

Unitarity Triangle Analysis and Searches for New Physics

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- Status of the Standard Model Analysis
- UT beyond the SM and constraints on NP:
 - Minimal Flavour Violation: UUT, $B \rightarrow \tau \nu$, ...
 - Beyond MFV: NP scales, B_s mixing, ...
- Conclusions and Outlook

<http://www.utfit.org>

INTRODUCTION

- SM UT analysis:
 - provide the best determination of CKM parameters
 - test the consistency of the SM (“direct” vs “indirect” determinations)
 - provide predictions for future experiments (ex. $\sin 2\beta$, Δm_s , ...)

INTRODUCTION

- NP UT analysis: given some hypothesis on NP (no corr. to SM tree, MFV, NMFV, ...)
 - disentangle possible NP effects: CKM determination, mixing, decay amplitudes
 - provide best CKM determination
 - extract allowed range for NP contributions to $\Delta F=2$ & $\Delta F=1$
 - obtain bounds on NP parameters

CP-CONSERVING INPUT

- $|V_{ub}|/|V_{cb}| \sim R_b$ (tree-level)
 - incl $b \rightarrow cl\nu \Rightarrow |V_{cb}| = (41.54 \pm 0.44 \pm 0.58) 10^{-3}$,
 $b \rightarrow ul\nu \Rightarrow |V_{ub}| = (39.9 \pm 1.5 \pm 4.0) 10^{-4}$
(HFAG + flat error for model spread)
 - excl: $B \rightarrow D^{(*)}l\nu \Rightarrow |V_{cb}| = (39.0 \pm 0.9) 10^{-3}$,
 $b \rightarrow \pi(\rho)l\nu \Rightarrow |V_{ub}| = (35.0 \pm 4.0) 10^{-4}$
using LQCD form factors

Lubicz & Tarantino

CP-CONSERVING INPUT

- $|V_{td}|/|V_{cb}| \sim R_+$ from $B_d - \bar{B}_d$ and $B_s - \bar{B}_s$ mixing

(loop mediated):

$$- \Delta m_d = (0.507 \pm 0.005) \text{ ps}^{-1}$$

$$- \Delta m_s = (17.77 \pm 0.12) \text{ ps}^{-1}$$

$$- f_{B_s} \sqrt{B_{B_s}} = (275 \pm 13) \text{ MeV}$$

$$- \xi = 1.24 \pm 0.03$$

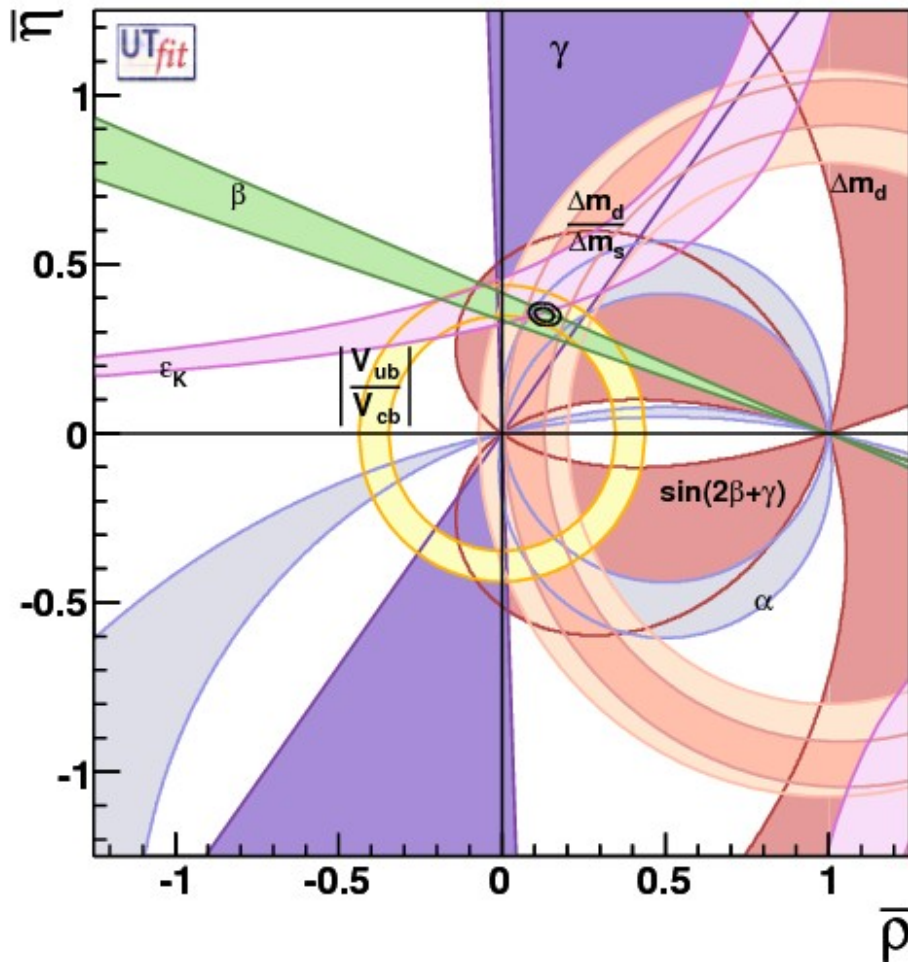
Laiho, Lunghi, V.d.Water

CP-VIOLATING INPUT

- $\sin 2\beta$ from $B \rightarrow J/\Psi K$ + th error from CPS:
 $\sin 2\beta = 0.655 \pm 0.027$
 - γ combined: $(72 \pm 11)^\circ \cup (-108 \pm 11)^\circ$
 - α combined: $(91 \pm 6)^\circ$
 - ε_K corrected for measured phase, $\text{Im } A_0$
and LD contributions
- Buras, Guadagnoli, Isidori
- $-F_K = 156.0 \pm 1.3 \text{ MeV}, B_K = 0.731 \pm 0.036$

Lubicz @ Lattice09

RESULTS OF THE SM FIT



$$\bar{\rho} = 0.130 \pm 0.020$$

$$\bar{\eta} = 0.352 \pm 0.013$$

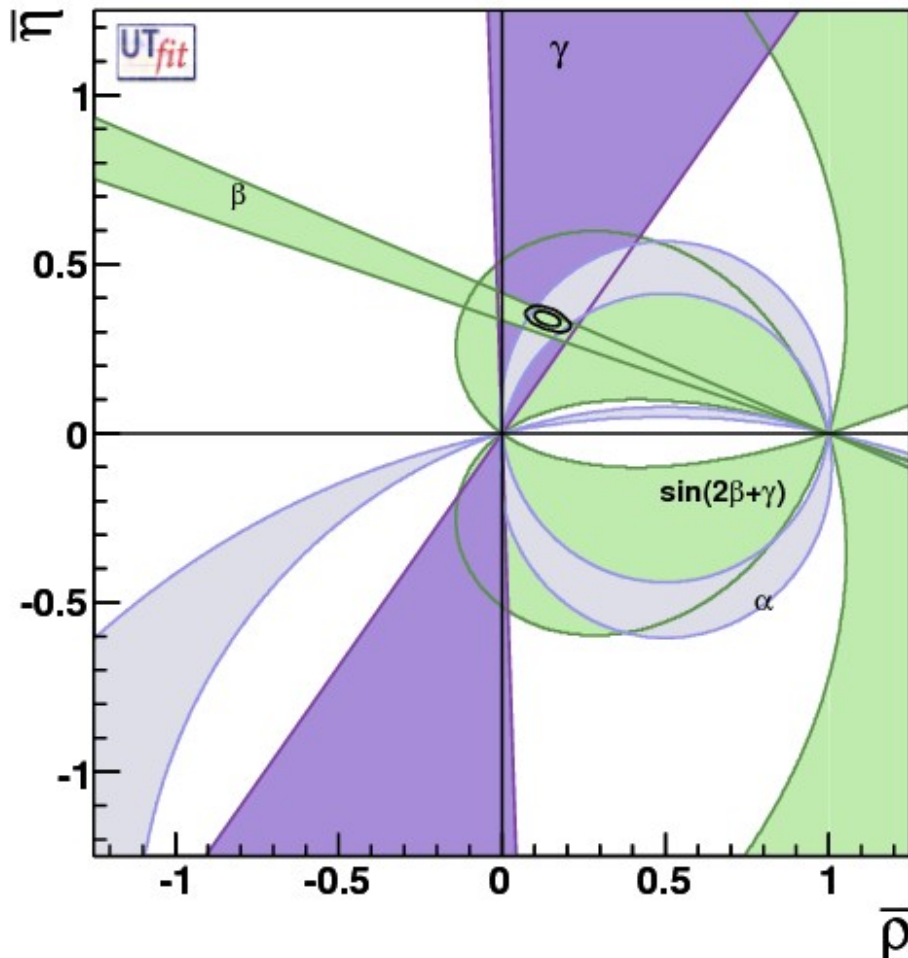
$$\beta = (22 \pm 1)^\circ$$

$$\gamma = (70 \pm 3)^\circ$$

$$\alpha = (88 \pm 3)^\circ$$

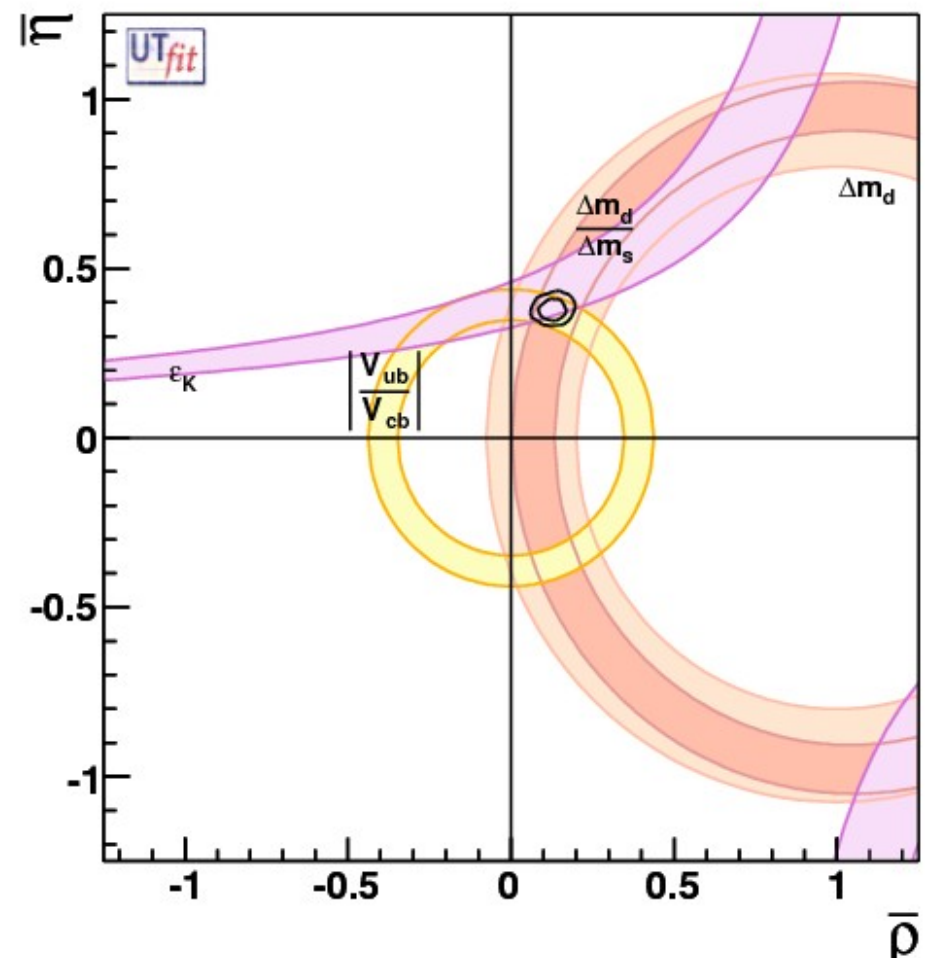
Here and in the following: **PRELIMINARY** winter 2010 results

ANGLES vs NON-ANGLES



$$\bar{\rho} = 0.139 \pm 0.029$$

$$\bar{\eta} = 0.338 \pm 0.015$$



$$\bar{\rho} = 0.129 \pm 0.028$$

$$\bar{\eta} = 0.380 \pm 0.021$$

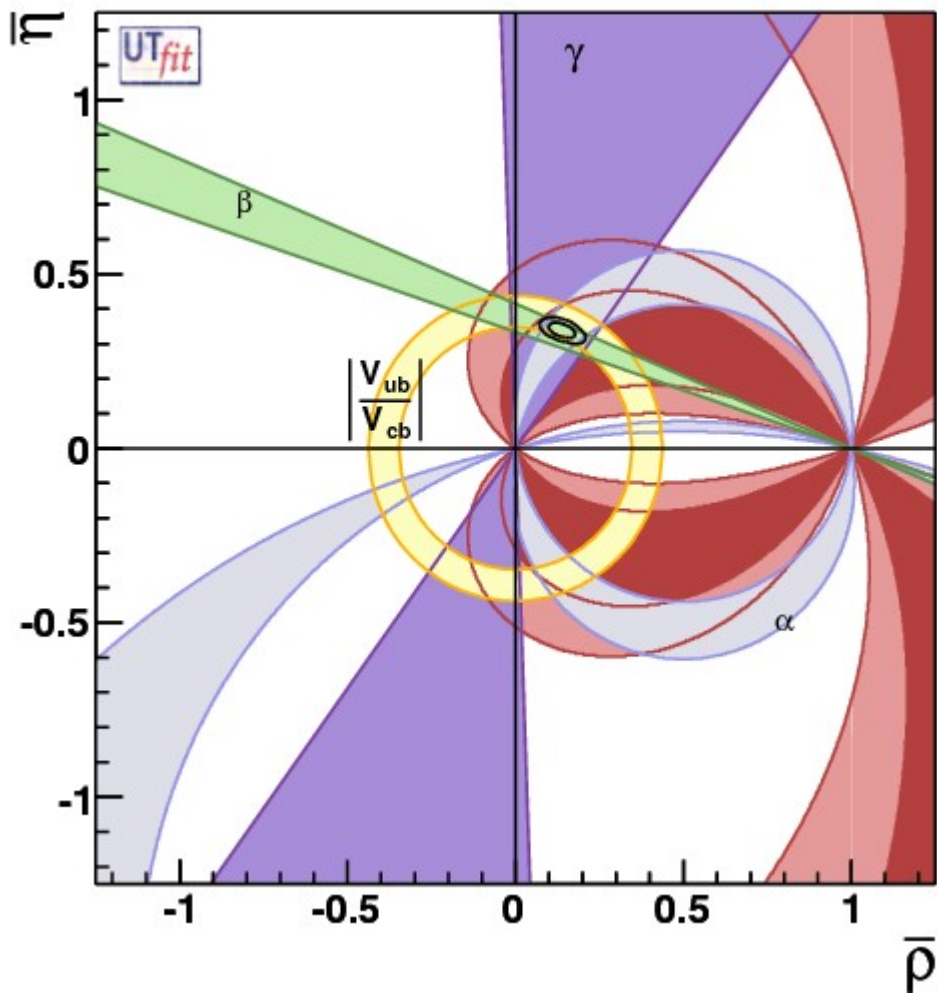
FIT PREDICTIONS vs INPUTS

| | Prediction | Measurement | Pull (σ) |
|-------------------------------------|---------------------------------|-----------------------------------|-------------------|
| $\sin 2\beta$ | 0.750 ± 0.033 | 0.655 ± 0.027 | 2.2 |
| α | $(86 \pm 4)^\circ$ | $(91 \pm 6)^\circ$ | <1 |
| γ | $(69 \pm 3)^\circ$ | $(74 \pm 11)^\circ$ | <1 |
| Δm_s | $(17.6 \pm 1.4) \text{ps}^{-1}$ | $(17.77 \pm 0.12) \text{ps}^{-1}$ | <1 |
| $ V_{ub} $ | $(36 \pm 1) 10^{-4}$ | $(38.8 \pm 2.2) 10^{-4}$ | 1.2 |
| B_K | (0.853 ± 0.070) | (0.731 ± 0.036) | 1.5 |
| $\text{BR}(B \rightarrow \tau \nu)$ | $(82 \pm 8) 10^{-6}$ | $(174 \pm 34) 10^{-6}$ | 2.6 |

THE UUT & MFV MODELS

- Consider MFV models ($Y_{u,d}$ only source of flavour & CPV) D'Ambrosio et al.,...
- Can define a Universal Unitarity Triangle using only observables unaffected by MFV-NP: R_b & angles Buras et al.
- UUT results starting point for model-dependent studies

UUT RESULTS



$$\bar{\rho} = 0.138 \pm 0.029$$

$$\bar{\eta} = 0.339 \pm 0.015$$

$$\beta = (22 \pm 1)^\circ$$

$$\gamma = (68 \pm 5)^\circ$$

$$\alpha = (90 \pm 5)^\circ$$

$B \rightarrow \tau \nu$ in MFV models

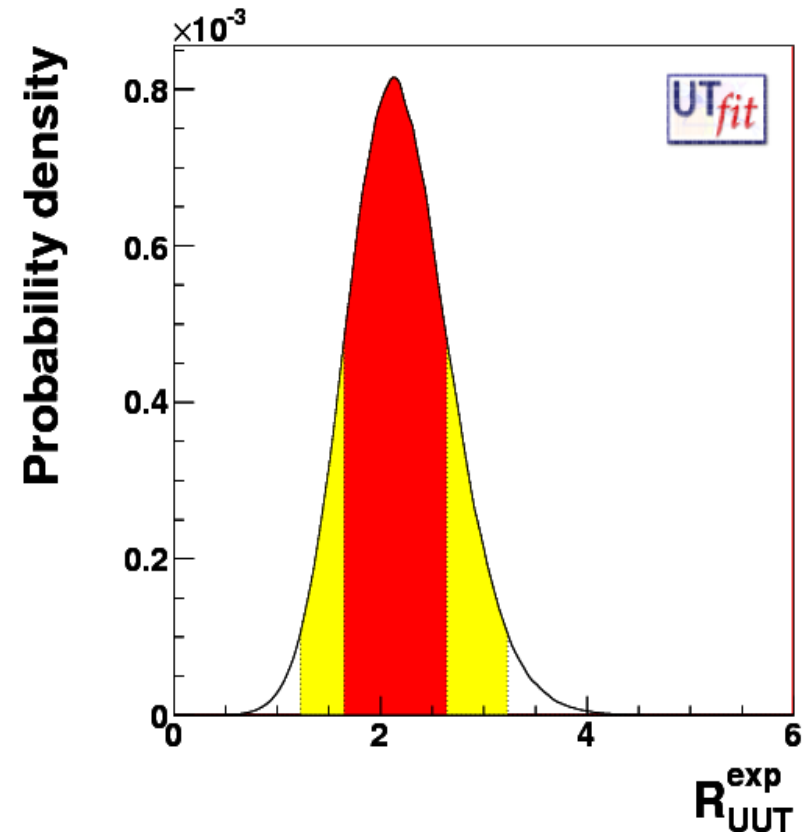
Define \overline{BR} as the prediction obtained assuming NO NP effect in the decay amplitude

We obtain:

$$R_{UUT}^{\text{exp}} = 2.1 \pm 0.5$$

where $R_{UUT}^{\text{exp}} = BR_{\text{exp}} / \overline{BR}_{UUT}$

to be compared with the $|V_{ub}|$ - and f_B -independent th. calculation of R_{UUT} in specific MFV models



Two Higgs Doublet Model II

$$R_{2\text{HDM}} = \left(1 - \tan^2 \beta \frac{m_B^2}{m_{H^+}^2} \right)^2 \rightarrow \text{bounds on } \tan\beta/m_{H^+}$$

Two regions selected:

1. small $\tan\beta/m_{H^+}$: $R < 1$ disfavoured at $\sim 2\sigma$

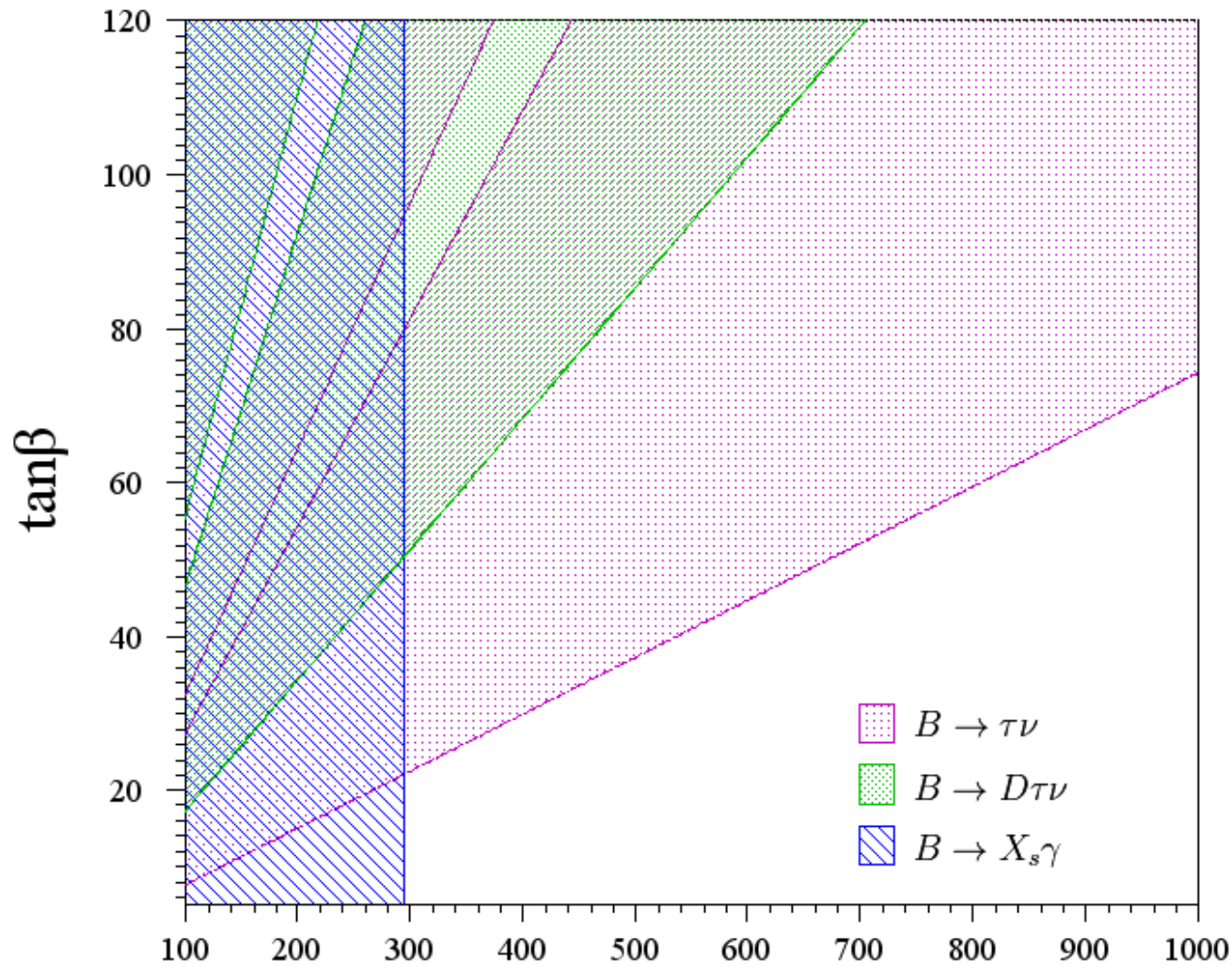
2. "fine-tuned" region for $\tan\beta/m_{H^+} \sim 0.3$:

positive correction, $R \sim R_{\text{exp}}$ can be obtained

incompatible with semileptonic decays

$\text{BR}(B \rightarrow D\tau\nu) / \text{BR}(B \rightarrow D\ell\nu) = (49 \pm 10)\%$ Iijima @ LP09
Kamenik & Mescia

$B \rightarrow X_s \gamma$ gives a lower bound on m_{H^+} : $m_{H^+} > 295 \text{ GeV}$



$$\tan\beta < 7.4 \frac{m_{H^+}}{100 \text{ GeV}}$$

$$BR(B_s \rightarrow \mu^+ \mu^-) = (4.3 \pm 0.9) \times 10^{-9} \quad ($$

$$([2.5, 6.2] \times 10^{-9} \text{ @95\% prob.})$$

MFV-MSSM at large $\tan\beta$

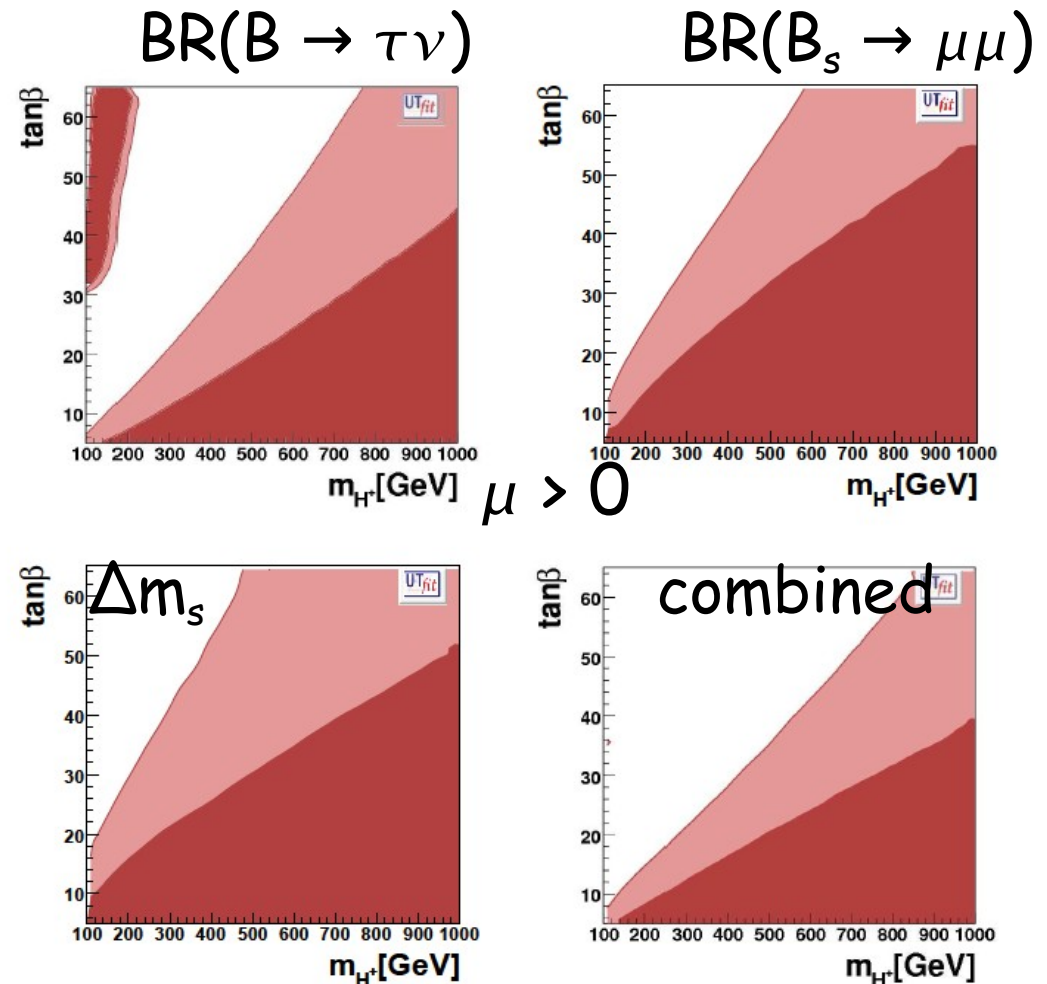
- * additional constraints exclude the "fine-tuned" region at very large $\tan\beta$
- * bound similar to 2HDM

$$\tan\beta < 7.3 m_{H^+} / (100 \text{ GeV})$$

In addition:

$$\text{BR}(B_s \rightarrow \mu\mu) < 26 \times 10^{-9}$$

@95% prob.



The case $\mu < 0$ is similar...

$$\epsilon_0 = -\frac{2\alpha_s\mu}{3\pi M_{\tilde{g}}} H_2 \left(\frac{M_{\tilde{q}_L}^2}{M_{\tilde{g}}^2}, \frac{M_{\tilde{d}_R}^2}{M_{\tilde{g}}^2} \right)$$

$$R_{B\tau\nu} = \left[1 - \left(\frac{m_B^2}{m_{H^\pm}^2} \right) \frac{\tan^2 \beta}{(1 + \epsilon_0 \tan \beta)} \right]^2$$

* for $\mu < 0$ the region of positive interference at very large $\tan\beta$ is enlarged

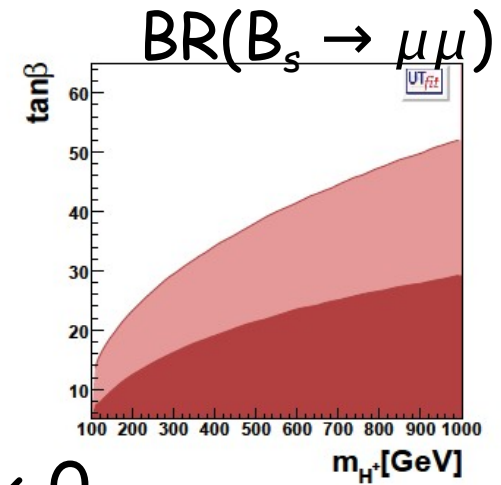
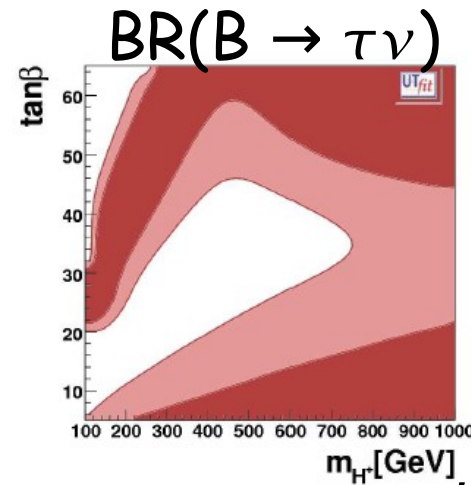
* yet the combined bound is stronger than for $\mu > 0$

$\tan\beta < 38$ @95% prob.

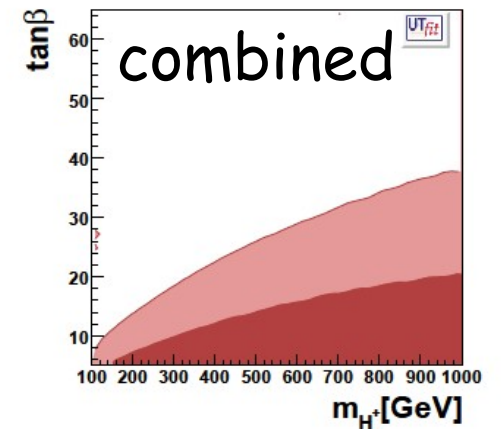
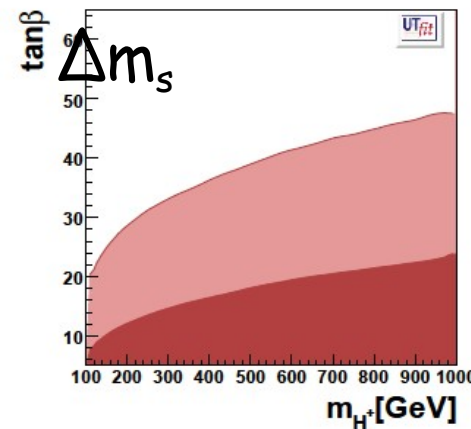
In this case:

$BR(B_s \rightarrow \mu\mu) < 17 \times 10^{-9}$

@95% prob.



$\mu < 0$



UTfit beyond MFV

1. fit simultaneously for the CKM and the NP parameters (generalized UT fit)

- add most general loop NP to all sectors

- use all available experimental info

- find out NP contributions to $\Delta F=2$ transitions

Soares, Wolfenstein; Deshpande, Dutta, Oh; Silva, Wolfenstein; Cohen et al.; Grossman, Nir, Worah; Laplace et al; Ciuchini et al; Ligeti; CKMFitter; UTfit; Botella et al.; Agashe et al.; ...

2. perform a $\Delta F=2$ EFT analysis to put bounds on the NP scale

- consider different choices of the FV and CPV couplings

1. Parameterization of generic NP contributions to the mixing amplitudes

K mixing amplitude (2 real parameters):

$$\text{Re} A_K = C_{\Delta m_K} \text{Re} A_K^{SM} \quad \text{Im} A_K = C_\varepsilon \text{Im} A_K^{SM}$$

B_d and B_s mixing amplitudes (2+2 real parameters):

$$A_q = C_{B_q} e^{2i\phi_{B_q}} A_q^{SM} e^{2i\phi_q^{SM}} = \left(1 + \frac{A_q^{NP}}{A_q^{SM}} e^{2i(\phi_q^{NP} - \phi_q^{SM})} \right) A_q^{SM} e^{2i\phi_q^{SM}}$$

Observables:

$$\begin{aligned} \Delta m_{q/K} &= C_{B_q/\Delta m_K} (\Delta m_{q/K})^{SM} & \varepsilon_K &= C_\varepsilon \varepsilon_K^{SM} \\ A_{CP}^{B_d \rightarrow J/\psi K_s} &= \sin 2(\beta + \phi_{B_d}) & A_{CP}^{B_s \rightarrow J/\psi \phi} &\sim \sin 2(-\beta_s + \phi_{B_s}) \\ A_{SL}^q &= \text{Im}(\Gamma_{12}^q / A_q) & \Delta \Gamma^q / \Delta m_q &= \text{Re}(\Gamma_{12}^q / A_q) \end{aligned}$$

A FEW REMARKS

- The value of $C_{\varepsilon K}$ extracted from the analysis potentially contains a mixture of $\Delta F=1$ & $\Delta F=2$ NP contributions

- For the B_s analysis, we use an improved theoretical prediction for $\Delta\Gamma$:

$$\Delta\Gamma_s/\Gamma_s = 0.14 \pm 0.02$$

Ciuchini et al., in preparation;
see also Lenz & Nierste

and allow for NP penguin effects in Γ_{12}

RESULTS OF GENERALIZED UTA

$$\bar{\rho} = 0.130 \pm 0.038$$

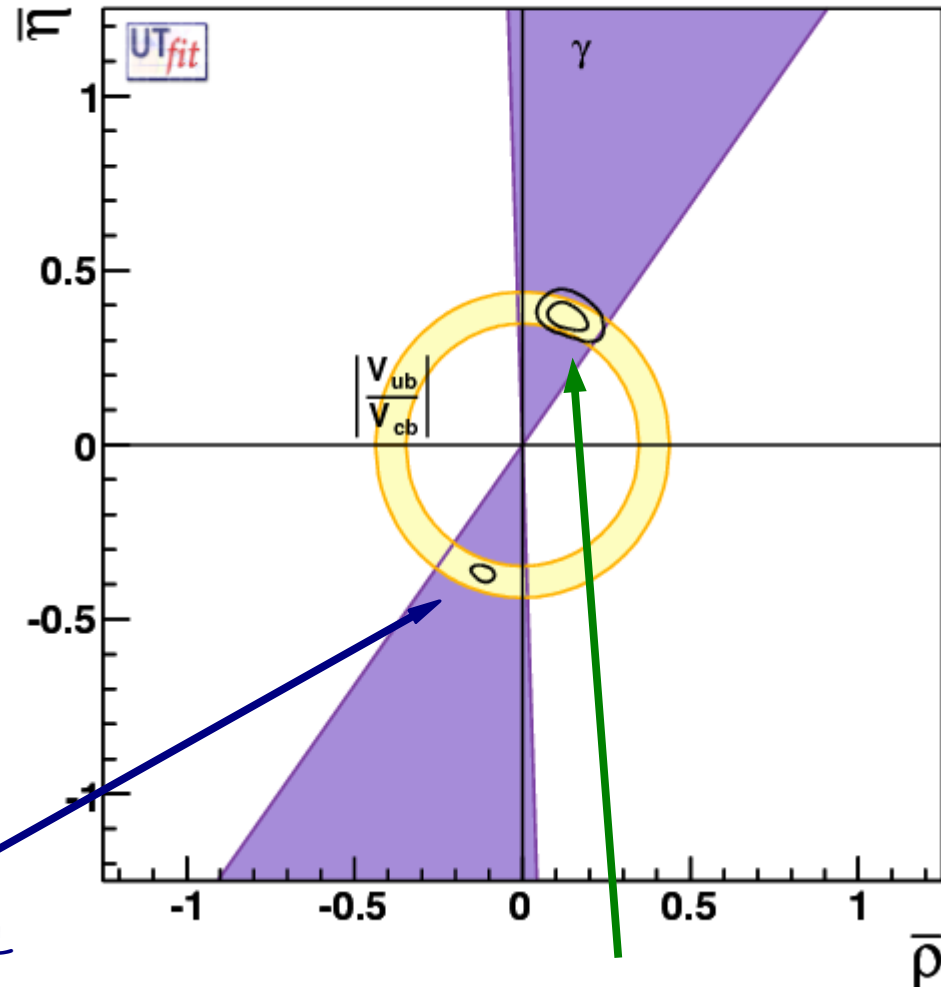
$$\bar{\eta} = 0.367 \pm 0.025$$

in the SM was:

$$\rho = 0.130 \pm 0.020$$

$$\eta = 0.352 \pm 0.013$$

degeneracy of γ broken by A_{SL}



Accuracy improved by α (assuming no huge NP contribution to EWP)

Results for NP parameters:

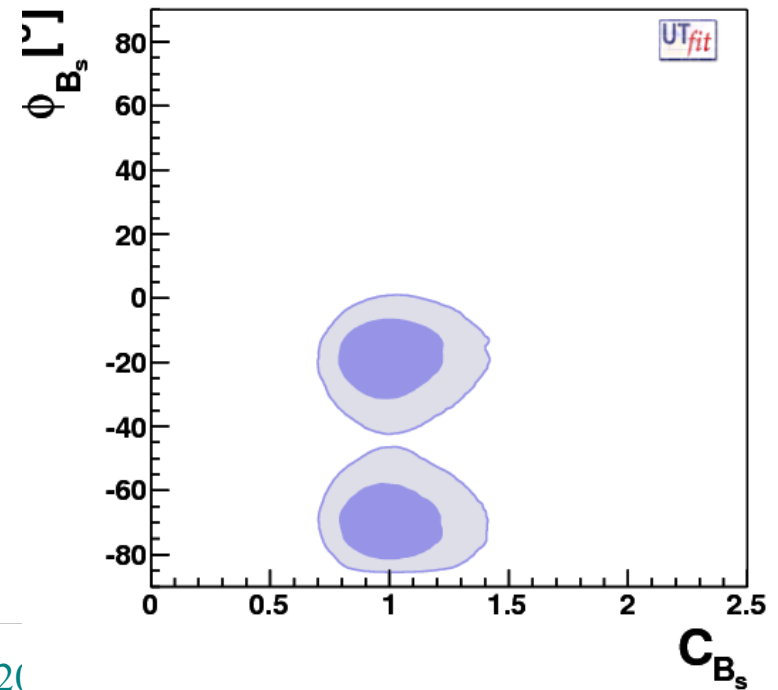
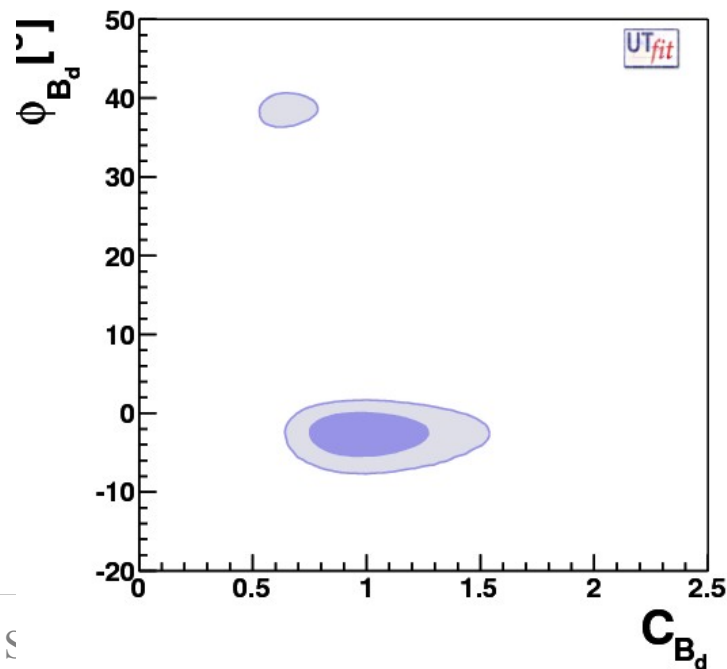
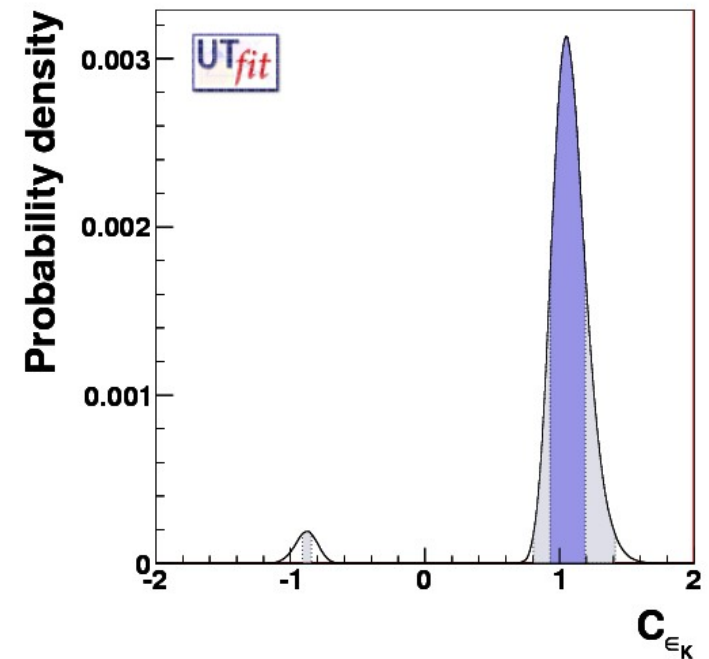
$$C_{\epsilon K} = 1.06 \pm 0.13 [0.81, 1.41]$$

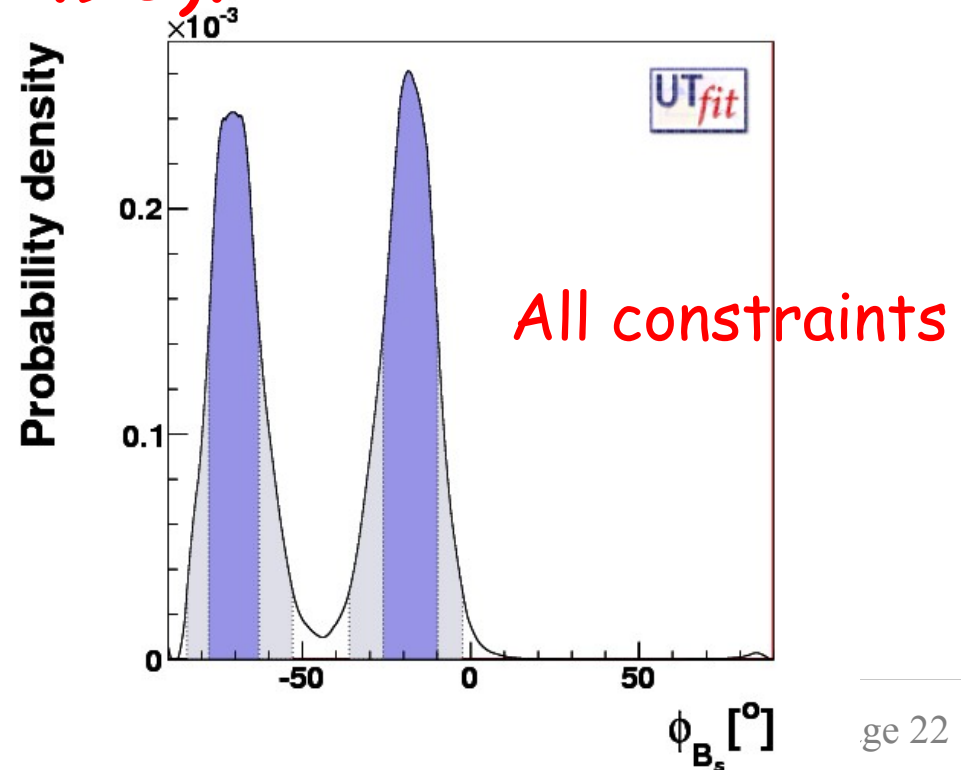
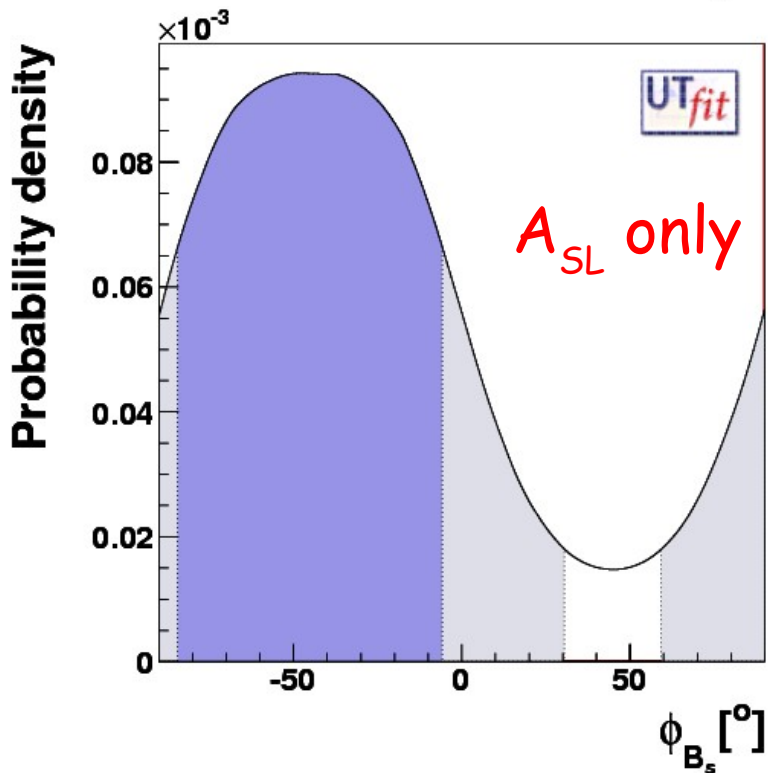
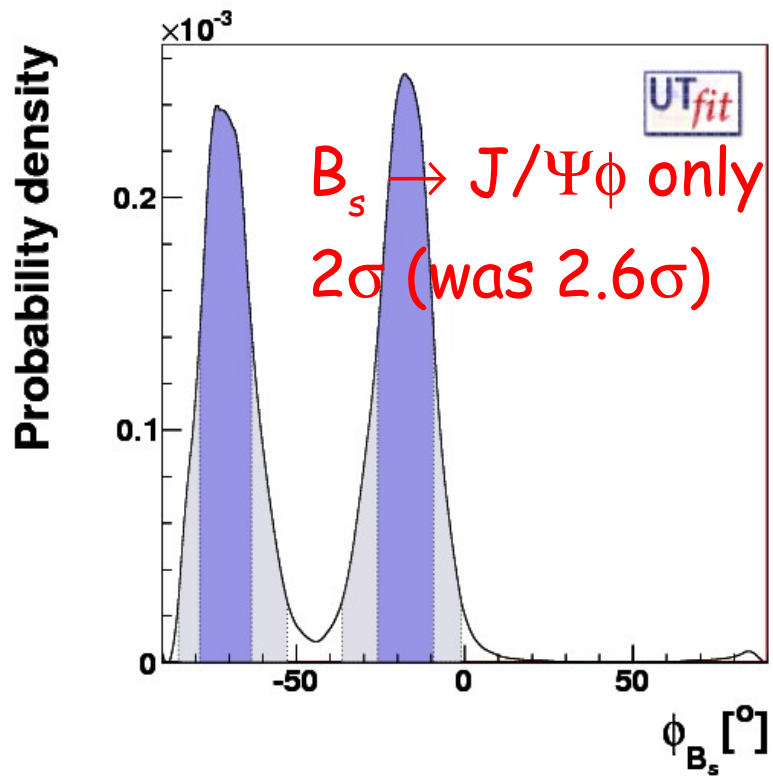
$$C_{B_d} = 0.99 \pm 0.18 [0.64, 1.41]$$

$$\phi_{B_d} = (-2.6 \pm 1.8)^\circ [-7.3, 1.4]^\circ$$

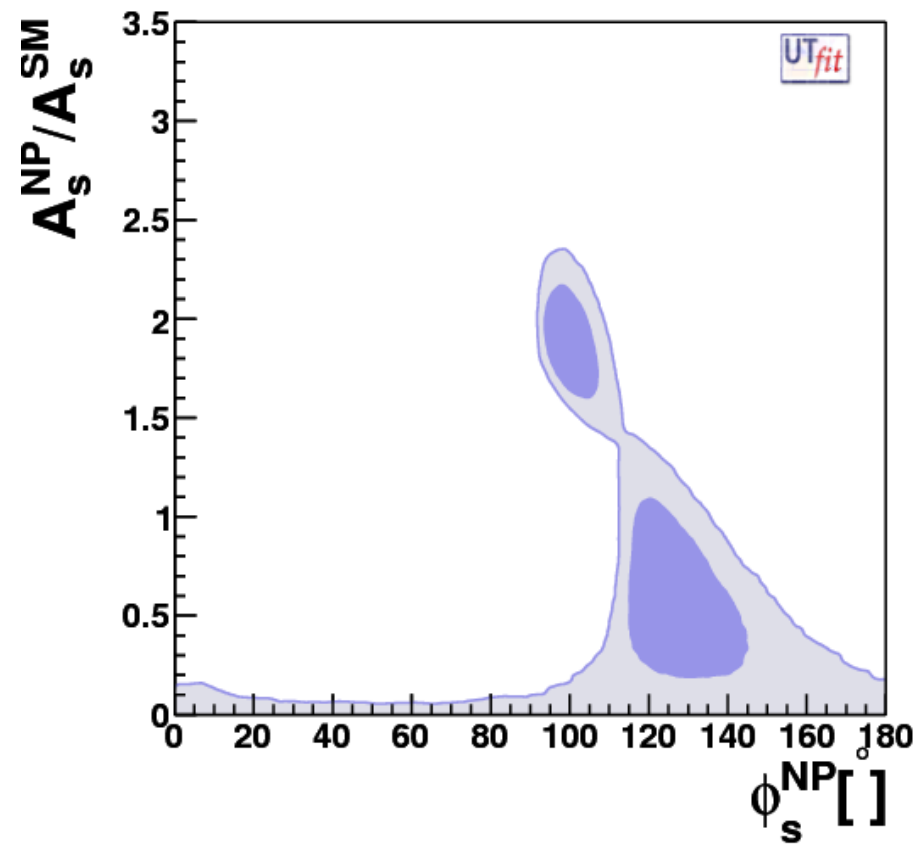
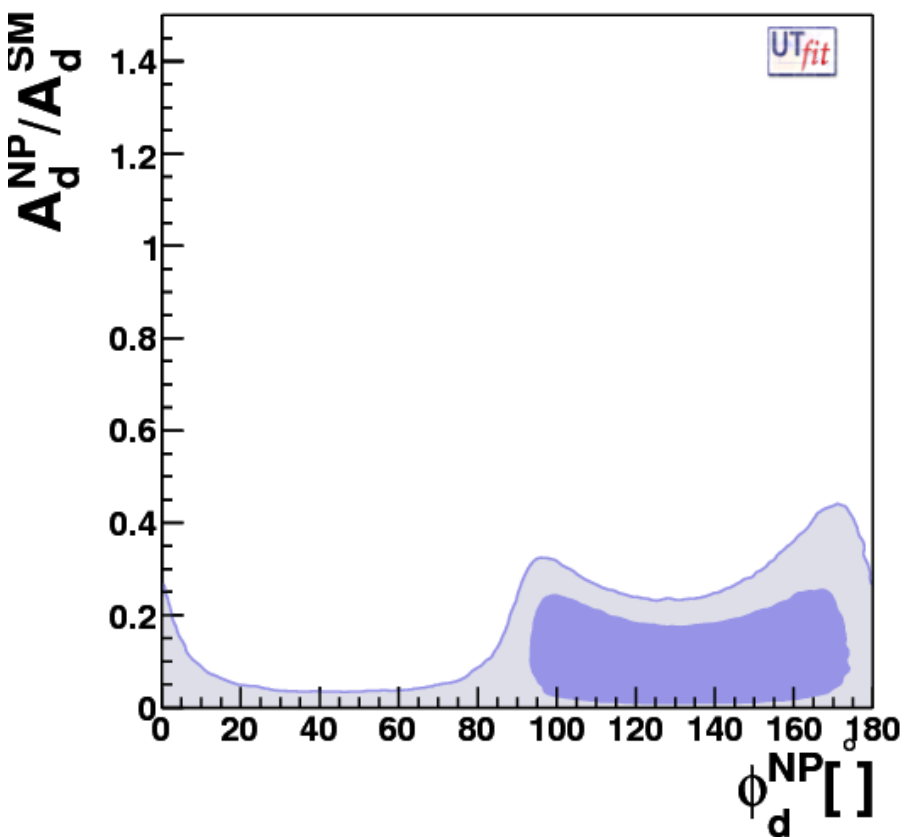
$$C_{B_s} = 1.00 \pm 0.14 [0.75, 1.33]$$

$$\phi_{B_s} = (-18 \pm 8)^\circ \cup (-70 \pm 8)^\circ [-36, -2]^\circ \cup [-85, -53]^\circ$$





We now use the combined Tevatron likelihood including frequentistic analysis of systematic errors (w. ~ 20 parameters varied at $\pm 5\sigma$). Using all data we are at 2.4σ (was 2.9σ).



Ratio of NP/SM contributions is $< 40\%$ @ 95% prob.
 in B_d mixing, and $\sim 60\%$ in B_s mixing (but 2σ range is very
 large)

See also Lunghi & Soni, Buras et al.

- Large NP contributions to $b \leftrightarrow s$ transitions are natural in nonabelian flavour models, given the large breaking of flavour $SU(3)$ due to the top quark mass

Pomarol, Tommasini; Barbieri, Dvali, Hall; Barbieri, Hall; Barbieri, Hall, Romanino; Berezhiani, Rossi; Masiero et al; ...

- GUTs can naturally connect the large mixing in ν oscillations with a large $b \leftrightarrow s$ mixing

Baek et al.; Moroi; Akama et al.; Chang, Masiero, Murayama; Hisano, Shimizu; Goto et al.; ...

- Might show up also in $\Delta F=1$ transitions ($b \rightarrow s\gamma$, $b \rightarrow sl^+l^-$, $B \rightarrow K\pi$, $B_s \rightarrow K^{*0}K^{*0}$, ...) and/or LFV ($\tau \rightarrow \mu\gamma$, $\mu \rightarrow e\gamma$)

2. EFT analysis of $\Delta F=2$ transitions

The mixing amplitudes $A_q e^{2i\phi_q} = \left\langle \bar{M}_q \left| H_{eff}^{\Delta F=2} \right| M_q \right\rangle$

$$H_{eff}^{\Delta B=2} = \sum_{i=1}^5 C_i(\mu) Q_i(\mu) + \sum_{i=1}^3 \tilde{C}_i(\mu) \tilde{Q}_i(\mu)$$

$$Q_1 = \bar{q}_L^\alpha \gamma_\mu b_L^\alpha \bar{q}_L^\beta \gamma^\mu b_L^\beta \quad (\text{SM/MFV})$$

$$Q_2 = \bar{q}_R^\alpha b_L^\alpha \bar{q}_R^\beta b_L^\beta$$

$$Q_3 = \bar{q}_R^\alpha b_L^\beta \bar{q}_R^\beta b_L^\beta$$

$$Q_4 = \bar{q}_R^\alpha b_L^\alpha \bar{q}_L^\beta b_R^\beta$$

$$Q_5 = \bar{q}_R^\alpha b_L^\beta \bar{q}_L^\beta b_R^\beta$$

$$\tilde{Q}_1 = \bar{q}_R^\alpha \gamma_\mu b_R^\alpha \bar{q}_R^\beta \gamma^\mu b_R^\beta$$

$$\tilde{Q}_2 = \bar{q}_L^\alpha b_R^\alpha \bar{q}_L^\beta b_R^\beta$$

$$\tilde{Q}_3 = \bar{q}_L^\alpha b_R^\beta \bar{q}_L^\beta b_R^\beta$$

7 new operators beyond MFV involving quarks with different chiralities

H_{eff} can be recast in terms of the $C_i(\Lambda)$ computed at the NP scale

- $C_i(\Lambda)$ can be extracted from the data (one by one)
- the associated NP scale Λ can be defined from

$$C_i(\Lambda) = \frac{LF_i}{\Lambda^2}$$

tree/strong interact. NP: $L \sim 1$
 perturbative NP: $L \sim a_s^2, a_w^2$

Flavour structures:

MFV

- $F_1 = F_{SM} \sim (V_{tq} V_{tb}^*)^2$
- $F_{i \neq 1} = 0$

next-to-MFV

- $|F_i| \sim F_{SM}$
- arbitrary phases

generic

- $|F_i| \sim 1$
- arbitrary phases

present lower bound on the NP scale (TeV):

| Process | C_4 (GeV ⁻²) | Λ_{GEN} (TeV) | Λ_{NMFV} (TeV) |
|--------------|----------------------------|------------------------------|-------------------------------|
| ϵ_K | $1.8 \cdot 10^{-17}$ | $24 \cdot 10^4$ | 62 |
| B_d | $1. \cdot 10^{-13}$ | $3.2 \cdot 10^3$ | 20 |
| B_s | $1.5 \cdot 10^{-11}$ | $2.6 \cdot 10^2$ | 8 |

- * $\Delta F=2$ chirality-flipping operators are RG enhanced and thus probe larger NP scales
- * when these operators are allowed, the NP scale is easily pushed beyond the LHC reach
- * suppression of the $1 \leftrightarrow 2$ transitions strongly weakens the lower bounds

CONCLUSIONS

- SM analysis displays good overall consistency but some tension in $\sin 2\beta$ and $B \rightarrow \tau \nu$
- The two tensions pull $|V_{ub}|$ in opposite directions
- Models predicting a suppression of $B \rightarrow \tau \nu$ disfavored by present data: 2HDM & MFV-MSSM @ large $\tan\beta$

CONCLUSIONS

- General UTA provides a precise determination of CKM parameters and NP contributions to $\Delta F=2$ amplitudes
- ε_K and B_d mixing give strong constraints on NP contributions, naively pushing the NP scale of several models beyond the LHC reach

CONCLUSIONS

- CPV in B_s mixing off from SM at $\sim 2.5\sigma$
- Requires new sources of flavour & CPV, natural in many extensions of the SM
- Theoretical uncertainty negligible now (and in the future if central value confirmed)
- Wait for confirmation from Tevatron/LHCb, look for other NP signals in $b \rightarrow s$ transitions