

The EFTofLSS: Successes and next moves

Based on several works with L. Senatore, P. Zhang, Y. Donath, M. Lewandowski, et al.

PyBird: *PYthon code for Biased tracers In ReDshift space*

PyOwl: *PyBird, but nightly*

PyFowl: *PYthon code For Observables in Weak Lensing*

New Physics from Galaxy Clustering at GGI, Florence, 2025/9/29

D'où venons-nous ?

Que sommes-nous ?

Où allons-nous ?



Gauguin 1898

What is this for?

- We are getting a wealth of LSS data (DESI, Euclid, Sphex...) What's the end goal of the experimental effort?
- **Understanding the universe!** Practically, this means **measuring known** and **detecting unknown physics** across a wide range of energy scales (neutrino masses, PNG, dynamical DE, light mediators, relics, ?)
- As with any experimental information, we struggle for
 - **Precision**: more volume, more galaxies, more time
 - **Accuracy**: less and less systematics, and a **reliable interpretation of the data**
- **EFTofLSS** gives us just that – analytical understanding in terms of a **description of the system based on good old effective theory**

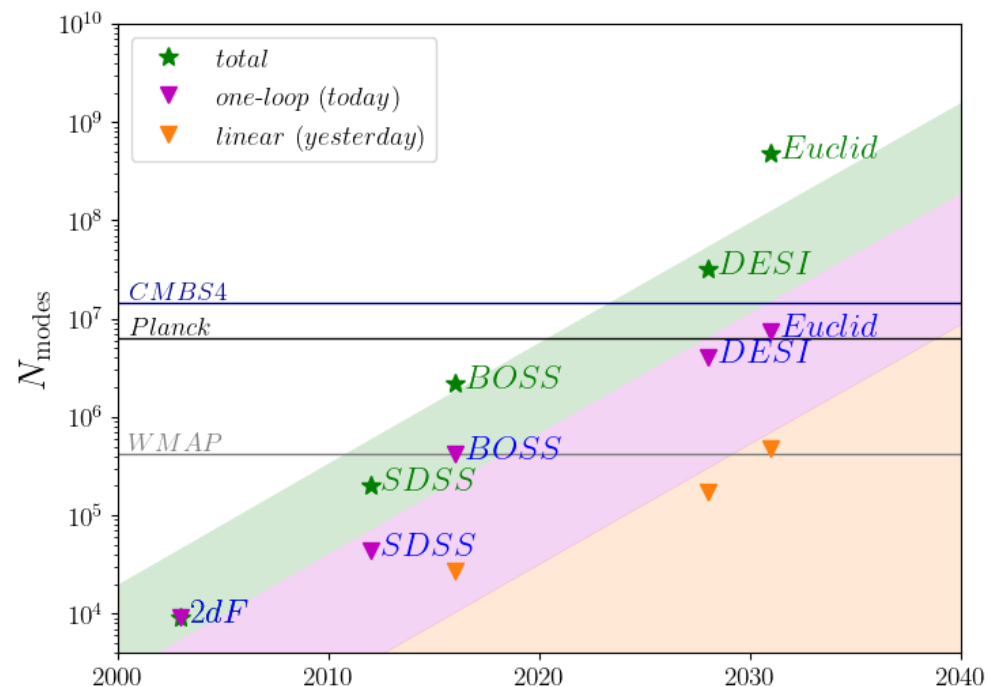
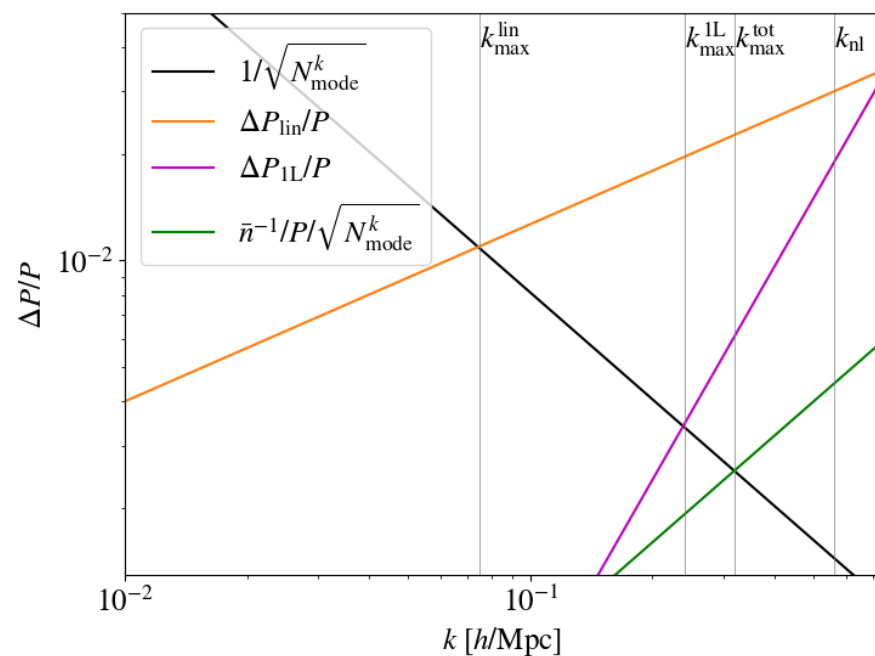
The EFTofLSS, in a nutshell

- **Basic dofs and symmetries:** CDM smoothed density and momentum (or displacement), translations, rotations, diffs
- **Physical cutoff(s)** of the theory: $k_{\text{NL}}, k_{\text{R}}$
- **Expansion parameter(s):** basically $k/k_{\text{NL}}, k/k_{\text{R}}$
- **Do perturbation theory, where possible.** Renormalise unknown effects of small scales. Resum non-perturbative effects. Congratulations, you have (unobservable) CDM!
- Add biasing scheme, go to redshift space, do projection integrals, to **get to specific observables**. Now you can build a likelihood!
- **Estimate the scale cuts.**
- Add eventual known systematics.
- Do the data analysis, report the results

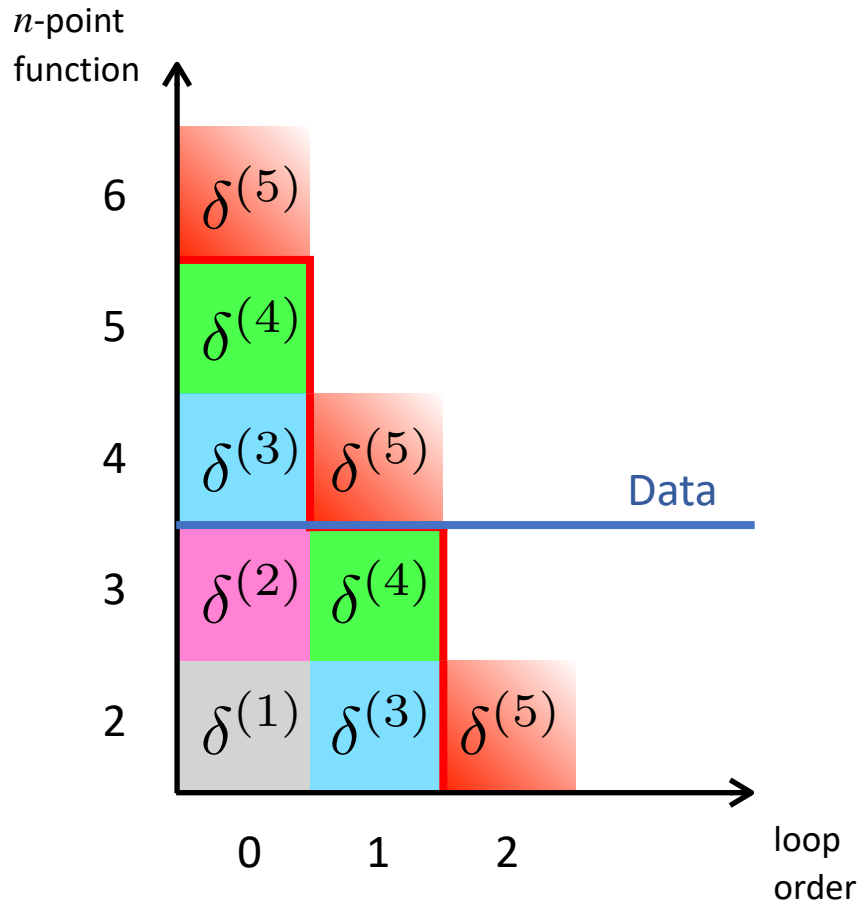
Why the pain?

$$N_{\text{modes}}(\text{CMB}) \sim \ell_{\text{max}}^2$$

$$N_{\text{modes}}(\text{LSS}) \sim V_{\text{survey}} k_{\text{max}}^3$$

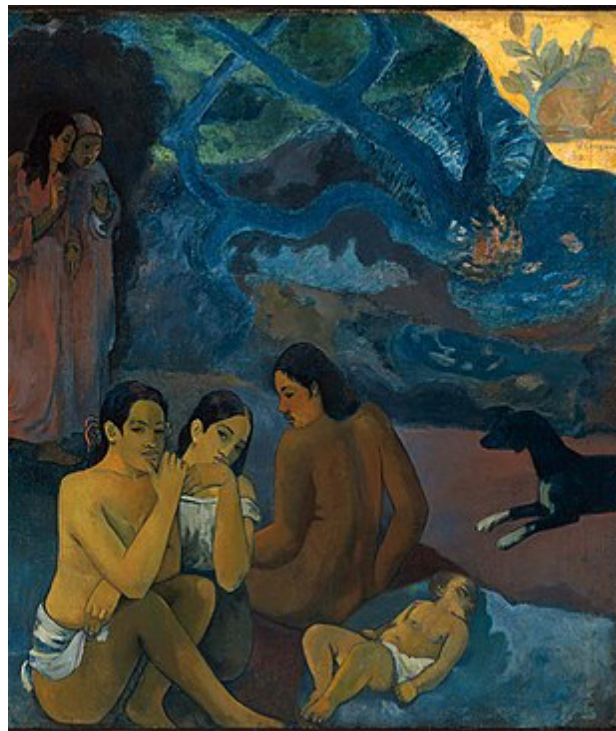


What have we computed?

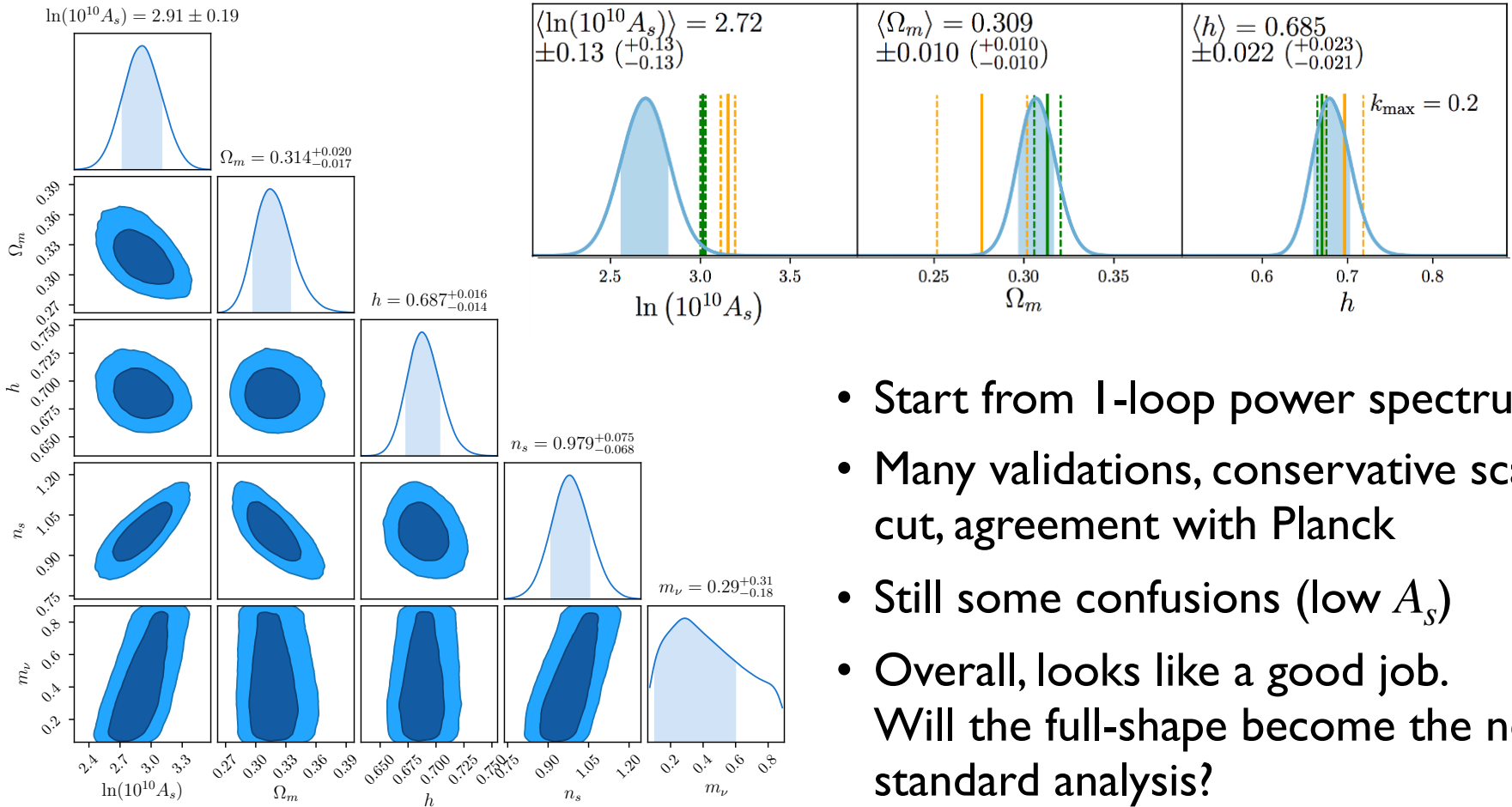


	Bias	Ren.	RSD
$\delta^{(3)}$	✓	✓	✓
$\delta^{(4)}$	✓	✓	✓
$\delta^{(5)}$	✓	✓	✗

D'où venons-nous ?



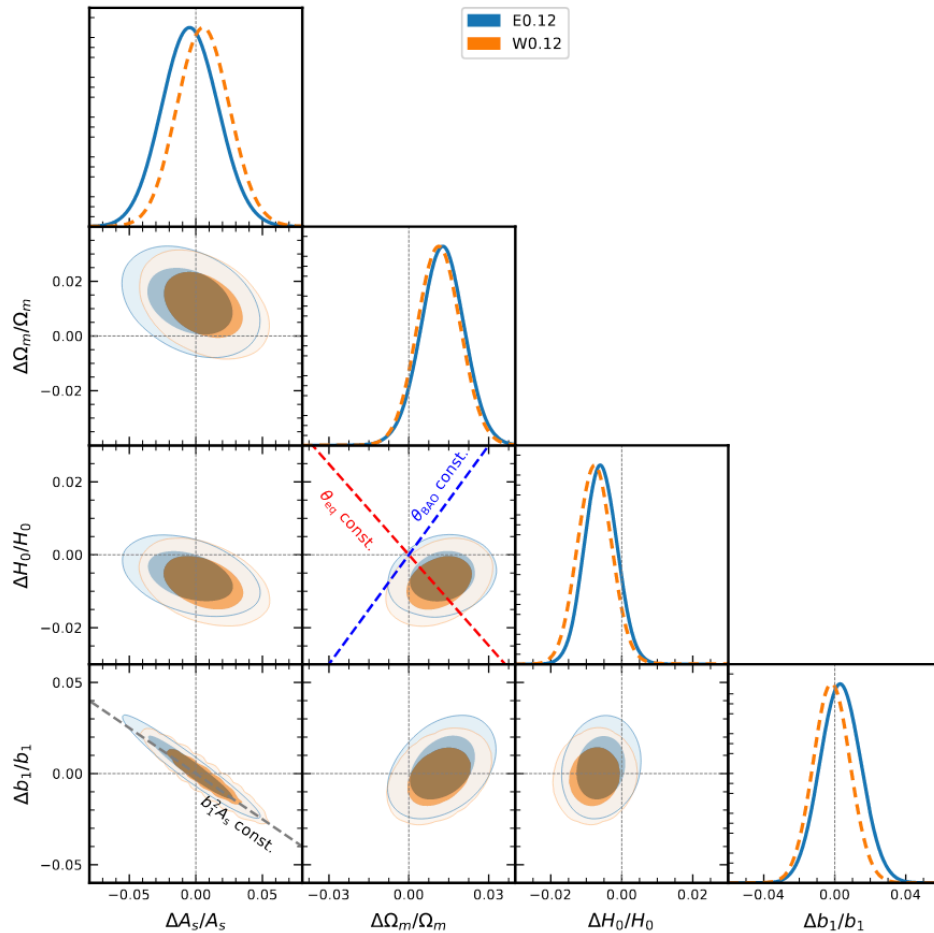
The story: 1-loop 2-point function



- Start from 1-loop power spectrum
- Many validations, conservative scale-cut, agreement with Planck
- Still some confusions (low A_s)
- Overall, looks like a good job. Will the full-shape become the new standard analysis?

D'Amico, Gleyzes, Kokron, Markovic,
 Senatore, Zhang, Beutler, Gil-Marín 2019
 Colas, D'Amico, Senatore, Zhang 2020

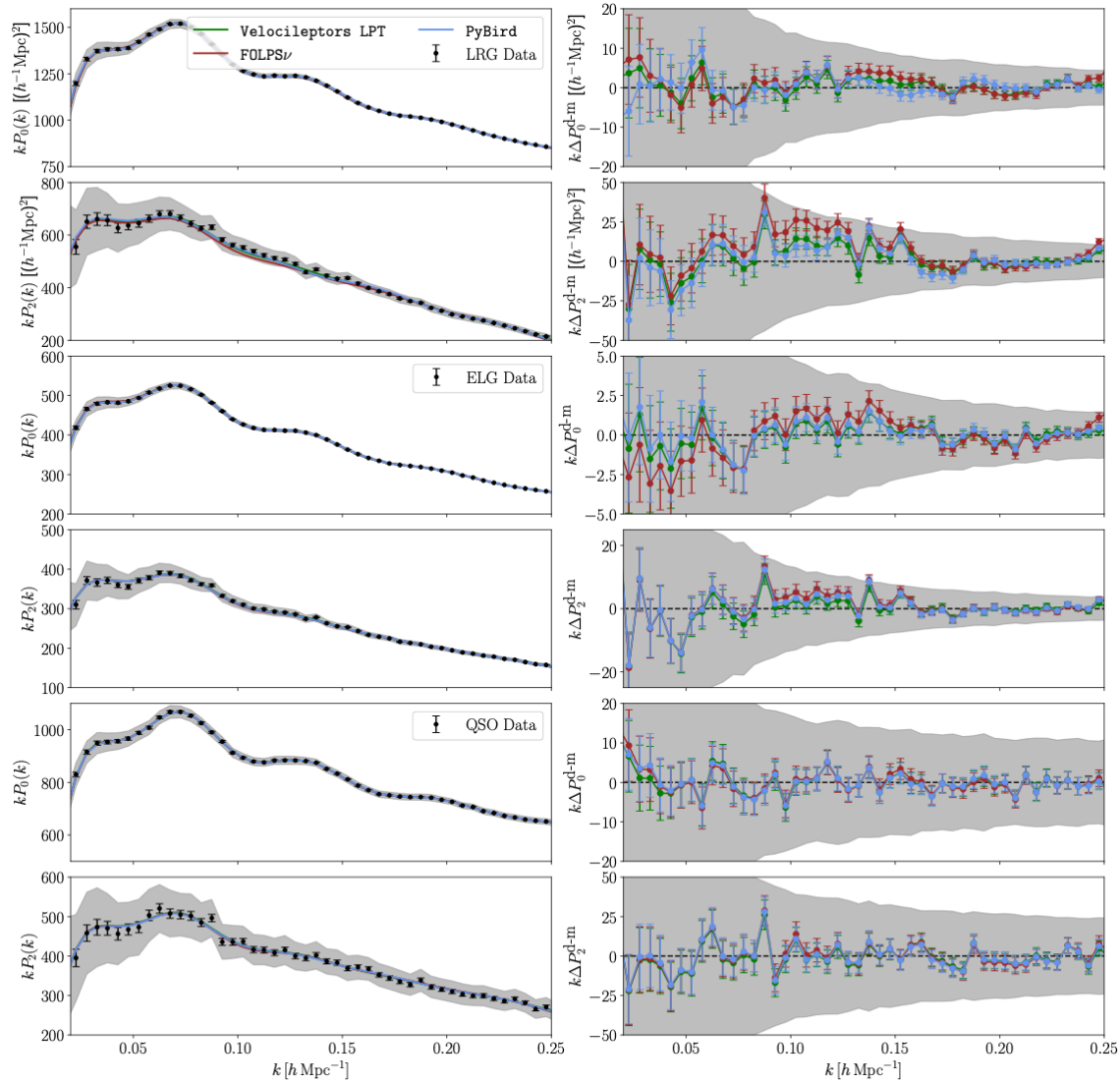
The blinded PT-challenge



- Total volume: $566 (h^{-1}\text{Gpc})^3$, ~ 100 times BOSS
- 2 independent blind analyses, on 3 cosmological parameters: we agreed between us and with the truth!
- This helped establish the reliability of the method
- Now many codes available: PyBird (and PyBird-Jax), CLASS-PT, Velocileptors, Class-OneLoop, PBj

Nishimichi, D'Amico, Ivanov, Senatore, Simonovic,
Takada, Zaldarriaga, Zhang 2020
Chen, Vlah, Castorina, White 2021

DESI validation

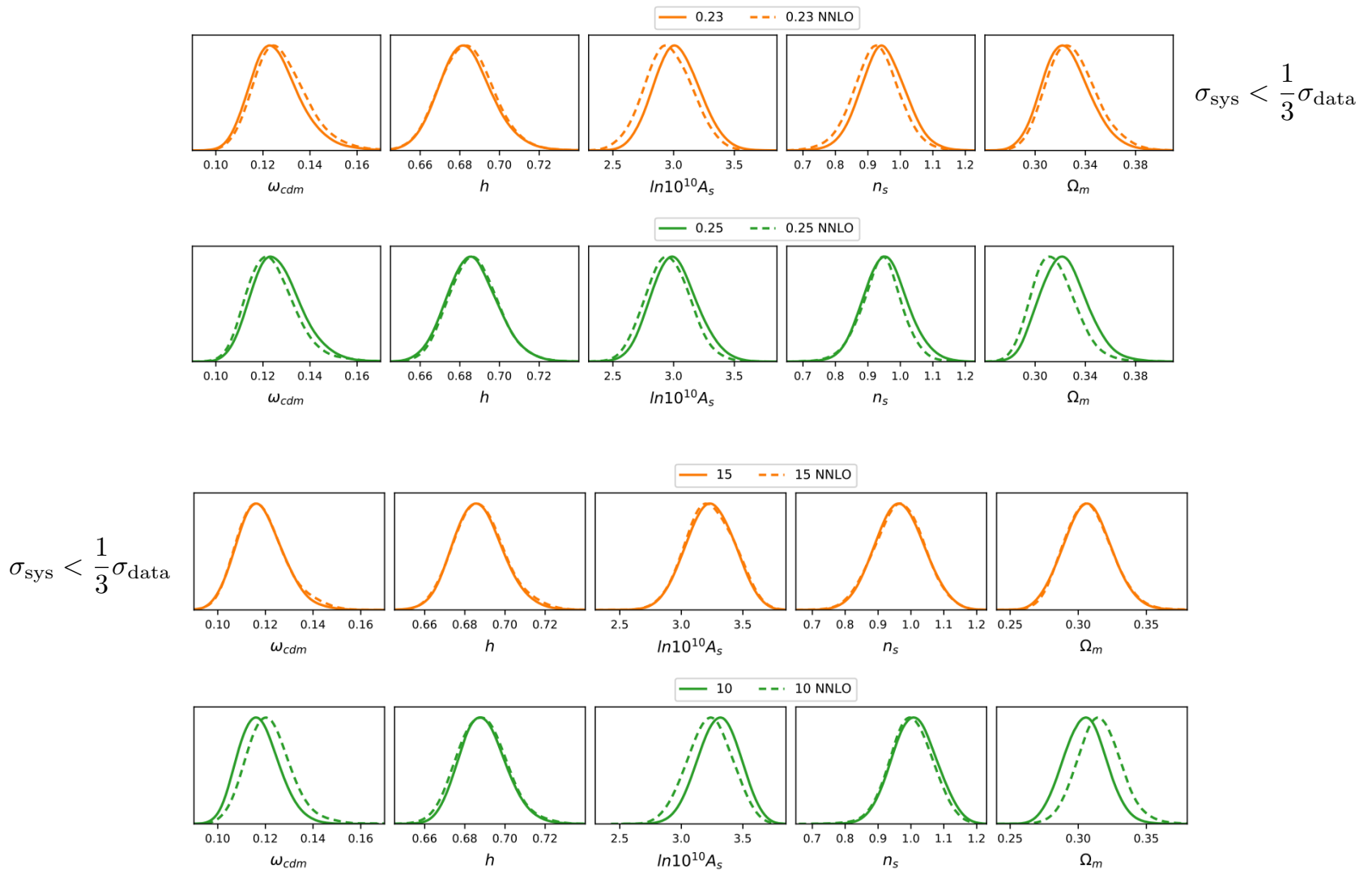


Maus et al. 2024

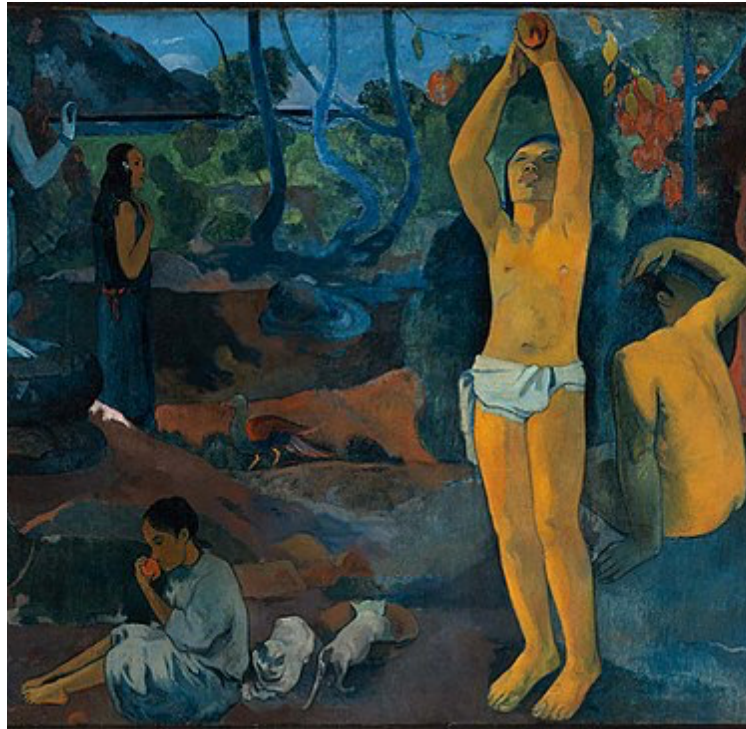
The scale-cut problem

- Crucial issue: **where to stop the fit?**
Usual tradeoff between *accuracy* and *precision*: smaller scales have smaller errors, but perturbative approach starts to fail
- Important observation: **effective field theories set their own scale cut, given a required precision**
- The theory is a controlled approximate description of reality.
 - If we compute up to order n , we are making a (theoretical) error of order $\Delta^{(n+1)}$
 - Data have an error σ_{data} . We choose how much accurate we want to be (30%? 10%? 1%?).
Then, let's require $\Delta^{(n+1)} < \alpha \sigma_{\text{data}}$
- **NNLO estimate** are more or less reliable. For the power spectrum
$$P_{2L}^{\mu=0}(k) \sim c_e \frac{k^2}{k_M^2} P_{1L}^{\mu=0}(k) \qquad P_{2L}(k, \mu) \sim b_1 (c_{r,4} b_1 + c_{r,6} \mu^2) \mu^4 \frac{k^4}{k_R^4} P_{11}(k)$$
- Of course, we can also use simulations (if they are reliable)

The scale-cut problem

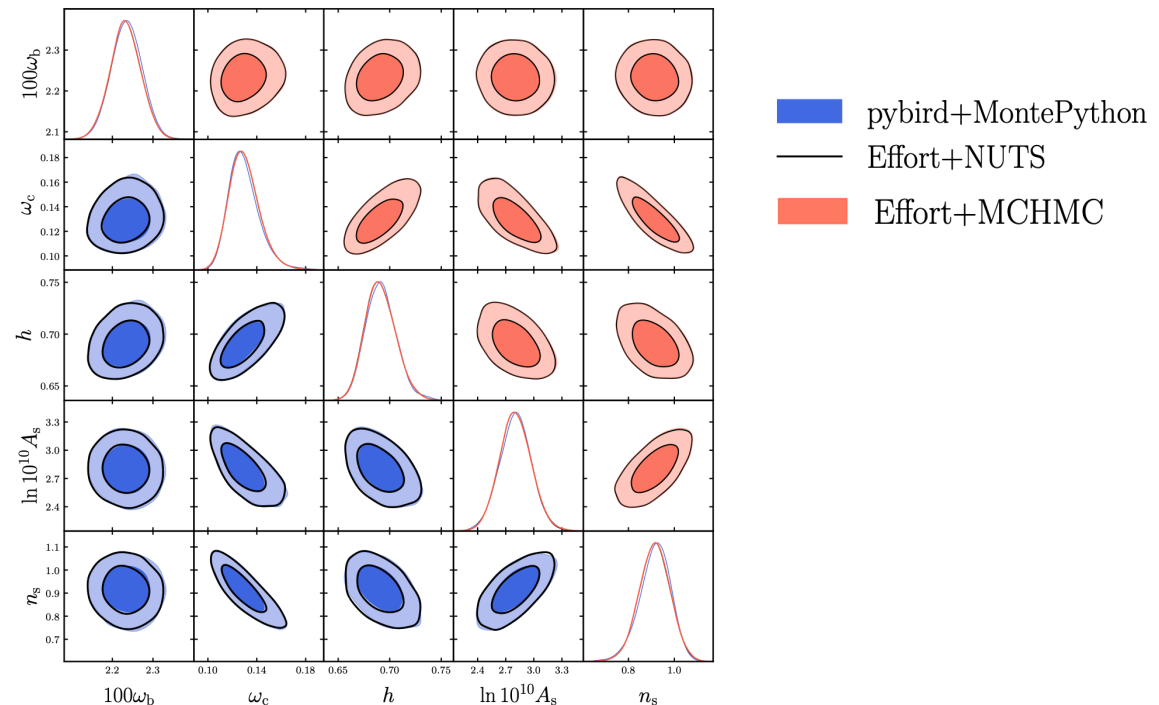


Que sommes-nous ?



Towards differentiable cosmology: *Effort.jl*

- Would be really nice to be super fast and use analytical gradients: easy minimisation, likelihood profile, use of HMC
- Possible solution: **emulate the theory, in a smart way**
For 2-pt very easy. One can afford to build a slower, more precise EFT code, since the NN is trained only once.
Each evaluation takes $\sim O(100 \mu\text{sec})$ ($15 \mu\text{sec}$ without AP)



Beyond 2-pt: the 1-loop bispectrum in LSS

- Lots of work to develop the pipeline for 1-loop bispectrum
 - Biased tracers to 4th order in perturbations
 - Redshift distortions up to 4th order
 - Counterterms up to 2nd order
 - Efficient way of computing loop integrals
 - Generalization to non-Gaussian initial conditions

GDA, Lewandowski, Senatore, Zhang (2022)

GDA, Donath, Lewandowski, Senatore, Zhang (2022a, 2022b)

Philcox, Ivanov, Cabass, Simonovic, Zaldarriaga (2022a, 2022b, 2022c)

Bakx, Ivanov, Philcox, Vlah (2025)

Theory Model

- Perturbation theory up to 4th order: 11 bias parameters

$$P_{11}^{r,h}[b_1] , \quad P_{13}^{r,h}[b_1, b_3, b_8] , \quad P_{22}^{r,h}[b_1, b_2, b_5] , \quad B_{321}^{r,h,(I)}[b_1, b_2, b_3, b_5, b_6, b_8, b_{10}] , \\ B_{211}^{r,h}[b_1, b_2, b_5] , \quad B_{321}^{r,h,(II)}[b_1, b_2, b_3, b_5, b_8] , \quad B_{411}^{r,h}[b_1, \dots, b_{11}] , \quad B_{222}^{r,h}[b_1, b_2, b_5]$$

- Stochastic and counterterms up to 2nd order: 30 parameters

$$P_{13}^{r,h,ct}[b_1, c_1^h, c_1^\pi, c_1^{\pi v}, c_3^{\pi v}] , \quad P_{22}^{r,h,\epsilon}[c_1^{\text{St}}, c_2^{\text{St}}, c_3^{\text{St}}] , \\ B_{321}^{r,h,(II),ct}[b_1, b_2, b_5, c_1^h, c_1^\pi, c_1^{\pi v}, c_3^{\pi v}] , \quad B_{321}^{r,h,(I),\epsilon}[b_1, c_1^{\text{St}}, c_2^{\text{St}}, \{c_i^{\text{St}}\}_{i=4,\dots,13}] , \\ B_{411}^{r,h,ct}[b_1, \{c_i^h\}_{i=1,\dots,5}, c_1^\pi, c_5^\pi, \{c_j^{\pi v}\}_{j=1,\dots,7}] , \quad B_{222}^{r,h,\epsilon}[c_1^{(222)}, c_2^{(222)}, c_5^{(222)}]$$

Observational effects

- **Window**: use approximation (on linear term) from Gil-Marín et al. (2014)

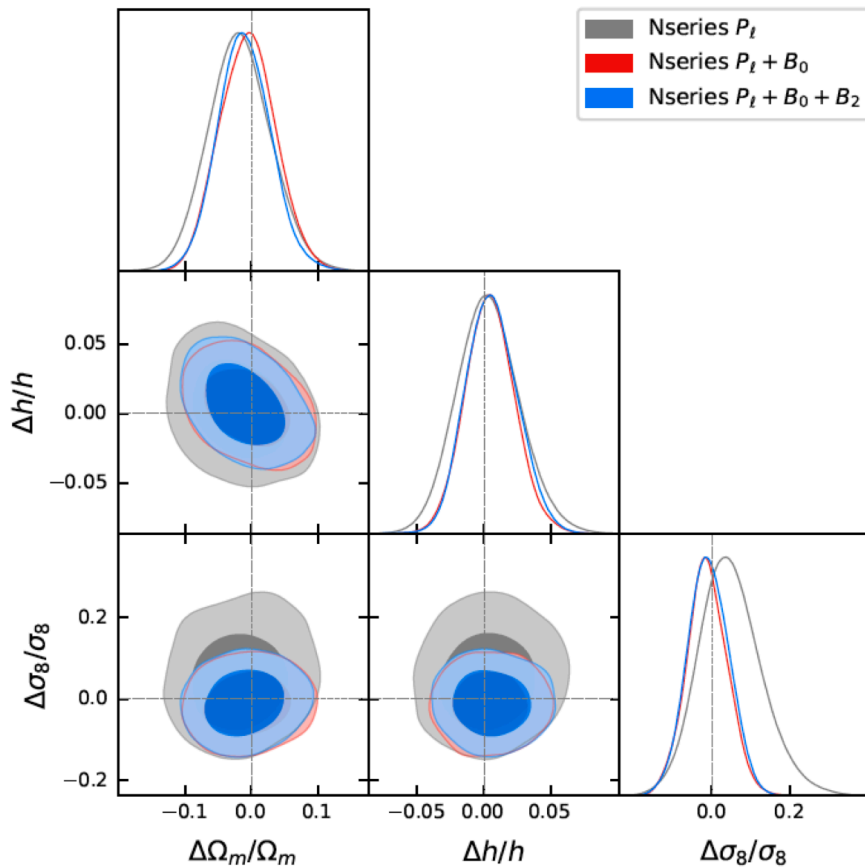
$$B_{211}^{r,h} = 2K_1^{r,h}(\vec{k}_1; \hat{z})K_1^{r,h}(\vec{k}_2; \hat{z})K_2^{r,h}(\vec{k}_1, \vec{k}_2; \hat{z})[W * P_{11}](\vec{k}_1)[W * P_{11}](\vec{k}_2) + 2 \text{ perms. } ,$$

$$[W * P_{11}](\vec{k}) = \int \frac{d^3k'}{(2\pi)^3} W(\vec{k} - \vec{k}')P_{11}(\vec{k}')$$

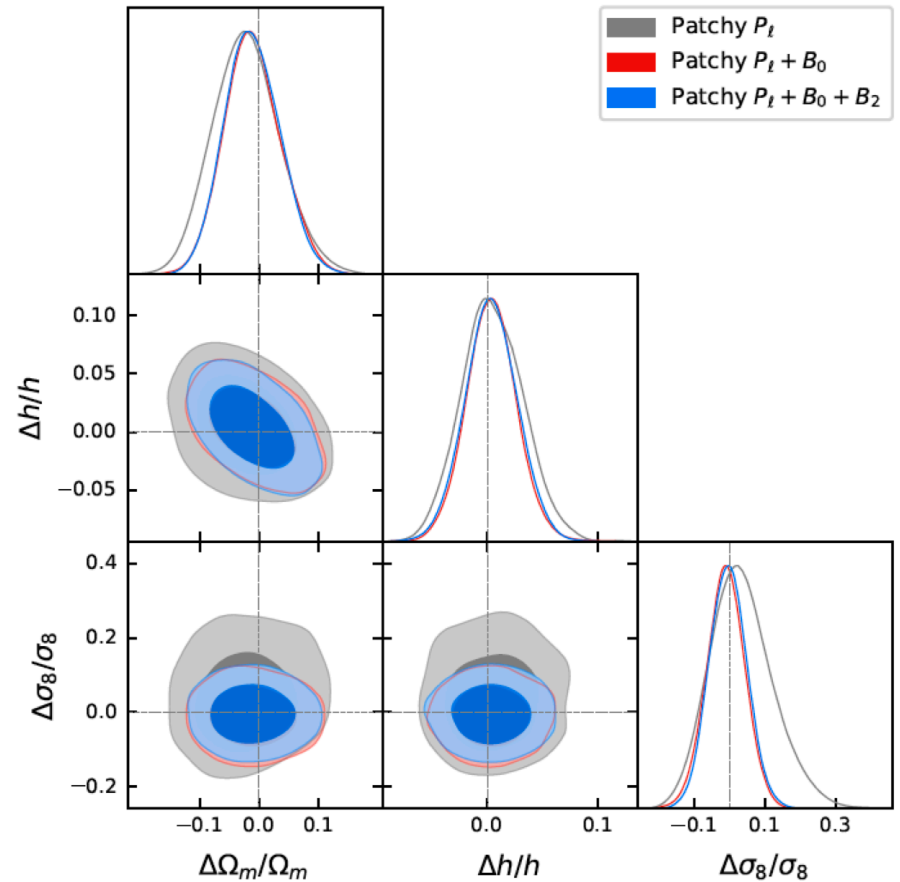
- **Binning** effect: it is performed (together with AP) exactly on the linear part, loop terms are small
- Effect of approximations are small

$\Delta_{\text{shift}}/\sigma_{\text{stat}}$	Ω_m	h	σ_8
$P_\ell + B_0$: base - w/ NNLO	-0.03	-0.09	-0.03
$P_\ell + B_0$: base - w/o B_0 window	0.11	-0.05	0.01
$P_\ell + B_0 + B_2$: base - w/o B_0, B_2 window	0.51	0.09	0.02

Theoretical error



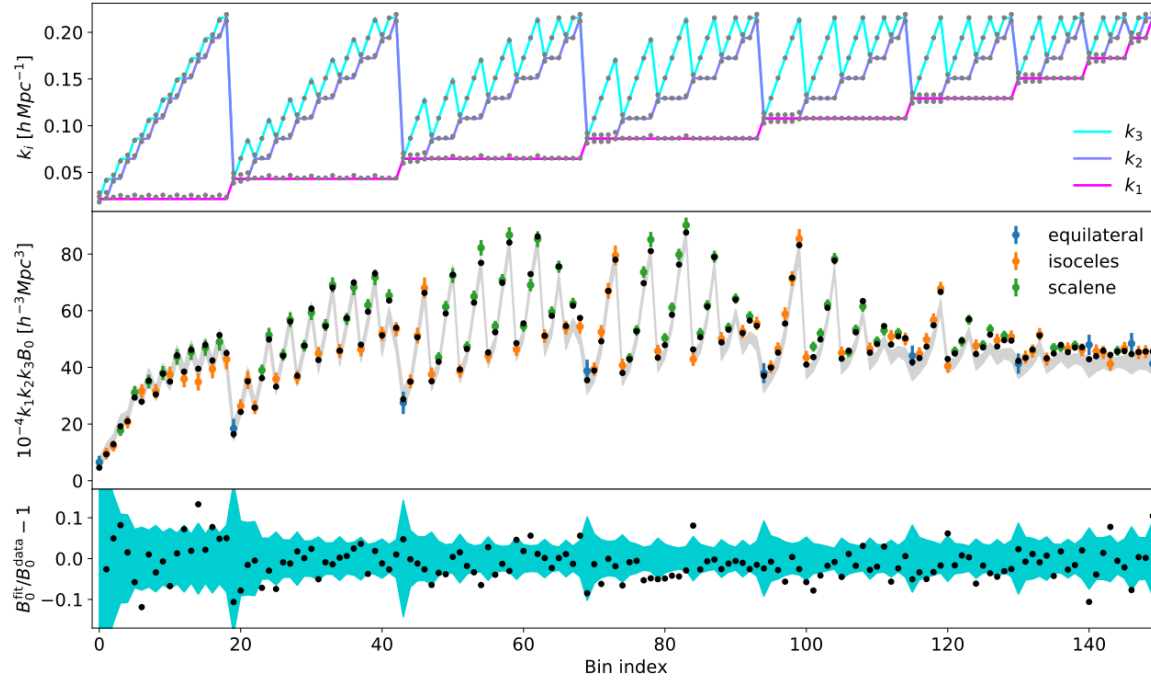
Nseries: 80 x BOSS volume



Patchy mocks: 2000 x BOSS volume

Fit with BOSS-volume covariance, safely within $\sigma_{data}/3$!

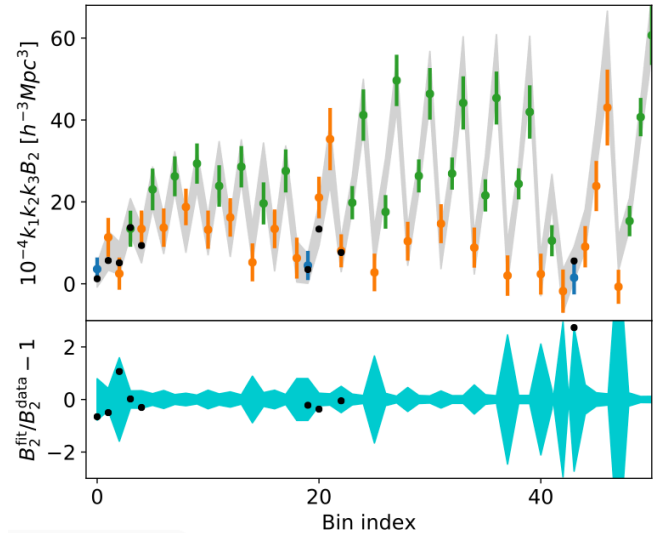
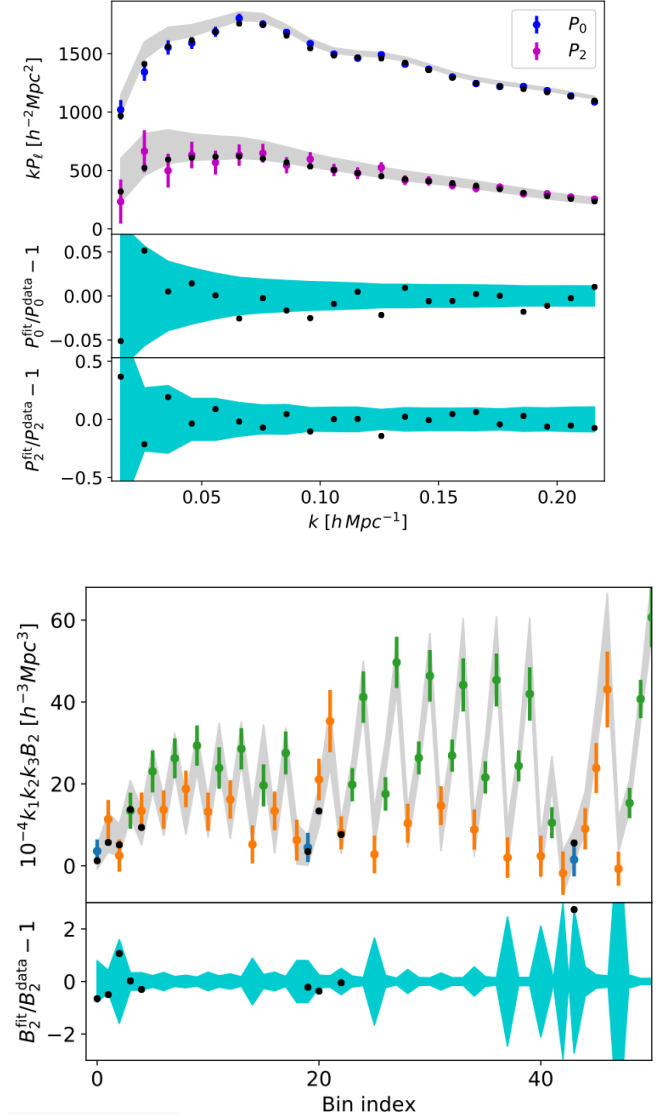
Best fit



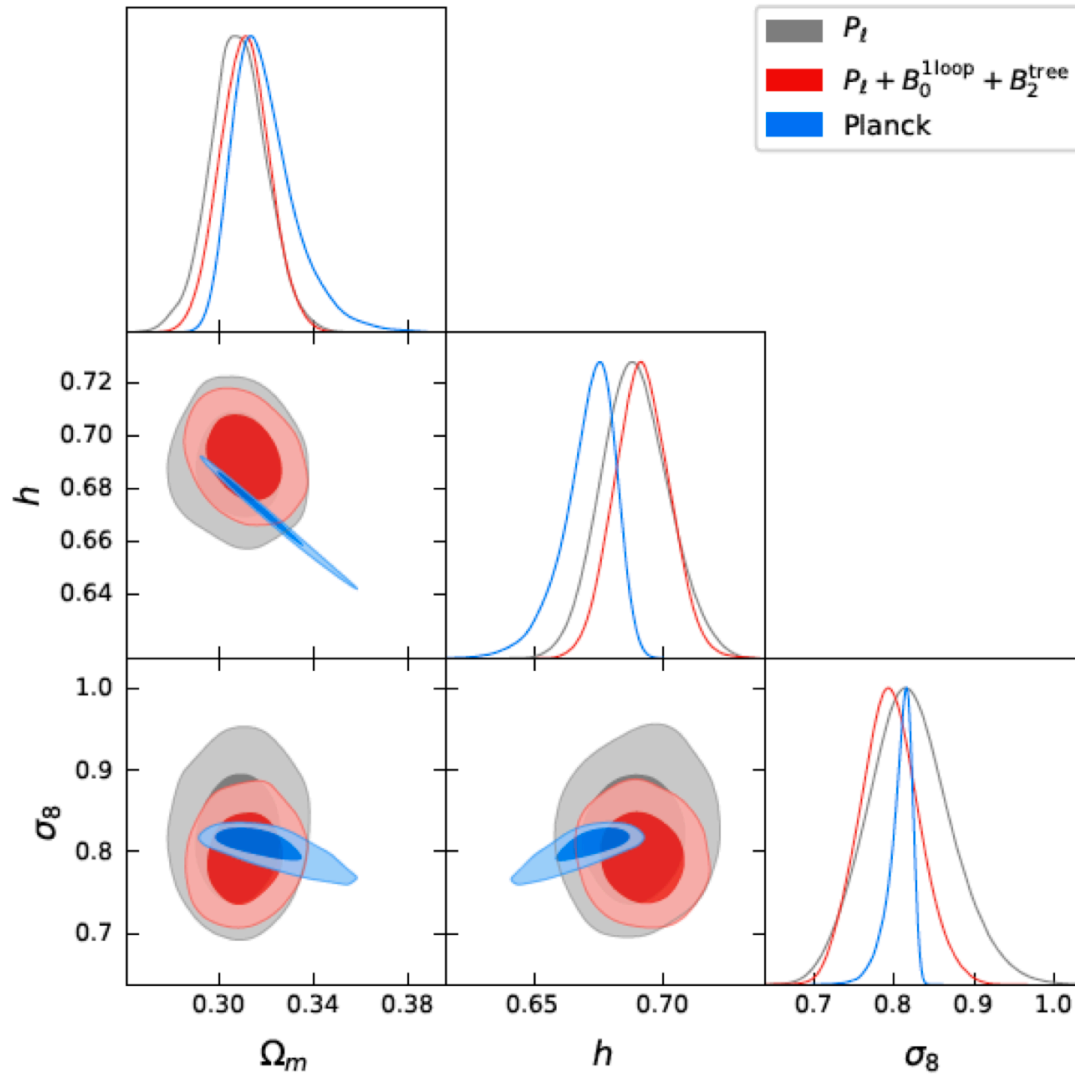
B_0 up to $k_{\text{max}} = 0.23 \text{ h/Mpc}$

B_2 tree-level, up to $k_{\text{max}} = 0.08 \text{ h/Mpc}$

Binned in triangles of 12 $k_f \sim 0.02 \text{ h/Mpc}$

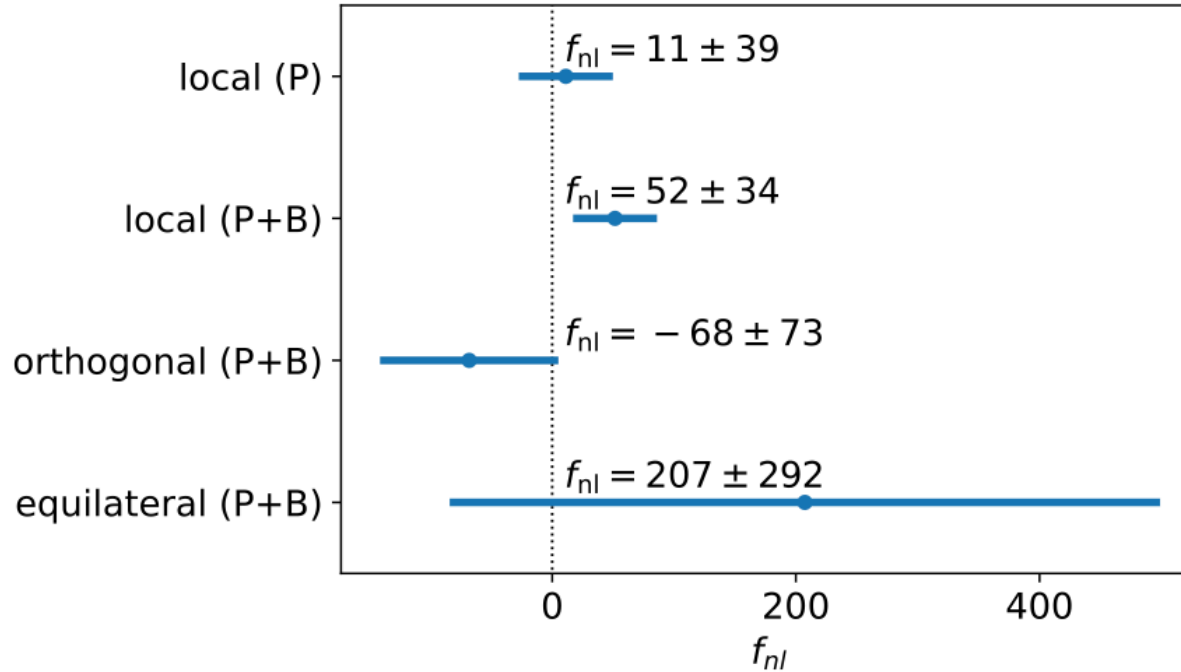


Results



- Improvements of
13% on Ω_m
18% on h
30% on σ_8
- Agreement with Planck

And with PNG!



	BOSS	WMAP	Planck
$f_{\text{NL}}^{\text{equil.}}$	207 ± 292	51 ± 136	-26 ± 47
$f_{\text{NL}}^{\text{orth.}}$	-68 ± 73	-245 ± 100	-38 ± 24
$f_{\text{NL}}^{\text{loc.}}$	52 ± 34	37.2 ± 19.9	-0.9 ± 5.1

Not only spectroscopy in the sky

- **Photometric surveys** will also be important now — e.g., for Euclid lensing is as important as clustering
- Why can't we apply the EFTofLSS to 3x2pt analysis? Turns out that, with a bit of effort, we can. Look also at DESxDESI 2x2pt (Chen et al. 2024)
- **Some peculiarities can be dealt with in an EFT-like way**: baryonic effects and intrinsic alignments. **Systematics are what they are**, hopefully improvements in the future
- Main practical challenges: determine the scale cuts, and keep the number of free parameters under control

Projection integrals

$$w^i(\theta) = \int \frac{dl l}{2\pi} J_0(l\theta) \frac{2}{\pi} \int dk k^2 \int d\chi_1 \int d\chi_2 f_{\delta_g}^i(\chi_1) f_{\delta_g}^j(\chi_2) P_{gg}(k, z(\chi_1), z(\chi_2))$$

$$\gamma_t^{ij}(\theta) = \int \frac{dl l}{2\pi} J_2(l\theta) \frac{2}{\pi} \int dk k^2 \int d\chi_1 \int d\chi_2 f_{\delta_g}^i(\chi_1) f_{\kappa}^j(\chi_2) P_{gm}(k, z(\chi_1), z(\chi_2))$$

$$\xi_+^{ij}(\theta) = \int \frac{dl l}{2\pi} J_0(l\theta) \int d\chi \frac{f_{\kappa}^i(\chi) f_{\kappa}^j(\chi)}{\chi^2} P_{mm}\left(\frac{l}{\chi}; z(\chi)\right)$$

$$\xi_-^{ij}(\theta) = \int \frac{dl l}{2\pi} J_4(l\theta) \int d\chi \frac{f_{\kappa}^i(\chi) f_{\kappa}^j(\chi)}{\chi^2} P_{mm}\left(\frac{l}{\chi}; z(\chi)\right)$$

Modeling choices

- In general, we use unequal-time PS at linear level; and Limber approximation, so equal-time PS, at loop level
- RSD and magnification term are linear, in Limber approx.

- P_{mm}

$$P_{mm,\text{lin}}(k, z_1, z_2) = D(z_1)D(z_2)P_{11}(k)$$

$$P_{mm,1L}(k, z) = D^4(z) (P_{22}(k) + P_{13}(k)) + 2c_s(z)^2 D^2(z) P_{11}(k) \frac{k^2}{k_{\text{NL}}^2}$$

- P_{gg}

$$P_{gg,\text{lin}}(k, z_1, z_2) = D(z_1)D(z_2)b_1(z_1)b_1(z_2)P_{11}(k)$$

$$P_{gg,1L}(k, z) = D^4(z) (P_{gg,22}(k; b_1, c_2) + P_{gg,13}(k; b_1, b_3)) + 2D^2(z)b_1(z)c_{gg}(z)P_{11}(k) \frac{k^2}{k_{\text{M}}^2}$$

- P_{gm}

$$P_{gm,\text{lin}}(k, z_1, z_2) = D(z_1)D(z_2)b_1(z_1)P_{11}(k)$$

$$P_{gm,1L}(k, z) = D^4(z) (P_{gm,22}(k; b_1, c_2) + P_{gm,13}(k; b_1, b_3)) + 2c_{gs}(z)D^2(z)P_{11}(k) \frac{k^2}{k_{\text{M}}^2}$$

Intrinsic alignments

- Shear modifies the shape of lensed galaxies. But these have intrinsic ellipticity, whose distribution is correlated with LSS
- And shear is inferred from observed ellipticity of background galaxies: thus, **intrinsic alignment term** $\gamma_{ab}(\hat{n}) \rightarrow \gamma_{ab}(\hat{n}) + \gamma_{ab}^{IA}(\hat{n})$
- Can be computed in EFT, and in fact it has:
Vlah, Chisari, Schmidt 2019; Chen, Kokron 2023
- DES-Y3 uses TATT model (Blazek et al. 2019); **we use the simpler NLA model**, used in Y1 (Hirata, Seljak 2004)

$$f_{\kappa}^j(z) \rightarrow f_{\kappa}^j(z) - A \left(\frac{1+z}{1+z_0} \right)^{\alpha} \frac{C_1 \rho_{m,0}}{D(z)} \frac{n_{source}^j(z)}{\bar{n}_s^j} \frac{dz}{d\chi}$$

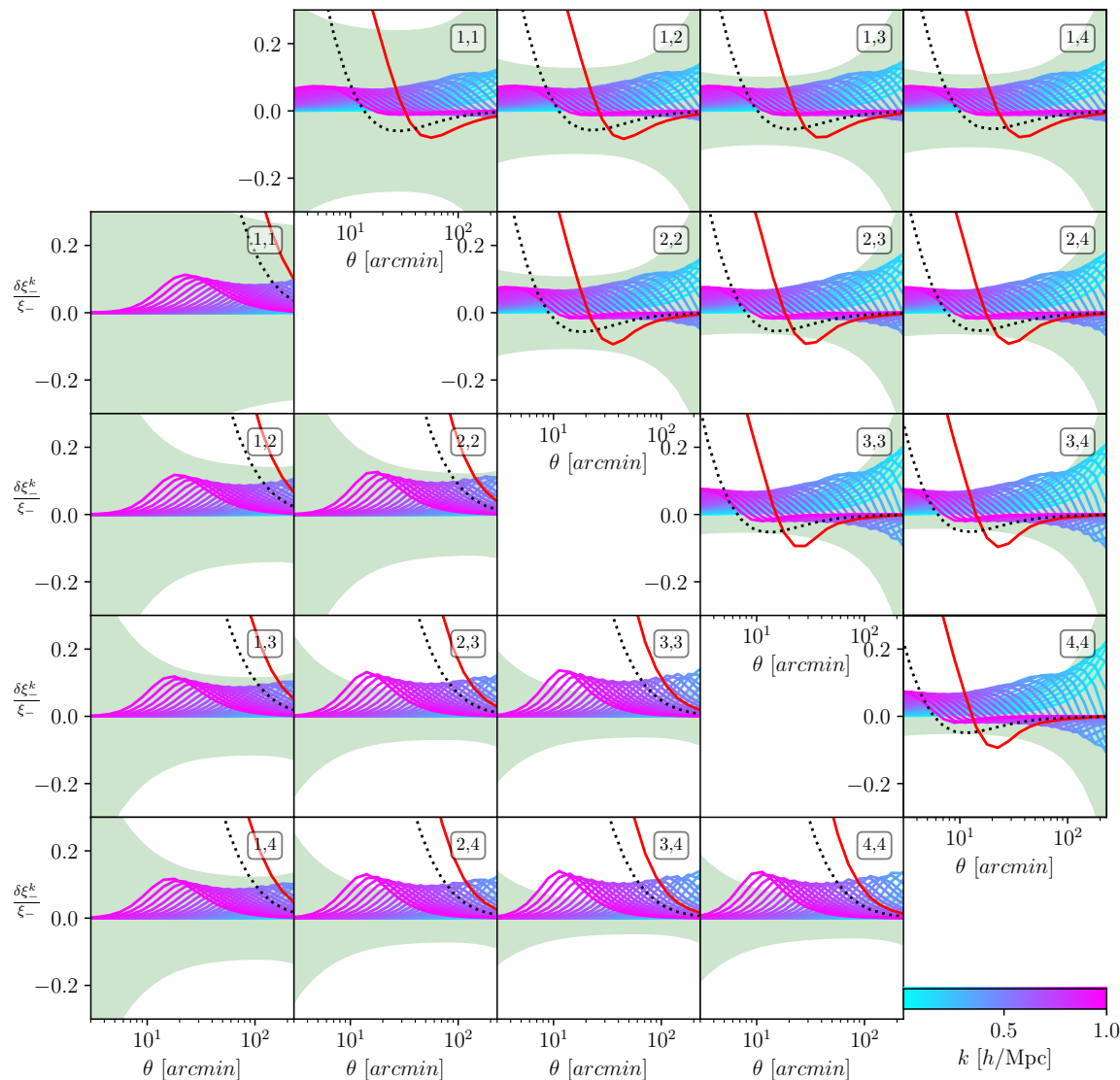
Integral over unknown bias functions

- Standard EFTofLSS expressions for the P_{mm} , P_{gm} , P_{gg} , but we have **integrals over parameters with unknown time dependence!**
- **Assume parameter is almost constant over the bin**, so each redshift bin combination now has its own parameters, e.g.

$$\xi_{+/-}^{ij}(\theta) \supset 2 \frac{c_{+/-}^{ij}}{k_{\text{NL}}^2} \int \frac{dl}{2\pi} l J_{0/4}(l\theta) \int d\chi \frac{f_{\kappa}^i(\chi) f_{\kappa}^j(\chi)}{\chi^2} D_+^2(z) P_{11} \left(\frac{l}{\chi} \right) \frac{l^2}{\chi^2}.$$

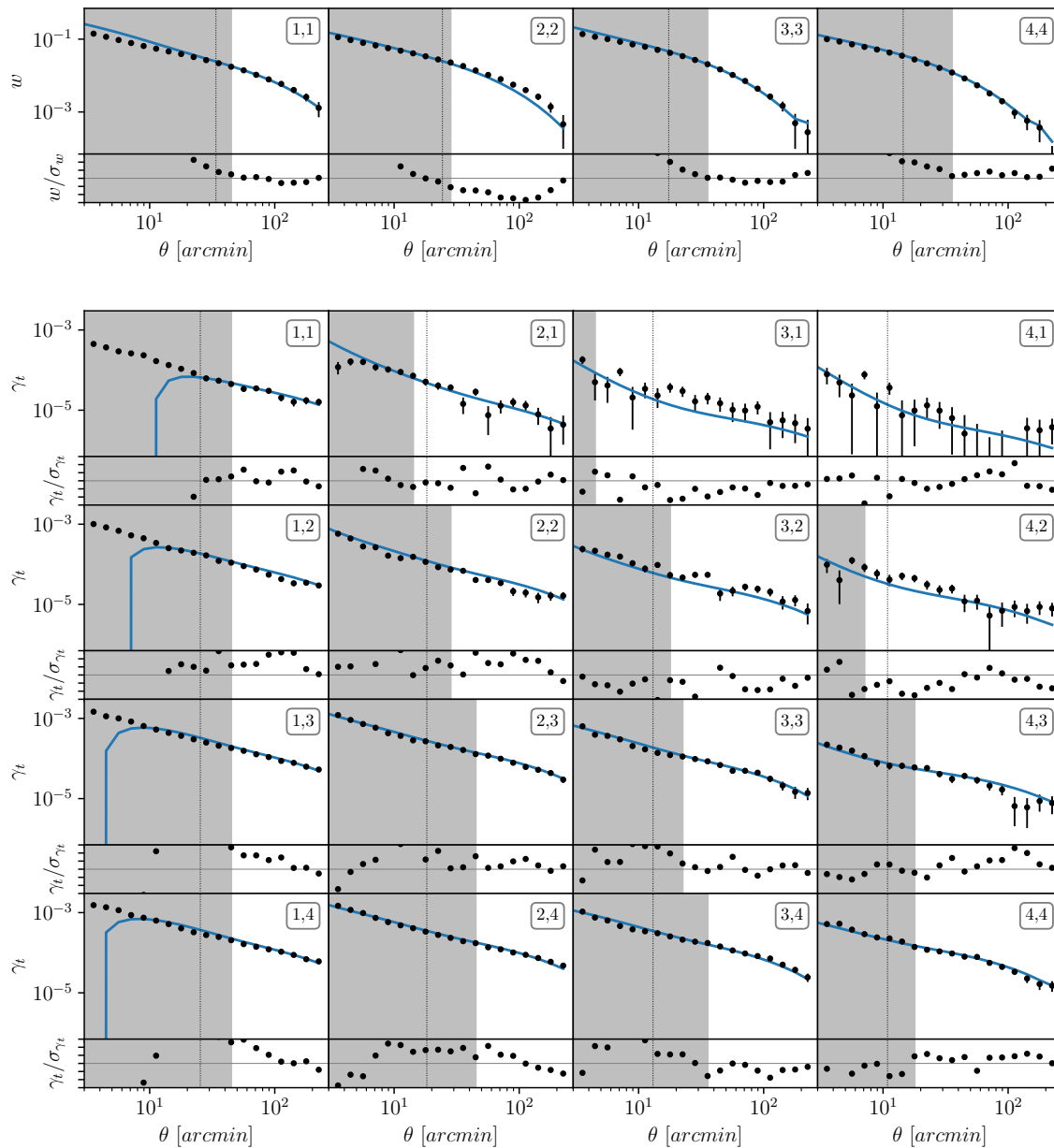
- **How to reduce their impact? Simple, we correlate them!**
After all, their time dependence is smooth and of order Hubble.

Scale cuts: NNLO and UV sensitivity

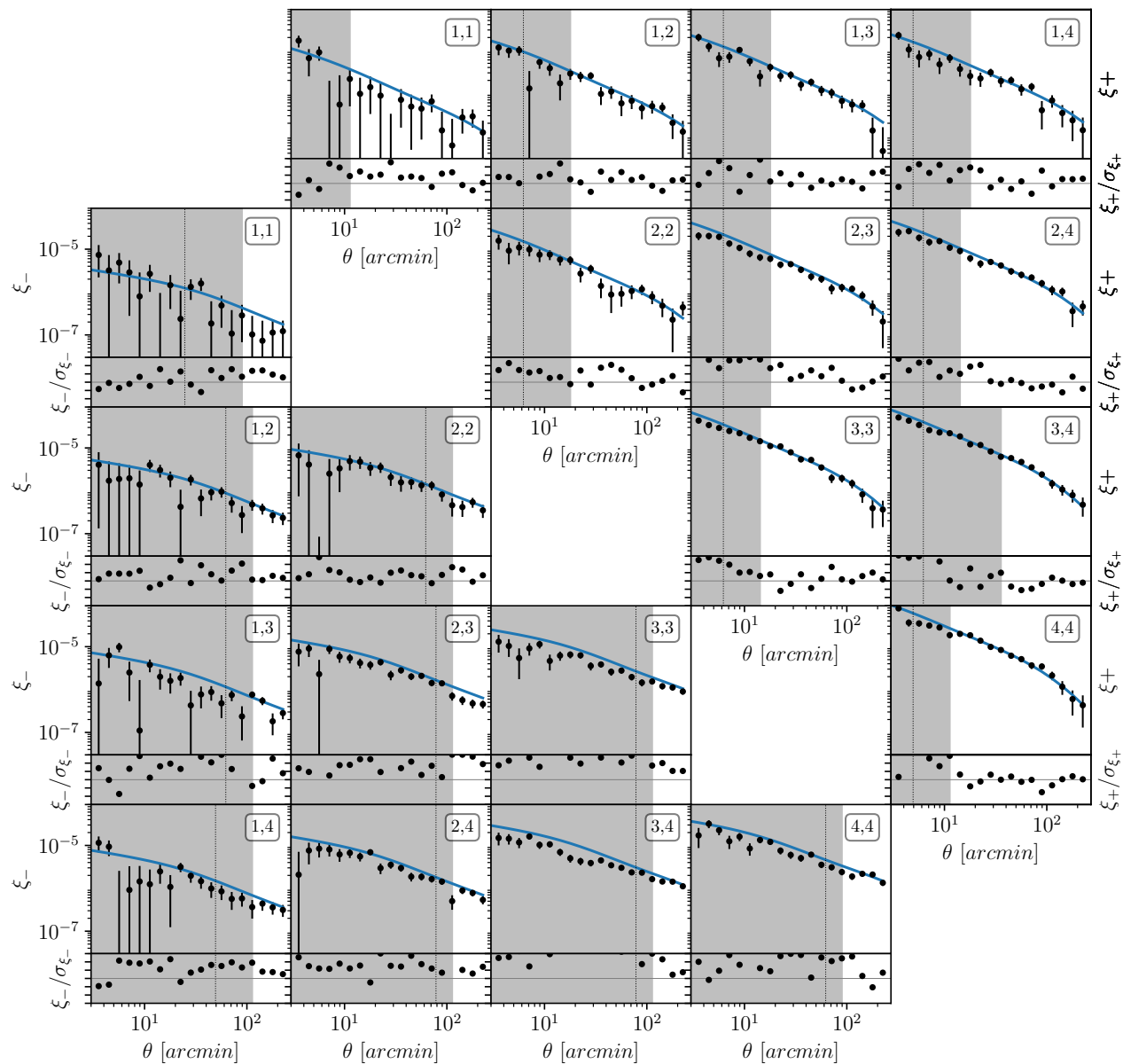


Results on Buzzard simulations

DES Y3 data and best-fit



DES Y3 data and best-fit



A new measurement of H_0

$$k_{\text{eq}} = H_0 \frac{\sqrt{2}\Omega_m}{\sqrt{\Omega_r}} \simeq 0.073 \Omega_m h^2 \text{ Mpc}^{-1}$$

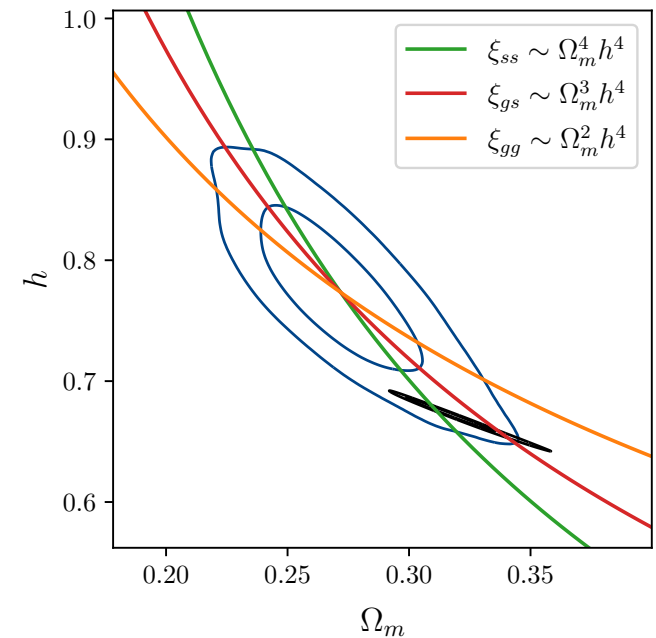
- Scale very small: but it enters the amplitude of power spectrum

$$P_{lin}(k, z) \approx D(z)^2 A_s \left(\frac{k}{k_*}\right)^{n_s-1} k (\Omega_m H_0^2)^{-2} \left(\frac{k_{\text{eq}}}{k}\right)^4 \log\left(\frac{k}{k_{\text{eq}}}\right)^2$$

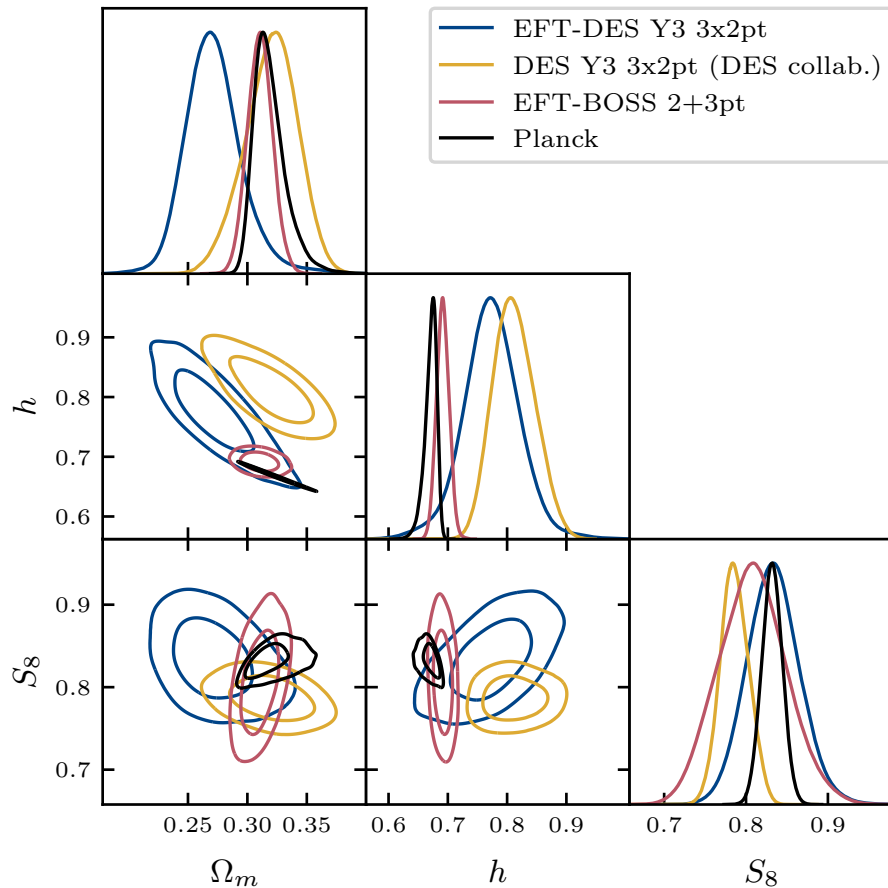
- This gives, for the scales of interest

$$P_{lin}(k, z) \propto D(z)^2 A_s (\Omega_m h^2)^2 k^\nu$$

- In fact, considering the lensing kernels, you can roughly predict the degeneracy Ω_m - h !



Flagship results



- We have larger errors than the official analysis. This is due to larger number of parameters and different scale cuts.
- But, we agree better with Planck and BOSS, especially in Ω_m - h
- Full posterior agreement at 1.7σ with BOSS and 2.3σ with Planck. Official analysis discrepant at 4.9σ with Planck.
- Reasonable precision on Ω_m , S_8 , not competitive on h .
- But, independent h measurement! It basically comes from k_{eq}

Where do we stand?

- Power spectrum at one loop: settled.
 - Several codes agreeing, and official analysis of DESI/Euclid
- Window function, real space $\xi(r)$, addition of BAO, all done
- Full EFT pipeline: self-consistent determination of k_{max} and analytical covariance available (Wadekar, Scoccimarro 2019)
- Prior on biases: reasonable perturbative prior, possible to include small-scale information

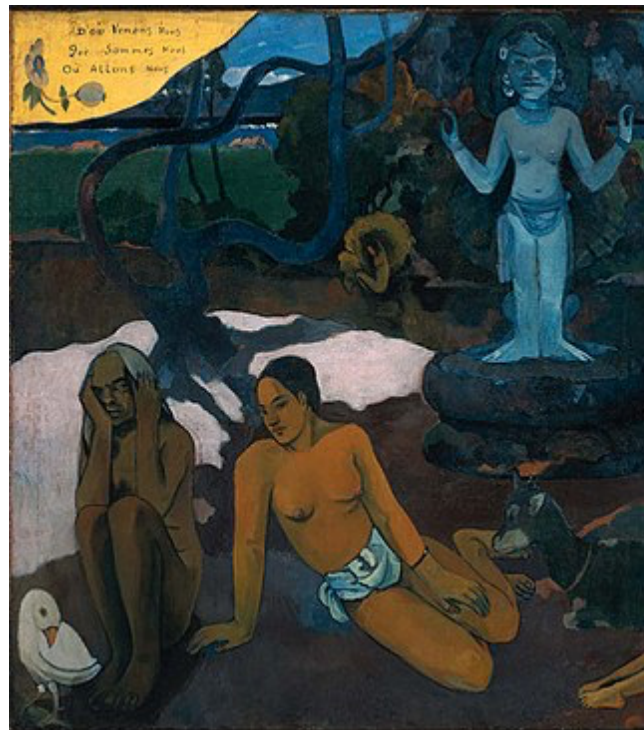
Where do we stand?

- Bispectrum at one loop: almost done
 - Done: one loop expression (twice), efficient implementation (twice), k_{max} determination
 - Done: study of IR resummation (Chen, Vlah, White, 2024)
 - Missing: a practical treatment of the survey mask
 - Missing: full analytical covariance. If too many triangles, can use mocks and compression, and approximate covariance available (Salvalaggio et al, 2024)
 - Not an official analysis yet, but soon

Where do we stand?

- 3x2pt: in progress
 - Definitely doable, needs more studying
 - Main difficulties (data side): systematics, and their modelling
 - Main difficulties (theory side)
 - Lots of parameters: can we reliably use some information on bias evolution?
 - Shear-shear: N-body are a good UV theory here, with 1-loop matter PS we stop early. Can do hybrid approach?
 - Does new physics show up at all? Which new physics can we probe by photometry?

Où allons-nous ?



A note on priors

- In Bayesian analysis, data update our belief on the model and its parameters. Must start from a probability measure on parameters. *There is no “uninformative prior”.*
- And what if data are not precise enough?

The status of LSS EFT analysis: two manifestly equivalent implementations of the manifestly correct theory on manifestly the same data produce manifestly different results. [#cosmology](#)

Priors on counterterms? The encarnation of the devil on Earth.



A note on priors

- What priors on EFT parameters?
We know they have to be perturbative correction, but not much more: except for a few, we center them at 0 with $\sigma=2$.
- Other choices are possible, as long as they cover the physically allowed region, but **EFT parametrization stays the same**.
“West Coast” and “East Coast” parameters are a *linear transformation of each other*. **Prior choice is, however, different**.
- Maximum likelihood points are unchanged. But credible intervals can shift due to projection.
- What should you put in the abstract on your paper?

$$\mathbb{E}_{\mathcal{P}}[\lambda_{\alpha_1} \dots \lambda_{\alpha_n}] = Z^{-1} \int \mathcal{M}(\boldsymbol{\lambda}) d^k \boldsymbol{\lambda} \lambda_{\alpha_1} \dots \lambda_{\alpha_n} \mathcal{P}(\boldsymbol{\lambda})$$

Simon, Zhang, Poulin, Smith (2022)
Carrilho, Moretti, Pourtsidou (2022)
Reeves, Zhang, Zheng (2025)

Fixing phase space issues

- Our solution, for bispectrum: adjusting the prior, measuring the effect on synthetic data fit to our data

$$\ln \mathcal{M} = -48 \left(\frac{b_1}{2} \right) + 32 \left(\frac{\Omega_m}{0.31} \right) + 48 \left(\frac{h}{0.68} \right)$$

$\sigma_{\text{proj}}/\sigma_{\text{stat}}$	Ω_m	h	σ_8	ω_{cdm}
1 sky, $\sim 100 V_{1\text{sky}}$	-0.1	-0.14	-0.21	-0.2
1 sky, $V_{1\text{sky}}$, adjust.	0.13	0.06	0.04	0.15
4 skies, $V_{4\text{skies}}$, adjust.	0.1	0.	-0.05	0.07

- On 3x2pt, similar thing:

$$\log \mathcal{M} = -55 \omega_{\text{cdm}} - 9 b_1$$

You are doing a morally wrong thing!

What to do?

- Just wait. Sooner or later data will just not care about the priors?
- Perturbativity prior

$$\mathcal{P}_P = \frac{1}{2N_{\text{bins}}^P} \sum_{i \in \text{bins}_P} \left(\frac{P_{1\text{-loop}}^h(k_i)}{\sigma_P^{\text{P.P.}}(k_i)} \right)^2 \quad \mathcal{P}_B = \frac{1}{2N_{\text{bins}}^B} \sum_{i \in \text{bins}_B} \left(\frac{B_{1\text{-loop}}^h(k_1^i, k_2^i, k_3^i)}{\sigma_B^{\text{P.P.}}(k_1^i, k_2^i, k_3^i)} \right)^2$$

$$\sigma_P^{\text{P.P.}}(k) \sim S^P(k) P_{1\text{-loop}}^{k_{\text{max}}} , \quad \sigma_B^{\text{P.P.}}(k_1, k_2, k_3) \sim S^B(k_1, k_2, k_3) B_{1\text{-loop}}^{k_{\text{max}}}$$

$$S_{1\text{-loop}}^P(k) \sim b_1^2 P_{11}(k) \left(\frac{k}{k_{\text{NL}}} \right)^{3+n(k)} \quad S_{1\text{-loop}}^B(k_1, k_2, k_3) \sim B_{211}^h(k_1, k_2, k_3) \sum_{i=1}^3 \left(\frac{k_i}{k_{\text{NL}}} \right)^{3+n(k_i)}$$

Bragança, Donath, Senatore, Zheng (2023)
GDA, Lewandowski, Senatore, Zhang (2025)

- Get some controlled UV information.

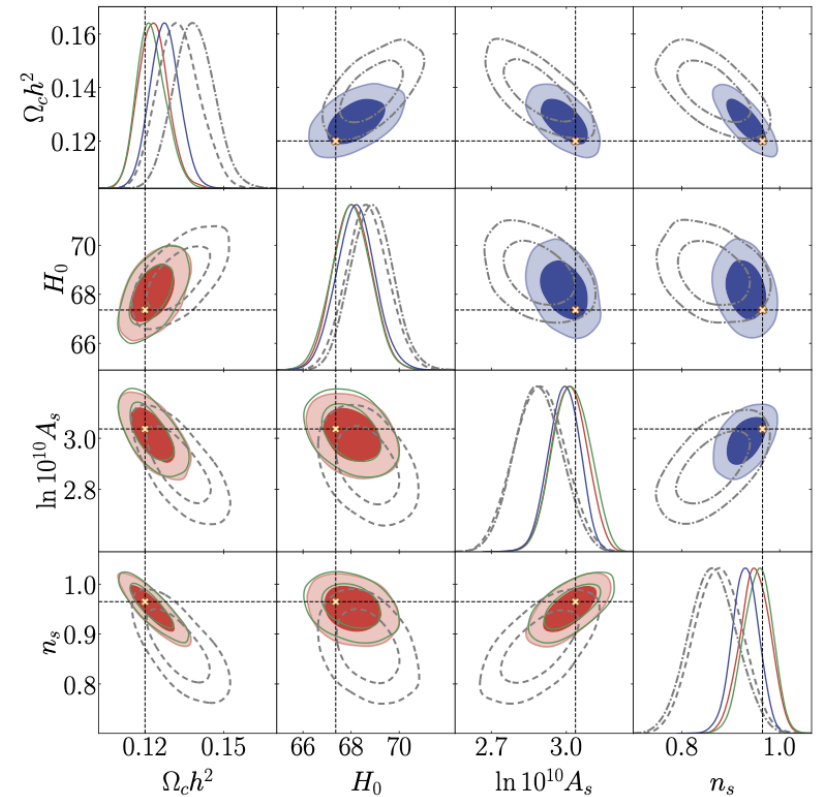
We used some old fitting formulas from simulations.

Now, the idea is to do a dedicated search. As a first step, fitting HOD.

But have to marginalise over them.

Small-scale information

- Can we get the best of both worlds?
- Simple, but effective, attempt:
translate HOD priors into EFT (Zhang, Bonici, GDA, et al. 2024; see also Ivanov et al. 2024)
- Interesting more general direction:
Hybrid SBI (Modi, Philcox 2024; Zhang, Modi, Philcox 2025)
- Can galaxies be reliably simulated?
In other words, can we claim detection of new physics based on small-scale information?



Where do we go now, technically?

- Higher loops?
 - 2-loop power spectrum requires 5th order. Several works on it, soon it will be done.
Integrals are complicated, but not an issue, really.
Determination of k_{max} easy.
 - But, will it be useful? Depending on the tracer and the data noise, it can be. Also, could construct $\mu \rightarrow 0$ variable and practically reach 0.5

Where do we go now, technically?

- Higher point functions / field-level?
 - Today could include up to tree-level 5-pt, though additional 4 biases.
 - But at least tree-level trispectrum has no more biases than 1-loop bispectrum.
 - Measurement and analysis seems daunting.
 - Field-level model and measurements simple. But, how about the likelihood?
 - But forward model super useful to help with standard analysis.

Getting new physics

- Λ CDM very well studied (and boring).
Cosmology's before LHC moment

Known new physics

- Neutrinos exist, they are light: must measure their mass, and we have surprises
- Inflationary dynamics is not free fields: need to push for primordial non-Gaussianities, not only local but single-field

Under the lamppost physics

- Dynamical dark energy: may or may not exist, no scalar field models solve the CC problem. Super interesting to look for, but must go beyond w_0 - w_a
- CDM, only billiard balls? No.
Here a question of scales (and naturalness): LSS exquisite probe for ultralights!

Getting new physics

- Can we just claim that “Something is rotten in the state of Λ CDM”?
- First time that we can use precise data at vastly different times, if not complementary length scales: CMB at $z=1100$, LSS at $z=0.3-2$, also SN, Ly- α , what else in the future?
- Could expose anomalies and tensions. Need some parametrisation though.
- **LSS bootstrap**: Impose symmetries, and parametrize the kernels with basic building blocks. Can match to specific models

$$G_2(\vec{k}_1, \vec{k}_2) = 2\beta(\vec{k}_1, \vec{k}_2) + d_1^{(2)} \gamma(\vec{k}_1, \vec{k}_2)$$

$$\begin{aligned} G_3(\vec{k}_1, \vec{k}_2, \vec{k}_3) = & 2\beta(\vec{k}_1, \vec{k}_2)\beta(\vec{k}_{12}, \vec{k}_3) + d_5^{(3)}\gamma(\vec{k}_1, \vec{k}_2)\gamma(\vec{k}_{12}, \vec{k}_3) - 2\left(d_{10}^{(3)} - h\right)\gamma(\vec{k}_1, \vec{k}_2)\beta(\vec{k}_{12}, \vec{k}_3) \\ & + 2(d_1^{(2)} + 2d_{10}^{(3)} - h)\beta(\vec{k}_1, \vec{k}_2)\gamma(\vec{k}_{12}, \vec{k}_3) + d_{10}^{(3)}\gamma(\vec{k}_1, \vec{k}_2)\alpha_a(\vec{k}_{12}, \vec{k}_3) + \text{cyc.} \end{aligned}$$

Fujita, Vlah (2020)

GDA, Marinucci, Pietroni, Vernizzi (2021)

Taule et al. (2024)

Marinucci, Pardede, Pietroni (2024)

Getting new physics

- Want to be more general? For instance, capture scale dependence
- Maybe we can use machine learning (agnostic parametrizations)
Already applied with success in particle physics

$$n(x|\vec{w}) = n(x|R)e^{f(x;\vec{w})}$$

$$t(\mathcal{D}) = -2\text{Min}_{\vec{w}}[N(\vec{w}) - N(R) - \sum_{x \in \mathcal{D}} f(x; \vec{w})]$$

End of the beginning or beginning of the end?

- EFTofLSS has been successful, stays strong in the wake of AI.
1-loop power spectrum, 1-loop bispectrum, 3x2pt.
What is next for it?
- Maybe more technical developments.
Higher loops? Field level analysis?
- Other observables? Ly- α , 21cm?
- Should we become friends with the machine?
Power analyses semianalytically with EFT forward model.
Controlled inclusion of small-scale information.
- Getting new physics
Must push for known physics, get tight constraints.
Solid modelling of beyond-standard models.
And surprises? Model-agnostic assessment of new physics.

Thank you!