



SISSA

Constraining mixed scenarios with galaxy clustering

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Which Dark Matter with galaxy clustering

What kind of new physics can we look for with galaxy clustering?

$$P(k, a) \sim D^2(a) T^2(k) P_0(k)$$

**Dynamics of
the Dark
Sector**

Primordial
features

Dynamics already constrained by larger and smaller scales (CMB, Lyman- α ,...)

+

galaxy clustering data getting to unprecedented precisions

↓

ideal place to look for small deviations from CDM

$$D(a)T(k) \sim aT_{\text{CDM}}(k) (1 + \Delta_{\text{NP}})$$

*Promising opportunity to constraints BSM scenarios.
If the dark sector is really dark, might be a unique window!*

Why mixed models

In many scenarios DM is CDM + something else (ultra-light axions [Lagüe et al 22, Rogers et al 23], warm thermal relics [Xu et al 21, Çelik&Schmidt 25], subcomponent with strong self-interactions [Garani et al 22],)

So, the DM is CDM+ χ . How strange can it be? Assuming that

1. χ non-relativistic
2. χ non-interacting

\Rightarrow fluid description [Shoji&Komatsu 10]

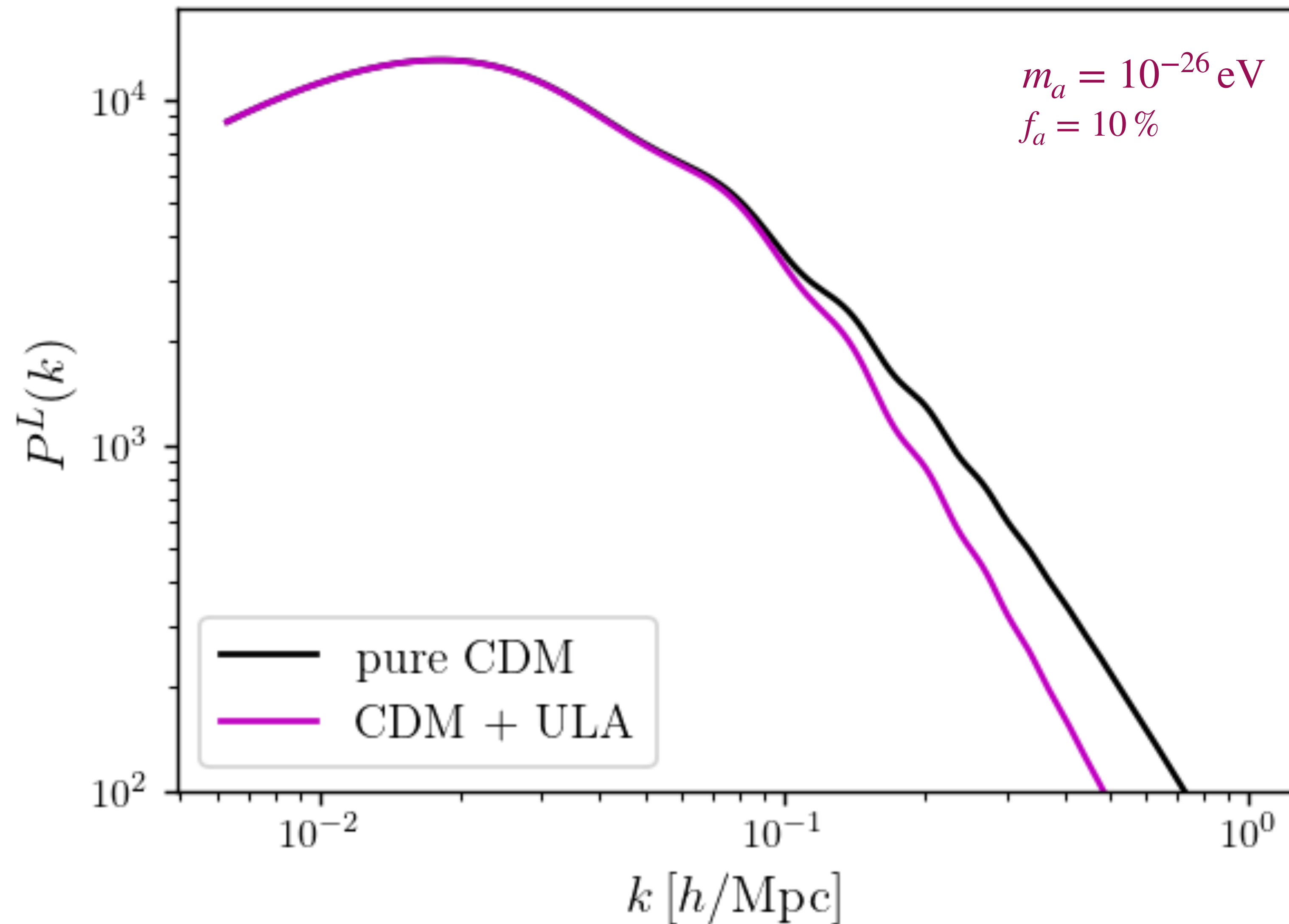
$$\ddot{\delta}_c + \mathcal{H}\dot{\delta}_c - \frac{3}{2}\mathcal{H}^2\delta_m = 0$$

$$\ddot{\delta}_\chi + \mathcal{H}\dot{\delta}_\chi - \frac{3}{2}\mathcal{H}^2\delta_m + c_s^2 k^2 \delta_\chi = 0$$

the deviation from CDM
ends up in this term

Phenomenologically, two new parameters $\left(f_\chi, \quad k_J \sim \frac{\mathcal{H}(a)}{c_s(a)} \right)$

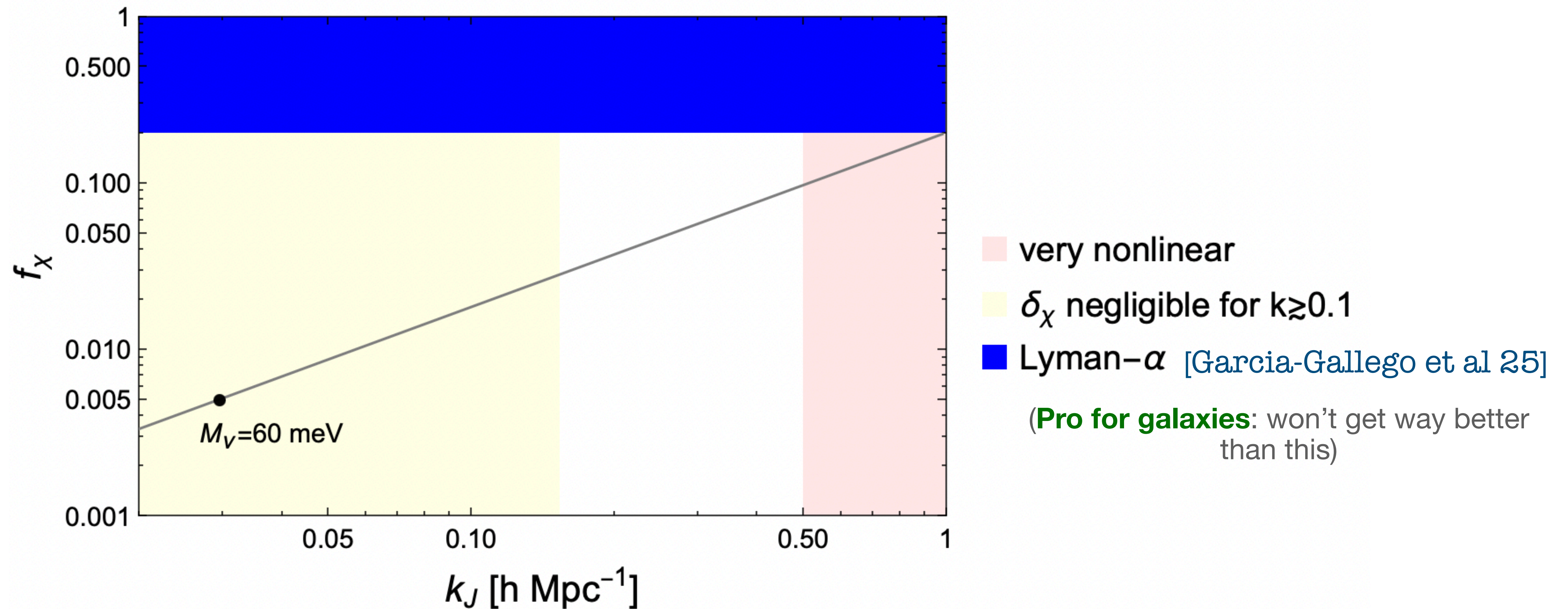
The growth of linear perturbations



⇒ full-shape analysis seems the ideal tool to probe these scenarios

Full-shape analysis seems the tool to probe these scenarios. Is it?

- **Motivation for the theoretical effort:** having a *first-principles*, robust, *controlled* computation
- How hard is the theoretical study? How hard is it to implement in a LSS inference code?
- Results would be generic or very model-dependent?
- By construction, marginalizing over unknowns (bias, counterterms...) Is there really a gain in constraining power after this?



- The fractions we can test is already small (helps theory)
- The “nontrivial” case reduces to $k \ll k_J$ [Çelik&Schmidt 25]

(**Stretch goal**: report generic constraints in this plot, and then read the corresponding particle physics model. **Unfortunately it is not that trivial...**)

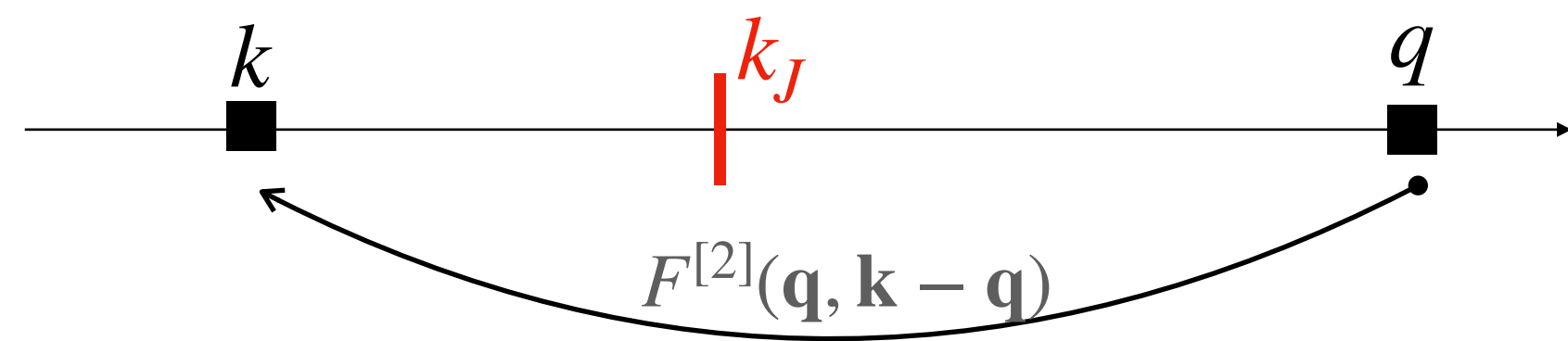
The steps one needs to take for a full, proper, EFTofLSS 1-loop analysis:

1. Have the BSM model implemented in an Einstein-Boltzmann solver
2. Study nonlinearities with perturbation theory
3. Bias expansion with two fluids
4. Include redshift-space distortions
5. Add this (at least in an approximate way) to an LSS code

Nonlinear Perturbation Theory with two components

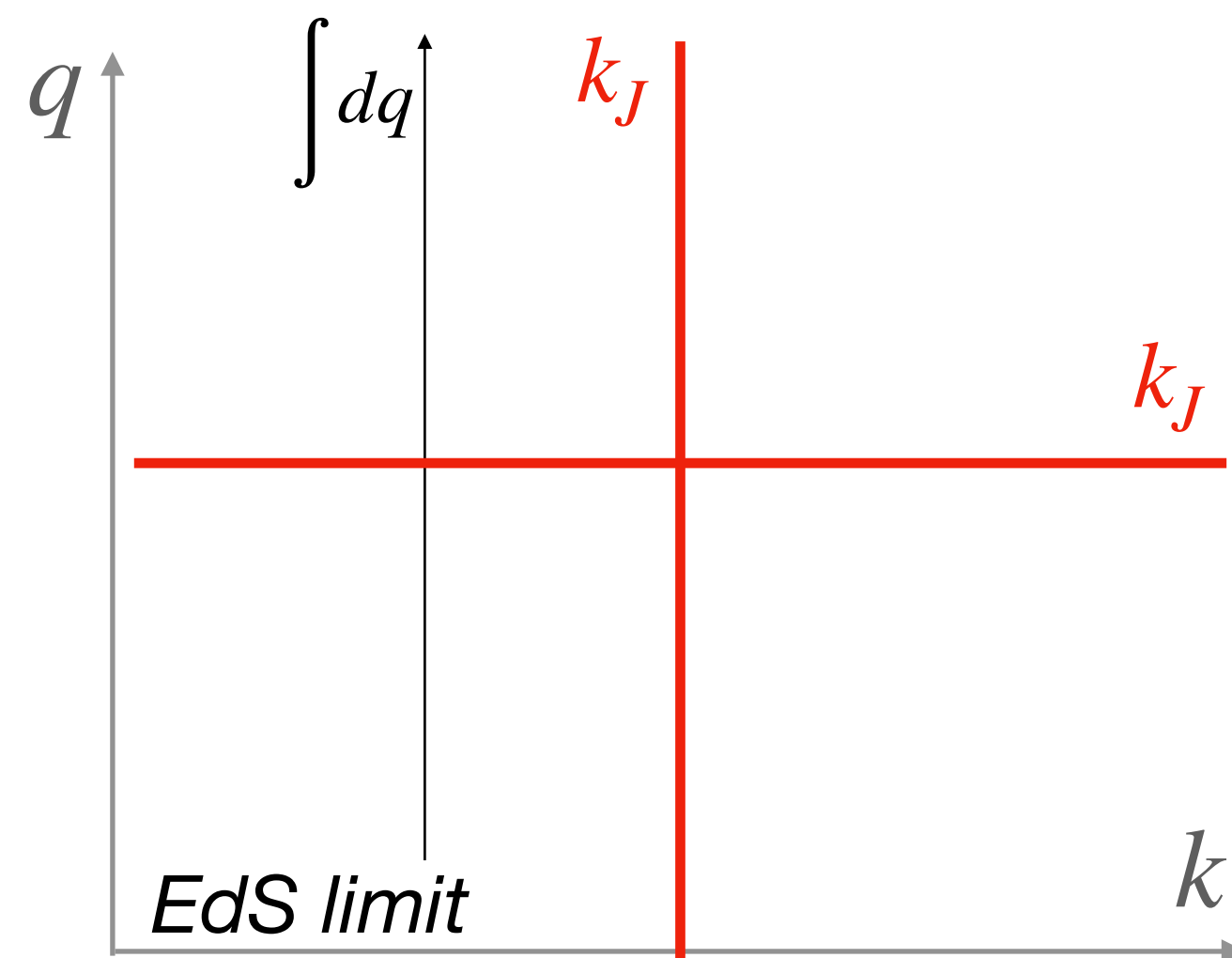
Nonlinearity couples the modes: all the scales are affected. In SPT approach

$$\delta(k) = \delta^L(k) + \int d^3q F_a^{[2]}(\mathbf{q}, \mathbf{k} - \mathbf{q}) \delta^L(\mathbf{q}) \delta^L(\mathbf{k} - \mathbf{q}) + \dots$$



PT Kernels (mode coupling)
For CDM is known, but for χ ?

Our approach: we can work in an $f_\chi \ll 1$ expansion. The problem is actually 2D now:



We can study analytically the 4 limits. We see that IR cancellation happens exactly as in EdS (even for $k \gg k_J$)

Could be guessed from first principles? [\[D'Amico et al 21\]](#)

Nonlinear Perturbation Theory with two components

Turns out that the symmetries are so constraining that kernels are very similar to the CDM ones

$$\int F_{\chi}^{[2]} \delta^L \delta^L \approx \int F_2 \delta_{\chi}^L \delta_{\chi}^L$$

or, in other words

$$\delta_{\chi}^{[n]}(\mathbf{k}) \simeq \frac{\mathcal{T}_{\chi}(k)}{\mathcal{T}_c(k)} \delta_c^{[n]}(\mathbf{k})$$

As already assumed in M_{ν} PT studies [[Aviles et al 21](#)]

Incredible advantage for code implementation: can recycle the FFTLog routines to compute 1-loop contributions

Side note: some claims on effects in the squeezed bispectrum [[Nascimento&Loverde 23](#), [Kamalinejad&Slepian 20](#), [Zhu&Castorina 20](#)]. In this fluid description, we find none: are we missing something?

Towards the analysis: bias and full-shape template

Start from the CDM-only “EFTofLSS” template for full-shape

$$P_g(k) = P_g^L(k) + P_g^{1\text{-loop}}(k) + P_g^{\text{ctr}}(k) + P_g^{\text{noise}}(k)$$

⇒ we allow galaxies trace also δ_a

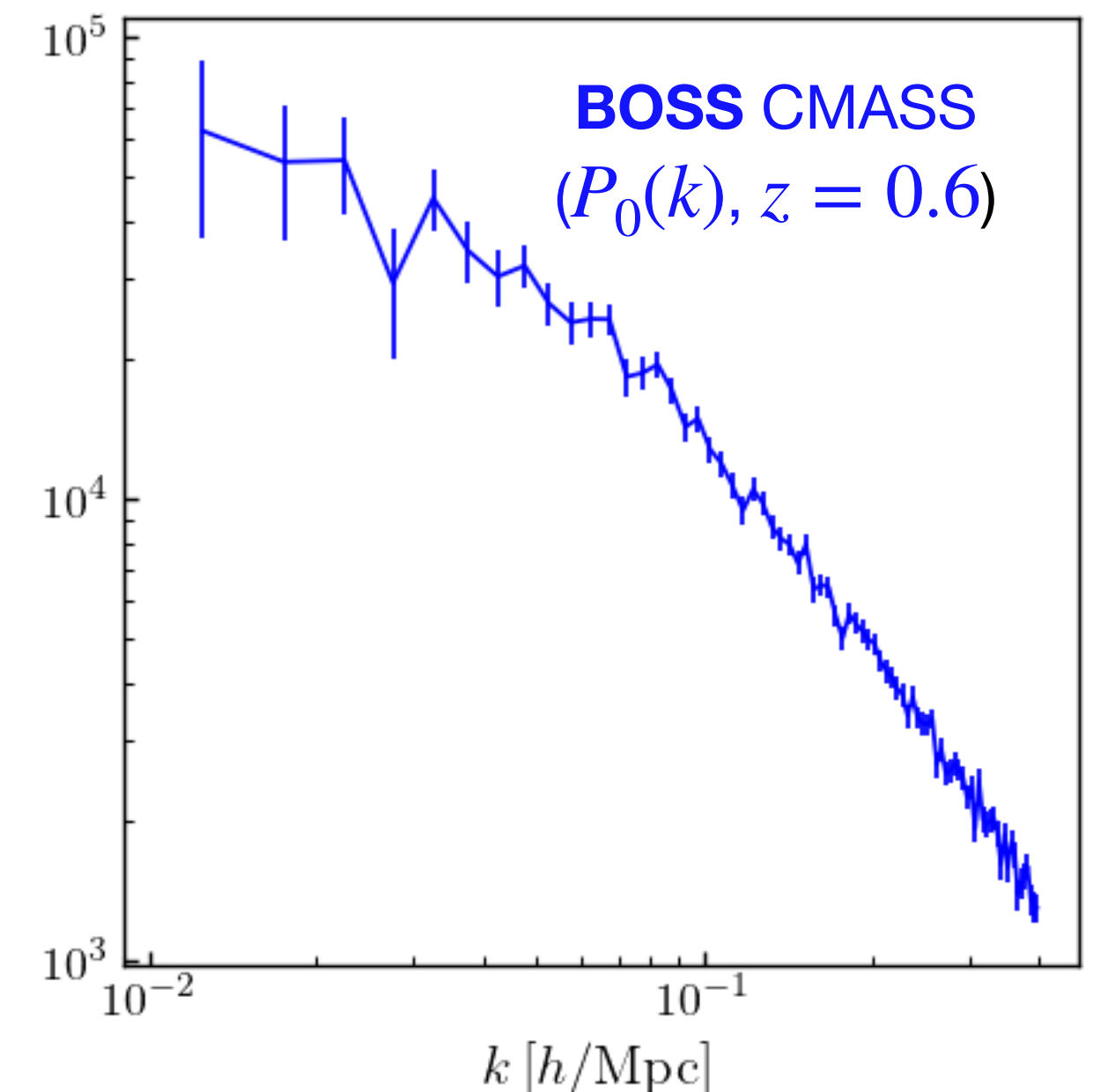
$$\delta_g = b_1 \delta_c + b_a \delta_a + \dots$$

Actually, the other modes can appear: for a proper treatment (see [Çelik&Schmidt 25])
Nonlinear biases? The number of operator explodes (see eg [Bottaro et al 23])

⇒ modify the code to implement new bias,
nonlinearities, counterterms...

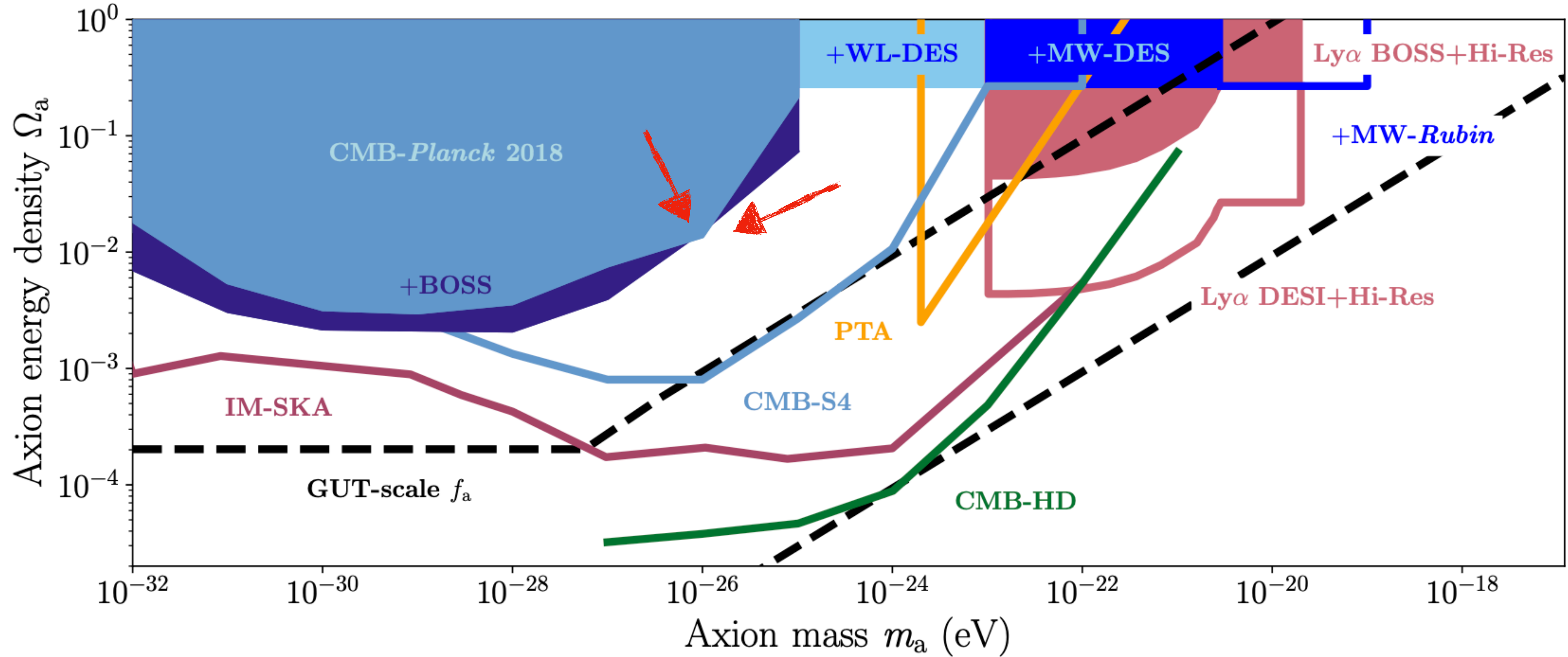
$$P_g^{\text{ctr}} \simeq -2c_c k^2 P_{cc} - 2c_a k^2 P_{ca} + \dots$$

⇒ finally, fit data!

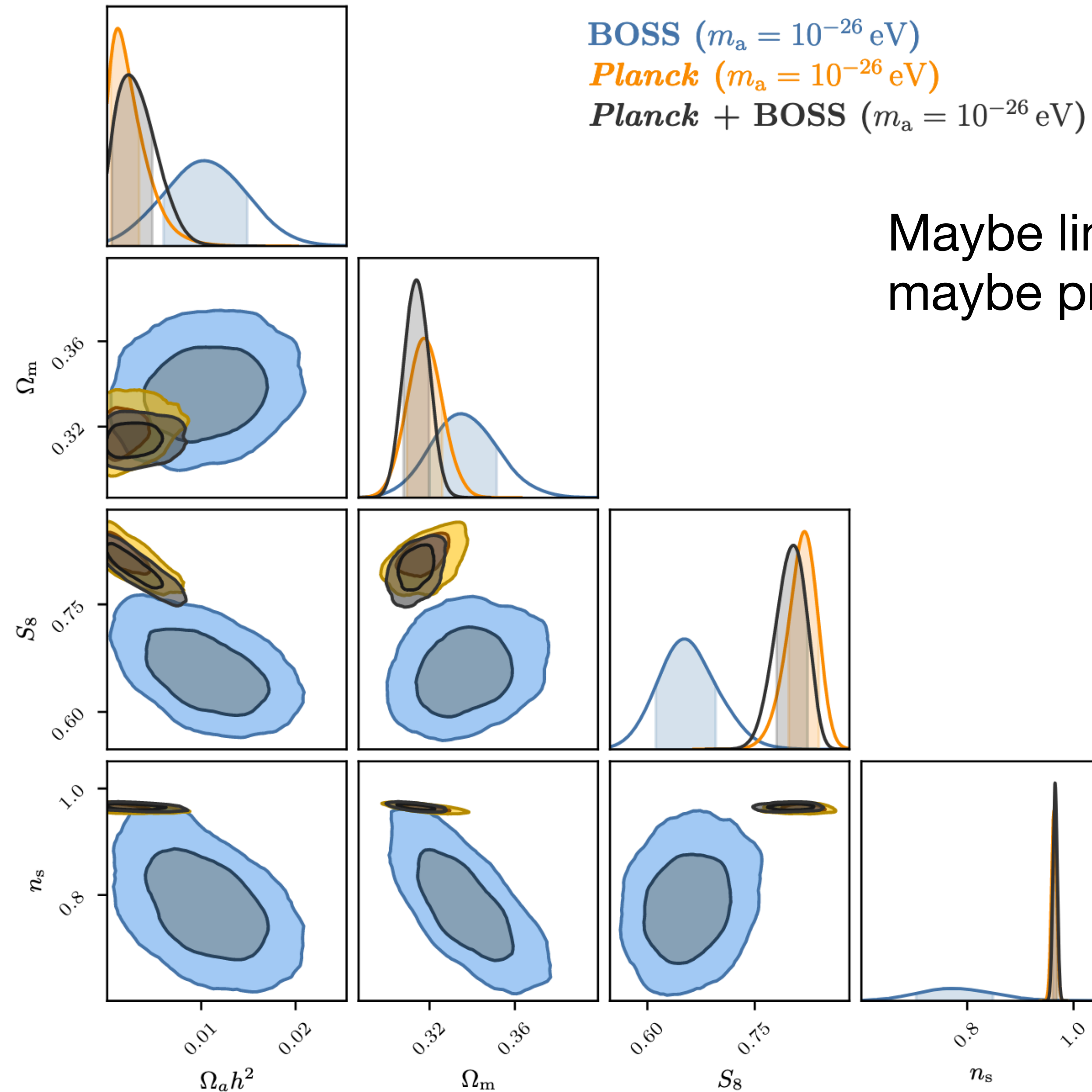


Our reference for ULAs

[Rogers et al. 23]



BOSS ($m_a = 10^{-26}$ eV)
Planck ($m_a = 10^{-26}$ eV)
Planck + BOSS ($m_a = 10^{-26}$ eV)

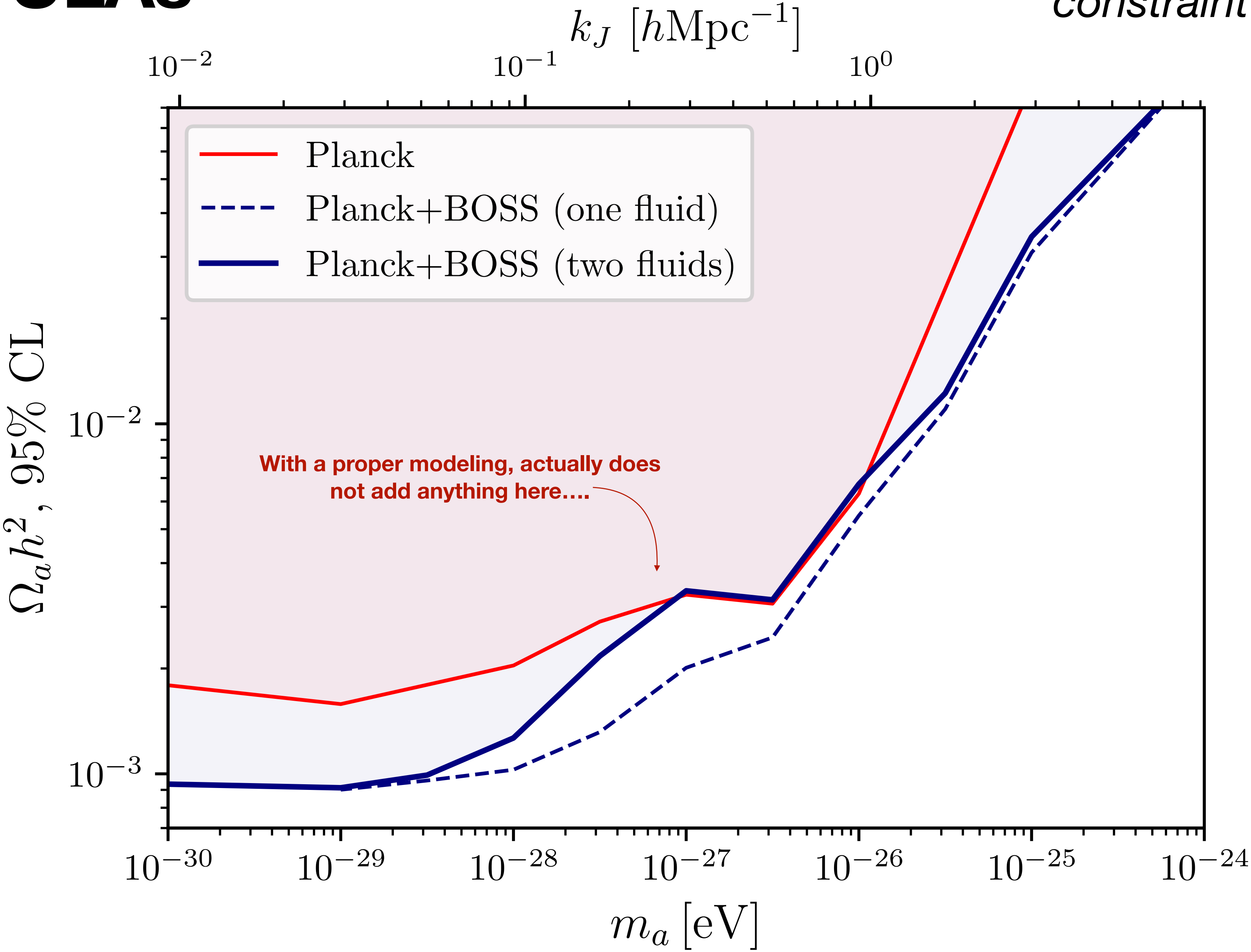


Maybe linked to the S_8 tension in BOSS,
maybe projection effects

What we are after to, changes a lot the
perspective:
just excluding, or do we have a detection?
In the latter case, proper modeling is imperative

Results: ULAs

We scan on m_a and then
constraint $\Omega_a h^2$



Hidden: the importance of Q_0 for BSM

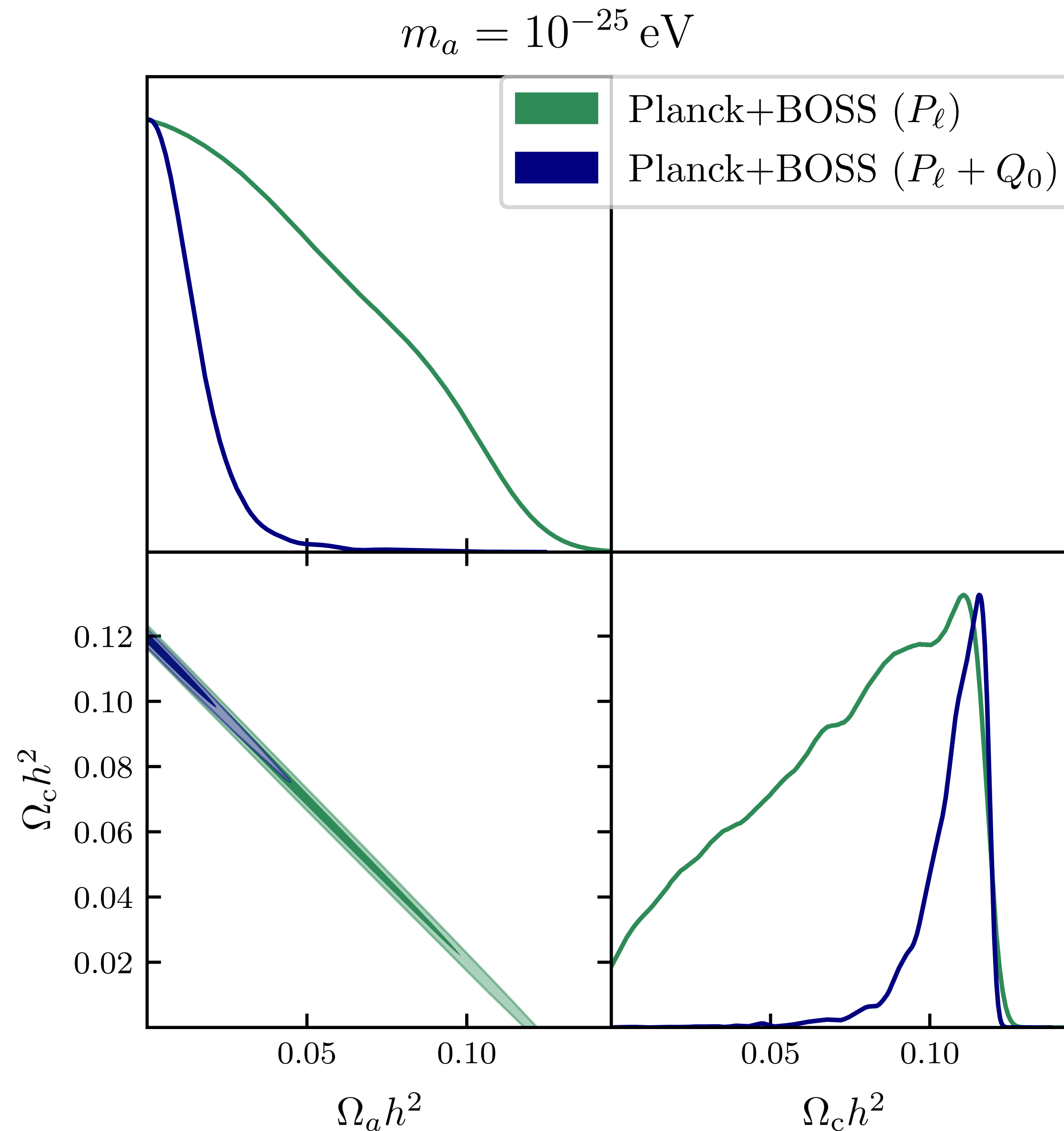
$$\check{Q}_0(k_i) = \check{P}_0 - \frac{1}{2}\check{P}_2 + \frac{3}{8}\check{P}_4$$

Being free from RSDs, one
can push the fit up to

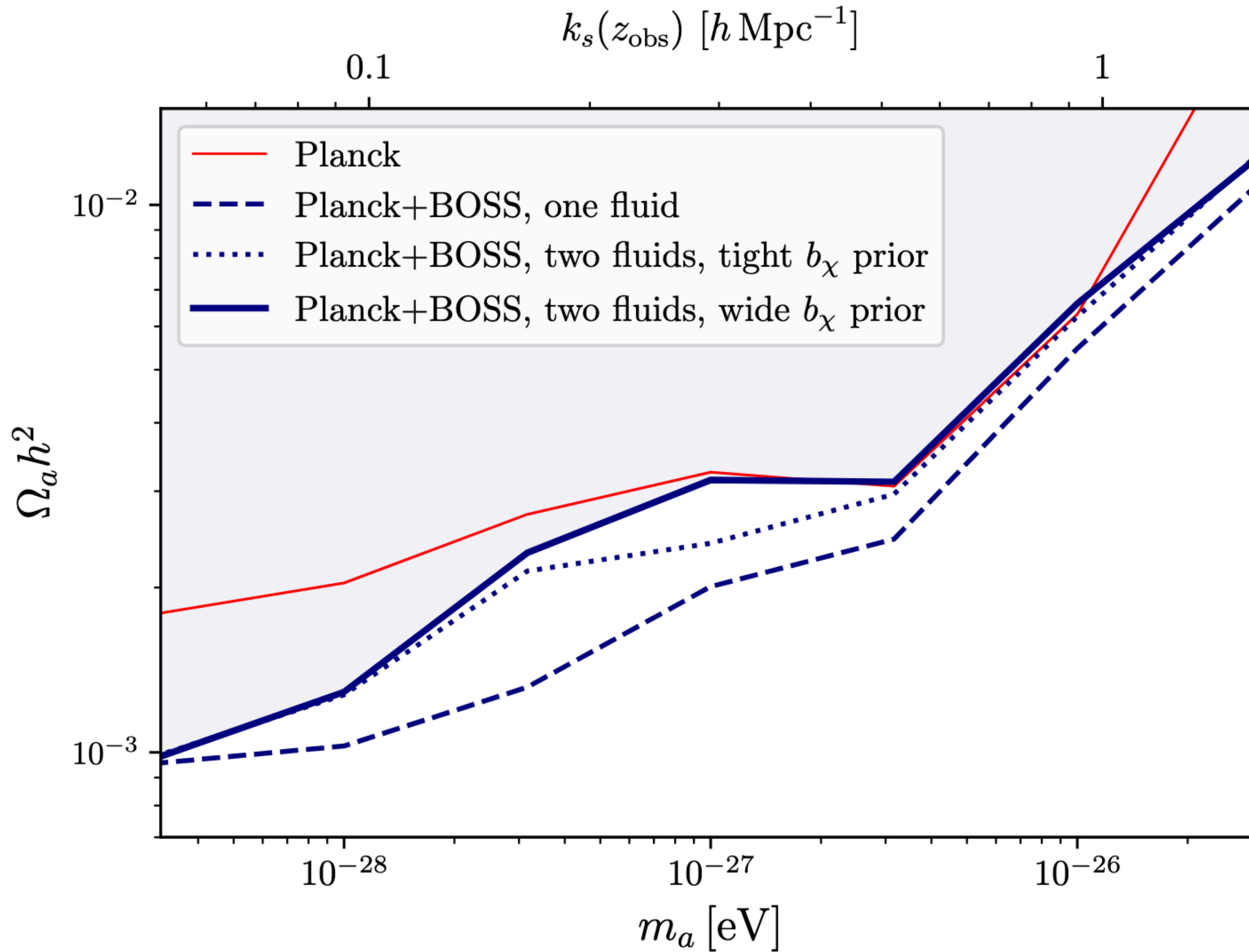
$$k_{\text{max}} = 0.4h\text{Mpc}^{-1}$$

[Ivanov et al. 22]

Particularly beneficial for
probing BSM suppression!



Do we know something more? For instance, $b_\chi \sim \mathcal{O}(f_\chi)$ [Çelik&Schmidt 25]



Some outlooks

- CMB + LSS is very powerful in constraining, even $\Omega_a \lesssim 0.01\Omega_m$!
- Theoretical modeling is important for controlled results. In constraining, not to overestimate [[Çelik&Schmidt 25](#)]. With a detection, totally new perspective
- The (f_χ, k_J) plot is not generic, but background dependent
- Not very clear how much information coming from background or shape suppression (neither for massive neutrinos, see [[Elbers et al 25](#)])