

Searching for New Physics at the LHC GGI Florence, October 29 2009

The Leptonic Sector in Composite Higgs Models



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F. del Aguila, A. Carmona, J.S. arXiv:0911.xxxx

The LHC is here!

- The LHC is finally ready to study the mechanism of EWSB
- Composite Higgs models are a likely candidate
 - Natural realization of EWSB
 - Potentially interesting collider implications
- Weak realization in 5D models with warped ED
 - Quark sector well studied
 - What about the lepton sector?



Composite Higgs Models

• Two scale symmetry breaking: Georgi, Kaplan, et al. 84-85



5D realization of CHM

 Models of gauge-Higgs unification in warped ED realize the composite Higgs idea

 $ds^{2} = \left(\frac{R}{z}\right)^{2} \left(\eta_{\mu\nu}dx^{\mu}dx^{\nu} - dz^{2}\right), \quad R \leq z \leq R'$

 Gauge symmetry broken on both branes contains a massless zero mode for A₅

 Composite Higgs

• Finite (calculable) potential generated at loop level

 Minimal realistic models have been constructed based on SO(5)/SO(4) symmetry breaking pattern Agashe, Contino, Pomarol '05

EW constraints on CHM

 T and Zbb protected by custodial (and LR) symmetry



Agashe, Delgado, May, Sundrum '03 Agashe, Contino, Da Rold, Pomarol '06

 Main constraint: S plus fermionic contribution to T and Zbb (1 loop) Carena, Pontón, J.S., Wagner '06-07

 $m_{\rm gauge} \gtrsim 3.5 {
m TeV}$

 Fermions are expected to be much lighter

 $m_{
m fermion}\gtrsim 0.5~{
m TeV}$

- Most models focus on the quark sector
 - top is expected to be also composite
 - Masses and mixing angles naturally hierarchical
- Why is the leptonic spectrum so different?



 $|V_{
m CKM}| = egin{pmatrix} 0.97 & 0.23 & 0.004 \ 0.23 & \sim 1 & 0.04 \ xx & xx & \gtrsim 0.7 \end{pmatrix}$

 $|U_{\rm PMNS}| = egin{pmatrix} 0.8 - 0.84 & 0.53 - 0.6 & 0 - 0.17 \ 0.29 - 0.52 & 0.51 - 069 & 0.61 - 0.76 \ 0.26 - 0.5 & 0.46 - 0.66 & 0.64 - 0.79 \end{pmatrix}$

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Goes well with hierarchical masses

 $|U_{\rm PMNS}| = egin{pmatrix} 0.8 - 0.84 & 0.53 - 0.6 & 0 - 0.17 \ 0.29 - 0.52 & 0.51 - 069 & 0.61 - 0.76 \ 0.26 - 0.5 & 0.46 - 0.66 & 0.64 - 0.79 \end{pmatrix}$

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 - The pattern changes for tiny neutrino Dirac masses

Agashe, Okui, Sundrum 08; Carena, Medina, Shah, Wagner 09



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- But:
 - Flavor?
 - Why neutrinos so different?

- Why is the leptonic spectrum so different?
 - Global symmetries (A4, T') in simple RS models

Csaki, Delauny, Grojean, Grossman 08; Chen, Mahanthappa, Yu 09

- Use see-saw to generate the right neutrino scale
- Discrete symmetries predict the mixing pattern
- Added bonus: flavor protection (very useful in models with warped extra dimensions)

 $|U_{\rm PMNS}| = egin{pmatrix} 0.8 - 0.84 & 0.53 - 0.6 & 0 - 0.17 \ 0.29 - 0.52 & 0.51 - 069 & 0.61 - 0.76 \ 0.26 - 0.5 & 0.46 - 0.66 & 0.64 - 0.79 \end{pmatrix}$

 $U_{\rm HPS} = \begin{pmatrix} \sqrt{\frac{2}{3}} & \frac{1}{\sqrt{3}} & 0\\ -\frac{1}{\sqrt{6}} & \frac{1}{\sqrt{3}} & -\frac{1}{\sqrt{2}}\\ -\frac{1}{\sqrt{6}} & \frac{1}{\sqrt{3}} & \frac{1}{\sqrt{2}} \end{pmatrix} = \begin{pmatrix} 0.82 & 0.58 & 0\\ -0.41 & 0.58 & -0.71\\ -0.41 & 0.58 & 0.71 \end{pmatrix}$

Tri-Bi-Maximal mixing

Harrison, Perkins, Scott 02

Roadmap

- 5D realization of CHM: SO(5)/SO(4) (x U(1))
- Use identical structure for the quark and lepton sectors for MCHM₅ Contino, Da Rold, Pomarol 07
- Impose a global $A_4 \times G_{FN}$ symmetry
- Compute spectrum: lepton masses and mixings
- Compute constraints: PMNS, EWPT, LFV
- Study collider implications (in progress)

• Lepton sector in MCHM₅: based on 5's of SO(5)

$$\zeta_1 = \begin{pmatrix} \widetilde{X}_1 & \nu_1 \\ \widetilde{\nu}_1 & e_1 \end{pmatrix} \oplus \nu_1'$$

Recall

 $SO(4) \sim SU(2)_L \times SU(2)_R$ $5 = (2,2) \oplus 1$

Lepton sector in MCHM₅: based on 5's of SO(5)

$$Y=rac{1}{2}-rac{1}{2}=0$$
 $\zeta_1=\left(egin{array}{c} \widetilde{X}_1&
u_1\ \widetilde{
u}_1&e_1\end{array}
ight)\oplus
u_1'$
 $Y=-rac{1}{2}-rac{3}{2}=0$
 $\zeta_3=\left(egin{array}{c}
u_3& ilde{e}_3\ e_3& ilde{Y}_3\end{array}
ight)\oplus e_3'$

Different quantum numbers governed by an extra U(1)

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$$\zeta_3 = \begin{pmatrix} \nu_3 & \tilde{e}_3 \\ e_3 & \tilde{Y}_3 \end{pmatrix} \oplus e'_3 \qquad \qquad \zeta_\alpha = \begin{pmatrix} \nu_\alpha & \tilde{e}_\alpha \\ e_\alpha & \tilde{Y}_3 \end{pmatrix} \oplus e'_\alpha$$

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- Impose a global A₄ symmetry
- A₄: even permutations of 4 elements
 - 3 one-dimensional plus 1 three-dimensional representations:

(1), (1'), (1'') = $(1')^*$, (3)

Leptons in the MCHM₅ Impose a global A, symmetry (3)(3) $\zeta_1 = \begin{pmatrix} \widetilde{X}_1 & \nu_1 \\ \widetilde{\nu}_1 & e_1 \end{pmatrix} \oplus \nu'_1 \qquad \qquad \zeta_2 = \begin{pmatrix} \widetilde{X}_2 & \nu_2 \\ \widetilde{\nu}_2 & e_2 \end{pmatrix} \oplus \nu'_2$ (1, 1', 1'')(3) $\zeta_lpha = egin{pmatrix} u_lpha & ilde e_lpha \ e_lpha & ilde Y_3 \end{pmatrix} \oplus e'_lpha \ e'_lpha \end{pmatrix}$ $egin{array}{ccc} \zeta_3 = \left(egin{array}{ccc} u_3 & \widetilde{e}_3 \ e_3 & \widetilde{Y}_3 \end{array} ight) \oplus e_3' \ \end{array}$

 $\zeta_{1,2,3}$ family independent localization

The second s

• Global A_4 broken by brane scalars to Z_2 and Z_3 (to ensure TBM mixing)

 $UV: \langle \phi \rangle = (\tilde{v}, 0, 0) \qquad IR: \langle \phi' \rangle = (v', v', v')$

Write the most general terms allowed by the symmetries

$$-\mathcal{L}_{UV} = \frac{x_{\eta}}{2\Lambda} \eta \overline{\psi}_{\nu'_{2}} \overline{\psi}_{\nu'_{2}} + \frac{x_{\nu}}{2\Lambda} \phi \overline{\psi}_{\nu'_{2}} \overline{\psi}_{\nu'_{2}} + x_{l} \chi_{l_{1}} \psi_{l_{3}} + \text{h.c.} + \dots$$

$$-\mathcal{L}_{IR} = \left(\frac{R}{R'}\right)^{4} \left[\frac{y_{b}^{\alpha}}{\Lambda'} \left\{ \left(\phi'^{\dagger} \chi_{l_{3}}\right)^{\alpha} \psi_{l_{\alpha}} + \left(\phi'^{\dagger} \chi_{\tilde{l}_{3}}\right)^{\alpha} \psi_{\tilde{l}_{\alpha}} \right\} + \frac{y_{s}^{\alpha}}{\Lambda'} \left(\phi'^{\dagger} \chi_{e'_{3}}\right)^{\alpha} \psi_{e'_{\alpha}}$$

$$+ \frac{y_{b}}{\Lambda'} \left\{ \eta'^{*} \chi_{l_{1}} \psi_{l_{2}} + \eta'^{*} \chi_{\tilde{l}_{1}} \psi_{\tilde{l}_{2}} \right\} + \frac{y_{s}}{\Lambda'} \eta'^{*} \chi_{\nu'_{1}} \psi_{\nu'_{2}} \right] + \text{h.c.} + \dots$$

After A₄ breaking

$$-\mathcal{L}_{UV} \rightarrow \frac{1}{2} \psi_{\nu_{2}'} \hat{\theta}_{M}^{\dagger} \psi_{\nu_{2}'} + x_{l} \chi_{l_{1}} \psi_{l_{3}} + \text{h.c.} + \dots$$

$$-\mathcal{L}_{IR} = \left(\frac{R}{R'}\right)^{4} \left[\sqrt{3} \frac{v'}{\Lambda'} \left(\chi_{l_{3}}^{\alpha} y_{b}^{\alpha} \psi_{l_{\alpha}} + \chi_{\tilde{l}_{3}}^{\alpha} y_{b}^{\alpha} \psi_{\tilde{l}_{\alpha}} + \chi_{e_{3}'}^{\alpha} y_{s}^{\alpha} \psi_{e_{\alpha}'}\right) + y_{b} \frac{v'_{\eta}}{\Lambda'} \left(\chi_{l_{1}} \psi_{l_{2}} + \chi_{\tilde{l}_{1}} \psi_{\tilde{l}_{2}}\right) + y_{s} \frac{v'_{\eta}}{\Lambda'} \chi_{\nu_{1}'} \psi_{\nu_{2}'}\right] + \text{h.c.} + \dots$$

 Z_2

 Z_3

 $\hat{\theta} = U_{\rm HPS} \begin{pmatrix} \epsilon_t + \epsilon_s & 0 & 0 \\ 0 & \epsilon_s & 0 \\ 0 & 0 & \epsilon_t - \epsilon_s \end{pmatrix} U_{\rm HPS}^{\dagger}$

After A₄ breaking

- Yukawas suppressed by $v'/\Lambda' \sim 0.1-0.01$
- Charged lepton Yukawas: diagonal and hierarchical
- Neutrino Yukawas: proportional to the identity
- The only source of mixing is the (UV localized) neutrino Majorana mass

 $\mathcal{M}_{M}^{
u} \propto U_{\mathrm{HPS}} egin{pmatrix} \epsilon_{t}+\epsilon_{s} & 0 & 0 \ 0 & \epsilon_{s} & 0 \ 0 & 0 & \epsilon_{t}-\epsilon_{s} \end{pmatrix} U_{\mathrm{HPS}}^{\dagger}$

• Scales:

Neutrino Majorana mass

 $M_M^{\nu} \sim R^{-1} \frac{\tilde{v}}{\Lambda} f_{R;UV}^2 \sim 10^{12-15} \text{ GeV}$

For moderately IR localized v_R

• Dirac (Yukawa) masses

 $M_D \sim v \frac{v'}{\Lambda'} f_{L;IR} \times f_{R;IR} \sim \begin{cases} 1 - 30 \text{ GeV}(\nu) \\ 10^{-4} - 1 \text{ GeV}(l) \end{cases}$

 u_R, au_R need a sizable IR localization

Leptonic spectrum in the MCHM₅

Neutrino masses generated through see-saw

 $M_{\rm eff}^{\nu} = -M_D^{\nu} (M_M^{\nu})^{-1} (M_D^{\nu})^{\rm T}$

 $\propto (M_M^{\nu})^{-1} = U_{\rm HPS} M_{\rm diag} U_{\rm HPS}^{\dagger}$

TBM mixing

- Charged lepton masses naturally hierarchical (suppressed with respect to v) and diagonal in current eingenstate basis: <u>No tree level LFV</u>
- Small corrections if fermion KK modes and nonlinear Higgs effects taken into account

Can higher dimension operators destabilize this pattern?

- Corrections to TBM mixing and LFV
- Easy to classify to all orders
 - $\langle \phi \rangle^3 \sim \langle \phi \rangle \qquad \qquad \langle \phi' \rangle^2 \sim 1 + \langle \phi' \rangle$
 - New structure for the Majorana mass
 - Diagonal but non-universal neutrino Yukawas

New in GHU!

- Non-diagonal charged lepton Yukawas

 As expected, there is more flavor violation in composite Higgs models than in simple RS

Csaki, Falkowski, Weiler 08

 But there is an extra flavor protection from custodial + LR symmetry + splitting in two sectors (neutrino and charged lepton sectors)

Agashe 09

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Motivated by Zbb protection

• Extra flavor protection: protected couplings to the Z

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 $\mu \to e\gamma$

• Fix the IR scale to 1.5 TeV ($m_{\kappa\kappa}$ ~3.5 TeV)

- Check constraints:
 - Lepton masses
 - EWPT:
 - Lepton mixing:
 - Tree level LFV:
 - One loop LFV:

 $\delta g/g \lesssim 0.2\%$ $U_{\rm PMNS}$ $\mu \to 3e, \quad \mu \leftrightarrow e, \dots$

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- Check constraints:
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Requires IR localized ν_R, τ_R $\delta g/g \lesssim 0.2\%$ $U_{\rm PMNS}$ $\mu \rightarrow 3e, \quad \mu \leftrightarrow e, \dots$ $\mu \rightarrow e\gamma$

Sizable but OK for $v/\Lambda \lesssim 0.1, c_3 \gtrsim 0.55$

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New light leptonic resonances!



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Conclusions

• Fully realistic composite Higgs model from 5D:

- Correct EWSB and quark spectrum
- Correct lepton spectrum with discrete symmetries
 - Hierarchical charged lepton masses
 - Correct neutrino masses and mixing
 - Double layer of flavor protection
 - A4 symmetry
 - Custodial + LR symmetry+ structure of MCHM5
- New quark and lepton resonances at the LHC!
- LFV close to observable limits