Introduction to accelerator physics

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Accelerators - A Window on Nature

- Particle accelerators provide the source for most high energy physics experiments
 - Provide high luminosity, high energy beams for colliders
 - Provide high brightness beams for secondary particle production
 - Also key technology for life sciences, engineering, chemistry
- How do they work?
 - How can we get to high energy?
 - How can we keep the beam in the accelerator?
 - How can we get to high luminosity?



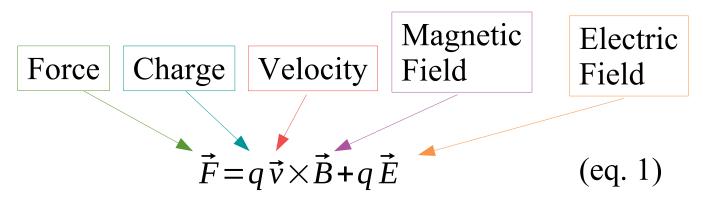
Accelerator Components

- Most accelerators share similar components
- Main components of an accelerator
 - Bending dipoles
 - Focussing quadrupoles
 - Acceleration RF cavities
- Important equipment I will not discuss
 - Vacuum
 - Diagnostics
 - Targets for secondary particle production



Lorentz force law

Fundamental equation for particles moving through fields

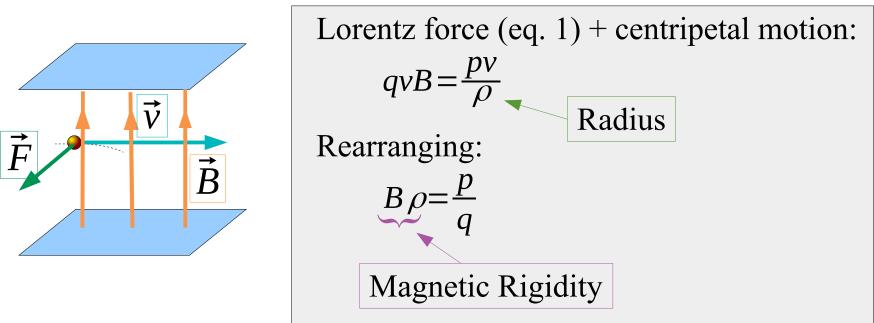


- Magnetic force is perpendicular to velocity
 - Magnetic field conserves energy
- Electric force is weaker by factor velocity
 - Magnets are better for bending and focussing



Magnetic Rigidity and Bending

- Simplest magnet "dipole"
 - Uniform magnetic field perpendicular to beam direction



- Constant force \rightarrow constant curvature \rightarrow circular motion
- Magnetic rigidity parameterises momentum
- Charge-to-mass ratio important when accelerating multiple particle species



Worked example – LHC

- If we wanted to accelerate, say, 7 TeV particles, what bending radius is required?
- Maximum dipole field around 8.3 T

$$B\rho = \frac{p}{q}$$

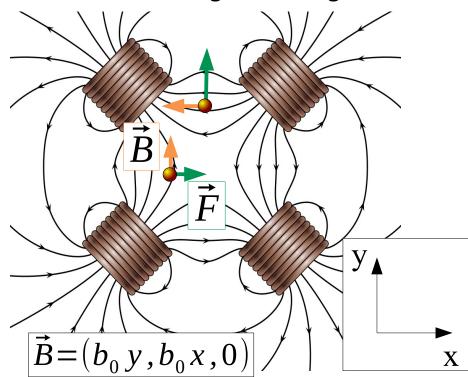
$$\rho = \frac{p}{qB} = \frac{7}{0.3 \times 8.3} = 2.8 \, km$$

- Nb: LHC radius ~ 4.1 km
 - Need space for detectors, etc



Quadrupole magnets

- If we only had bending magnets, particles would soon be lost from the accelerator
- Need to keep the particles in the accelerator using focussing elements
 - Usually use quadrupoles
- Field stronger away from beam centre
 - Like a spring or pendulum
 - Simple harmonic motion
- "F" quad focuses in x and defocuses in y
- "D" quad focuses in y and defocuses in x
- Overall focussing by alternating "F" and "D"
 - Just reverse the field





Quadrupole field – horizontal (1)

- For a particle moving near to the z-axis $\vec{F} = q \vec{v} \times \vec{B} + q \vec{E}$
- Considering only p_x for now

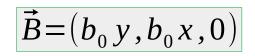
$$\frac{dp_x}{dt} = q \frac{dz}{dt} B_y$$

Use the chain rule

$$\frac{dp_x}{dt} = \frac{dp_x}{dz}\frac{dz}{dt}$$

Combining these equations:

$$\frac{dp_x}{dz} = q b_0 x$$



$$\vec{v} \times \vec{B} = \begin{pmatrix} v_z B_y \\ -v_z B_x \\ 0 \end{pmatrix}$$



Quadrupole field – horizontal (2)

$$\frac{dp_x}{dz} = q b_0 x \qquad \textcircled{\circ}$$

Definition of x-component of momentum

$$p_x = m \gamma v_x = m \gamma \frac{dz}{dt} \frac{dx}{dz} = p_z \frac{dx}{dz}$$

Substitute this definition into

$$p_z \frac{d^2 x}{dz^2} = q b_0 x$$

Rearrange and wrap up constant terms in focussing strength k

$$\frac{d^2x}{dz^2} - kx = 0$$



Quadrupole field – vertical

• Lorentz force law with quadrupole field definition $dp_{...}$

$$\frac{dp_y}{dt} = -q b_0 v_z y$$

Use chain rule and eliminate vz

$$p_z \frac{d^2 y}{dz^2} = -q b_0 y$$

Rearrange and wrap up constant terms in defocussing strength k

$$\frac{d^2 y}{dz^2} + k y = 0$$





Motion is governed by $\frac{d^2 x}{dz^2} - kx = 0 \qquad \frac{d^2 y}{dz^2} + ky = 0$

• This is simple harmonic motion – solutions are of form $x = x_0 \cos(\sqrt{k} z) + \frac{dx_0}{dz} \frac{1}{\sqrt{k}} \sin(\sqrt{k} z)$

• Taking derivative

$$\frac{dx}{dz} = -x_0 \sqrt{k} \sin(\sqrt{k} z) + \frac{dx_0}{dz} \cos(\sqrt{k} z)$$

For y

$$y = y_0 \cosh(\sqrt{k} z) + \frac{dy_0}{dz} \frac{1}{\sqrt{k}} \sinh(\sqrt{k} z)$$
$$\frac{dy}{dz} = y_0 \sqrt{k} \sinh(\sqrt{k} z) + \frac{dy_0}{dz} \cosh(\sqrt{k} z)$$



Transfer Matrix

Just thinking about x, the particles move according to

$$x_1 = x_0 \cos(\sqrt{k} z) + \frac{dx_0}{dz} \sin(\sqrt{k} z)$$
$$\frac{dx_1}{dz} = -x_0 \sqrt{k} \sin(\sqrt{k} z) + \frac{dx_0}{dz} \sqrt{k} \cos(\sqrt{k} z)$$

• We can rewrite this as a matrix

$$\begin{pmatrix} x_1 \\ \frac{dx_1}{dz} \end{pmatrix} = \begin{pmatrix} \cos(\sqrt{k}z) & \frac{1}{\sqrt{k}}\sin(\sqrt{k}z) \\ -\sqrt{k}\sin(\sqrt{k}z) & \cos(\sqrt{k}z) \end{pmatrix} \begin{pmatrix} x_0 \\ \frac{dx_0}{dz} \end{pmatrix}$$

This matrix is known as the quadrupole's transfer matrix $\underline{u}_1 = M_{01} \underline{u}_0$





- Exercise what is the transfer matrix for a drift space, that is a region with no fields at all?
 - What is the force acting on the particle?
 - What is x(z) in terms of dx_0/dz and x_0
 - What is dx/dz in terms of dx_0/dz
 - Now write that as a matrix



- Exercise what is the transfer matrix for a drift space?
 - What is the force acting on the particle?
 - No force
 - What is x(z) in terms of dx₀/dz and x₀ $x = x_0 + \frac{dx_0}{dz} z$



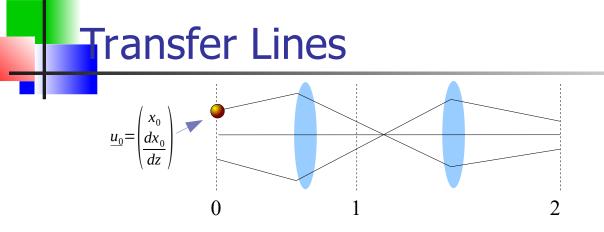
What is dx/dz in terms of dx₀/dz

$$\frac{dx}{dz} = \frac{dx_0}{dz}$$

Now write that as a matrix

$$\begin{pmatrix} x \\ \frac{dx}{dz} \end{pmatrix} = \begin{pmatrix} 1 & z \\ 0 & 1 \end{pmatrix} \begin{pmatrix} x_0 \\ \frac{dx_0}{dz} \end{pmatrix}$$





- Transfer matrix defines transport through a region
- Transfer matrices can be combined by multiplication
- Say we have transfer matrices like:

$$\underline{u_1} = \boldsymbol{M_{01}} \underline{u_0}$$
$$\underline{u_2} = \boldsymbol{M_{12}} \underline{u_1}$$

Then

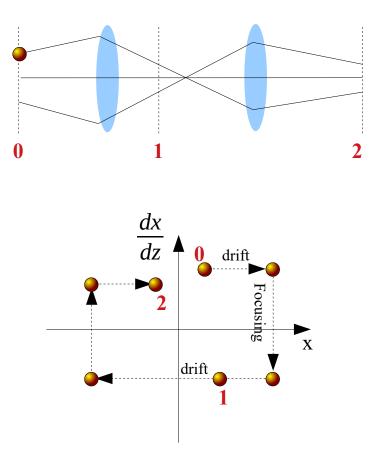
$$\underline{u}_2 = \boldsymbol{M}_{12} \boldsymbol{M}_{01} \underline{u}_0$$

• i.e. we can define a combined transfer matrix like $M_{02} = M_{12} M_{01}$



Phase space

 Another instructive way to look at beam optics is by considering the phase space





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Symplecticity

- There is a general rule for what transfer matrices are allowed by equations of motion
 - "Symplectic condition"
- Formally a matrix \mathbf{M} is *symplectic* if it satisfies $\mathbf{M}^{\mathrm{T}} \mathbf{S} \mathbf{M} = \mathbf{I}^{\mathsf{T}}$

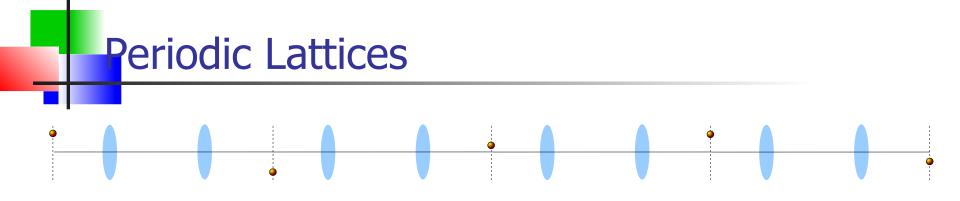
Where

$$\boldsymbol{S} = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix}$$

It can be shown that any *symplectic* matrix **M** can be written as

 $\boldsymbol{M} = \boldsymbol{I} \cos \mu + \boldsymbol{J} \sin \mu$ $\boldsymbol{J} = \begin{pmatrix} \alpha & \beta \\ -\gamma & -\alpha \end{pmatrix} \text{ with } \gamma \beta - \alpha^2 = 1 \text{ and } \boldsymbol{I} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$





Following *n* identical cells or turns in a ring with one-turn matrix
 M

$$\underline{u}_n = \boldsymbol{M}^n \underline{u}_0$$

Rewrite

$$M = I \cos \mu + J \sin \mu$$

$$J = \begin{pmatrix} \alpha & \beta \\ -\gamma & -\alpha \end{pmatrix} \text{ with } \gamma \beta - \alpha^2 = 1 \text{ and } I = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$

So
$$J^2 = -I$$

And

$$M^n = I \cos(n\mu) + J \sin(n\mu)$$



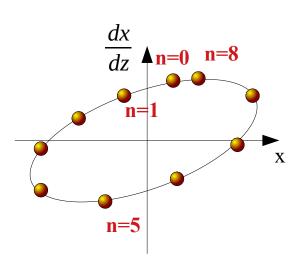
Periodic Lattices

What does this mean?

 $\boldsymbol{M}^{\boldsymbol{n}} = \boldsymbol{I}\cos(n\,\mu) + \boldsymbol{J}\sin(n\,\mu)$

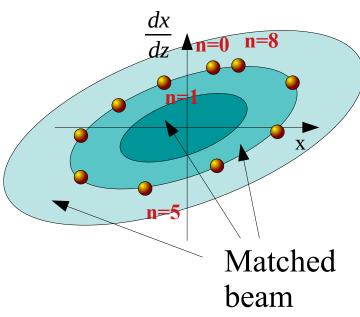
- Particles move around an ellipse in phase space if Trace(M) < 2
- µ is the "phase advance"
 - Sometimes use "tune" ... $2\pi v = \mu$
- α, β and γ are "Twiss parameters"
 - Tell us the alignment of the ellipse
- Each particle sits on ellipse area ε the particle's amplitude



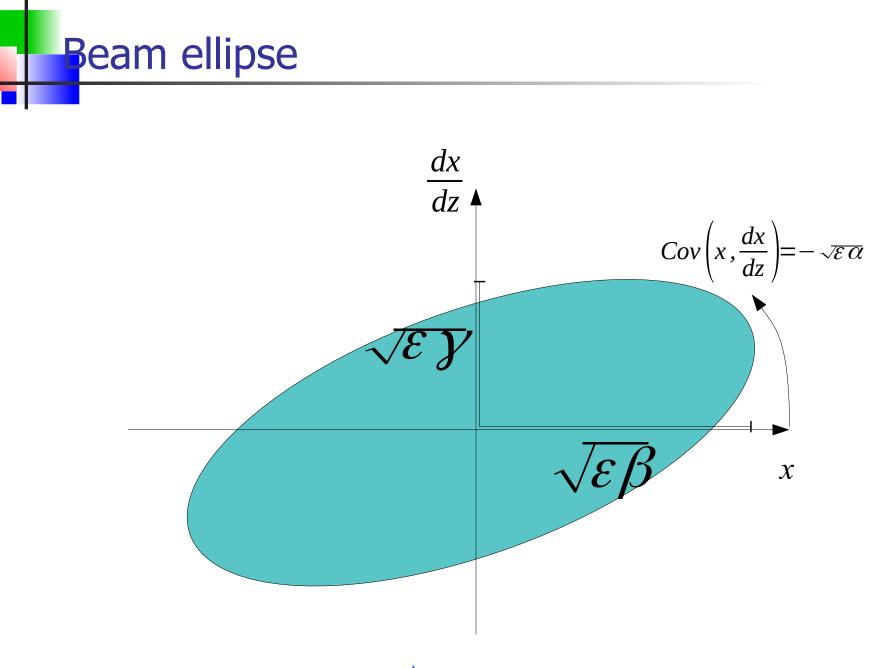


Periodic Lattices and beams

- Beam is composed of many particles
 - Particles occupy a region in phase space
- "Emittance" is area occupied by the entire beam
- Sometimes classify "RMS emittance"
 - Area occupied by ellipse 1 RMS distance from beam centre
- Low emittance is crucial for
 - High luminosity
 - Low losses









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- What is behaviour of particles in phase space if
 - Trace(M) < 2</p>
 - Trace(M) = 2
 - Trace(M) > 2



- What is behaviour of particles in phase space if
 - Trace(M) < 2
 - Motion is an ellipse
 - Trace(M) = 2
 - x → +/- x
 - Trace(M) > 2
 - Motion is a hyperbola

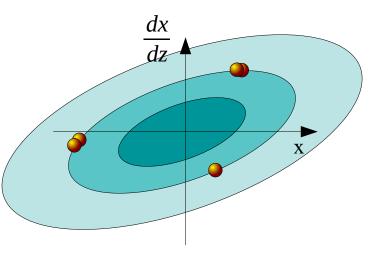


Emittance Growth

- Ideally emittance is conserved, but this is not always the case
- Long list of effects that can cause emittance growth
 - Beam mismatch
 - Scattering off residual gas
 - Scattering off particles in the same beam
 - Scattering off particles in other beams (e.g. in collider)
 - Space charge
 - Resonances



Resonances



- Reminder:-
 - Tune v is number of SHM oscillations per turn
 - Phase advance $\mu = 2\pi v$ is "angle" advanced per turn
- The beam does not behave well when

$$v_x = l + \frac{m}{n}$$
 integer

- Beam passes through the same field region every nth turn
- Imperfections in the field get amplified
- Resonance



Resonances $v_x = 4 + \frac{1}{2} = 4.33$ Can see poor performance for Only a very small area in phase space is transmitted (v_x, v_y) In fact, a 2D phenomenon in Bad transmission hot) ISIS Synchrotron (measured) Good transmission (cold) 3.95 3.90 3.85 \mathcal{V}_{v} 3.80 3.75 3.70

3.65

4.10

4.15

4.20

4.25

4.30

4.35

4.40

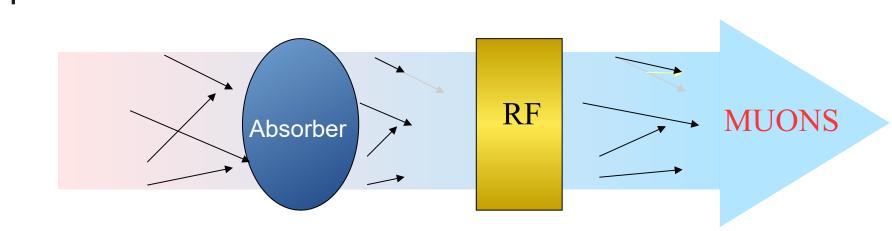
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Emittance Reduction (Cooling)

- Several techniques to reduce emittance
 - Synchrotron radiation cooling
 - Stochastic cooling
 - Laser cooling
 - Electron cooling
 - Ionisation cooling
- Fundamental principle is to remove "heat" from the beam using a neighbouring heat sink
 - Comoving electron beam \rightarrow electron cooling
 - Comoving laser \rightarrow laser cooling
 - Emission of synchrotron radiation
 - Photon emission caused by (principally) electrons bending in magnetic field



E.g. Ionisation Cooling



- Beam loses energy in absorbing material
 - Absorber removes momentum in all directions
 - RF cavity replaces momentum only in longitudinal direction
 - End up with beam that is more straight
- Multiple Coulomb scattering from nucleus ruins the effect
 - Mitigate with tight focussing
 - Mitigate with low-Z materials
 - Equilibrium emittance where MCS completely cancels the cooling



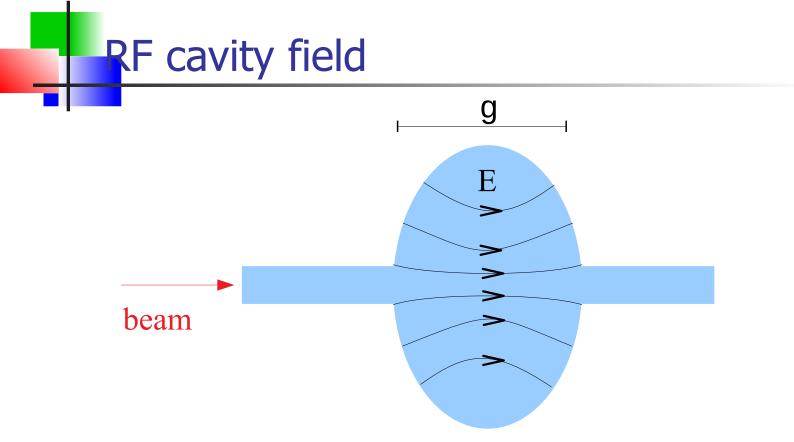
Longitudinal Dynamics and Acceleration



Longitudinal Dynamics

- So much for transverse motion (i.e. x and y planes)
- What about energy and acceleration?
- Electrostatic acceleration limited by breakdown potential
 - Change in energy is given by voltage differential
 - High voltage differentials cause breakdown (sparks)
 - Practically limits electrostatic acceleration to few MeV
- To accelerate beyond MeV require oscillating electric field
- RF Cavities





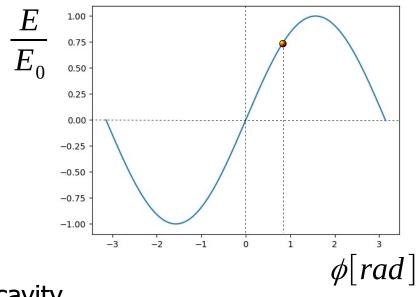
RF cavity holds a resonating EM wave

• Recall Lorentz force law
$$\vec{F} = q \vec{v} \times \vec{B} + q \vec{E}$$

Force is in direction of motion - energy changes!



RF cavity field



In RF cavity

$$\vec{E} = E_0 \sin(\omega t + \phi) \vec{\hat{z}}$$

• Energy change of particle crossing at ϕ

$$\delta W = \int F \, dz = q \, g \, E_0 \sin(\phi)$$

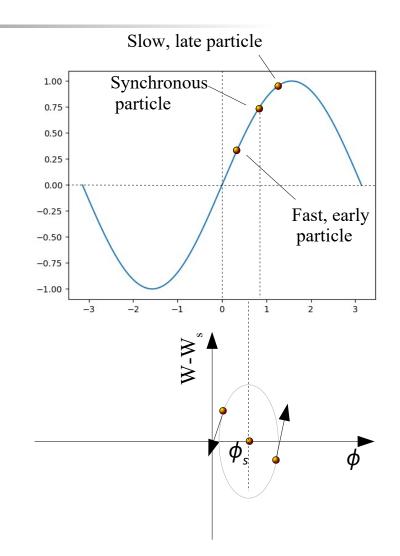
- g is the gap length
- Assumes gap is short so electric field doesn't change much
 - For longer gaps, can introduce an effective gap length g T
 - T is the "transit time factor" \rightarrow reduces the effective gap length

Phase stability

- Phase cavities so that a "synchronous" particle always crosses at phase φ_s
- Particle crossing at phase ϕ relative to synchronous particle

 $\delta W = q g T E_0 \sin(\phi + \phi_s)$

- Particle arriving early
 - Fast
 - t negative
 - Gets smaller energy kick
 - Ends up relatively slower
- Particle arriving late
 - Slow
 - t positive
 - Gets bigger energy kick
 - Ends up relatively faster
- Phase stability!

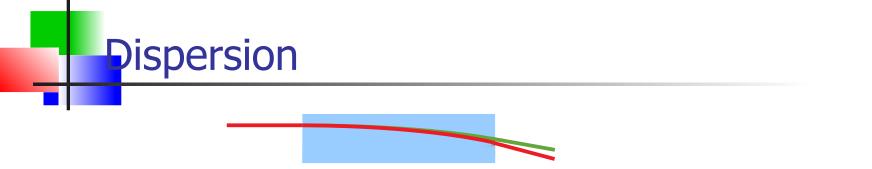




Dealing with momentum spread

- Momentum spread introduces a few effects
 - Dispersion
 - Chromaticity
 - Momentum compaction
- Dispersion:
 - Off-momentum particles follow a different trajectory
- Momentum compaction (rings):
 - Different path length yields different time of flight
- Chromaticity:
 - Off-momentum particles get a different focussing strength





Recall the definition of magnetic rigidity

$$B\rho = \frac{p}{q}$$

- Particles having different momentum (p) get different radius of curvature
 - Introduce dispersion D

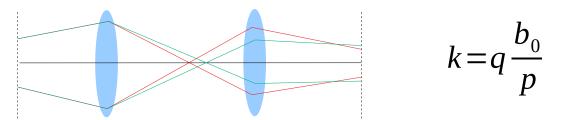
$$D = p \frac{dx}{dp}$$

• Which is another optical function that we must make periodic



Chromaticity

 Chromaticity arises because quadrupoles focus differently for different momenta



- This often limits the degree of focussing at a collision point
 - Limits luminosity
- Can deliberately enhance/reduce chromaticity by
 - Introduce a dispersion
 - Using a magnet with variable focussing strength across the aperture
 "sextupole"







- Dipoles are used to bend a beam rigidity is
- Quadrupoles are used to focus a beam:

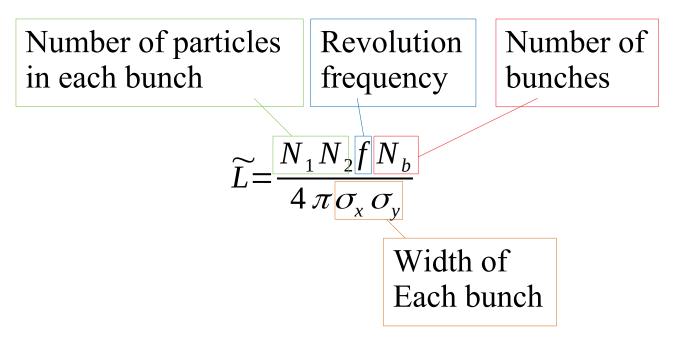
$$B\rho = \frac{p}{q}$$
$$k = q\frac{b_0}{p}$$

- Beam in each of x and y can be characterised by 3 Twiss parameters and an emittance
- Lattices can be characterised by a phase advance
- RF cavities are used to accelerate the beam
- Introducing momentum spread, one can also define a dispersion (and its derivative with respect to z)



Finally... luminosity

- Luminosity defines the number of interactions in a collider per unit time for a given cross section
- Luminosity will increase if
 - Beam is narrower
 - Current is higher





What dictates luminosity?

$$\widetilde{L} = \frac{N_1 N_2 f N_b}{4 \pi \sigma_x \sigma_y}$$

- Typically
 - Number of particles → wake field or space charge
 - Revolution frequency \rightarrow ring circumference
 - Number of bunches \rightarrow Bunch merge
 - Beam width $\rightarrow \sqrt{\varepsilon\beta}$
 - Emittance (cooling?)
 - Twiss beta (final focus and chromaticity)
- Colleagues will discuss in more detail in following lectures







Consider a circular beam of radius *a* having uniform density

$$\rho(r) = q \frac{I}{\beta_{rel} c \pi a^2} \qquad r < a$$
Quote field around a cylinder of charge/current
$$E(r) = \frac{1}{2 \pi \varepsilon_0} \frac{I}{\beta_{rel} c} \frac{r}{a^2}$$

$$B_{\phi}(r) = \frac{1}{2 \pi \varepsilon_0} \frac{I}{c^2} \frac{r}{a^2}$$
Apply Lorentz force law
$$\vec{F} = q \vec{v} \times \vec{B} + q \vec{E}$$

$$F_r = q \vec{v} \times \vec{B} + q \vec{E} = \frac{1}{2 \pi \varepsilon_0} \frac{r}{a^2} (\frac{I}{\beta_{rel} c} - \frac{I}{\beta_{rel} c} \beta_{rel}^2) = \frac{1}{2 \pi \varepsilon_0} \frac{r}{a^2} (\frac{I}{\gamma^2 \beta_{rel} c})$$

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Force is defocusing

$$\frac{d^2x}{dz^2} - (k - K_{sc})x = 0 \quad \text{with}$$

Treat SC as a perturbation

$$M_{p} = M M_{sc}$$
$$M = I \cos \mu + J \sin \mu$$
$$M_{sc} = \begin{pmatrix} 1 & 0 \\ -K_{SC} & 1 \end{pmatrix}$$

Change of beam size (β)

Change of phase advance

- Drive the beam onto resonances \rightarrow ruin the acceptance
- Phase advance \rightarrow look at Trace of M_n

$$Tr(\boldsymbol{M}_{\boldsymbol{p}}) = 2\cos(\mu) + \alpha\sin(\mu) - \alpha\sin(\mu) + \beta K\sin(\mu)$$



$$K_{sc} = \frac{1}{2 \pi \varepsilon_0} \frac{1}{a^2} \left(\frac{I}{\gamma_{rel}^3 \beta_{rel}^2 c}\right)$$

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$$\boldsymbol{I} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$

$$\mathbf{J} = \begin{pmatrix} \alpha & \beta \\ -\gamma & -\alpha \end{pmatrix}$$

- Consider just the *trace(M_p)* $Tr(M_p) = 2\cos(\mu) + \beta K \sin(\mu)$
- Consider compound angle formula $\cos(\mu + \delta\mu) = \cos(\mu)\cos(\delta\mu) + \sin(\mu)\sin(\delta\mu)$ $\cos(\mu + \delta\mu) \simeq \cos(\mu) + \sin(\mu)\sin(\delta\mu)$
- Looking at the tune

$$K_{sc} = \frac{1}{2 \pi \varepsilon_0} \frac{1}{a^2} \left(\frac{I}{\gamma_{rel}^3 \beta_{rel}^2 c} \right)$$

$$\delta v = \frac{r_0 N}{2 \pi \varepsilon \beta_{rel}^2 \gamma_{rel}^3}$$

 $\delta v = \frac{\delta \mu}{2\pi} = \frac{\beta K}{4\pi}$

$$\sigma(x) = \sqrt{\beta \epsilon}$$



