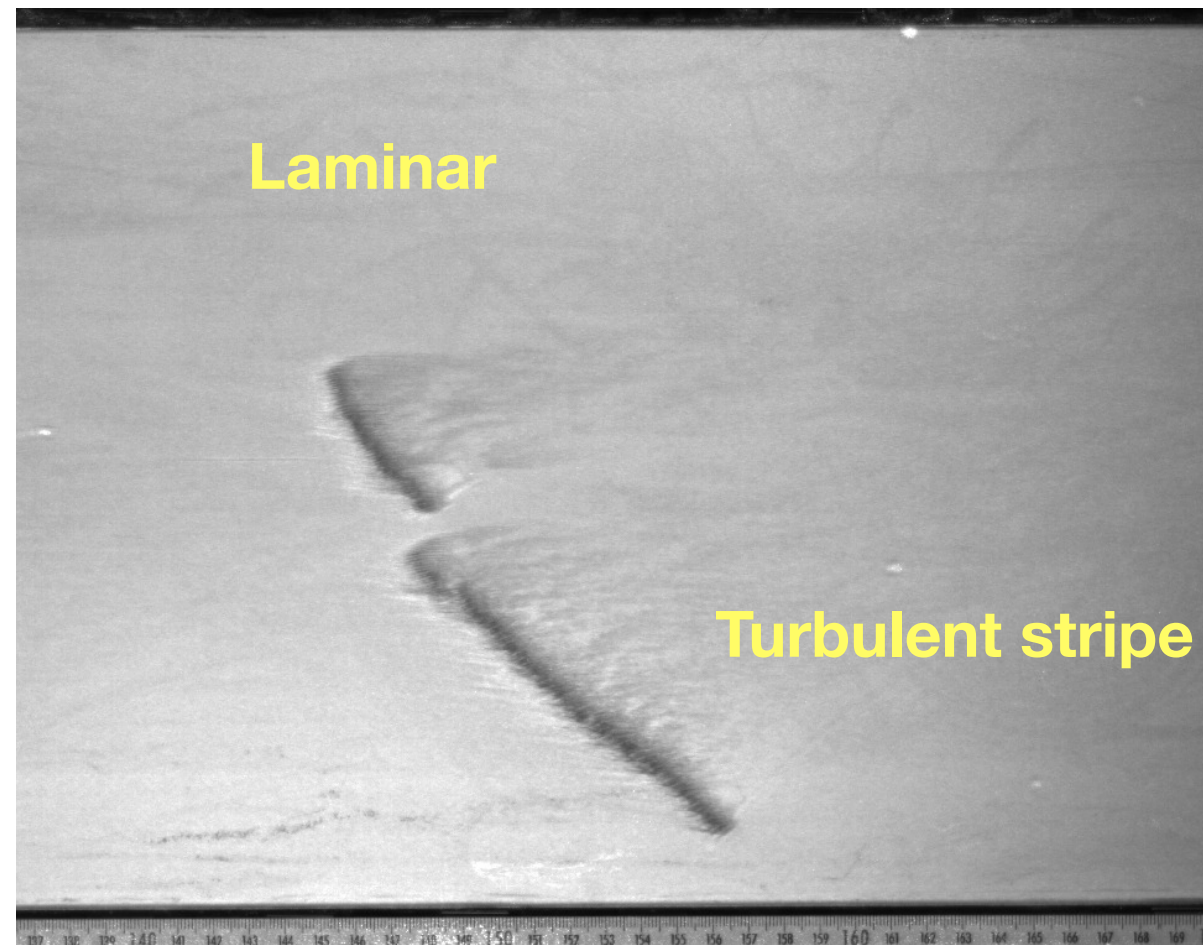
$$E \left[\sum_n c_n \phi_n \right] \neq \sum_n E [\phi_n]$$



Phenomenological importance

Linear perts ultimately decay

Nonlinearities

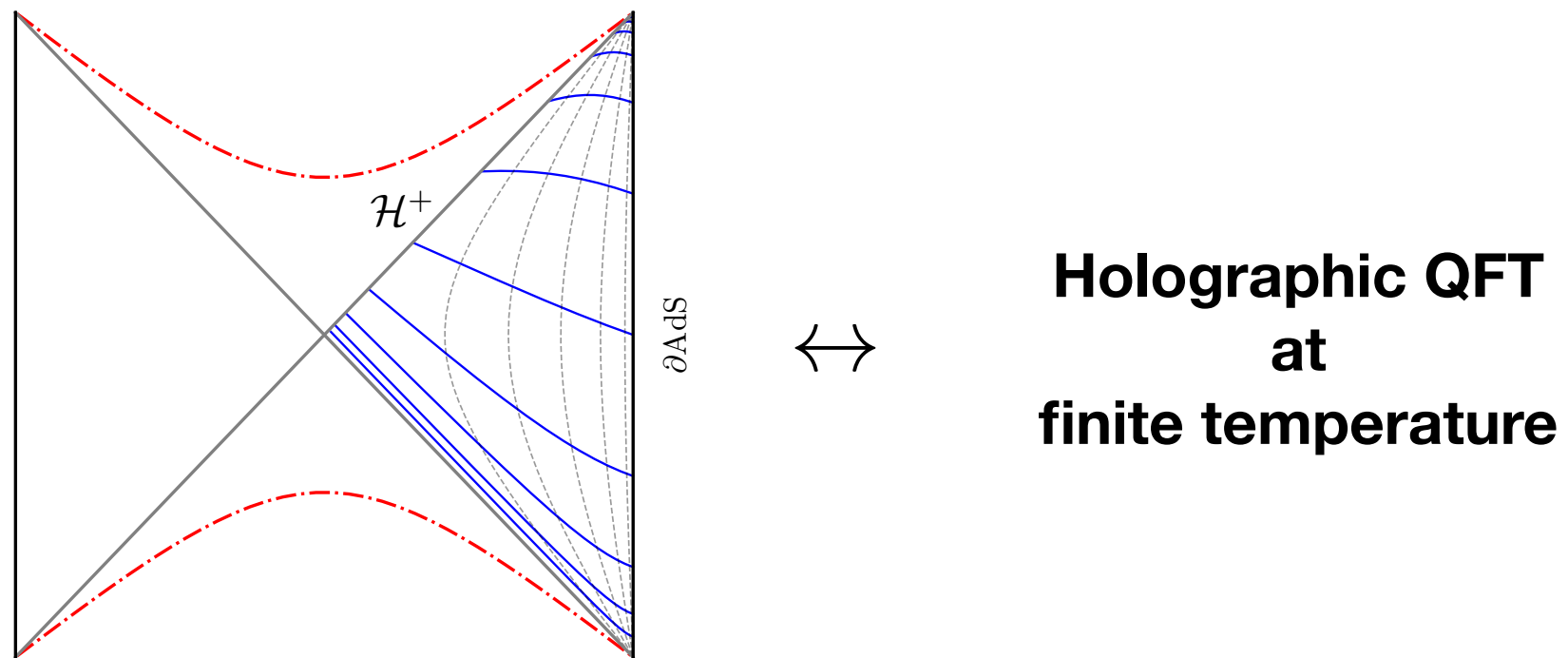


*From 'Onset of turbulence in plane Poiseuille flow'
C Paranjape, PhD Thesis (2019)*

~ Sustained nonlinear structures can form,
provided suitably large amplitude initial data

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$$E_\psi = \int_{\Sigma_\tau} (T_\psi)^\mu{}_\tau n_\mu d\Sigma_\tau$$

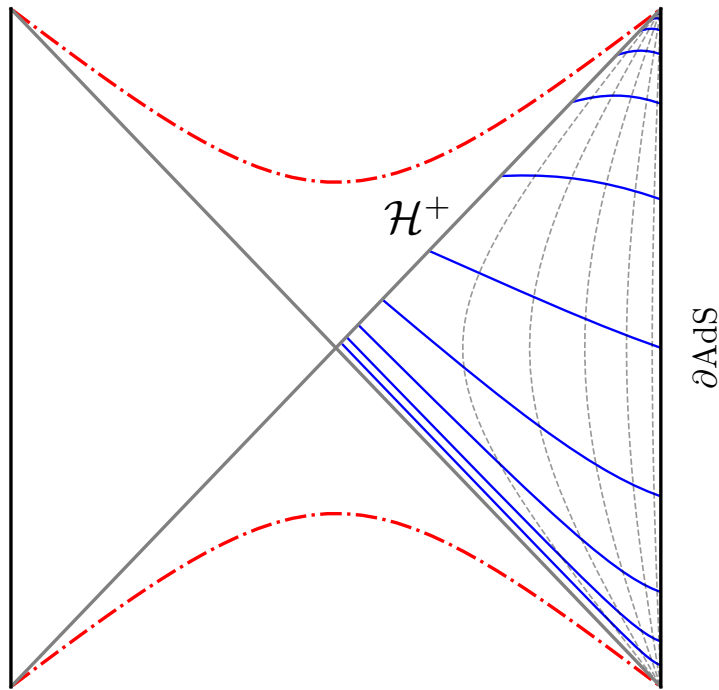
$$\langle a, b \rangle = \dots$$

\mathcal{H} is non-normal with respect to $\langle \cdot, \cdot \rangle$

Analytic example (BTZ)

$$\phi_n(\tau, z) = e^{-(2n+\Delta)\tau} (z+1)^{-n-\frac{\Delta}{2}} {}_2F_1\left(-n, -n; 1-2n-\Delta; 1-z^2\right)$$

$$\omega_n = -i(\Delta + 2n),$$



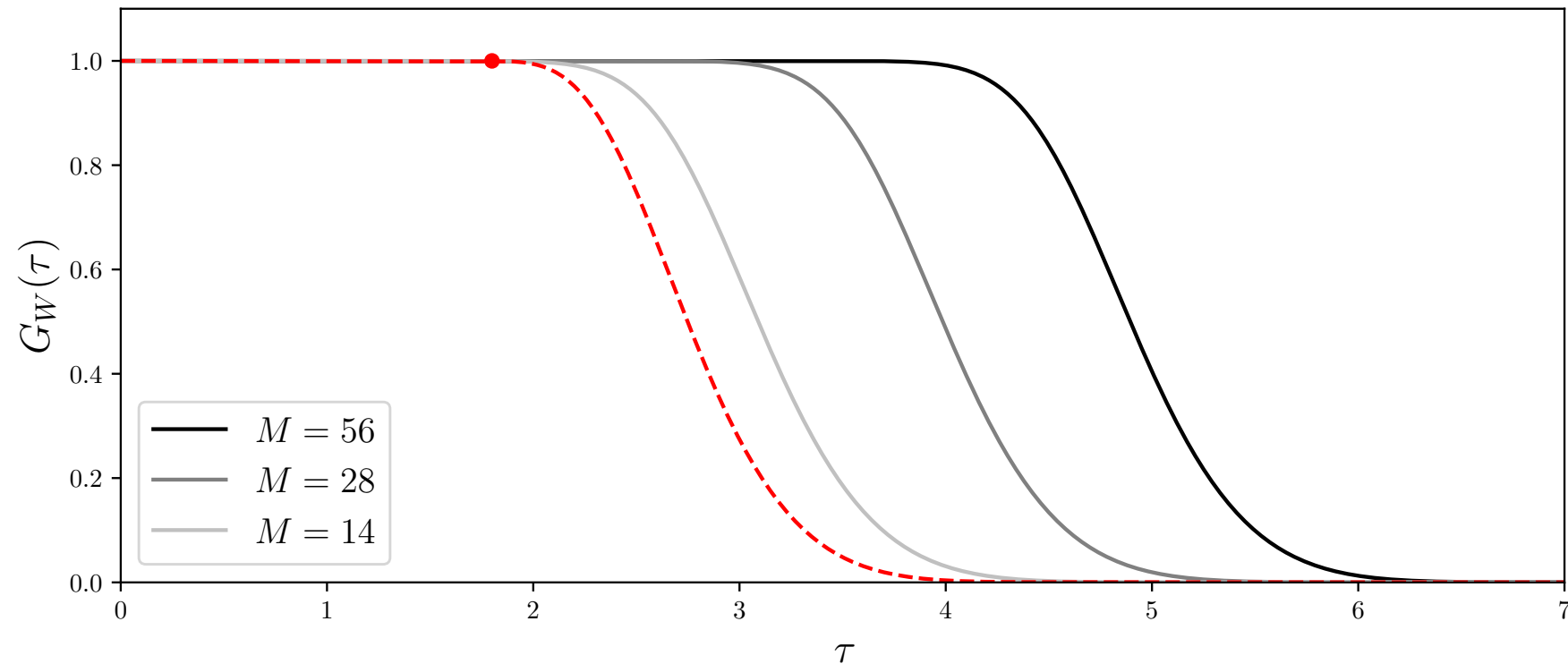
2 modes

$$\phi(\tau, z) = a_1 \frac{e^{-2\tau}}{1+z} + a_2 \frac{e^{-4\tau}(2+z^2)}{3(1+z)^2},$$

$$E(\tau) = \frac{1}{4}|a_1|^2 e^{-4\tau} + \frac{1}{6}(a_1^* a_2 + a_1 a_2^*) e^{-6\tau} + \frac{1}{8}|a_2|^2 e^{-8\tau}.$$

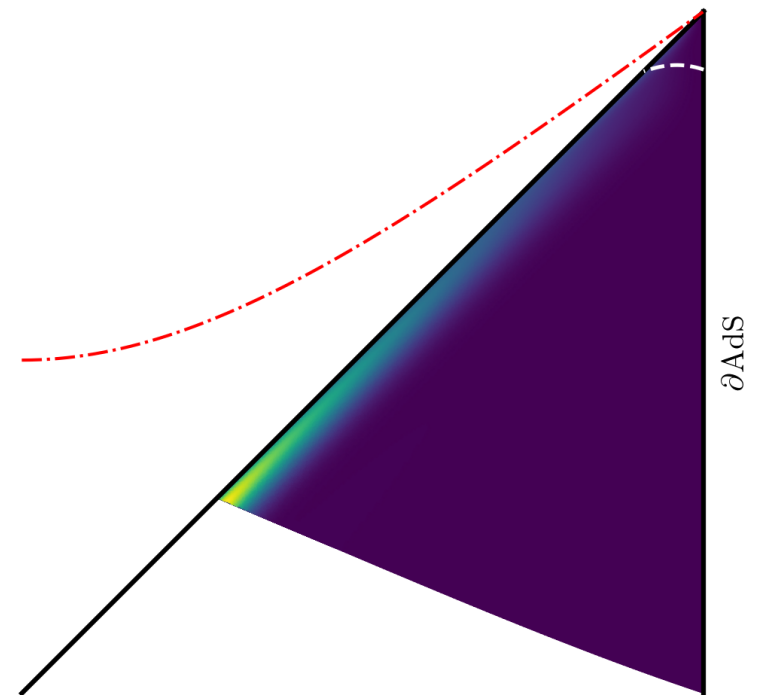
$$\frac{E(\tau)}{E(0)} = 1 - 24 \frac{|a_1 + a_2|^2}{6|a_1|^2 + 4(a_1^* a_2 + a_1 a_2^*) + 3|a_2|^2} \tau + O(\tau)^2$$

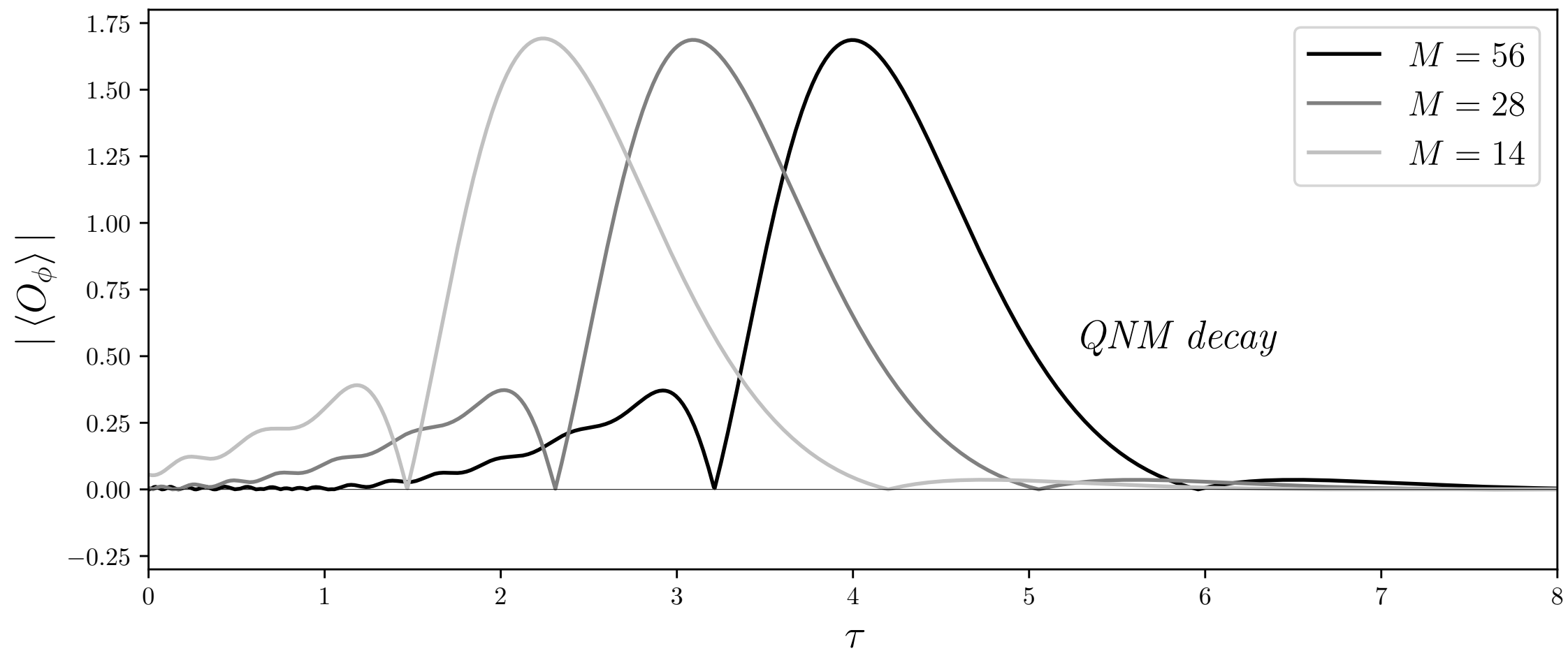
M modes

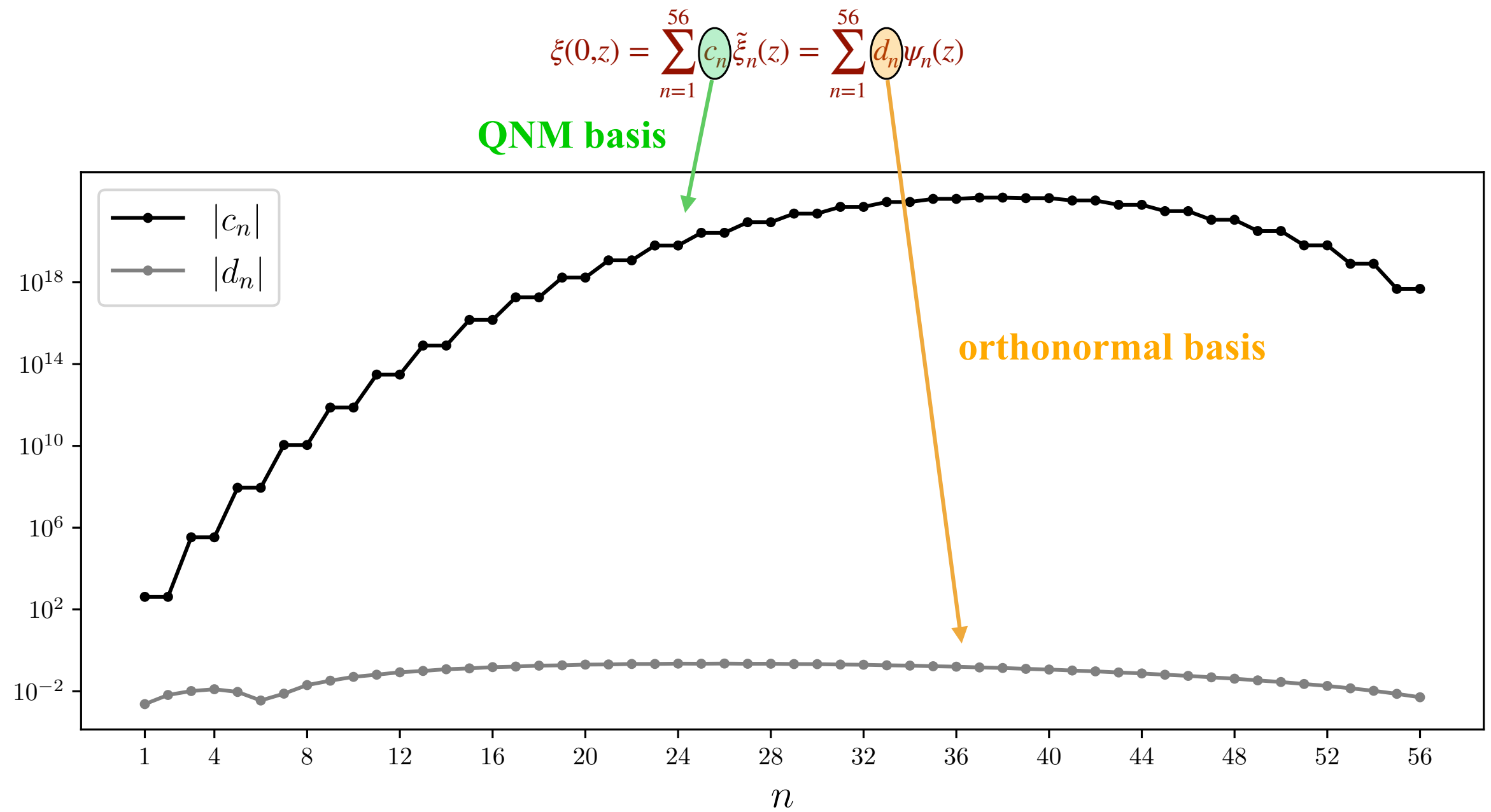


- More modes allows for longer periods of sustained E
- States which take an arbitrarily long time to thermalise
- Non-hydrodynamic in nature
- Holds for Schwarzschild- AdS_{d+1} , dS_{d+1} , Schwarzschild

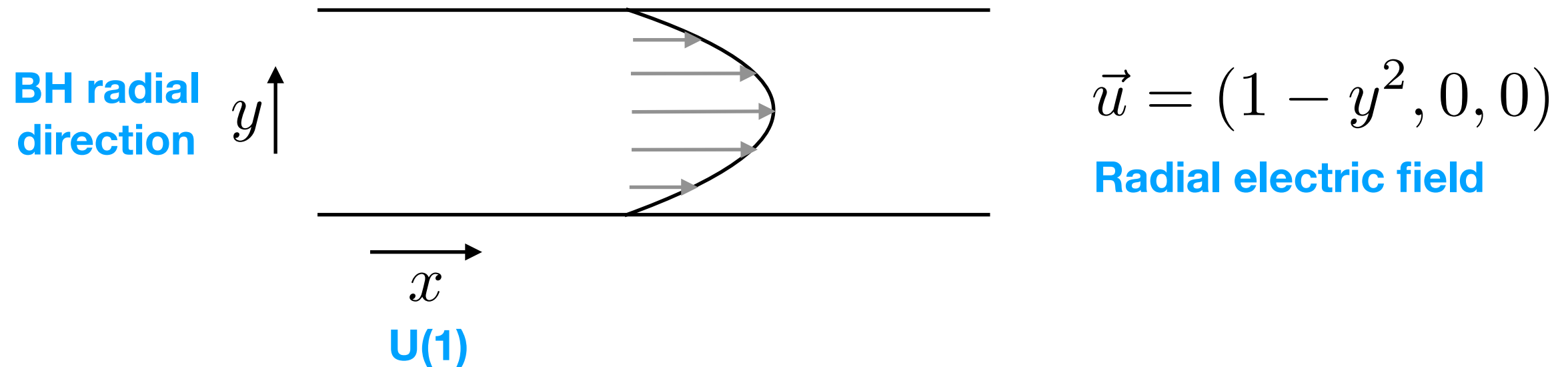
$$\tau \sim \log M$$







But what about *growth*?



$$\Phi = \phi(y) e^{-i\omega t + i\alpha x}$$

Charged scalar QNMs

**RN-AdS w/
charged scalar**

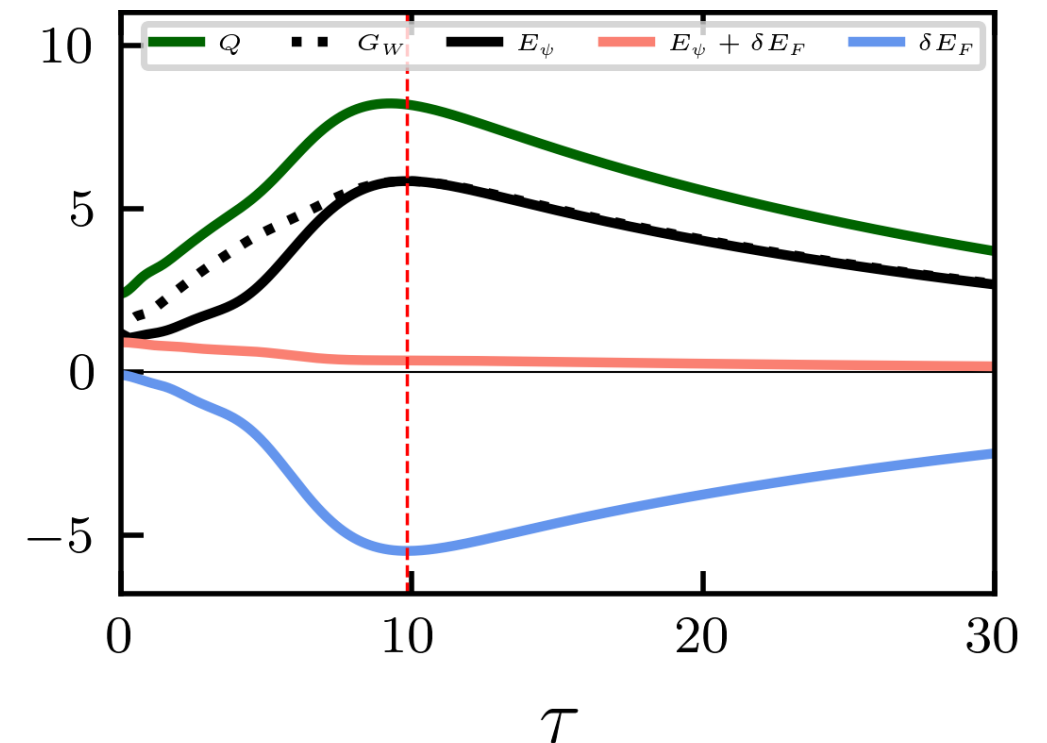
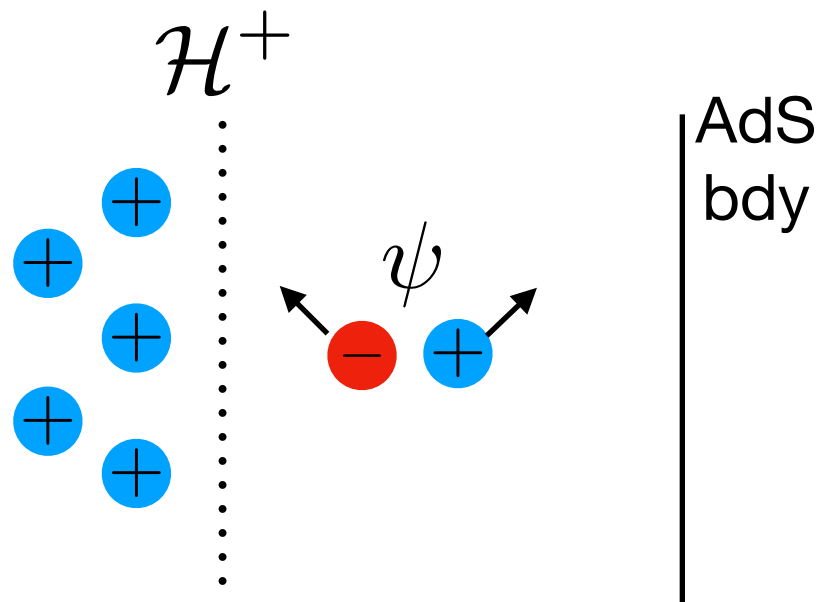


**Holographic
superfluid**

Energy growth from superradiance

Physical mechanism

Classical wave analogue of pair production

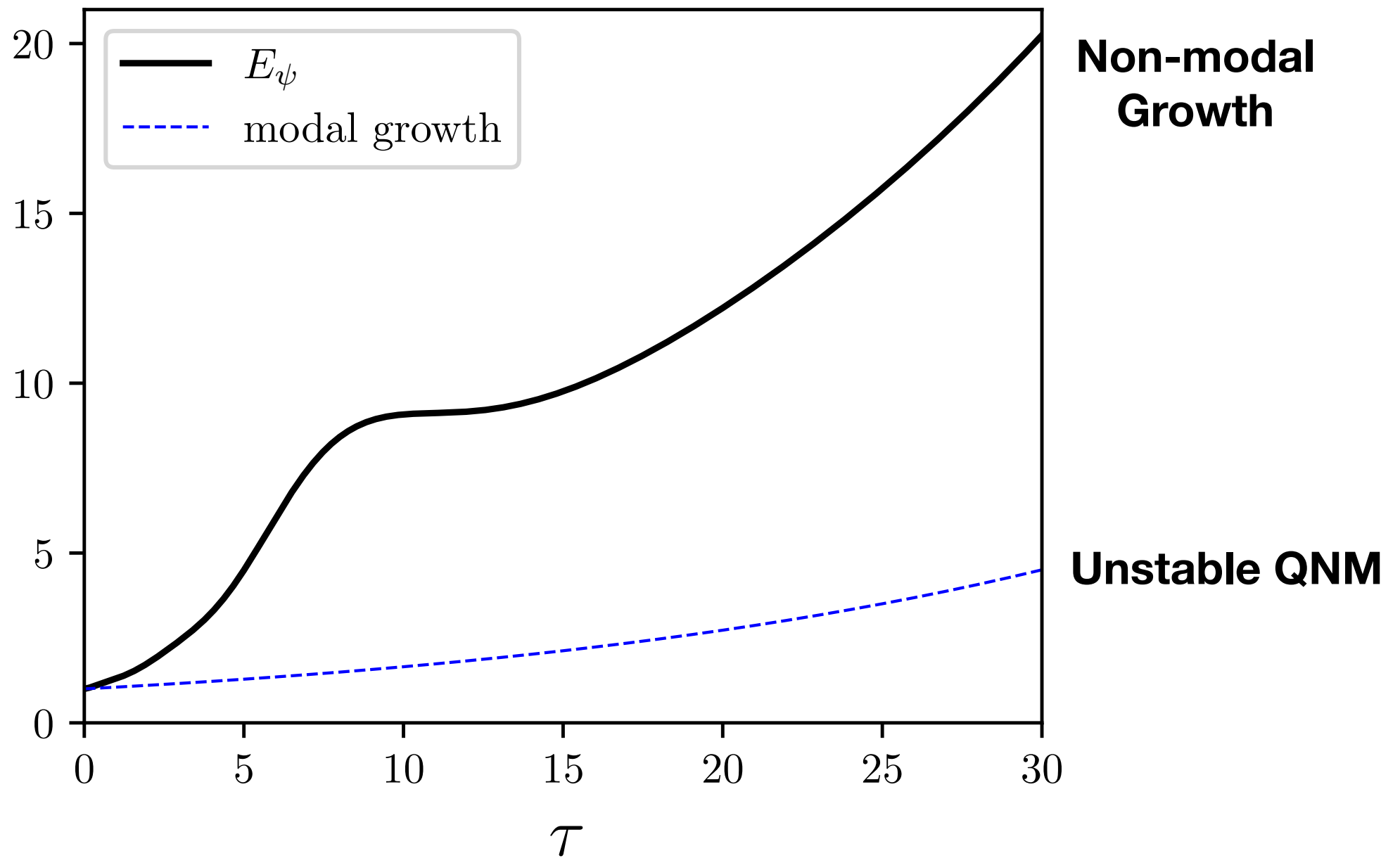


$T < T_c$ Runaway process (to broken phase)

$T > T_c$ Process still occurs, but transiently

“Transient superradiance”

Even when $T < T_c$ the non-modal instability is faster than the QNM



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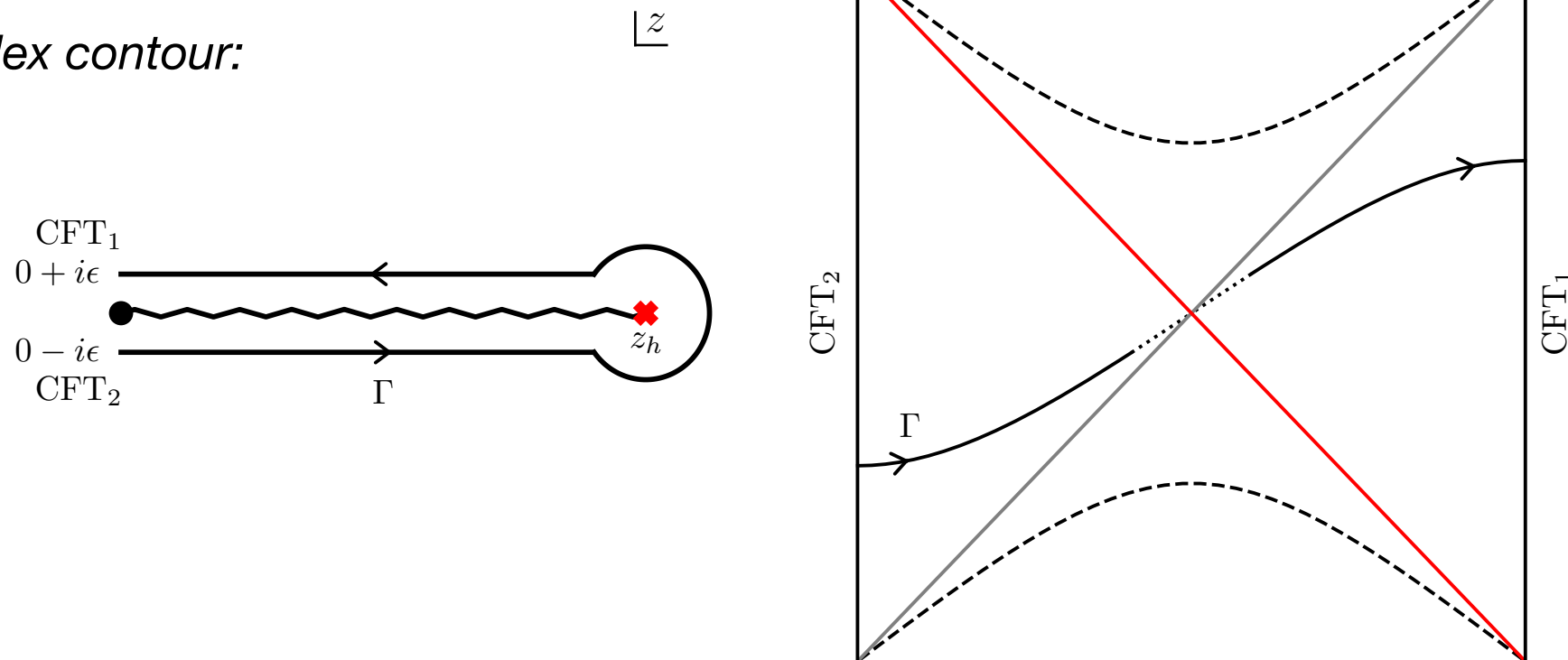
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Aside: QNM orthogonality relations

Constructions: de Sitter [Jafferis, Lupsasca, Lysov, Ng, Strominger] (2013)
Kerr [Green, Hollands, Sberna, Toomani, Zimmerman] (2022)
AdS black holes [Arnaudo, Carballo, BW] (2025)

Main ingredients:

- *Complex contour:*



- *Discrete symmetry operators:* $\langle a, b \rangle_{\text{ortho.}} \equiv \langle \mathcal{CPT} a, b \rangle_{\text{KG}}$

Orthogonality... but with respect to a bilinear (not inner product)

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Summary

- In **non-normal** systems, eigenvalue analysis misses transient effects
- Perturbations of BHs are non-normal
- Govern perturbations of strongly coupled plasmas
- We identified transient effects:
 - States which take an arbitrarily long time to thermalise
 - Holographic superconductors are non-modally unstable for $T > T_c$
(just like transition to turbulence for water in a pipe)

Thank you for your attention!