

TOWARDS RELATIVISTIC ATOMIC PHYSICS AND POST-MINKOWSKIAN GRAVITATIONAL WAVES

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LUSANNA - THE CHRONOGEOMETRICAL STRUCTURE OF SPECIAL AND GENERAL
RELATIVITY: A REVISITATION OF CANONICAL GEOTHEODYNAMICS

42 KARPACZ WINTER SCHOOL INT. J. GEOM. METHODS IN MOD. PHYS. 4 (2007) 79
GR-9C/0604120

ALBA-LUSANNA - GENERALIZED RADAR 4-COORDINATES AND EQUAL TIME CAUCHY SURFACES
FOR ARBITRARY ACCELERATED OBSERVERS

INT. J. MOD. PHYS. D16 (2007) 1149 GR-9C/0501090

ALBA-LUSANNA-PAURI - CENTERS OF MASS AND ROTATIONAL KINEMATICS FOR THE
RELATIVISTIC N-BODY PROBLEM IN THE REST-FRAME INSTANT FORM

J. MATH. PHYS. 43 (2002) 1677 hep-th/0102087

CRATER-LUSANNA - THE REST-FRAME DARWIN POTENTIAL FROM THE LIENARD-WIECHERT
SOLUTION IN THE RADIATION GAUGE

ANN. PHYS. ~~400~~ 289 (2001) 87

ALBA-CRATER-LUSANNA - HAMILTONIAN RELATIVISTIC 2-BODY PROBLEM: CENTER OF MASS
AND ORBIT RECONSTRUCTION

J. PHYS. A 40 (2007) 9535 GR-9C/0610200

ALBA-LUSANNA-CRATER - TOWARDS RELATIVISTIC ATOMIC PHYSICS. I. THE REST-FRAME

INSTANT FORM OF DYNAMICS AND A CANONICAL TRANSFORMATION FOR A SYSTEM OF
CHARGED PARTICLES PLUS THE ELECTROMAGNETIC FIELD

2008 ARXIV 0806.2383

II COLLECTIVE AND RELATIVE RELATIVISTIC VARIABLES FOR A SYSTEM OF CHARGED
PARTICLES PLUS THE ELECTROMAGNETIC FIELD IN PREPARATION

III CLOCK SYNCHRONIZATION, QUANTIZATION AND RELATIVISTIC ENTANGLEMENT
IN PREPARATION

ACES { ESA FIRENZE (2008) FTP://CACCIAPODI1:ln73rn0@FTP.ESTEC.ESA.INT/
ESA MÜNICH GEODESY (2008) <http://www.iapg.bv.tu.munich.de/12735-naces-programme.htm>

ALBA-LUSANNA - CHARGED PARTICLES AND THE ELECTROMAGNETIC FIELD IN NON-INERTIAL
FRAMES IN MINKOWSKI SPACETIME 0812.3057

GRAVITY LUSANNA - THE REST-FRAME INSTANT FORM OF PERIC GRAVITY GRG 33 (2001) 1579

DE PIETRI-LUSANNA-PARTUCCI-RUSSO - ... TETRAD GRAVITY GRG 34 (2002) 877

ALBA-LUSANNA - YORK MAP GRG 39 (2007) 2149 GR-9C/0604086

LUSANNA-PAURI - DYNAMICAL 3-SPACE AND HOLE ARGUMENT

GRG 38 (2006) 187 and 229 GR-9C/0403081, 0407007

HISTORY AND PHILOSOPHY OF MODERN PHYSICS 37 (2006) 692 GR-9C/0604087

MOTIVATIONS FOR INTRODUCING THE LIGHTCONE

SPACE PHYSICS - LIGHT PROPAGATION IN THE SOLAR SYSTEM

CLOCK SYNCHRONIZATION AND INSTANTANEOUS 3-SPACE - ACES (POST-NEWTONIAN NULL GEODESICS FOR LIGHT), CAUCHY PROBLEM, PREDICTABILITY

"PHOTONS" OF PROTOCOLS FOR ENTANGLEMENT → MASSLESS PARTICLES ON NULL GEODESICS

2-LEVEL ATOMS - RELATIVISTIC POINT PARTICLES WITH STRUCTURE ($\neq$ SPIN)

INERTIAL AND TIDAL EFFECT IN NON-INERTIAL FRAMES

INTERACTION OF GRAVITATIONAL WAVES WITH PARTICLES AND LIGHT

ATOM INTERFEROMETERS

2-BODY PROBLEM IN POST-MINKOWSKIAN GRAVITY

COSMOLOGY - HOW MUCH DARK MATTER AND DARK ENERGY CAN BE EXPLAINED BY EINSTEIN GR?

WE NEED A FULLY RELATIVISTIC THEORETICAL FORMULATION OF ATOMIC PHYSICS (EFFECTIVE THEORY BELOW THE THRESHOLD OF PAIR PRODUCTION) INTERACTING WITH A DYNAMICAL ELECTROMAGNETIC FIELD AND ITS EXTENSION TO GRAVITY

A) INERTIAL AND NON-INERTIAL FRAMES IN MINKOWSKI SPACETIME

ISOLATED SYSTEMS, POINCARÉ GENERATORS, RELATIVISTIC CENTER OF MASS, INERTIAL AND NON-INERTIAL REST-FRAME INSTANT FORM OF DYNAMICS

B) LINEARIZATION TO POST-MINKOWSKIAN GRAVITY IN ASYMPTOTICALLY MINKOWSKIAN SPACETIMES (ONLY NON INERTIAL FRAMES DUE TO EQUIVALENCE PRINCIPLE)

GRAVITATIONAL WAVES WITH ASYMPTOTIC BACKGROUND



C) RELATIVISTIC QUANTUM MECHANICS AND RELATIVISTIC ENTANGLEMENT

FAILURE OF ZEROth POSTULATE → NON-LOCALITY AND SPATIAL NON-SEPARABILITY

RELATIVISTIC ASPECTS PRESENT AT $1/c^2$ LEVEL WHICH CANNOT BE IGNORED IN SPACE PHYSICS

NOT 1+3 BUT 3+1 AND RADAR COORDINATES

CLASSICAL AND QUANTUM THEORIES

PDE PREDICTABILITY $\rightarrow$ CAUCHY PROBLEM

$\downarrow$
INSTANTANEOUS 3-SPACE

$\parallel$
CLOCK SYNCHRONIZATION

1) GALILEI SPACE-TIME

$\left\{ \begin{array}{l} \text{INERTIAL FRAMES} \\ \text{NON-INERTIAL FRAMES} \end{array} \right.$

2) MINKOWSKI SPACE-TIME

3+1 SPLITTINGS
AND THEIR GAUGE EQUIVALENCE

$\left\{ \begin{array}{l} \text{INERTIAL FRAMES} \\ \text{NON-INERTIAL FRAMES} \end{array} \right.$

ISOLATED SYSTEMS (PARTICLES, STRINGS, FIELDS, FLUIDS)

RELATIVISTIC CENTER OF MASS AND POINCARÉ GROUP

QUANTUM REL. MECHANICS

REL. ATOMIC PHYSICS

REL. ENTANGLEMENT

QUANTIZATION IN NON-INERTIAL FRAMES

3) EINSTEIN SPACE-TIMES

3+1 SPLITTINGS

ONLY NON-INERTIAL FRAMES
EQUIVALENCE PRINCIPLE

ACES - CLOCK SYNCHRONIZATION IN SOLAR SYSTEM (PN GRAVITY)

YORK GAP - INERTIAL AND TIDAL EFFECTS IN CANONICAL GRAVITY

COSMOLOGY - FRW PARADIGM

$\left\{ \begin{array}{l} \text{GR} \rightarrow \text{DARK MATTER (RELATIVISTIC INERTIAL EFFECT?)} \\ \text{GR} \rightarrow \text{DARK ENERGY (BACKREACTION? BUCHERT, ELUS, ...)} \end{array} \right.$

$\rightarrow$ CAUCHY PROBLEM, CLOCK SYNCHRONIZATION

3+1
1+3
VIEWPOINTS - BOTH NON FACTUAL

NON-RELATIVISTIC PHYSICS

(NEWTON'S THEORY IN GALILEI SPACETIME)

ABSOLUTE TIME
ABSOLUTE 3-SPACE (EUCLIDEAN)



ABSOLUTE SIMULTANEITY AND SPATIAL DISTANCE

GALILEI LAW OF INERTIA $\rightarrow$ INERTIAL OBSERVERS

GALILEI RELATIVITY PRINCIPLE - INVARIANCE IN FORM OF $\vec{F} = m\vec{a}$ IN INERTIAL FRAMES
(CENTERED ON INERTIAL OBSERVERS) WITH CARTESIAN
3-COORDINATES - KINEMATICAL GALILEI GROUP

NEWTON'S GRAVITY - A-A-A-D INTERACTION WITH GALILEI EQUIVALENCE PRINCIPLE

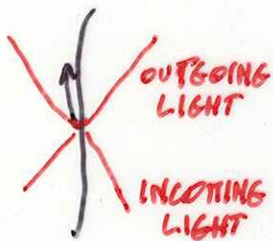
$$m\vec{a} = \vec{F}_g = m_g \vec{g} \quad m = m_g$$

INERTIAL FRAMES - IDEAL LIMIT

RIGID NON-INERTIAL FRAMES CENTERED - INERTIAL FORCES $\propto$ INERTIAL MASS
ON ACCELERATED OBSERVERS

$$\vec{a} = \underbrace{\frac{\vec{F}}{m}}_{\text{LINEAR}} - \underbrace{\vec{\omega}(t)}_{\text{JACOBI}} + \underbrace{2 \frac{d\vec{x}}{dt} \times \vec{\omega}(t)}_{\text{CORIOLIS}} + \underbrace{\vec{\omega}(t) \times [\vec{x} \times \vec{\omega}(t)]}_{\text{CENTRIFUGAL}}$$

SPECIAL RELATIVITY (MINKOWSKI SPACETIME)



ABSOLUTE FLAT SPACETIME

$$ds^2 = [c^2 dt^2 - d\vec{x}^2] = \eta_{\mu\nu} dx^\mu dx^\nu$$

$$\eta_{\mu\nu} = \epsilon(t, -, -, -), \epsilon = \pm 1$$

CARTESIAN 4-COORDINATES

ONLY THE
CONFORMAL STRUCTURE
(THE LIGHT CONE) IS
INTRINSICALLY GIVEN

NO NOTION OF

INSTANTANEOUS 3-SPACE, SPATIAL DISTANCE
SIMULTANEITY (SYNCHRONIZATION OF DISTANT CLOCKS NEEDED)
CAUCHY 3-SURFACE FOR MAXWELL EQS
1-WAY VELOCITY OF LIGHT (2 SYNCHRONIZED CLOCKS NEEDED)

LIGHT POSTULATES - THE 2-WAY (OR ROUND TRIP; ONLY 1 CLOCK NEEDED) VELOCITY OF
LIGHT IS A) ISOTROPIC B) CONSTANT $c = 2.997924580 \cdot 10^8 \text{ m.s}^{-1}$

METROLOGY - c and standard of time

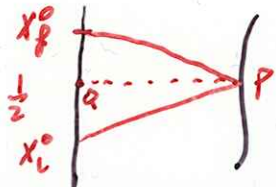
$$L = \frac{1}{2} c \Delta t$$

$\rightarrow$ but curved 3-space?

LAW OF INERTIA $\rightarrow$ INERTIAL OBSERVERS

SPECIAL RELATIVISTIC RELATIVITY PRINCIPLE - INVARIANCE IN FORM OF PHYSIC'S LAWS
IN INERTIAL FRAMES (CENTERED ON INERTIAL OBSERVERS) WITH CARTESIAN
4-COORDINATES - KINETICAL POINCARÉ GROUP

IDENTIFIED BY EINSTEIN'S $\frac{1}{2}$ CONVENTION FOR CLOCK SYNCHRONIZATION



$$x_P^0 \stackrel{\text{def}}{=} x_Q^0 = x_L^0 + \frac{1}{2}(x_R^0 - x_L^0) = \frac{1}{2}(x_R^0 + x_L^0)$$

$\Downarrow$

A (INERTIAL) B (NON-INERTIAL) 3 SPACES = EUCLIDEAN $x^0 = ct = \text{CONST}$ HYPERPLANES

1-WAY = 2-WAY = c



INERTIAL OBSERVERS = IDEAL LIMIT $\rightarrow$ ACCELERATED OBSERVERS

$$\begin{aligned} \text{ACCELERATION LENGTH} & \left\{ \begin{array}{ll} L = c^2/a & \text{TRANSLATIONAL} \quad (\sim 1 \text{ LY}) \\ L = \frac{c^2}{\omega} & \text{ROTATIONAL} \quad (\sim 10^{10} \text{ AU}) \end{array} \right. \end{aligned}$$

LOCALITY HYPOTHESIS - an accelerated observer is a succession of instantaneously comoving inertial observers

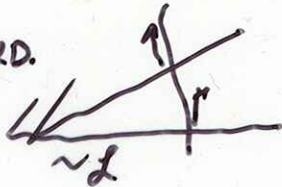
$\lambda > L$ PROBLEMS

1+3 POINT OF VIEW

ONLY THE OBSERVER'S WORLDLINE AND 4-VELOCITY ARE GIVEN

NO KNOWN METHOD TO DEFINE GLOBAL INSTANTANEOUS 3-SPACES (CAUCHY SURFACES)

FERMI COORD.



ROTATING DISK



$$\omega R \ll c \Rightarrow g_{00}(Y, \vec{Y} = \omega R \hat{n} \approx c \hat{n}) = 0$$

$\downarrow$ $x^\mu \rightarrow Y^\mu(x), \eta_{\mu\nu} \rightarrow g_{\mu\nu}(Y)$
SAGNAC EFFECT

COORDINATE SINGULARITIES

PARAMETRIZED MINKOWSKI THEORIES

WORLD-LINE OF AN ACCELERATED OBSERVER (TIRELIKE)

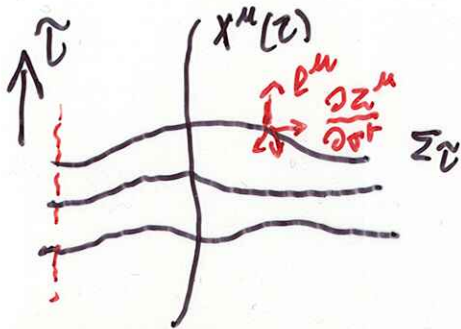
→ 3+1 SPLITTING OF MINKOWSKI SPACETIME

(= NICE FOLIATION WITH SPACELIKE 3-SURFACES =

= CLOCK SYNCHRONIZATION CONVENTION =

= DEFINITION OF THE INSTANTANEOUS 3-SPACE)

⇒ GLOBAL NON-INERTIAL FRAME CENTERED ON THE ACCELERATED OBSERVER



ASYMPTOTIC INERTIAL OBSERVER

FRAME- AND OBSERVER- DEPENDENT, LORENTZ-SCALAR

RADAR 4-COORDINATES (BONDI) $(\tau; \sigma^A)$

[DIFFERENTLY FROM FERMİ COORDINATES, THEY CAN BE OPERATIONALLY DEFINED]

$$\begin{cases} x^\mu \mapsto \sigma^A(x) = (\tau; \sigma^A) \\ \sigma^A \mapsto x^\mu = z^\mu(\tau, \sigma^A) \end{cases}$$

x^μ CARTESIAN 4-COORD. IN AN INERTIAL FRAME

EMBEDDING OF Σ_τ

$$\eta_{\mu\nu} \rightarrow g_{AB}[z(\tau, \sigma^A)] = \frac{\partial z^\mu(\sigma^A)}{\partial \sigma^A} \eta_{\mu\nu} \frac{\partial z^\nu(\sigma^B)}{\partial \sigma^B}$$

NORMAL TO Σ_τ

$$\frac{\partial z^\mu(\tau, \vec{\sigma})}{\partial \tau} = \underbrace{N(\tau, \vec{\sigma})}_{\text{LAPSE}} l^\mu(\tau, \vec{\sigma}) + \underbrace{N^r(\tau, \vec{\sigma})}_{\text{SHIFT}} \frac{\partial z^\mu(\tau, \vec{\sigma})}{\partial \sigma^r}$$

$$l^\mu = \frac{1}{\sqrt{g}} \varepsilon^\mu \alpha_{\beta\gamma} z_1^\alpha z_2^\beta z_3^\gamma$$

$\vec{l} = z^\mu_{,\tau}(\tau, \vec{\sigma})$ TANGENT TO Σ_τ

MÖLLER ADMISSIBILITY CONDITIONS FOR THE 3+1 SPLITTINGS

$$g_{\tau\tau}(\tau, \vec{\sigma}) > 0, \quad g_{\tau r}(\tau, \vec{\sigma}) < 0, \quad \det g_{rs}(\tau, \vec{\sigma}) < 0$$

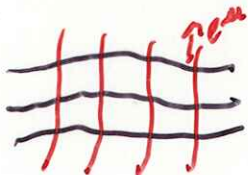
↓
NO RIGID REL. ROTATIONS

$$\begin{vmatrix} g_{\tau\tau} & g_{\tau s} \\ g_{s\tau} & g_{ss} \end{vmatrix}(\tau, \vec{\sigma}) > 0$$

$$N(\tau, \vec{\sigma}) > 0 \quad \text{NO INTERSECTION OF 3-SPACES}$$

Σ_τ SPATIAL INFINITY
SPACELIKE HYPERPLANE

2 CONGRUENCES OF ACCELERATED OBSERVER ASSOCIATED TO EVERY 3+1 SPLITTING



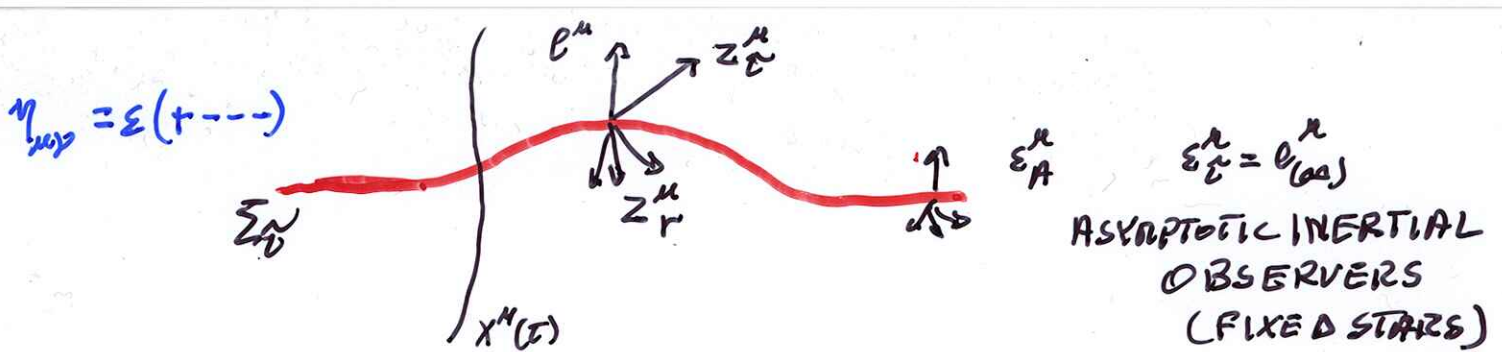
NORMAL = 4-VELOCITY SURFACE-FORMING

EULERIAN OBSERVERS



$$\frac{\partial z^\mu}{\partial \sigma^r} \frac{1}{\sqrt{g_{rr}}} = 4\text{-VELOCITY}$$

NON-SURFACE-FORMING (ROTATING DISK)



$$Z^\mu(\tau, \sigma^u)$$

$$\begin{cases} Z^\mu_t(\tau, \sigma^u) = \frac{\partial Z^\mu(\tau, \sigma^u)}{\partial \tau} = \left[(1+n) e^\mu + h^{rs} n_s Z^\mu_r \right] (\tau, \sigma^u) \\ Z^\mu_r(\tau, \sigma^u) = \frac{\partial Z^\mu(\tau, \sigma^u)}{\partial \sigma^r} \end{cases} \quad \begin{matrix} \text{LAPSE} & \text{SHIFT} \end{matrix}$$

$$e^\mu(\tau, \sigma^u) = \frac{1}{\sqrt{\gamma(\tau, \sigma^u)}} \epsilon^\mu \alpha_{\beta\gamma} [Z^\alpha_1 Z^\beta_2 Z^\gamma_3](\tau, \sigma^u)$$

$$\begin{cases} g_{\tau\tau} = Z^\mu_t \eta_{\mu\nu} Z^\nu_t = \epsilon \left[(1+n)^2 - h^{rs} n_r n_s \right] \\ g_{\tau r} = Z^\mu_t \eta_{\mu\nu} Z^\nu_r = -\epsilon n_r \\ g_{rs} = Z^\mu_r \eta_{\mu\nu} Z^\nu_s = -\epsilon h_{rs} \end{cases}$$

$$h_{rs} \text{ (ttt), INVERSE } h^{rs} \left\{ \begin{array}{l} \gamma = \det h_{rs} \\ \sqrt{|\det g_{AB}|} = (1+n) \sqrt{\gamma} \end{array} \right.$$

$$\begin{aligned} & \text{I} \quad d\tilde{t} > 0 \quad \boxed{1+n(\tau, \sigma^u) > 0} \\ & \Rightarrow \lambda_i = \gamma^{1/3} e^{2\frac{z}{4} a} \delta_{ai} R_a \\ & \epsilon g_{AB} = \begin{pmatrix} \boxed{\epsilon g_{\tau\tau}(\tau, \sigma^u) > 0} & -n_s(\tau, \sigma^u) \\ -n_r(\tau, \sigma^u) & -h_{rs} = -[R^T(\lambda_1, \lambda_2, \lambda_3) R]_{rs} \end{pmatrix} \end{aligned}$$

\(\rightarrow n r < c\)

\(\lambda_i(\tau, \sigma^u) > 0\)

\(\pi\phi\text{LLER}\)

$$h_{rs} = \gamma^{1/3} \sum_u e^{2\frac{z}{4} a} \delta_{ai} R_a V_{ru}(\theta^u) V_{su}(\theta^u) = \sum_u e_{(u)r} e_{(u)s} \rightarrow \text{COTRIADS ON } \Sigma_{\tilde{c}}$$

$$3 K_{rs} = \frac{1}{2(1+n)} (n_{rs} + n_{sr} - \partial_\tau h_{rs})$$

$\gamma, R_a, \theta^u, 1+n, n_r$ INERTIAL GAUGE VARIABLES

$\checkmark_{\text{GR}} \rightarrow$ NONLINEAR GRAVITON IN GR



$$Z^\mu(\tau, \sigma^u) = X^\mu(\tau) + \epsilon^\mu_r R^r_s (F(\sigma) \alpha_L(\tau)) \sigma^s$$

$$\pi\phi\text{LLER} \quad \frac{dF(\sigma)}{d\sigma} \neq 0, \quad 0 < F(\sigma) < \frac{1}{A\sigma}$$

$$F(\sigma) = \frac{1}{1 + \frac{\omega^2 \sigma^2}{c^2}}$$

DIFFERENTIAL ROTATIONS (REL. STARS)

$$h_{rs} \approx \delta_{rs} + [\Omega_{(s)u}^r(\tau, \sigma) + \Omega_{(u)s}^r(\tau, \sigma)] \sigma^u + \sum_w \Omega_{(u)w}^r(\tau, \sigma) \Omega_{(s)w}^v(\tau, \sigma) \sigma^u \sigma^v$$

$$\Omega = -\frac{1}{2} R^{-1} \partial_\tau R$$

$$\Omega_{(u)} = R^{-1} \partial_\tau R$$

ISOLATED SYSTEM (PARTICLES, STRINGS, FLUIDS, FIELDS) WITH LAGRANGIAN DESCRIPTION (CONFIGURATIONS WITH THE 10 POINCARÉ GENERATORS FINITE)

COUPLING TO EXTERNAL GRAVITATIONAL FIELD $g_{\mu\nu}(x)$ $\mathcal{L}(\text{matter}, g_{\mu\nu})$

$$g_{\mu\nu}(x) \mapsto g_{AB}[z^\mu(\tau, \vec{\sigma})] \Rightarrow S = \int d\tau d^3\sigma \tilde{\mathcal{L}}(\text{matter}, z(\tau, \vec{\sigma}))$$

$$\phi(\tau, \vec{\sigma}) = \hat{\phi}(z(\tau, \vec{\sigma})), x^\mu_L(\tau) = z^\mu(\tau, \vec{\eta}_L(\tau))$$

ACTION PRINCIPLE FOR MATTER AND EMBEDDING $z^\mu(\tau, \vec{\sigma})$

FLUIDS - ACTION OF BROWN (CQG 10(1993)1579)

S INVARIANT UNDER FRAME-PRESERVING DIFFEOMORPHISMS

(SPECIAL RELATIVISTIC GENERAL COVARIANCE)

4 FIRST-CLASS CONSTRAINTS (DEPARAMETRIZED ANALOGUES OF THE SUPER-HAM. AND SUPER-MOMENTUM CONSTRAINTS OF GR)

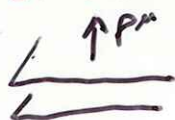
$$\left\{ \begin{aligned} \mathcal{H}^\mu(\tau, \vec{\sigma}) &= S^\mu(\tau, \vec{\sigma}) - \rho^\mu(\tau, \vec{\sigma}) T^{\hat{t}\hat{t}}(\tau, \vec{\sigma}) - z^\mu_\tau(\tau, \vec{\sigma}) T^{\hat{t}\hat{r}}(\tau, \vec{\sigma}) \approx 0 \\ &\quad \begin{array}{l} \text{4-MOMENTUM} \\ \text{OF } z^\mu(\tau, \vec{\sigma}) \end{array} \quad \begin{array}{l} \text{MATTER} \\ \text{ENERGY} \\ \text{DENSITY} \end{array} \quad \begin{array}{l} \text{MATTER} \\ \text{3-MOMENTUM} \\ \text{DENSITY} \end{array} \\ \{ \mathcal{H}^\mu(\tau, \vec{\sigma}_1), \mathcal{H}^\mu(\tau, \vec{\sigma}_2) \} \approx 0 \end{aligned} \right.$$

$\Rightarrow z^\mu(\tau, \vec{\sigma})$ GAUGE VARIABLES $\Rightarrow$ GAUGE EQUIVALENCE OF ADMISSIBLE NON-INERTIAL FRAMES (I.E. OF CLOCK SYNCHRONIZATION CONVENTIONS AND OF CHOICE OF 3-COORDINATES ON THE INSTANTANEOUS 3-SPACES)

REST-FRAME INSTANT FORM OF DYNAMICS

INERTIAL 3+1 SPLITTING CENTERED ON AN INERTIAL OBSERVER (DESCRIBING THE COLLECTIVE VARIABLES OF THE ISOLATED SYSTEM)

WHOSE INSTANTANEOUS 3-SPACES ARE ORTHOGONAL TO THE 4-MOMENTUM OF THE ISOLATED SYSTEM (INTRINSICALLY DEFINED BY ITS CONFIGURATIONS)



$$z^\mu(\tau, \vec{\sigma}) = x^\mu(\tau) + \varepsilon^\mu_\tau(P) \sigma^\tau$$

NON-INERTIAL REST FRAMES



ASYMPTOTIC INERTIAL OBSERVERS = FIXED STARS

ISOLATED SYSTEMS IN THE REST-FRAME (INSTANT FORM)

STANDARD DESCRIPTION - POINCARÉ GENERATORS $P^\mu, J^{\mu\nu}$ FINITE

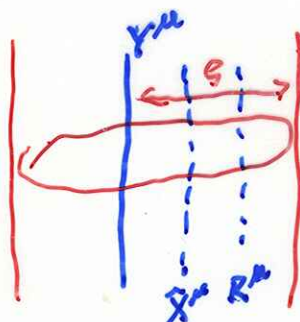
NON-LOCAL QUANTITIES - THEY KNOW THE WHOLE INSTANTANEOUS 3-SPACE $P^\mu = \int_{\Sigma_t} d^3x \dots$
 LORENTZ BOOSTS! INTERACTION-DEPENDENT $[P^\mu, X^\nu] = \delta^\mu_\nu P^0$
 $\neq$ GALILEI BOOSTS

$P^\mu, J^{\mu\nu} \rightarrow$ 3 NOTIONS OF COLLECTIVE VARIABLES

$\begin{cases} \hat{X}^\mu = (\hat{X}^0; \vec{\hat{X}}) & \text{NON-COVARIANT CANONICAL CENTER OF MASS (CENTER OF SPIN)} \\ \Upsilon^\mu = (\hat{X}^0; \vec{\Upsilon}) & \text{COVARIANT NON-CANONICAL FOKKER-PRYCE CENTER OF INERTIA} \\ R^\mu = (\hat{X}^0; \vec{R} = -\frac{\vec{K}}{\pi c}) & \text{NON-COVARIANT NON-CANONICAL MÖLLER CENTER OF ENERGY} \end{cases}$

$\rightarrow$ NEWTON CENTER OF MASS
 $\hookrightarrow \Delta$

MÖLLER
WORLD-TUBE
OF NON-COVARIANCE
(LORENTZ SIGNATURE)



MÖLLER RADIUS

$$S \approx \frac{\sqrt{-W^2}}{P^2 c} = \frac{|\vec{S}|}{\pi c}$$

$S \sim$ COMPTON WAVELENGTH
OF ISOLATED SYSTEM
REMNANT OF ENERGY CONDITIONS
OF GENERAL RELATIVITY

REST-FRAME INSTANT FORM

WIGNER INSTANTANEOUS 3-SPACES



$\tau = ct$ REST TIME
(LORENTZ SCALAR)

Y^μ ORIGIN OF 3-COORD.

$$Z^\mu(\tau, \vec{h}) = Y^\mu(\tau) + \varepsilon^\mu_\tau(\vec{h}) \tau$$

$$P^\mu = \pi c (\sqrt{1 + \vec{h}^2}; \vec{h}) = L^\mu_\nu (P, \vec{P}) \bar{P}^\nu$$

$$\bar{P}^\mu = \pi c (1; \vec{0}), \pi c = \sqrt{P^2}$$

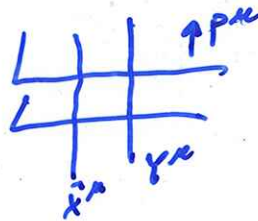
$\downarrow$ WIGNER
BOOST
FOR TIME-LIKE ORBITS

POSITIVE ENERGY PARTICLES $X^\mu_i(\tau) = Z^\mu(\tau, \vec{\eta}_i(\tau)) = Y^\mu(\tau) + \varepsilon^\mu_\tau(\vec{\eta}_i) \eta^\tau_i(\tau)$

$\hat{X}^\mu(\tau), \Upsilon^\mu(\tau), R^\mu(\tau)$ KNOWN FUNCTIONS OF $\tau, \pi, \vec{S}$, and of the FROZEN
JACOBI DATA $\vec{Z}, \vec{h}$ OF THE EXTERNAL CENTER OF MASS

$\vec{X}_{NW} = \frac{\vec{Z}}{\pi c}$ NEWTON-WIGNER POSITION

ISOLATED SYSTEM



- 1) NON-COVARIANT DECOUPLED EXTERNAL CENTER OF MASS
(DESCRIBED BY THE FROZEN JACOBI DATA) CARRYING A POLE-DIPOLE
STRUCTURE : M INTERNAL MASS, $\vec{S}$ REST SPIN
(UNIVERSAL BREAKING OF LORENTZ COVARIANCE BUT DECOUPLED)

EXTERNAL
POINCARÉ GENERATORS

$$P^0 = \sqrt{1+h^2} \pi c \sqrt{1+h^2}$$

$$\vec{P} = \pi c \vec{h}$$

$$\vec{J} = \vec{Z} \times \vec{h} + \vec{S}$$

$$\vec{K} = \sqrt{1+h^2} \vec{Z} + \frac{\vec{S} \times \vec{h}}{1+\sqrt{1+h^2}}$$

- 2) INSIDE THE INSTANTANEOUS WIGNER 3-SPACE THERE ARE THE WIGNER-COVAR
VARIABLES OF THE ISOLATED SYSTEM (INTERNAL SPACE)
AND A UNFAITHFUL INTERNAL REALIZATION OF THE POINCARÉ GROUP

$\mathcal{L} = cT$ LORENTZ-SCALAR REST TIME

$$\left\{ \begin{array}{l} \pi c = c \int d^3\sigma T^{00}(\mathcal{L}, \vec{\sigma}) \\ \vec{P} \approx 0 \quad \text{REST-FRAME CONDITION} \\ \vec{J} \approx \vec{S} \\ \vec{K} \approx 0 \quad \text{(INTERACTION-DEPENDENT) ELIMINATION OF THE INTERNAL 3-CENTER OF MASS} \end{array} \right. \quad \left\{ \begin{array}{l} \vec{P} = \int d^3\sigma T^{0i}(\mathcal{L}, \vec{\sigma}) \\ \vec{S} = \frac{1}{2} \varepsilon^{\mu\nu} \int d^3\sigma \sigma^\mu T^{0\nu}(\mathcal{L}, \vec{\sigma}) \\ K^i = - \int d^3\sigma \sigma^i T^{00}(\mathcal{L}, \vec{\sigma}) \end{array} \right.$$

POSITIVE-ENERGY SCALAR FREE PARTICLES $\vec{\eta}_L(\mathcal{L}), \vec{K}_L(\mathcal{L})$

$$\pi c = \sum_L \sqrt{m_L^2 c^2 + \vec{K}_L^2}, \quad \vec{P} = \sum_L \vec{K}_L \approx 0, \quad \vec{S} = \sum_L \vec{\eta}_L \times \vec{K}_L, \quad \vec{K} = - \sum_L \vec{\eta}_L \sqrt{m_L^2 c^2 + \vec{K}_L^2} \approx 0$$

$\Rightarrow \vec{\eta}_L, \vec{K}_L, L=1, \dots, N$ FUNCTIONS ONLY OF RELATIVE VARIABLES $\vec{S}_a, \vec{T}_a, a=1, \dots, N$

2-BODY CASE - ORBIT RECONSTRUCTION

$$X_L^\mu(\mathcal{L}) = Y^\mu(\mathcal{L}) + (-)^{L+1} \varepsilon_L^\mu(\vec{h}) \frac{\sqrt{m_L^2 c^2 + \vec{\eta}_L^2}}{\sum_K \sqrt{m_K^2 c^2 + \vec{\eta}_K^2}} \vec{S}_{12}(\mathcal{L}) \xrightarrow{\mathcal{L} \rightarrow \infty} \left\{ \vec{X}_L(\omega) + (-)^{L+1} \frac{m_{L+1}}{m_1 + m_2} \vec{F}_L(\omega) = \vec{X}_{(N)L} \right.$$

COVARIANT WORLD LINES - DERIVED NON-CANONICAL QUANTITIES

$\{X_L^\mu, X_S^\nu\} \neq 0$ PREDICTIVE COORDINATES
VIOLATION OF MICROCAUSALITY

DERIVED P_L^μ SATISFY $P_L^2 = m_L^2 c^2$ ALSO IN PRESENCE OF INTERACTIONS

THE SIMPLEST INTERACTING 2-BODY SYSTEM

$$\left\{ \begin{array}{l} M = \sqrt{m_1^2 + \vec{K}_1^2 + \phi(\vec{S}^2)} + \sqrt{m_2^2 + \vec{K}_2^2 + \phi(\vec{S}^2)} \quad \vec{S} = \vec{\eta}_1 - \vec{\eta}_2 \\ \vec{P} = \vec{K}_1 + \vec{K}_2 \approx 0 \\ \vec{J} = \vec{S} = \vec{\eta}_1 \times \vec{K}_1 + \vec{\eta}_2 \times \vec{K}_2 \\ \vec{K} = -\vec{\eta}_1 \sqrt{m_1^2 + \vec{K}_1^2 + \phi(\vec{S}^2)} - \vec{\eta}_2 \sqrt{m_2^2 + \vec{K}_2^2 + \phi(\vec{S}^2)} \approx 0 \end{array} \right.$$

THE INTERNAL POINCARÉ ALGEBRA IS SATISFIED

$$\Pi \approx M = \sqrt{m_1^2 + \vec{\Pi}^2 + \phi(\vec{S}^2)} + \sqrt{m_2^2 + \vec{\Pi}^2 + \phi(\vec{S}^2)} \quad \text{HAMILTONIAN FOR THE RELATIVE MOTION}$$

$$m_0 = \sqrt{m_1^2 + \vec{\Pi}^2} + \sqrt{m_2^2 + \vec{\Pi}^2}$$

$$\vec{K} \approx 0$$

$$\Rightarrow \vec{q} \approx \vec{q}(\vec{S}, \vec{\Pi}) = \frac{m_1^2 - m_2^2}{2} \left(\frac{1}{m_0^2} - \frac{1}{M^2} \right) \vec{S}$$

$$\Rightarrow \vec{\eta}_L \approx \frac{(-)^{L+1}}{2} \left(1 - \frac{m_1^2 - m_2^2}{m^2} \right) \vec{S} \xrightarrow{c \rightarrow \infty} (-)^{L+1} \frac{m_i}{m_1 + m_2} \vec{S}$$

IN THE PAPER THERE ARE THE ORBITS FOR COULOMB-LIKE POTENTIALS

$$\phi(\vec{S}^2) = -2\mu \frac{e^2}{\sqrt{\vec{S}^2}}$$

N.B. FOR $\Pi = \sqrt{m_1^2 + \vec{K}_1^2} + \sqrt{m_2^2 + \vec{K}_2^2} + \frac{e^2}{4\pi |\vec{\eta}_1 - \vec{\eta}_2|}$ (REAL COULOMB POTENTIAL)

IT IS NOT KNOWN THE FORM OF THE INTERNAL BOOST $\vec{K}$

(IT IS NOT KNOWN THE LAGRANGIAN, FROM WHICH TO GET THE ENERGY-MOMENTUM TENSOR.)

RELATIVISTIC PSEUDOCCLASSICAL ATOMIC PHYSICS

GRASSMANN-VALUED ELECTRIC CHARGES

$$q_L^2 = 0, q_L q_j \neq 0 \quad L \neq j$$

- ULTRAVIOLET REGULARIZATION OF SELF-ENERGIES ~~3~~
- INFRARED REGULARIZATION ~~3~~
- POTENTIAL IN CAUCHY PROBLEM AND IN THE RADIATION GAUGE FOR 1-PHOTON EXCHANGE ~~3~~

⇒ DARWIN AND SALPETER POTENTIALS ADDED TO THE COULOMB ONE
(TILL NOW FROM BETHE-SALPETER EQ. FOR RELATIVISTIC BOUND STATES)

WAY OUT OF HAAG THEOREM WITH A CLASSICAL CANONICAL TRANSF.

$$\begin{aligned} \vec{\eta}_L, \vec{k}_L + \text{COULOMB} \\ \vec{A}_L, \vec{E}_L \text{ (INTERPOLATING FIELDS)} \end{aligned} \rightarrow \begin{aligned} \vec{\tilde{\eta}}_L, \vec{\tilde{k}}_L \text{ (COULOMB DRESSED)} + \\ + \text{COULOMB} + \text{DARWIN} \\ \vec{A}_L^{\text{RAD}}, \vec{E}_L^{\text{RAD}} \text{ (IN OR OUT FREE FIELDS)} \end{aligned}$$

ONLY IN THE GLOBAL REST-FRAME

MISSING - RELATIVISTIC STATISTICAL MECHANICS WITH LONG-RANGE FORCES

FIND MICROCANONICAL PARTITION FUNCTION IN REST FRAME

LIKE IT HAS BEEN DONE FOR NEWTONIAN GRAVITY

D.H.E. GROSSE ET AL. PRL 89 (2002) 031101

A) MASSLESS PARTICLES IN THE REST-FRAME INSTANT FORM

B) 2-LEVEL ATOM (JAYNES-CUMMING MODEL) - PASSIVE RELATIVISTIC PARTICLE
WITH STRUCTURE $(\vec{\eta}, \vec{k})$ POINT, $(\vec{S}, \vec{\pi})$ DIPOLE

$$H \approx \sqrt{m^2 c^2 + 2\omega \theta^\dagger \theta + \vec{k}^2} + A(\theta^\dagger \alpha + \alpha^\dagger \theta) \vec{S} \cdot \vec{\pi}_\perp (\vec{E}_1 \vec{\eta}(\omega)) + \frac{1}{2c} \int d\vec{r} (\vec{\pi}_\perp^2 + \vec{B}^2)(\vec{r}, \vec{r})$$

$$\vec{\pi}^2 \approx 2\omega \theta^\dagger \theta$$

$$\alpha^\dagger \alpha + \theta^\dagger \theta \approx 0$$

$\theta^\dagger, \theta, \alpha^\dagger, \alpha$ GRASSMANN

- QUANTIZATION - KG FIELD + 2-LEVEL STRUCTURE
- TRY TO SIMULATE ABCD BORDE METHOD FOR ATOM INTERFEROMETER
- COUPLING TO GRAVITY (BORDE', KASEVICH et al. ...)

$$\begin{aligned}
M_C &= \sum_i \sqrt{m_i^2 c^2 + (\vec{K}_i - \frac{q_i}{c} \vec{A}_\perp(\vec{r}_i, \vec{\eta}_i))}^2 + \sum_{i \neq j} \frac{q_i q_j}{4\pi |\vec{\eta}_i - \vec{\eta}_j|} \\
&\quad + \frac{1}{2} \int d^3\sigma [\vec{\Pi}_\perp^2 + \vec{B}^2](\vec{r}, \vec{\sigma}) \\
\vec{P} &= \sum_i \vec{K}_i + \frac{1}{c} \int d^3\sigma (\vec{\Pi}_\perp \times \vec{B})(\vec{r}, \vec{\sigma}) \approx 0 \\
\vec{J} = \vec{S} &= \sum_i \vec{\eta}_i \times \vec{K}_i + \frac{1}{c} \int d^3\sigma [\vec{\sigma} \times (\vec{\Pi}_\perp \times \vec{B})](\vec{r}, \vec{\sigma}) \\
\vec{K} &= - \sum_i \vec{\eta}_i \sqrt{m_i^2 c^2 + (\vec{K}_i - \frac{q_i}{c} \vec{A}_\perp(\vec{r}_i, \vec{\eta}_i))}^2 + \\
&\quad + \frac{1}{c} \sum_{i \neq j} q_i q_j \left(\frac{1}{A_{ij}} \frac{\partial C(\vec{\eta}_i - \vec{\eta}_j)}{\partial \vec{\eta}_j} - \vec{\eta}_j C(\vec{\eta}_i - \vec{\eta}_j) \right) \\
&\quad + \sum_i \frac{q_i}{c} \int d^3\sigma \vec{\Pi}_\perp(\vec{r}, \vec{\sigma}) C(\vec{\sigma} - \vec{\eta}_i) - \\
&\quad - \frac{1}{2c} \int d^3\sigma \vec{\sigma} (\vec{\Pi}_\perp^2 + \vec{B}^2)(\vec{r}, \vec{\sigma}) \approx 0
\end{aligned}$$

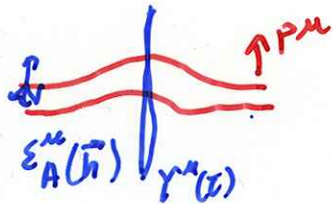
$q_i^2 = 0$
 $q_i q_j \neq 0$
 $\vec{\Pi}_\perp = \frac{1}{c} \vec{E}_\perp = \frac{\partial \vec{A}_\perp}{\partial t}$
 $C(\vec{\sigma}) = \frac{1}{4\pi |\vec{\sigma}|}$

CRATER - LL - ANN. PHYS. 289(2001)37
 hep-th/0001046
 ALBA, CRATER, LL
 INT. J. MOD. PHYS. A16(2001)3365
 hep-th/0103109

2 PARTICLES WITH DARWIN POTENTIAL

$$\begin{aligned}
M_C &= \sqrt{m_1^2 c^2 + \vec{K}_1^2} + \sqrt{m_2^2 c^2 + \vec{K}_2^2} + \frac{q_1 q_2}{4\pi |\vec{\eta}_1 - \vec{\eta}_2|} + V_{\text{DARWIN}} \\
\vec{P} &= \vec{K}_1 + \vec{K}_2 \approx 0 \\
\vec{J} = \vec{S} &= \vec{\eta}_1 \times \vec{K}_1 + \vec{\eta}_2 \times \vec{K}_2 \\
\vec{K} &= - \sum_i \vec{\eta}_i \sqrt{m_i^2 c^2 + \vec{K}_i^2} - \\
&\quad - \frac{q_1 q_2}{2c} \left[\vec{\eta}_1 \frac{\vec{K}_1}{\sqrt{m_1^2 c^2 + \vec{K}_1^2}} \cdot \left(\frac{1}{2} \frac{\partial K_{12}}{\partial \vec{S}} - 2 \vec{A}_{1S2}(\vec{K}_1, \vec{S}) \right) + \vec{\eta}_2 \frac{\vec{K}_2}{\sqrt{m_2^2 c^2 + \vec{K}_2^2}} \cdot \left(\frac{1}{2} \frac{\partial K_{12}}{\partial \vec{S}} - 2 \vec{A}_{1S2}(\vec{K}_2, \vec{S}) \right) \right] \\
&\quad - \frac{q_1 q_2}{2c} \left(\sqrt{m_1^2 c^2 + \vec{K}_1^2} \frac{\partial}{\partial \vec{K}_1} + \sqrt{m_2^2 c^2 + \vec{K}_2^2} \frac{\partial}{\partial \vec{K}_2} \right) K_{12}(\vec{K}_1, \vec{K}_2, \vec{S}) - \\
&\quad - \frac{q_1 q_2}{4\pi c} \int d^3\sigma \left(\frac{\vec{\Pi}_{1S1}(\vec{\sigma} - \vec{\eta}_1, \vec{K}_1)}{|\vec{\sigma} - \vec{\eta}_1|} + \frac{\vec{\Pi}_{1S2}(\vec{\sigma} - \vec{\eta}_2, \vec{K}_2)}{|\vec{\sigma} - \vec{\eta}_2|} \right) - \\
&\quad - \frac{q_1 q_2}{c} \int d^3\sigma \left[\vec{\Pi}_{1S1}(\vec{\sigma} - \vec{\eta}_1, \vec{K}_1) \cdot \vec{\Pi}_{1S2}(\vec{\sigma} - \vec{\eta}_2, \vec{K}_2) + \vec{B}_{S1}(\vec{\sigma} - \vec{\eta}_1, \vec{K}_1) \cdot \vec{B}_{S2}(\vec{\sigma} - \vec{\eta}_2, \vec{K}_2) \right] \approx 0
\end{aligned}$$

NON-INERTIAL REST-FRAME INSTANT FORM



$$Z^\mu(\tau, \sigma^\mu) = \gamma^\mu(x) + u^\mu(\vec{h}) g(\tau, \sigma^\mu) + \epsilon_T^\mu(\vec{h}) [\sigma^\tau + g^\tau(\tau, \sigma^\mu)]$$

$$\xrightarrow{\sigma \rightarrow \infty} \gamma^\mu(x) + \epsilon_T^\mu(\vec{h}) \sigma^\tau \quad \text{ASYMPTOTIC WIGNER 3-SPACES}$$

↓ FOKKER-PRICE

$$H_D = M_0 + \int d^3\sigma (\partial_\sigma g Z_\tau + \partial_\tau g^\tau Z_\tau)(\tau, \vec{\sigma}) + (\text{em constraints})$$

INVARIANT
BASE

INERTIAL POTENTIALS

$$\int d^3\sigma [(1+n) T_{\tau\tau} - n^\tau T_{\tau\tau}]$$

ASYMPTOTIC INTERNAL POINCARÉ GENERATORS

$$P_0, \quad \vec{S}^\tau = \dots$$

$$\left\{ \begin{array}{l} P^\mu = \int d^3\sigma [A^\mu T_{\tau\tau} - (\delta_\tau^\mu + \partial_\tau g^\mu) h^{\tau s} T_{\tau s}] (\tau, \sigma^\mu) \approx 0 \\ K^\mu = \int d^3\sigma [(\sigma^\mu + g^\mu) Z^\tau + g Z^\tau] (\tau, \sigma^\mu) \approx 0 \end{array} \right.$$

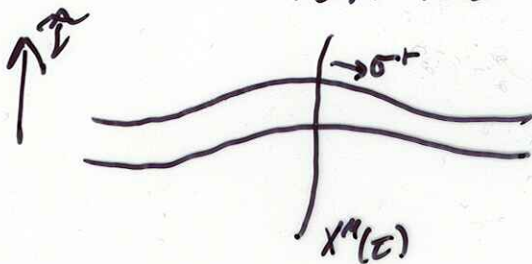
REST FRAME
CONDITION

ELIMINATION
OF INTERNAL COM

MISSING - FIND RIEMANN NORMAL COORDINATES

RADAR

FOR THE 3+1 SPLITTING



NO COORDINATE SINGULARITIES (LIKE IN 1+3):

SEE FERMI COORDINATES AND ROTATING DISK)

EQUATIONS OF MOTION OF CHARGED PARTICLES IN NON-INERTIAL FRAMES IN THE RADIATION GAUGE

$$m_i c \frac{d}{d\tau} \frac{\dot{\eta}_i^\mu(\tau) + n^\mu(\tau, \vec{\eta}_i(\tau))}{\sqrt{[1+n(\tau, \vec{\eta}_i(\tau))]^2 - h_{\mu\nu}(\tau, \vec{\eta}_i(\tau)) [\dot{\eta}_i^\mu(\tau) + n^\mu(\tau, \vec{\eta}_i(\tau))] [\dot{\eta}_i^\nu(\tau) + n^\nu(\tau, \vec{\eta}_i(\tau))]} =$$

$$= \frac{q_i}{c} h^{\mu\nu}(\tau, \vec{\eta}_i(\tau)) \left[E_{\perp\mu}(\tau, \vec{\eta}_i(\tau)) + \varepsilon_{\mu\nu\lambda} \dot{\eta}_i^\lambda(\tau) B_\nu(\tau, \vec{\eta}_i(\tau)) \right] + \left[-\frac{1}{c} \frac{\partial}{\partial \eta_i^\mu} W(\vec{\eta}_1(\tau), \dots, \vec{\eta}_N(\tau)) + F_{L\mu}(\tau) \right]$$

NON-INERTIAL COULOMB POTENTIAL
LORENTZ FORCE
INERTIAL FORCES

in
hrs and
 γ
are hidden
the
GRAVITON
effects
in GR

$$\vec{F}_{LP} = \frac{m_i c}{\sqrt{(1+n)^2 - h_{\mu\nu}(\dot{\eta}_i^\mu + n^\mu)(\dot{\eta}_i^\nu + n^\nu)}} \left\{ \frac{1}{1+n} \left[\frac{\partial h_{\mu\sigma}}{\partial \eta_i^\mu} (\dot{\eta}_i^\sigma + n^\sigma)(\dot{\eta}_i^\mu + n^\mu) - \frac{\partial n}{\partial \eta_i^\mu} + \frac{\partial n^\sigma}{\partial \eta_i^\mu} h_{\sigma\mu} (\dot{\eta}_i^\mu + n^\mu) \right] - \left(\frac{\partial h_{\mu\sigma}}{\partial \tau} + \dot{\eta}_i^\nu \frac{\partial h_{\mu\sigma}}{\partial \eta_i^\nu} \right) (\dot{\eta}_i^\mu + n^\mu) \right\}$$

$$W = \frac{1}{24\pi} \sum_{i \neq j} \sum_{\mu} \left\{ \frac{h_{\mu\sigma} (1+n)}{2\sqrt{\gamma}} \frac{q_i q_j}{|\vec{\eta}_i(\tau) - \vec{\eta}_j(\tau)|} + \int d^3\sigma \sum_{\mu} \left\{ \frac{q_i q_j}{8\pi |\vec{\sigma} - \vec{\eta}_i(\tau)|} \left[\frac{h_{\mu\sigma} (1+n)}{2\sqrt{\gamma}} E^{\mu\sigma}(\vec{\sigma} - \vec{\eta}_i(\tau)) - \frac{1}{4\pi} \frac{(\vec{\sigma} - \vec{\eta}_i(\tau))^2}{|\vec{\sigma} - \vec{\eta}_i(\tau)|^3} \frac{\partial}{\partial \sigma^3} \left[\frac{h_{\mu\sigma} (1+n)}{2\sqrt{\gamma}} \right] (\tau, \vec{\sigma}) + i \leftrightarrow j \right\} + \right. \\ \left. + \sum_j \int d^3\sigma \frac{q_j}{4\pi |\vec{\sigma} - \vec{\eta}_j(\tau)|} \frac{\partial}{\partial \sigma^3} \left[\frac{h_{\mu\sigma} (1+n)}{2\sqrt{\gamma}} \right] \pi_\mu^\dagger + n^\mu F_{\mu\sigma} \right\} (\tau, \vec{\sigma})$$

$$E^{\mu\sigma}(\vec{\sigma}) = \frac{1}{|\vec{\sigma}|^3} \left(\delta^{\mu\sigma} - 3 \frac{\sigma^\mu \sigma^\sigma}{\sigma^2} \right) = \frac{4\pi}{3} \delta^{\mu\sigma} \delta^3(\vec{\sigma}) + \vec{E}^{\mu\sigma}(\vec{\sigma})$$

NON-RELATIVISTIC LIMIT (WITHOUT EM FIELDS) $R(\tau, \sigma) \rightarrow R(\tau) + O(c^{-2})$

$$m_i \frac{d^2 \vec{\eta}_i(\tau)}{d\tau^2} = \frac{q_i}{c} [\vec{E} + \frac{d\vec{\eta}_i(\tau)}{d\tau} \times \vec{B}] - \frac{\partial}{\partial \eta_i^\mu} W + \vec{F}_{L(\text{lin})}$$

$$\vec{F}_{L(\text{lin})} = m_i \left[\underbrace{-\frac{d\vec{v}(\tau)}{d\tau} \cdot \vec{\eta}_i(\tau) \times \vec{v}(\tau)}_{\text{LINEAR ACCELERATION}} + \underbrace{\vec{\eta}_i(\tau) \times \frac{d\vec{\omega}(\tau)}{d\tau}}_{\text{JACOBI}} + 2 \underbrace{\frac{d\vec{\eta}_i(\tau)}{d\tau} \times \vec{\omega}(\tau)}_{\text{CORIOLIS}} + \underbrace{\vec{\omega}(\tau) \times [\vec{\eta}_i(\tau) \times \vec{\omega}(\tau)]}_{\text{CENTRIFUGAL}} \right]$$

$$W = - \sum_{i \neq j} \frac{q_i q_j}{4\pi |\vec{\eta}_i(\tau) - \vec{\eta}_j(\tau)|} + \frac{1}{c} \int d^3\sigma \sum_i \frac{q_i}{4\pi |\vec{\sigma} - \vec{\eta}_i(\tau)|} \left\{ -\vec{v}(\tau) \cdot \vec{\sigma} \times \vec{B}(\tau, \vec{\sigma}) - \vec{\sigma} \cdot [\vec{\sigma} \cdot \vec{\omega}(\tau) \vec{\omega}(\tau) - \vec{\omega}(\tau) \cdot \vec{B}(\tau, \vec{\sigma})] (\tau, \vec{\sigma}) \right\}$$

QUANTUM MECHANICS OF SCALAR AND SPINNING PARTICLES IN NON-INERTIAL FRAMES

ALBA, LL - INT. J. MOD. PHYS. 21, 2781 (2006) [hep-th/0602060]

NON-REL. CASE ALBA - INT. J. MOD. PHYS 21, 3917 (2006) [hep-th/0604060]

MULTITEMPORAL QUANTIZATION - QUANTIZE THE PARTICLES

BUT NOT THE INERTIAL APPEARANCES [$z^{\mu}(x, 0^{\mu})$ C-NUMBER]

ATOMIC SPECTRA - OK

←————→
OPEN PROBLEM

QUANTIZATION OF KG FIELD IN NON-INERTIAL FRAMES

(STANDARD MODEL OF PARTICLES IN NON-INERTIAL FRAMES)

TORRE-VARADARAJAN - CLAS. Q. GRAV. 16, 2651 (1999)

ARAGEORGIS, EARTMAN, RUETSCHKE - STUDIES HIST. PHIL. MODERN PHYSICS
33, 151 (2002)

FIND HOW TO AVOID THE TORRE-VARADARAJAN NO-GO THEOREM (NO UNITARY
EVOLUTION IN THE TOPOLOGA-SCHWINGER FORMALISM)

1) USE INSTANTANEOUS 3-SPACES ADMITTING FOURIER TRANSFORM
→ FOCK SPACE (INSTEAD OF DETECTORS REPLACING PARTICLES LIKE IN
QFT IN CURVED SPACETIMES)

OR 2) HAMILTONIAN EVOLUTION DEFINED ONLY IN A 3+1 SPLITTING

a) 3+1 SPLITTINGS WITH FT ON Σ_t  BOGOLIOBOV TRANSF.

b) FIND 3+1 SPLITTINGS WITH HILBERT-SCHMIDT BOGOLIOBOV TRANSF.

$$a'_k = \sum_e \bar{z}_{ke} [a_{ke} a_e + \beta_{ke} a_e^\dagger] \quad \sum_{ke} |\beta_{ke}|^2 < \infty$$

c) UNITARY EQUIVALENCE OF HS 3+1 SPLITTINGS



NON-RELATIVISTIC QUANTUM MECHANICS

2-BODY PROBLEM

$$\vec{X}_{(n)i}, \vec{P}_{(n)i} \rightarrow \vec{X}_{(n)}, \vec{P}_{(n)}, \vec{r}_{(n)}, \vec{q}_{(n)}$$

$$\begin{aligned} m_1, m_2 \\ m = m_1 + m_2 \\ \mu = m_1 m_2 / m \end{aligned}$$

GALILEI GENERATORS

$$\begin{cases} E = H = H_{com} + H_{rel}, & H_{com} = \frac{\vec{P}_{(n)}^2}{2m}, & H_{rel} = \frac{\vec{q}_{(n)}^2}{2\mu} + V(\vec{r}_{(n)}) \\ \vec{P} = \vec{P}_{(n)} = \sum_i \vec{P}_{(n)i} \\ \vec{L} = \sum_i \vec{X}_{(n)i} \times \vec{P}_{(n)i} = \vec{X}_{(n)} \times \vec{P}_{(n)} + \vec{S}_{(n)}, & \vec{S}_{(n)} = \vec{r}_{(n)} \times \vec{q}_{(n)} \\ \vec{K} = -m \vec{X}_{(n)} + \vec{P}_{(n)} t \end{cases}$$

REST FRAME CENTERED IN $\vec{X}_{(n)}$

EXTERNAL C.O.P. $\vec{X}_{(n)}, \vec{P}_{(n)}$ DECOUPLED WITH POLE-DIPOLE STRUCTURE: $H_{rel}, \vec{S}_{(n)}$

→ EXTERNAL GALILEI GROUP

+ INTERNAL SPACE WITH $\vec{\eta}_{(n)i} = \vec{X}_{(n)i} |_{\vec{X}_{(n)}=0}, \vec{K}_{(n)i} = \vec{P}_{(n)i} |_{\vec{P}_{(n)}=0}$ $\vec{\eta}_{(n)1} - \vec{\eta}_{(n)2} = \vec{r}_{(n)}$

and unfaithful Galilei realisation

$$H_{rel}, \vec{P} \approx 0, \vec{S}_{(n)}, \vec{K} \approx 0$$

$$\begin{cases} \vec{\eta}_i \rightarrow \vec{\eta}_{(n)i} \\ \vec{K}_i \rightarrow \vec{K}_{(n)i} \end{cases}$$

$H = H_1 \otimes H_2$	$=$	$H_{com} \otimes H_{rel}$
COMPOSITE SYSTEM SEPARABILITY (ZEROth POSTULATE of QM, ZEN)	UNITARY TRANSF.	$\Psi(t, \vec{X}_{(n)}) \Phi(t, \vec{r}_{(n)})$
$\Psi(t, \vec{X}_{(n)1}) \Psi(t, \vec{X}_{(n)2})$	DICTATED BY SEPARATION OF VARIABLES IN SCHR. EQ	

$$\begin{cases} i \frac{\partial}{\partial t} \Psi(t, \vec{X}_{(n)}, \vec{r}_{(n)}) = (\hat{H}_{com} + \hat{H}_{rel}) \Psi(t, \vec{X}_{(n)}, \vec{r}_{(n)}) \\ i \frac{\partial}{\partial t} (e^{-i \frac{\vec{P}_{(n)}^2}{2m} t} \Psi_P(\vec{X}_{(n)})) = \hat{H}_{com} (e^{-i \frac{\vec{P}_{(n)}^2}{2m} t} \Psi_P(\vec{X}_{(n)})) \\ i \frac{\partial}{\partial t} (e^{-i \varepsilon_n t} \Phi_{nem}(\vec{r}_{(n)})) = \hat{H}_{rel} (e^{-i \varepsilon_n t} \Phi_{nem}(\vec{r}_{(n)})) \end{cases} \quad \Psi_P(\vec{X}_{(n)}) \sim e^{i \vec{X}_{(n)} \cdot \vec{P}_{(n)}}$$

HAMILTON-JACOBI FOR CENTER OF MASS

FROZEN JACOBI DATA $\vec{X}_{(n)}(0), \vec{P}_{(n)}$ FOR COO
INTERNAL SPACE $\vec{r}_{(n)}, \vec{q}_{(n)}$

$$\begin{cases} \hat{P} \Psi_P(\vec{P}_{(n)} | t, \vec{r}_{(n)}) \\ i \frac{\partial}{\partial t} \Psi_P(\vec{P}_{(n)} | t, \vec{r}_{(n)}) = \hat{H}_{rel} \Psi_P(\vec{P}_{(n)} | t, \vec{r}_{(n)}) \end{cases}$$

$$H_{com} \xrightarrow{e^{i \vec{S}}} 0$$

$$S = -H_{com} t$$

QUANTIZATION OF THE REST-FRAME INSTANT FORM

$$H = H_{com} \otimes H_{rel} \neq H_1 \otimes H_2$$

Tensor product of 2
Klein-Gordon

NO CORRELATION BETWEEN x_1^0 and x_2^0

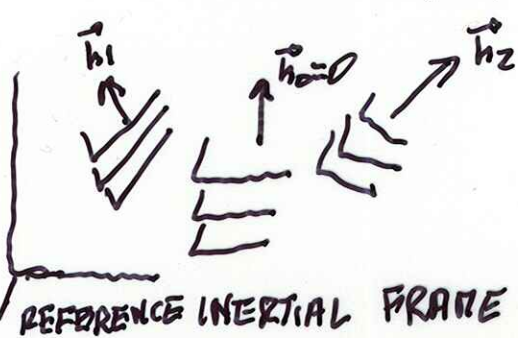
$$\frac{\partial}{\partial x_1^0} \langle \psi_{a1}, \psi_{a2} \rangle = 0$$

IN PAST OF THE STATES PARTICLE i
IN THE ABSOLUTE FUTURE OF PARTICLE j

IRRELEVANT FOR MATRIX $x_i^0 \rightarrow \pm \infty$

ALL THE IMPLICATIONS OF POINCARÉ GROUP
TAKEN INTO ACCOUNT

$$P^\mu = m c (\sqrt{1 + \vec{h}^2}; \vec{h})$$



GLOBAL (NOT LOCALLY DEFINED)

CLOSED UNIVERSE $\vec{h} \Rightarrow$ CONSTANT OF MOTION

NON-COVARIANT

1) H_{com}

FROZEN JACOBI DATA

$$\vec{z}, \vec{h}$$

OF EXTERNAL CENTER OF MASS $[\vec{z}, \vec{h}] = i \hbar \delta^3$

PREFERRED $\vec{h}$ -BASIS FOR DEFINITION OF INSTANTANEOUS WIGNER 3-SPACE

$$\vec{h} \psi_{\vec{R}}(\vec{h}) = \vec{R} \psi_{\vec{R}}(\vec{h}) \quad \psi_{\vec{R}}(\vec{h}) = \delta^3(\vec{h} - \vec{R})$$

$$\vec{x}_{NW} = \frac{\vec{z}}{m c} \quad \text{DEPENDS ON MASS EIGENVALUE}$$

$$\langle \psi_1 | \psi_2 \rangle = \int \frac{d^3 h}{1 + \vec{h}^2} \psi_1^*(\vec{h}) \psi_2(\vec{h})$$

$\vec{h} \approx$ CONSTANT OF MOTION

- $\vec{z}$ NOT MEASURABLE (WIGNER-ARAKI-YAMASE TH.)

CLOSED UNIVERSE

$$\psi_{\vec{R}}(\vec{z}) = \int d^3 h e^{-i \vec{h} \cdot \vec{z}} \psi_{\vec{R}}(\vec{h}) = e^{-i \vec{R} \cdot \vec{z}}$$

COM WAVE PACKETS

a) ALLOWED? (ALSO SUPERPOSITION OF 3-SPACES)

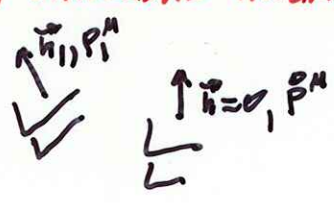
b) FORBIDDEN FROM SUPERSELECTION RULE?
(LIKE IN CANONICAL GRAVITY - GIULINI, MAROLF)

WHO WILL OBSERVE THE WAVE FUNCTION OF THE UNIVERSE?

FROZEN JACOBI DATA - NO EVOLUTION OF COM - NO INSTANTANEOUS SPREADING
OF COM WAVEPACKETS (WAY OUT OF HEISENBERG'S THEOREM)

2) UNIVERSAL INTERNAL SPACE H_{rel} DUE TO WIGNER COVARIANCE

$$P^\mu \approx m c (1; 0)$$



$$P_1^\mu = L^\mu_\nu(P_1, \vec{P}) \vec{P}^\nu = m c L^\mu_0(\vec{P}_1, \vec{P}), \quad h_1^i = L^i_0(P_1, \vec{P})$$

$$\eta_L^i \xrightarrow{\Lambda} R^i_s(P, \Lambda) \eta_L^s \quad \Lambda = L(P, \vec{P}) \quad \boxed{R(P, L(P, \vec{P})) = 1}$$

$$\vec{\eta}_L \rightarrow \vec{\eta}_i$$

$$H_{rel} = \bigotimes_i H_{\vec{p}_i} \bigg|_{\substack{\vec{p} \approx 0 \\ \vec{k} \approx 0}} = \bigotimes_a H_{\vec{s}_a}$$

INTERMEDIATE UNPHYSICAL HILBERT SPACE

$$\langle \phi_1, \phi_2 \rangle = \int d^3 \vec{s}_a \phi_1^*(t, \vec{s}_a) \phi_2(t, \vec{s}_a)$$

$$i \frac{\partial}{\partial t} \phi(t, \vec{s}_a) = \hat{H}_c \phi(t, \vec{s}_a)$$

$$\langle \psi_1, \psi_2 \rangle = \int d^3 \vec{q}_i \psi_1^*(t, \vec{q}_i) \psi_2(t, \vec{q}_i)$$

$$i \frac{\partial}{\partial t} \psi(t, \vec{q}_i) = \hat{H}_c \psi(t, \vec{q}_i)$$

INTERACTION DEPENDENT

GUPTA-BLEULER FOR REST-FRAME CONDITIONS

$$\langle \text{PHYS} | \vec{P} | \text{PHYS} \rangle = \langle \text{PHYS} | \vec{K} | \text{PHYS} \rangle = 0$$

$$\begin{cases} \hat{H}_c \phi_{nem}(\vec{s}_a) = \epsilon_n \phi_{nem} \\ \hat{L}_z \phi_{nem} = l(l+1) \phi_{nem} \\ \hat{S}_3 \phi_{nem} = m \phi_{nem} \end{cases}$$

ELEMENTARY SOLUTIONS

↓
WAVE PACKETS

$$\psi_{nem}(\vec{q}_i | t, \vec{s}_a) = \psi_{nem}(\vec{q}_i) e^{-i \epsilon_n t} \phi_{nem}(\vec{s}_a)$$

$$\xrightarrow{t \rightarrow \infty} e^{-i m c^2 t} \left(\text{NON-REL WAVE FUNCTION IN HAMILTON-JACOBI FOR COM} \right)$$

STANDARD ATOMIC PHYSICS - $O(\frac{1}{c^2})$ - NEITHER GALILEI NOR POINCARÉ GROUP

RELATIVISTIC ATOMIC PHYSICS - EFFECTIVE THEORY FOR $E < \text{THRESHOLD FOR PAIR PRODUCTION}$

FIXED NUMBER OF POSITIVE-ENERGY SCALAR (OR SPINNING) PARTICLES WITH COULOMB MUTUAL INTERACTION + TRANSVERSE ELECTRO-MAGNETIC FIELD (RADIATION GAUGE)

ISOLATED SYSTEM IN THE REST-FRAME INSTANT FORM

GRASSMANN-VALUED ELECTRIC CHARGES $Q_L^2 = 0, Q_L Q_J = Q_J Q_L \neq 0$ for $L \neq J$

UV REGULARIZATION OF SELF-ENERGIES

IR REGULARIZATION (NO SOFT PHOTON EMISSION)

POTENTIAL IN CAUCHY PROBLEM OF 1-PHOTON EXCHANGE

LIENARD-WIECHERT POTENTIALS = FUNCTIONS OF $\vec{q}_i, \vec{k}_i$ INDEPENDENT

OR RETARDED, ADVANCED, ... GREEN FUNCTIONS

CLASSICAL CANONICAL TRANSF

$\vec{q}_i, \vec{k}_i$ + COULOMB

$\vec{q}_i, \vec{k}_i$ (COULOMB DRESSED) + COULOMB + DARWIN

← BETHE SALPETER

$\vec{A}_\perp, \vec{E}_\perp$ INTERPOLATING FIELDS

$\vec{A}_\perp^{\text{RAD}}, \vec{E}_\perp^{\text{RAD}}$

IN, OUT FREE FIELDS (NO HADG THEOREM) QUANTUM 2

IMPLICATIONS FOR RELATIVISTIC ENTANGLEMENT

ISOLATED SYSTEMS - CLOSED UNIVERSE

CLOCK SYNCHRONIZATION (INSTANTANEOUS 3-SPACE, CAUCHY PROBLEM)

KINEMATICAL NON-LOCALITY - THE RELATIVISTIC CENTER OF MASS IS A GLOBAL QUANTITY NOT LOCALLY MEASURABLE

+ PREFERRED $\vec{h}$ -BASIS
(3-SPACE)

SEE NEWTON-WIGNER LOCALIZATION PROBLEM
THE NW POSITION OPERATOR CANNOT BE BUILT WITH
LOCAL OR QUASI-LOCAL OPERATORS

$\Rightarrow$ NON-COVARIANCE OF $\vec{P} = m\vec{c}\vec{\alpha}_{NW}$ IRRELEVANT

KINEMATICAL SPATIAL NON-SEPARABILITY

$$H \approx H_{com} \otimes H_{rel} \neq H_1 \otimes H_2 \otimes \dots$$

POSITION OR MOMENTUM BASES KINEMATICALLY
PREFERRED

PREFERRED ROBUST POINTER BASES
IN H_{rel} FROM DECOHERENCE

NO ZEROth POSTULATE ON COMPOSITE
SYSTEMS, NO EINSTEIN SEPARABILITY

ONLY RELATIVE VARIABLES (WEAK RELATIONISM $\neq$ ROVELLI)

THESE RELATIVISTIC PROPERTIES BEFORE QUANTUM NON-LOCALITY (ENTANGLEMENT)

NON-LOCALITY (FAKE A-A-A-D INTERACTION) - FROM KINEMATICS OF PONCARE GROUP
(NO SUPERLUMINAL SIGNALS)

NO EINSTEIN REALISM - LOCAL HIDDEN VARIABLES RULED OUT BY EXPERIMENTS
SUPERPOSITION OF MACROSCOPIC OBJECTS (DOUBLE SLIT
EXPERIMENT WITH FULLERENE MOLECULES) SHIELDED FROM
RELEVANT ENVIRONMENT

REALITY? - PROBABILISTIC RULES + PROBLEM OF OUTCOMES (COLLAPSE
AND EIGENVECTOR-EIGENVALUE LINK OR SPONTANEOUS COLLAPSE
OR EVERETT OR BOHN)

INTERPRETATION OF THE REDUCED DENSITY MATRIX
(IMPROPER MIXTURE)

DECOHERENCE - OPEN SUBSYSTEMS, BREAKING OF UNITARITY, PREFERRED ROBUST POINTER
BASES (BUT SPECIAL INTERACTIONS, SCHRÖDTER DECOMPOSITION, 3 SUBSYSTEMS?)
NO SOLUTION TO PROBLEM OF OUTCOMES

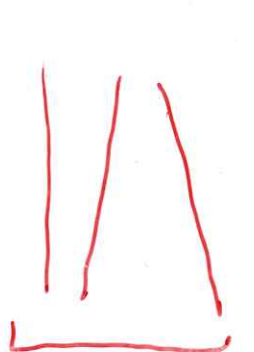
EMERGENCE OF QUASI-CLASSICAL MACROSCOPIC OBJECTS (ENHANCEMENT, GAUSSIAN MINIMAL UNCERTAINTY
WAVE PACKETS, WITH POSITIVE WIGNER FUNCTION) - OPEN PROBLEM

O. PESSOA JR - WE NEED AN ENVIRONMENT-INDUCED COLLAPSE (MANY-BODY PROBLEM)

ISOLATED SYSTEM IN REST-FRAME QM

$H_{com} \otimes H_{rel}$

RELEVANCE OF
TRAJECTORIES
OF PARTICLES
AND PHOTONS
(ABSENT IN
NRQM)



SYSTEM
MICROSCOPIC

ALICE
OBSERVER 1
MEASURING
APPARATUS

MACROSCOPIC
QUASI-CLASSICAL

BOB
OBSERVER 2
MEASURING
APPARATUS

MACROSCOPIC
QUASI-CLASSICAL

ENVIRONMENT

MACROSCOPIC
EITHER QUANTUM
OR QUASI-CLASSICAL

LOOK FOR THE SET OF CANONICAL RELATIVE VARIABLES MORE SUITABLE
TO THE PROTOCOL OF THE EXPERIMENT (BELL INEQUALITIES, TELEPORTATION)

LIKE IN MOLECULAR PHYSICS - CANONICAL JACOBI COORDINATES MOST SUITABLE
TO DESCRIBE THE STRONGEST CHEMICAL BONDS INSIDE THE
MOLECULE (THE OTHERS TREATED WITH PERTURBATION THEORY)

MAX. NUMBER OF CONSTANTS OF MOTION
PREFERRED MACROSCOPIC OBJECTS

STUDY $\langle \hat{X}_L^A(\omega) \rangle$ OF PREDICTIVE WORDLINES (NOT COMMUTING) TO RECOVER CLASSICAL TRAJECTORIES

NO MICROCAUSALITY BUT CLUSTER DECOMPOSITION
(CRITICAL TO LOCAL QFT - UNSHARP
MEASUREMENTS)

ENTANGLEMENT WITH EXTERNAL NON-LOCAL CENTER OF MASS NOT MEASURABLE

SUPERSELECTION RULE?

$$\Psi = \Psi_{com} \sum \Pi_a \Psi(\vec{r}_i, \vec{p}_a)$$

$$H_{rel} = \otimes_a H_{sa} \quad \text{OR} \quad \otimes_i H_{\vec{r}_i} \Big|_{\vec{p} \approx 0, \vec{r} \approx 0} \quad (\text{INTERACTION-DEPENDENT})$$

STUDY SIMPLE EXAMPLES

WITH MACROSCOPIC BODIES $\vec{r} \approx 0$ MAY BE APPROXIMATED
WITH $\vec{\eta}_{INT.C.O.M.} \approx 0$ AND WE MAY WORK IN THE
SEPARABLE UNPHYSICAL HILBERT SPACE $\otimes_L H_{\vec{r}_L}$
RESTRICTED AT THE END ($\vec{p} \approx 0, \vec{r} \approx 0$)

ATOMIC PHYSICS - RAY OF LIGHT + POLARIZATION (IN NRQM ONLY POLARIZATION)
SPINNING PARTICLES

TELEPORTATION TO SPACE STATION - DIFFICULT CONTROL OF POLARIZATION DUE TO ATMOSPHERE

GENERAL RELATIVITY

EINSTEIN SPACETIMES AND BEYOND

GEOMETRICAL VIEW OF THE GRAVITATIONAL FIELD

DESER

$$G_{\mu\nu} = \frac{8\pi G}{c^4} T_{\mu\nu}$$

EVERYTHING IS DYNAMICAL

CAUCHY SURFACE NEEDED

DOUBLE ROLE OF THE METRIC TENSOR $g_{\mu\nu}(x)$

A) POTENTIAL OF THE GRAVITATIONAL FIELD (LIKE IN EM AND YM)

B) DYNAMICAL CHRONOGEOMETRICAL STRUCTURE OF SPACETIME

$$ds^2 = g_{\mu\nu}(x) dx^\mu dx^\nu$$

$g_{\mu\nu}(x)$ TEACHES RELATIVISTIC CAUSALITY
TO ALL OTHER FIELDS



POINT-DEPENDENT TRAJECTORY OF MASSLESS PARTICLES
(PHOTONS, GLUONS ---) - NULL GEODESICS ($ds^2 = 0$)

ACES WINDOW ON IT! (2013 - GRAV. REDSHIFT OF THE GEOID)

NON-GEOMETRICAL PARTICLE VIEW

FEYNMAN

QFT, SM --- STRING THEORY

FIX THE BACKGROUND SPACE-TIME - ONLY KNOWN WAY TO DEFINE
FOCK SPACE AND QFT



PROPERTY B) IS LOST! GRAVITON - MASSLESS PARTICLE LIKE PHOTON,
TO RECOVER IT IS A UNSOLVED NON-PERTURBATIVE PROBLEM

EVEN NOT KNOWN HOW TO FORMULATE QFT (PARTICLE SM) IN

NON-INERTIAL FRAMES IN MINKOWSKI SPACETIME

(TORRE-VARADARASIAN NO-GO THEOREMS 1999

- NO UNITARY EVOLUTION)

RELATIVITY PRINCIPLE $\rightarrow$ PRINCIPLE OF GENERAL COVARIANCE

POINCARÉ GROUP $\rightarrow$ DIFFEOMORPHISM GROUP

EQUIVALENCE PRINCIPLE - GLOBAL INERTIAL FRAMES DO NOT EXIST

ONLY LOCAL FREE FALL

CHOICE OF THE CLASS OF SPACETIMES

AND THE PROBLEM OF TIME

CANONICAL GRAVITY REQUIRES GLOBALLY HYPERBOLIC SPACETIMES
(EXISTENCE OF GLOBAL TIME) AND ABSENCE OF SINGULARITIES FOR FINITE TIMES

A1) SPATIALLY COMPACT SPACETIMES WITHOUT BOUNDARY (LOOP QG, MACH PRINCIPLE)

NO BACKGROUND SPACETIME

NO POINCARÉ GROUP (PARTICLES?), NO EXTRA DIMENSIONS
AND NO FOCK SPACE

3+1 SPLITTING OF SPACETIME, QUANTIZATION OF 3-SPACE

H₂O FROZEN PICTURE, TIME EVOLUTION DEFINED ONLY LOCALLY
NO TIME PICTURE (ROVELLI, BARBOUR, ...)

AVOIDS SINGULARITIES (BOUNCES) IN QUANTUM LOOP COSMOLOGY

A2) QFT AND STRING THEORY

FIXED BACKGROUND SPACETIME (USUALLY MINKOWSKI), EXTRA DIMENSIONS

GRAVITY REDUCED TO SPIN 2 GRAVITON OVER MINKOWSKI SPACETIME

TO GET A VARYING LIGHTCONE IS A NON-PERTURBATIVE PROBLEM

HOW TO DEFINE THE LIGHTCONE WITH COMPACTIFIED EXTRA DIMENSIONS?

NOT YET EXISTING IN NON-INERTIAL FRAMES

IN BOTH CASES THE USE OF NUCLEOSYNTHESIS IS A BIG EXTRAPOLATION

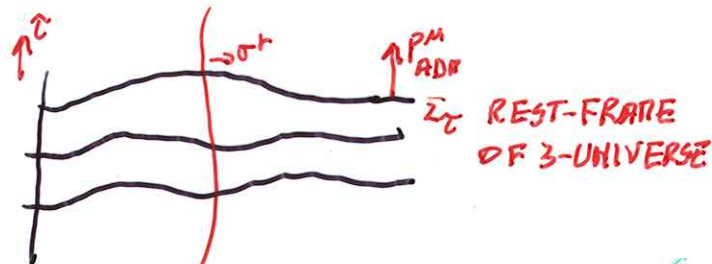
PARTICLE PHYSICS → LOOK FOR A FORMULATION OF
GENERAL RELATIVITY CONTAINING A
POINCARÉ GROUP

SPACETIMES ASYMPTOTICALLY FLAT AT SPATIAL INFINITY
WITHOUT SUPER-TRANSLATIONS

THE ASHTEKAR SPI GROUP OF ASYMPTOTIC SYMMETRIES

REDUCES TO THE ASYMPTOTIC ADM POINCARÉ GROUP

$$P_{ADM}^{\mu}, J_{ADM}^{\mu\nu}$$



~~GLOBAL~~

GLOBAL NON-INERTIAL FRAMES CENTERED ON ACCELERATED OBSERVERS

ASYMPTOTIC INERTIAL OBSERVERS → FIXED STARS
ASYMPTOTIC BACKGROUND

$G \rightarrow 0$ DEPARAMETRIZATION TO THE NON-INERTIAL REST-FRAME INSTANT FORM OF PARTICLE PHYSICS

$$P_{ADM}^{\mu} \rightarrow P^{\mu}, J_{ADM}^{\mu\nu} \rightarrow J^{\mu\nu}$$

$g_{\mu\nu}(x) \xrightarrow[\text{SPATIAL INFINITY}]{\text{}} \eta_{\mu\nu}$ (IN A DIRECTION-INDEPENDENT WAY)

NOT THE STANDARD BOUNDARY CONDITIONS OF FRW COSMOLOGY

IN ABSENCE OF MATTER CHRISTODOULOU-KLAINTERMANN SPACETIMES ARE OK

MODEL OF UNIVERSE AFTER THE LAST SCATTERING (OR RECOMBINATION) TIME, CORRESPONDING TO A 3-SPACE Σ_{CMB} WHERE THE EMISSION OF THE PHOTONS OF CMB HAPPENED ($\sim 3 \cdot 10^5$ YEARS AFTER BIG BANG)

NO FROZEN PICTURE $H_0 \approx \dot{E}_{ADM} \neq 0$

WEAK ADM ENERGY $\dot{E}_{ADM} = \int d^3\vec{\sigma} \dot{E}_{ADM}(\vec{\sigma}, \vec{\sigma}) \approx E_{ADM} = \int_{\Sigma_{ADM}} \dots$
STRONG ADM ENERGY

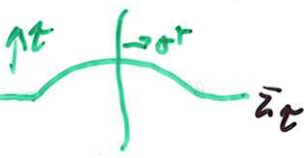
USE ADM ACTION FOR EINSTEIN THEORY WITH THE 4-METRIC EXPRESSED IN TERMS OF TETRADS (REQUIRED BY

$${}^4g^{\mu\nu}(x) \approx E_{(a)}^{\mu}(x) \gamma^{(a)(b)} E_{(b)}^{\nu}(x)$$

THE COUPLING TO FERMIONS)

NO DEGENERACIES DUE TO KILLING SYMMETRIES

1) GO TO RADAR 4-COORDINATES ADAPTED TO THE 3+1 SPLITTING AND HAVING AN ACCELERATED OBSERVER AS ORIGIN



$$x^A \mapsto \sigma^A(t; \sigma^r)$$

$$\sigma^A \mapsto x^\mu = z^\mu(t, \vec{\sigma}) \quad \text{EMBEDDING OF } \Sigma_t$$

$$z^\mu_A = \frac{\partial z^\mu}{\partial \sigma^A}$$

$${}^4g_{AB} = z^\mu_A z^\nu_B g_{\mu\nu}$$

$${}^4E^\mu_A = z^\mu_A {}^4E^A_\alpha$$

2) IN THE TANGENT SPACE OF EACH POINT OF Σ_t BOOST TETRADS AND COTETRADS TO TETRADS AND COTETRADS ADAPTED TO Σ_t WITH STANDARD WIGNER BOOST



COTETRADS

$$E^{(\alpha)}_A = L^{(\alpha)}_{(a)}(\varphi_{(a)}) \begin{pmatrix} N \\ 0 \end{pmatrix} + L^{(\alpha)}_{(a)}(\varphi_{(a)}) \begin{pmatrix} n_{(a)} \\ {}^3e_{(a)r} \end{pmatrix}$$

$\begin{matrix} \text{EA} \\ \text{E}^{(a)}_{(a)} \end{matrix}$

$$N = 1 + n \quad \text{LAPSE}$$

$$n_{(a)} = {}^3e_{(a)r} n^r \quad \text{SHIFT}$$

$\begin{cases} {}^3e^r_{(a)} \text{ TRIADS ON } \Sigma_t \\ {}^3e_{(a)r} \text{ COTRIADS ON } \Sigma_t \end{cases}$

$\varphi_{(a)}$ BOOST PARAMETERS

$${}^4g_{tt} = \varepsilon \left[(1+n)^2 - \sum_a n_{(a)}^2 \right]$$

$${}^4g_{tr} = -\varepsilon n_r = -\varepsilon n_{(a)} {}^3e_{(a)r}$$

$${}^4g_{rs} = -\varepsilon {}^3g_{rs} = -\varepsilon \sum_a {}^3e_{(a)r} {}^3e_{(a)s}$$

$$\varepsilon = + \quad \gamma^{\mu\nu} = \begin{pmatrix} + & - & - & - \\ - & + & + & + \end{pmatrix}$$

$${}^3g_{rs} \quad (+++)$$

CANONICAL BASIS

$\varphi_{(a)}$	$1+n$	$n_{(a)}$	${}^3e_{(a)r}$
x_0	x_0	x_0	${}^3\pi^r_{(a)}$

16 FIELDS
+
16 MOMENTA

14 FIRST CLASS CONSTRAINTS (GENERATORS OF GAUGE TRANSF.)
↓
14 GAUGE VARIABLES (INERTIAL EFFECTS)

⇒ 2+2 PHYSICAL D.O.F. OF THE GRAY. FIELD (TIDAL EFFECTS)



CONGRUENCE OF EULERIAN OBSERVERS

VORTICITY = 0, $\sigma_{(a)(b)}$ SHEAR ($= {}^3e^r_{(a)} {}^3e^s_{(b)} \sigma_{rs}$)

EXPANSION θ , ACCELERATION ${}^3a_r = \nabla_r(1+n)$
 $\approx -s^3K$

↓
POLARIZATION OF GRAY. WAVES IN WEAK FIELD APPROX.

YORK CANONICAL BASIS

AFTER A SUITABLE CANONICAL TRANSF.

ADAPTED TO 10 OF THE 14 CONSTRAINTS SUCH THAT

$${}^3e_{(a)r} = R_{(a)(b)}(\alpha_{(a)}) V_{rb}(\vartheta^n) \tilde{\Phi}^{1/3} e^{\frac{1}{2} \sum_a \gamma_{aa} R_a}$$

$${}^3g_{rs} = \sum_a {}^3e_{(a)r} {}^3e_{(a)s} = \tilde{\Phi}^{2/3} \sum_a V_{ra}(\vartheta^n) V_{sa}(\vartheta^n) e^{2 \sum_a \frac{1}{2} \gamma_{aa} R_a}$$

DIAGONALIZED WITH ROTATION MATRIX

$\varphi_{(a)}$	$\alpha_{(a)}$	$1+n$	$\bar{\eta}_{(a)}$	ϑ^n	$\tilde{\Phi}$	R_a
≈ 0	≈ 0	≈ 0	≈ 0	$\pi_n^{(b)}$	$\pi_{\tilde{\Phi}}$	π_a

TIDAL EFFECTS

ABSENT WITH A TIMELIKE KILLING VECTOR

$\leftrightarrow \sigma_{(a)(a)}$

0 GAUGE VARIABLES (INERTIAL EFFECTS)

ONLY GAUGE POTENTIAL (LORENTZ SIGNATURE)

$\varphi_{(a)}, \alpha_{(a)}$

NEWMAN-PENROSE $O(3,1)$ FREEDOM IN THE CHOICE OF TETRAIDS

GAUGE FREEDOM OF EACH OBSERVER IN THE ORIENTATION OF 3 GYROSCOPES AND IN THE CHOICE OF THE LAW OF TRANSPORT ALONG THE WORLDLINE

$\varphi_{(a)}(\xi, \bar{\tau}) \approx 0, \alpha_{(a)}(\xi, \bar{\tau}) \approx 0$ SCHWINGER TIME GAUGES

ϑ^n

GAUGE FREEDOM IN THE CHOICE OF 3-COORDINATES ON $\Sigma_{\bar{\tau}}$

$\pi_n^{(b)} \leftrightarrow \sigma_{(a)(b)} |_{a \neq b}$

DETERMINED BY SUPERPOTENTIAL CONSTRAINTS

$\chi_{(a)}(\xi, \bar{\tau}) \approx 0$

$\bar{\eta}_{(a)}$

GAUGE FREEDOM IN GRAVITOMAGNETISM

3-COORD-DEPENDENT CORIOLIS ... INERTIAL EFFECTS

$$\tilde{\Phi} = \Phi^6 = |{}^3e_{(a)r}| = \sqrt{\det {}^3g_{rs}}$$

3-VOLUME ELEMENT DETERMINED BY THE SUPER-HAMILTONIAN CONSTRAINT (LICHNEROWICZ EQ.)

$$\chi(\xi, \bar{\tau}) \approx 0 \rightarrow \tilde{\Phi} \approx F(R_a, \pi_a, \vartheta^n, \pi_{\tilde{\Phi}})$$

$$\pi_{\tilde{\Phi}} \approx \frac{c^3}{12\pi G} {}^3K$$

YORK TIME - GAUGE FREEDOM IN THE CHOICE OF THE INST. 3-SPACE (CLOCK SYNCHRONIZATION)

$${}^3K = -\epsilon \bar{\nabla} \cdot \vec{n}$$

THIS INERTIAL EFFECT IS ABSENT IN NEWTONIAN GRAVITY

$\Sigma_{\bar{\tau}}$ $\uparrow \uparrow \uparrow$ ${}^3K_{rs}$ EXTRINSIC CURVATURE (HOW $\Sigma_{\bar{\tau}}$ IS EMBEDDED IN 4-SPACE)

$1+n$

GAUGE FREEDOM IN THE CHOICE OF THE LOCAL UNIT OF PROPER TIME

HAMILTONIAN

$$H_D = \frac{1}{c} \dot{E}_{AD\pi} + \int d^3\sigma \left[\eta \mathcal{H} + \bar{\eta}_{(\omega)} \mathcal{H}_{(\omega)} + \lambda_n \Pi_n + \lambda_{(\omega)} \Pi_{n(\omega)} \right] (\mathcal{E}_i \vec{\sigma})$$

$$\lambda_n \stackrel{\circ}{=} \partial_c \eta, \quad \lambda_{(\omega)} \stackrel{\circ}{=} \partial_c \bar{\eta}_{(\omega)}, \quad \text{REST-FRAME CONDITION} \quad \sum_{AD\pi} \dot{P}_{AD\pi} \approx 0$$

WEAK AD\pi ENERGY

$$\frac{1}{c} \dot{E}_{AD\pi} = M_{AD\pi} c = \int d^3\sigma \left[\frac{4\pi G}{c^3} \tilde{\Phi}^{-1} \sum_{\vec{a}} \Pi_{\vec{a}}^2 + \frac{c^3}{16\pi G} \tilde{\Phi} \sum_{a \neq b} \sigma_{(a)(b)}^2 - \frac{c^3}{4\pi G} \tilde{\Phi} ({}^3K)^2 \right] (\mathcal{E}_i \vec{\sigma})$$

TIDAL KINETIC TERM

GRAVITOMAGNETIC KINETIC TERM
MATTER CURRENT M_t

NEGATIVE ENERGY DENSITY IN EVERY GAUGE $[{}^3K = \epsilon \mathcal{G}]$

INERTIAL EFFECT
CLOCK SYNCH. GAUGE
FREEDOM

$$- \frac{c^3}{16\pi G} \int d^3\sigma \mathcal{G}(\mathcal{E}_i \vec{\sigma})$$

$\mathcal{G}[R_{\vec{a}}, \tilde{\Phi}, \eta^n]$ - THE $\mathcal{P}$ TERM IN 3R
3-COORD-DEPENDENT POTENTIAL

= 0 IN CIRC 3-SPACES ${}^3K(\mathcal{E}_i \vec{\sigma}) \approx 0$

$$\rightarrow \int d^3\sigma \mathcal{M}(\mathcal{E}_i \vec{\sigma})$$

MATTER DENSITY

$$\mathcal{G}(\mathcal{E}_i \vec{\sigma}) \propto \mathcal{M}(\mathcal{E}_i \vec{\sigma})$$

SUPERHAMILTONIAN CONSTRAINT (LICHNEROWICZ EQ)

$$\Phi = \tilde{\Phi}^{1/6}$$

$${}^3R[\Phi, \eta^n, R_{\vec{a}}] = \Phi^5 (-8\hat{\Delta}\Phi + {}^3\hat{R}[\Phi, R_{\vec{a}}]\Phi)$$

$$\Phi (\hat{\Delta}\Phi - \frac{1}{8} {}^3\hat{R}\Phi) + \frac{8\pi G^2}{c^6} \Phi^{-6} \sum_{\vec{a}} \Pi_{\vec{a}}^2 + \frac{1}{8} \Phi^6 \sum_{a \neq b} \sigma_{(a)(b)}^2 - \frac{1}{12} \Phi^6 ({}^3K)^2 = - \frac{2\pi G}{c^3} \mathcal{M} \quad \text{PASS DENSITY}$$

SUPERMOMENTUM CONSTRAINTS $\mathcal{H}_{(\omega)}(\mathcal{E}_i \vec{\sigma}) \approx 0$ DEPEND ON THE PASS CURRENT (GRAVITOMAGNETISM)

$$\text{AD\pi EQS} \quad \partial_c {}^3g_{rs} \stackrel{\circ}{=} \dots, \quad \partial_c {}^3K_{rs} \stackrel{\circ}{=} \dots$$

$$\text{A) CONTRACTED BIANCHI IDENTITIES} \quad \partial_c \tilde{\Phi} \stackrel{\circ}{=} \dots, \quad \partial_c \sigma_{(a)(b)}|_{a \neq b} \stackrel{\circ}{=} \dots$$

PRESERVATION IN TIME OF THE CONSTRAINTS (PROBLEM IN NUMERICAL GRAVITY)

$$\text{B) EVOLUTION EQS FOR GAUGE VARIABLES} \quad \left\{ \begin{array}{l} \partial_c \vartheta \stackrel{\circ}{=} \dots \quad \vartheta = -\epsilon {}^3K \\ \partial_c \vartheta^n \stackrel{\circ}{=} \dots \end{array} \right. \quad \text{RAYCHAUDHURI EQ.}$$

GAUGE FIXING

$$\left\{ \begin{array}{l} \vartheta^n x \dots \Rightarrow \partial_c \vartheta^n = \dots \quad \text{EQ FOR SHIFT } \bar{\eta}_{(\omega)} \\ \vartheta = -\epsilon {}^3K x \dots \Rightarrow \partial_c \vartheta = \dots \quad \text{EQ FOR LAPSE } 1+\eta \end{array} \right.$$

$$\text{C) EVOLUTION EQS FOR TIDAL VARIABLES} \quad R_{\vec{a}}, \Pi_{\vec{a}} (\Leftrightarrow \sigma_{(a)(b)})$$

$$\left\{ \begin{array}{l} \partial_c R_{\vec{a}} \stackrel{\circ}{=} \dots \rightarrow \Pi_{\vec{a}} \stackrel{\circ}{=} \dots \\ \partial_c \Pi_{\vec{a}} \stackrel{\circ}{=} \dots \end{array} \right. \rightarrow \partial_c^2 R_{\vec{a}} \stackrel{\circ}{=} \dots$$

ONLY FREE
CAUCHY DATA

D) MATTER HAMILTON EQS

RESEARCH PROGRAM

A) DARK ENERGY FROM BACKREACTION - BUCHERT'S AVERAGING APPROACH

HE ADDS A COSMOLOGICAL PERFECT FLUID (DUST), ASSUMES THAT IT IS IRROTATIONAL AND USES THE 3-SPACES $\bar{\Sigma}_{\text{DUST}} \perp$ FLUID 4-VELOCITY AS INST. 3-SPACES

N.B. IN FRW COSMOLOGY BASED ON THE COSMOLOGICAL PRINCIPLE THE FLUID HAS ZERO VORTICITY, SHEAR AND 3-ACCELERATION, ONLY THE EXPANSION IS NOT ZERO (CHOOSING OBSERVERS FOR CMB)

N.B. FLATNESS PROBLEM - IT CONCERNS THE INTRINSIC CURVATURE OF 3-SPACE, NOT ITS EXTRINSIC CURVATURE AS A 3-SPACE EMBEDDED IN SPACETIME

USE DUST AS MATTER - NO REASON FOR BEING IRROTATIONAL

ADD BY HAND 3 CONSTRAINTS $\omega_r(x, \vec{r}) \approx 0$ (VANISHING VORTICITY)

LOOK FOR CONSISTENCY WITH THE EINSTEIN'S EQS $\{\omega_r(x, \vec{r}), H_D\} \approx 0$

IF CONSISTENT (I.E. IF THERE IS A SUBSET OF SOLUTIONS OF EINSTEIN'S EQS WITH IRROTATIONAL DUST), THEN

a) THERE IS A PREFERRED 3+1 SPLITTING IDENTIFIED BY MATTER (CMB?)

b) IF THE PREFERRED 3-SPACES ARE NOT CTC ($3K \neq 0$), THERE IS A NEGATIVE TERM IN $\check{E}_{\text{ADM}}$

c) FIND THE GAUGE SELECTING THE PREFERRED 3+1 SPLITTING FOR INSTANCE BY IMPOSING $\ell^A(x, \vec{r}) \approx u^A_{\text{DUST}}(x, \vec{r})$

d) REDO BUCHERT'S CALCULATIONS

AVERAGED UNIVERSE EXPANDS AT AN ACCELERATED PACE IF

$$\frac{2}{3} \langle (3K - \langle 3K \rangle)^2 \rangle - 2 \langle \sigma^2 \rangle > \frac{1}{2} \langle S \rangle$$

$\downarrow$ $\sum_{a \neq b} \sigma_{(a)(b)}^2 + \alpha \sum_a \pi_a^2$ $\rightarrow$ MATTER DENSITY
 $\uparrow$ MATTER CURRENT

3) DARK MATTER

STUDY THE 2-BODY PROBLEM, AT LEAST IN THE WEAK FIELD APPROXIMATION,
TO FIND POST-MINKOWSKIAN GRAVITY IN 3-ORTHOGONAL GAUGES

[${}^3g_{ij}(t, \vec{r})$ DIAGONAL - IS IT A HARTONIC GAUGE FOR SOME ${}^3K(t, \vec{r})$?]

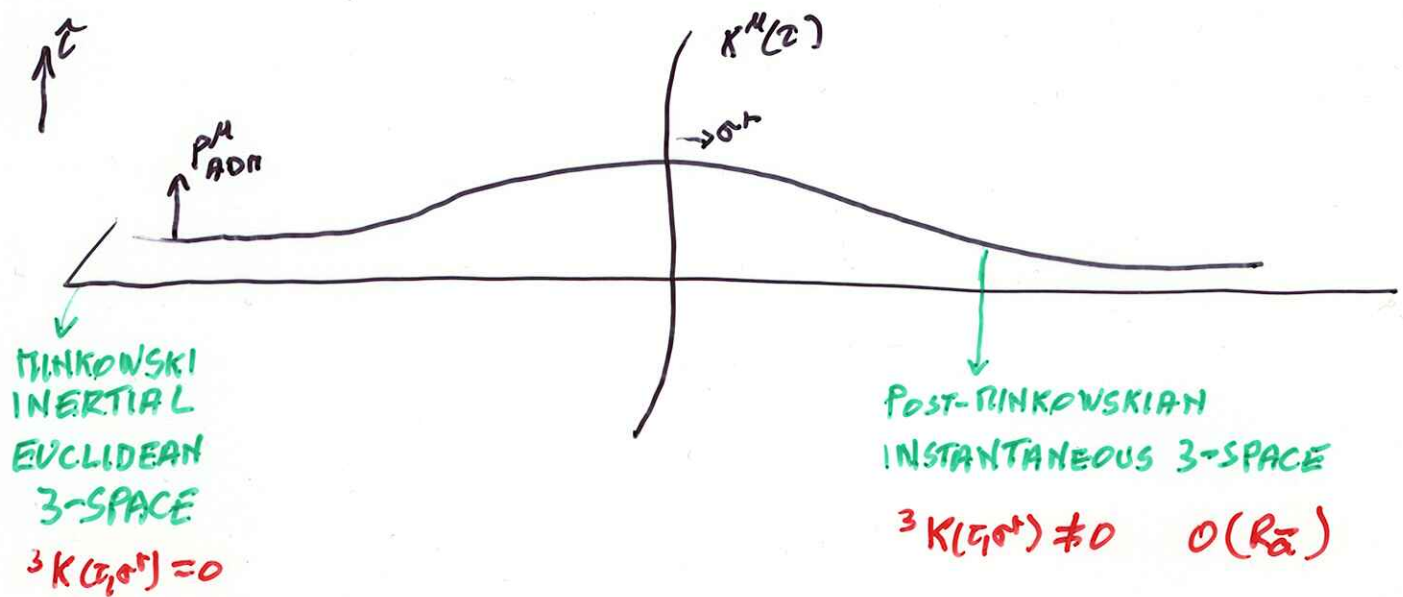
- FIND THE POST-MINKOWSKIAN NEWTON AND GRAVITOMAGNETIC POTENTIALS BETWEEN THE TWO PARTICLES FOR SOME ${}^3K(t, \vec{r})$
 - TRY TO SOLVE THE TIDAL EVOLUTION EQS. $\partial_t^2 R_{\alpha\beta}(t, \vec{r}) \stackrel{?}{=} \dots$ IN THE WEAK FIELD APPROXIMATION [IF $R_{\alpha\beta}(t, \vec{r}) \sim e^{i[\omega t - \vec{k} \cdot \vec{r}]}$ $\Rightarrow \omega^2 - \vec{k}^2 \neq 0$ GAUGE FUNCTION OF MATTER AND ${}^3K(t, \vec{r})$]
 - FIND THE POST-MINKOWSKIAN GRAVITATIONAL WAVES WITH ASYMPTOTIC BACKGROUND, BOTH THE HOMOGENEOUS PART AND THE MATTER-DEPENDENT PART (LIKE WITH THE LIENARD-WIECHERT SOLUTION IN ELECTROMAGNETISM) FOR SOME ${}^3K(t, \vec{r})$.
 - FIND THE EMBEDDING $Z^4(t, \vec{r})$ OF Σ_t IN SPACETIME FOR SOME ${}^3K(t, \vec{r})$ - IT DESCRIBES THE DYNAMICAL NON-INERTIAL FRAMES
 - FIND THE (3K -DEPENDENT) TIMELIKE GEODESICS OF THE PREVIOUS SOLUTIONS AND STUDY THE MOTION OF A TEST PARTICLE CIRCULATING AROUND THE 2-BODY SYSTEM (SIMULATING A GALAXY). FIND THE POST-MINKOWSKIAN VIRIAL THEOREM, REPLACING THE KEPLER ONE ($v^2 \sim \frac{M}{R}$) USED FOR THE ROTATION CURVES OF GALAXIES.
- HOW MUCH DARK MATTER CAN BE DESCRIBED AS A RELATIVISTIC INERTIAL EFFECT INDUCED BY ${}^3K(t, \vec{r})$ TOGETHER WITH GRAVITOMAGNETISM (COOPERSTOCK-TIEU)?

THE RELATIVISTIC EQ. OF MOTION OF THE TEST PARTICLE CAN BE PUT IN THE FORM $m \vec{a} = \vec{F}_{\text{TIDAL+INERTIAL}}$

IN NEWTON EUCLIDEAN 3-SPACE, THE MOND APPROACH WRITES THEM IN THE FORM $m \vec{f}(\vec{a}) = \vec{F}_{\text{GRAVITY}}$ WITH $\vec{f}(\vec{a}) \neq \vec{a}$ AT LOW ACCELERATIONS.

$\Rightarrow$ THE INERTIAL COMPONENT OF $\vec{F}_{\text{TIDAL+INERTIAL}}$ SHOULD BE EFFECTIVE FOR LOW ACCELERATIONS.

POST-MINKOWSKIAN GRAVITATIONAL WAVES IN 3-ORTHOGONAL GAUGES



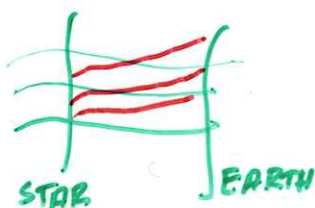
$$\left\{ \begin{array}{l} {}^3g_{rs}(x,0^+) \approx \tilde{\Phi}_{(0)}^{2/3}(x,0^+) \left[1 + 2 \sum_{\bar{a}} \gamma_{\bar{a}\bar{a}} R_{\bar{a}}(x,0^+) \right] \delta_{rs} \\ {}^4g_{rs}(x,0^+) \\ \varepsilon {}^4g_{zz}(x,0^+) \approx 1 + 2 \eta_{(0)}(x,0^+) \\ \varepsilon {}^4g_{zr}(x,0^+) \approx \tilde{\Phi}_{(0)}^{1/3}(x,0^+) \left[1 + \sum_{\bar{a}} \gamma_{\bar{a}\bar{a}} R_{\bar{a}}(x,0^+) \right] \bar{n}_{(r)(0)}(x,0^+) \end{array} \right.$$

$$\tilde{\Phi}_{(0)}^{(0)} = \sqrt{\det {}^3g_{rs}} \quad \text{3-VOLUME ELEMENT}$$

FINAL COMMENTS

- A) IN GR THE LIGHT-POSTULATES ON THE 2-WAY VELOCITY OF LIGHT ARE ASSUMED LOCALLY VALID

THE FITTING PROBLEM, THE EVALUATION OF THE LUMINOSITY DISTANCE, THE 1-WAY VELOCITY OF LIGHT REQUIRE 3+1 SPLITTINGS AND THE STUDY OF THE NULL GEODESICS OF SOLUTIONS OF EINSTEIN'S EQS



PROBLEMS OF LIGHT PROPAGATION, LENSING, ...

↓
ALSO FOR BEPI-COLOMBO!

- B) INHOMOGENEOUS AND ANISOTROPIC UNIVERSES REQUIRE SOLUTIONS OF EINSTEIN'S EQS **WITHOUT KILLING VECTORS**. INSTEAD THE COSMOLOGICAL (AND COPERNICAN) PRINCIPLE LEADS TO FRW COSMOLOGICAL MODEL WITH 6 KILLING VECTORS.

SPACE OF METRICS - STRATIFIED MANIFOLD WITH SINGULARITIES (GONE OVER CONE STRUCTURE - FISHER, MARSDEN ...)

NO REAL CONTROL OVER POSSIBLE DISCONTINUITIES IN THE GAUGE PROBLEM OF COMPARISON OF SPACETIMES WITHOUT KILLING VECTORS WITH FRW SPACETIMES

- C) HORIZON PROBLEMS (CURED WITH INFLATION) - RELATIVISTIC ENTANGLEMENT INTRODUCES "NON-LOCALITY AND SPATIAL NON-SEPARABILITY". AFTER A 3+1 SPLITTING OF SPACETIME CAUSALLY DISCONNECTED REGIONS HAVE A NON-SEPARABILITY IN THE INSTANTANEOUS 3-SPACES AND THE CAUCHY DATA HAVE TO BE GIVEN FOR RELATIVE VARIABLES (THE NON-LOCAL AND NON-COVARIANT RELATIVISTIC CENTER OF MASS OF THE 3-UNIVERSE IS DECOUPLED).

- d) OBSERVABLES FOR GENERAL RELATIVITY (IN ABSENCE OF KILLING VECTORS)

{ TIDAL $R_{\alpha\beta}, T_{\alpha\beta} \rightarrow 4$ EIGENVALUES OF WEYL TENSOR (4 ALGEBRAIC INVARIANTS INDEPENDENT FROM 4-COORDINATES AND TETRAIDS)
 MANY INERTIAL EFFECTS } $\rightarrow 4$ -SCALARS BUT TETRAD-DEPENDENT
 ELECTRO-MAGNETIC FIELD } (NEWMAN-PENROSE FORMALISM - FIND CANONICAL FORMULATION)
 MATTER (THINK TO GPS SATELLITES) - INTRINSICALLY COORDINATE-DEPENDENT

