# **Axion Cold Dark Matter in Standard** and Non-Standard Cosmologies

Paolo Gondolo University of Utah

Visinelli, Gondolo, arxiv:0903.4377, Phys. Rev. D 80, 035024 (2009) Visinelli, Gondolo, arxiv:09 | 2.00 | 5, Phys. Rev. D 8 | , 063508 (20 | 0)



Luca Visinelli

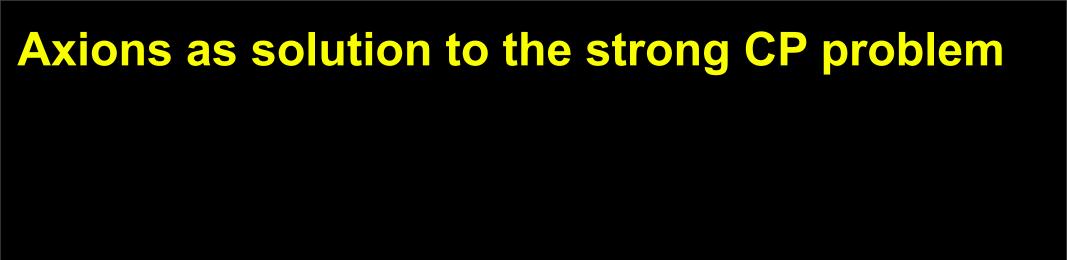
## **Axion cold dark matter**

### When are axions 100% of cold dark matter?

Study axion parameter space imposing

$$\Omega_a = \Omega_{\text{CDM}} = 0.1131 + 0.0034$$

And update cosmological constraints and include anharmonicities



The strong CP problem

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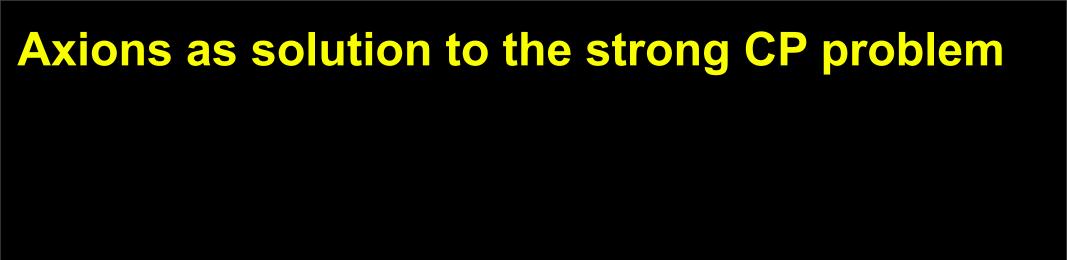
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Why  $\theta$  should be so small is the strong CP problem



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Introducing a U(I)<sub>PQ</sub> symmetry replaces  $\theta_{\text{total}} = \theta + \arg \det M_{\text{quark}} \quad \Rightarrow \quad \theta(x) = a(x)/f_a$ 

static CP-violating angle

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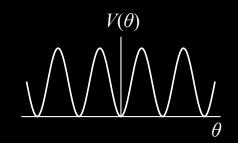
$$heta_{
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Before QCD phase transition,  $\langle \theta \rangle$  can be anything

After QCD phase transition, instanton effects generate

$$V(\theta) = m_a^2 f_a^2 (1 - \cos \theta) \label{eq:Vtheta}$$
 and  $\langle \theta \rangle = 0$  dynamically



## **Axions as dark matter**

Hot

Cold

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Produced thermally in early universe

Important for  $m_a > 0.1 eV$  ( $f_a < 10^8$ ), mostly excluded by astrophysics

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#### Cold

Produced by coherent field oscillations around mimimum of  $V(\theta)$  (Vacuum realignment)

Produced by decay of topological defects (Axionic string decays)

## Axion cold dark matter parameter space

 $f_a$  Peccei-Quinn symmetry breaking scale

N Peccei-Quinn color anomaly

 $N_d$  Number of degenerate QCD vacua

Kim-Shifman-Vainshtain-Zakharov Dine-Fischler-Srednicki-Zhitnistki

Couplings to quarks, leptons, and photons

 $H_{\rm I}$  Expansion rate at end of inflation

 $heta_i$  Initial misalignment angle

Harari-Sikivie-Hagmann-Chang Davis-Battye-Shellard

Axionic string parameters

Assume  $N = N_d = 1$  and show results for KSVZ and HSHC string network

Thus 3 free parameters  $f_a$ ,  $\theta_i$ ,  $H_{\rm I}$  and one constraint  $\Omega_a = \Omega_{\rm CDM}$ 

## Cold axion production in cosmology

## Vacuum realignment

- ullet Initial misalignment angle  $heta_i$
- Coherent axion oscillations start at temperature  $T_1$

$$3H(T_1)=m(T_1)$$

Hubble expansion parameter non-standard expansion histories differ in the function H(T)

T-dependent axion mass axions acquire mass through instanton effects at  $T < \Lambda \approx \Lambda_{\rm QCD}$ 

• Density at  $T_1$  is  $n_a(T_1) = \frac{1}{2} m_a(T_1) f_a^2 \chi \langle \theta_i^2 f(\theta_i) \rangle$ 

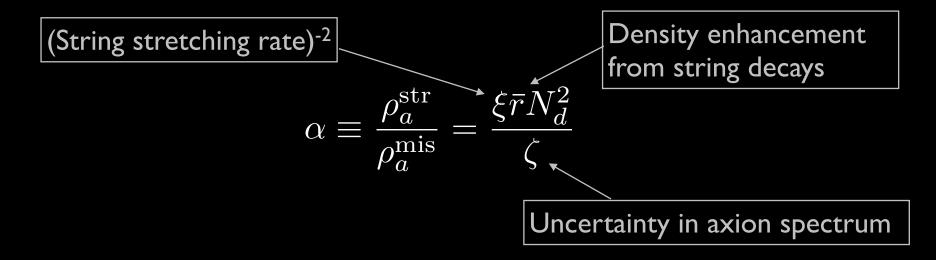
Anharmonicity correction  $f(\theta)$  axion field equation has anharmonic terms  $\ddot{\theta}+3H(T)\dot{\theta}+m_a^2(T)\sin\theta=0$ 

ullet Conservation of comoving axion number gives present density  $\Omega_a$ 

## Cold axion production in cosmology

## Axionic string decays

Energy density ratio (string decay/misalignment)



Slow oscillating strings (Davis-Battye-Shellard)

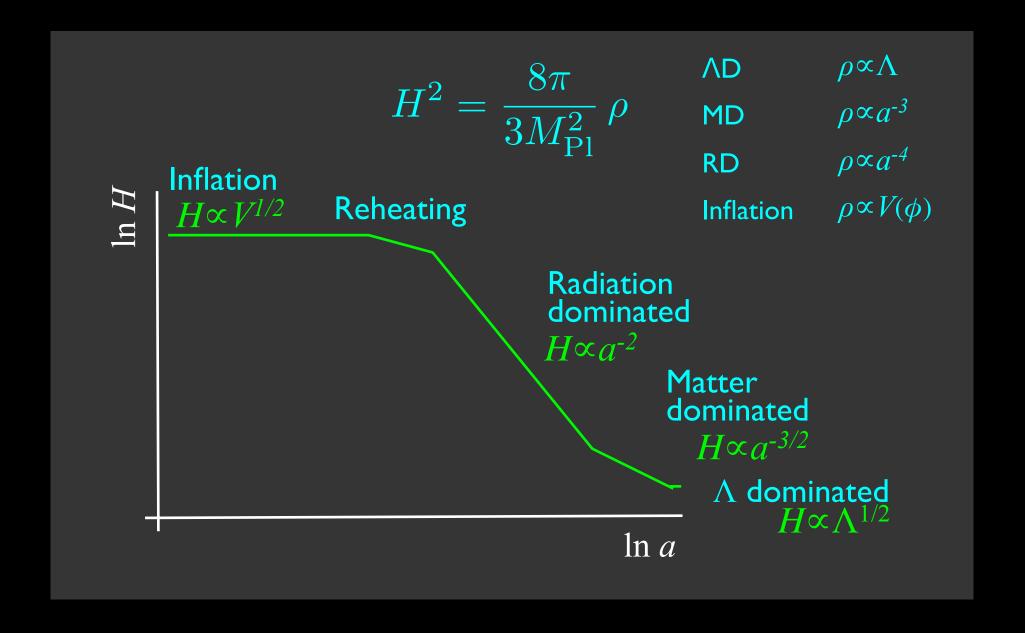
$$\bar{r} = \frac{1-\beta}{3\beta-1} \ln(t_1/\delta)$$

Fast-oscillating strings (Harari-Hagmann-Chang-Sikivie)  $\bar{r} = \frac{1-\beta}{3\beta-1} 0.8$ 

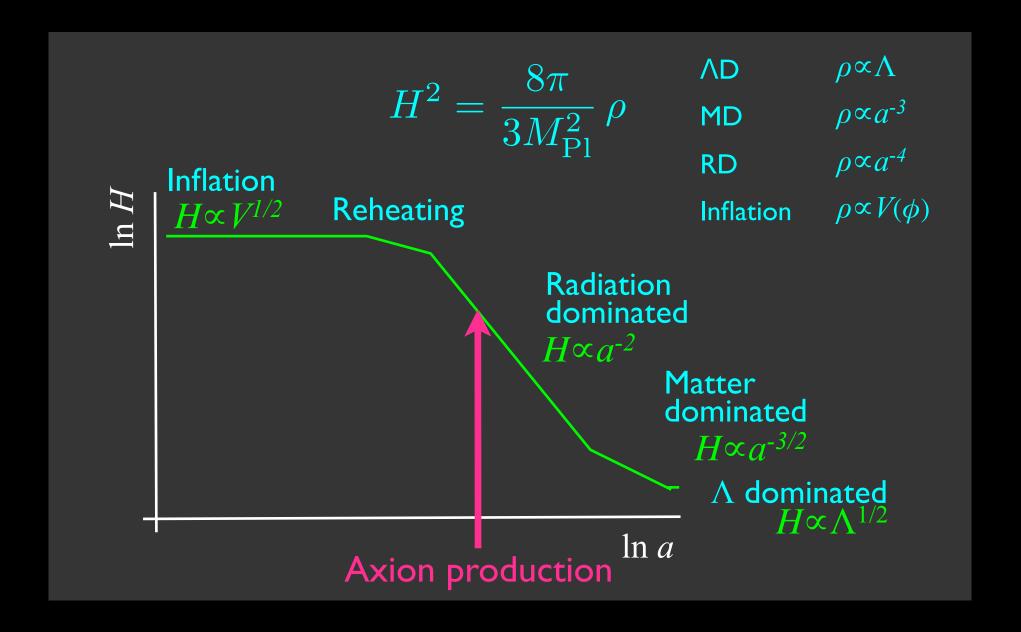
$$\bar{r} = \frac{1-\beta}{3\beta - 1} \, 0.8$$

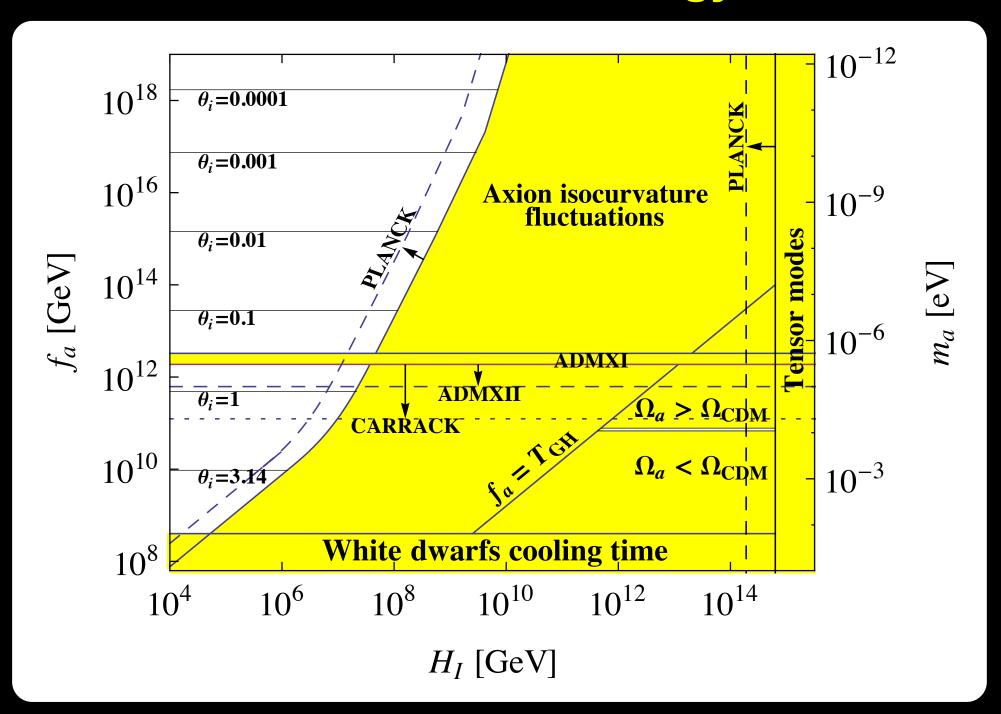
with  $a(t) \propto t^{\beta}$ 

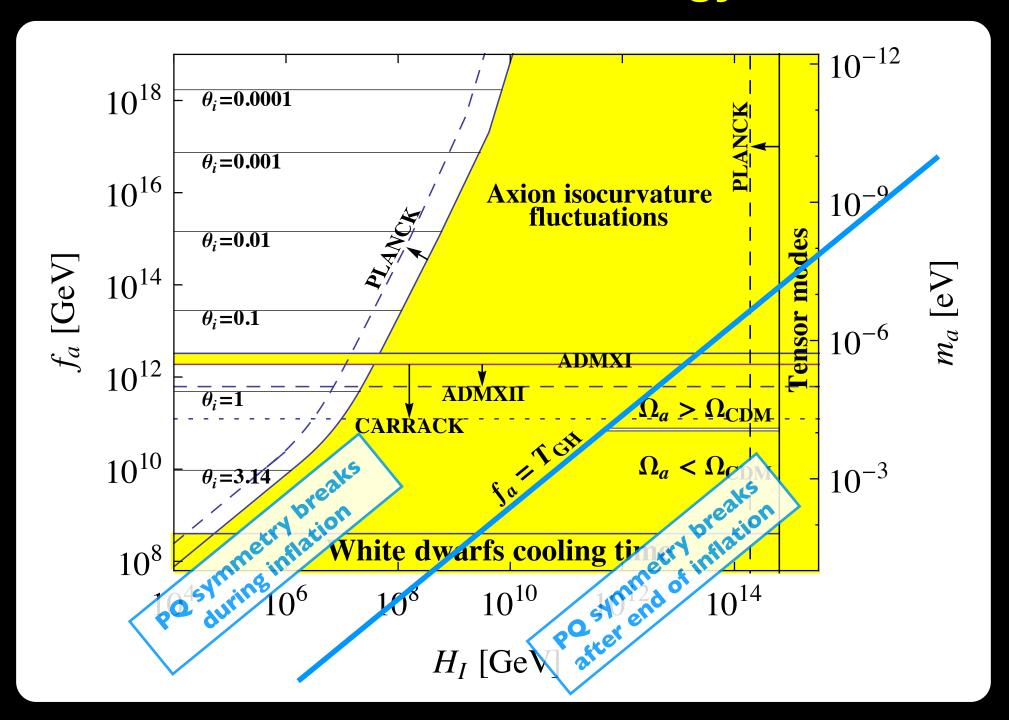
## **Standard cosmology**



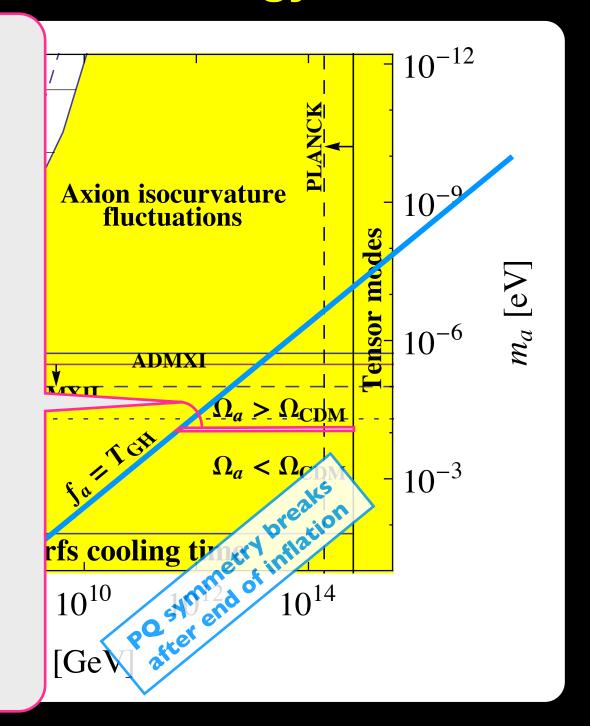
# Standard cosmology





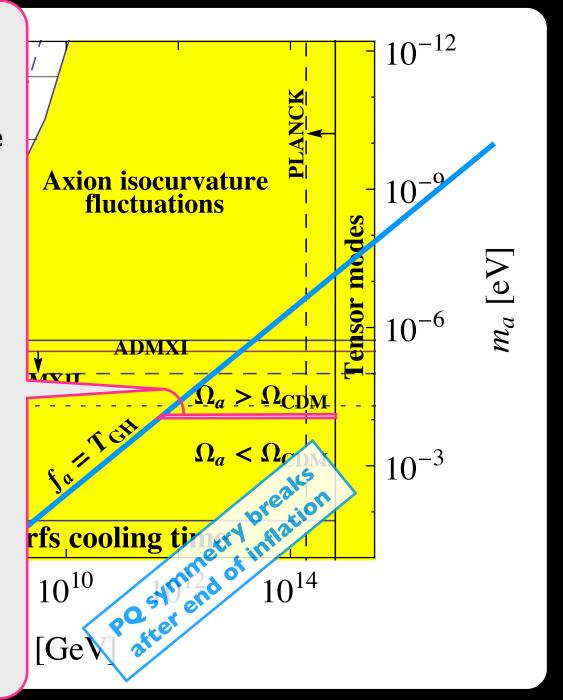


PQ symmetry breaks after end of inflation



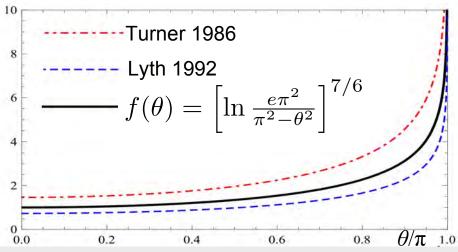
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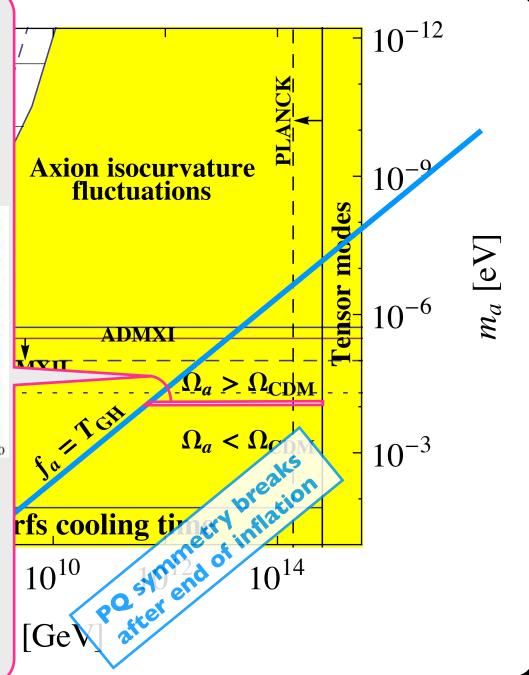
• Average  $\theta_i$  over Hubble volume



# PQ symmetry breaks after end of inflation

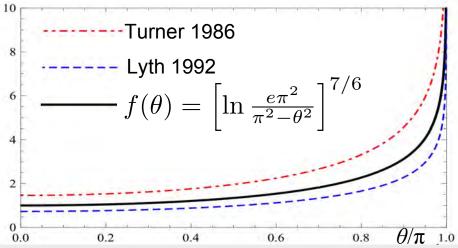
- Average  $\theta_i$  over Hubble volume
- Anharmonicities are important



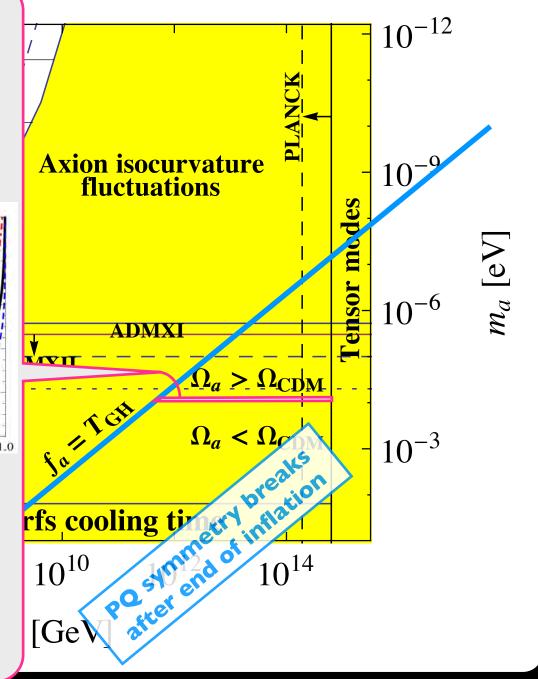


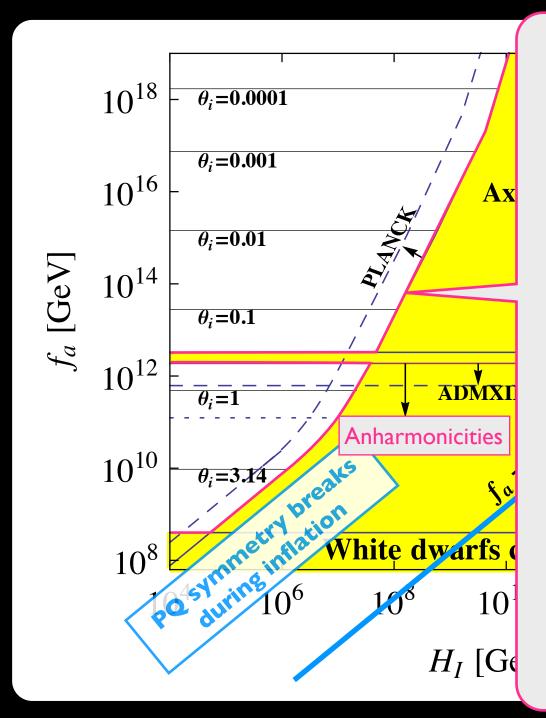
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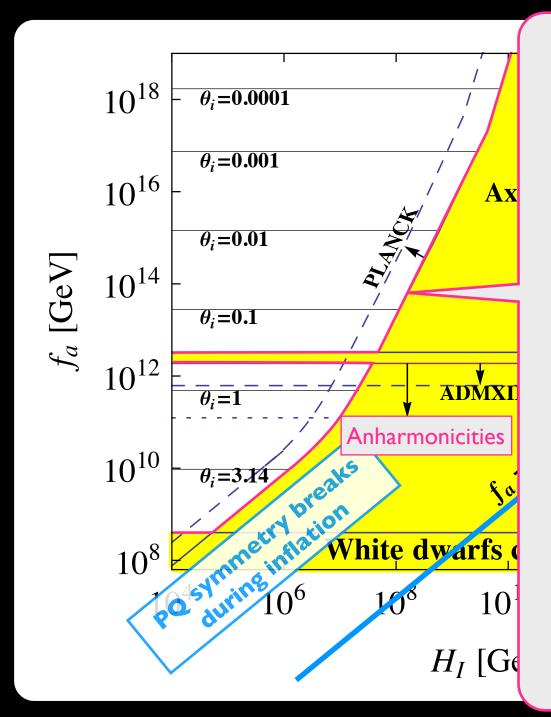


String decay contribution is~16% of vacuum realignment



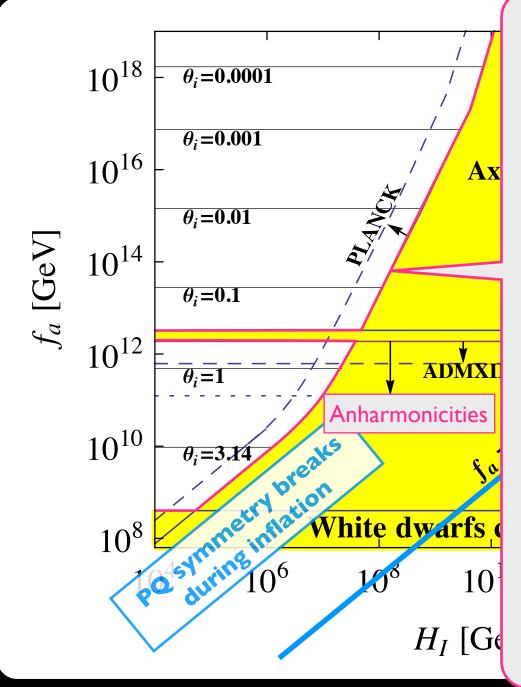


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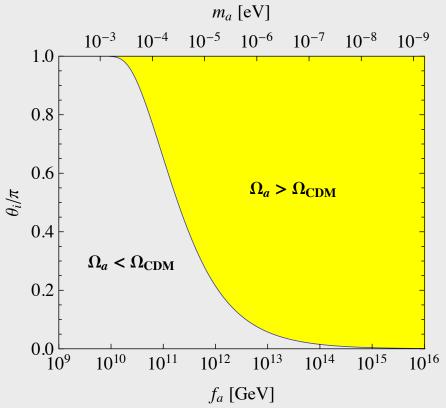
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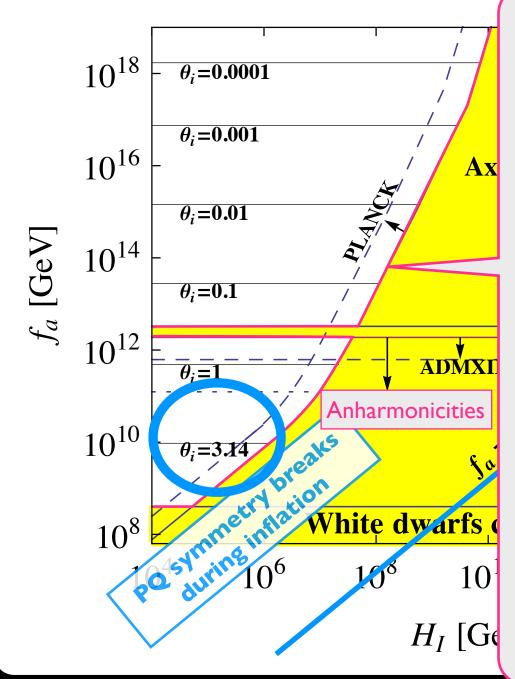
 Constrained by non-adiabatic fluctuations



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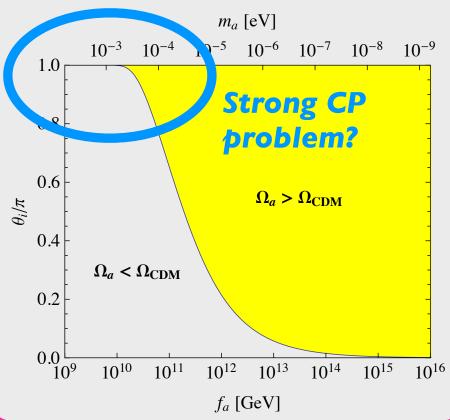
- Constrained by non-adiabatic fluctuations
- Single value of  $\theta_i$  throughout Hubble volume

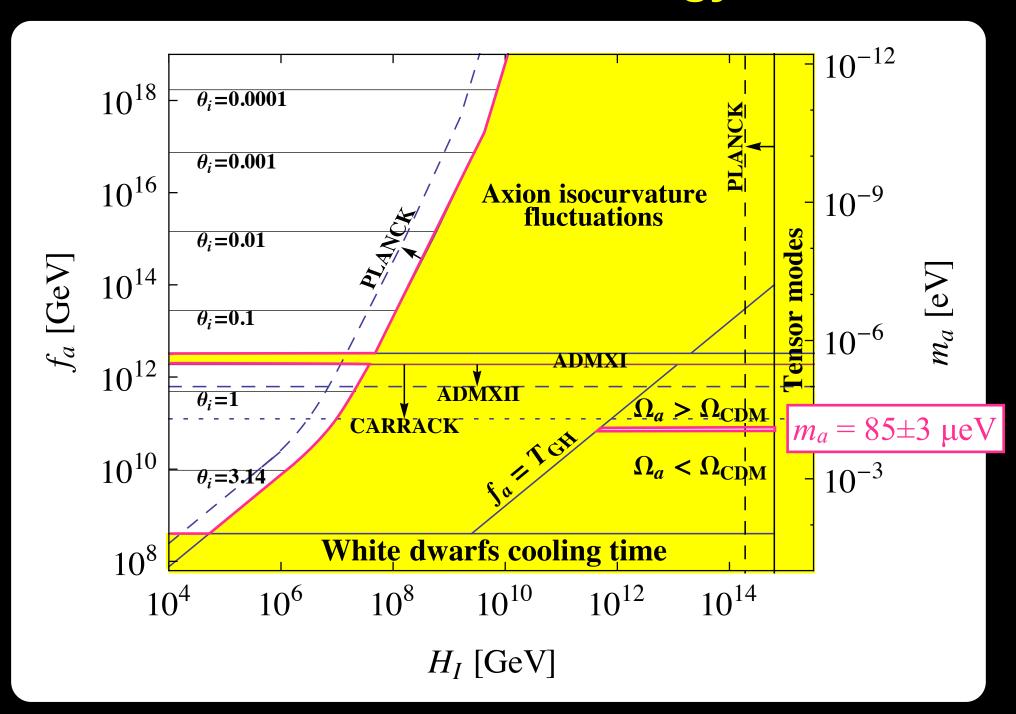




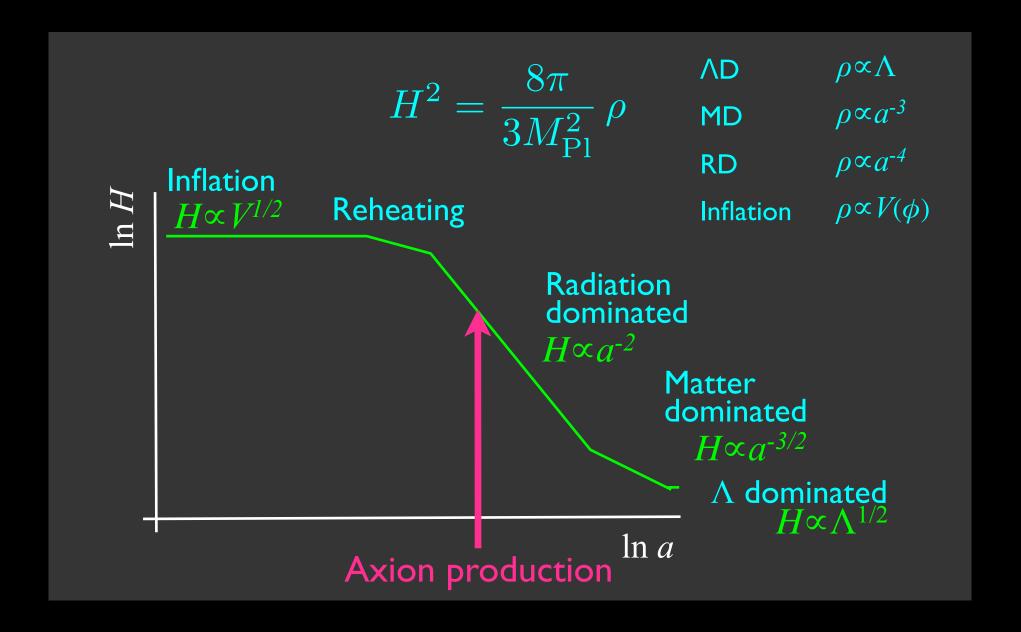
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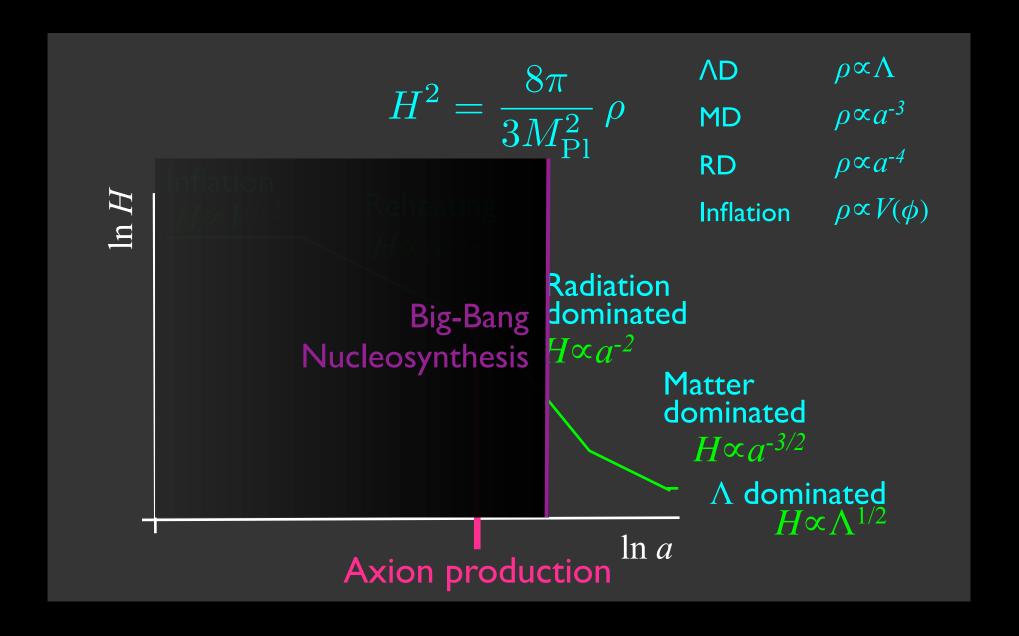




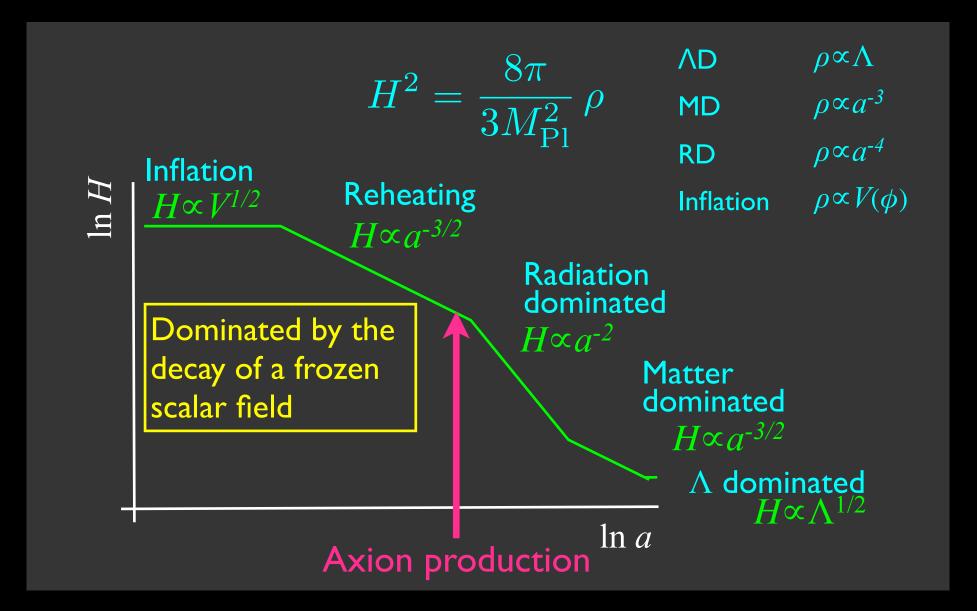
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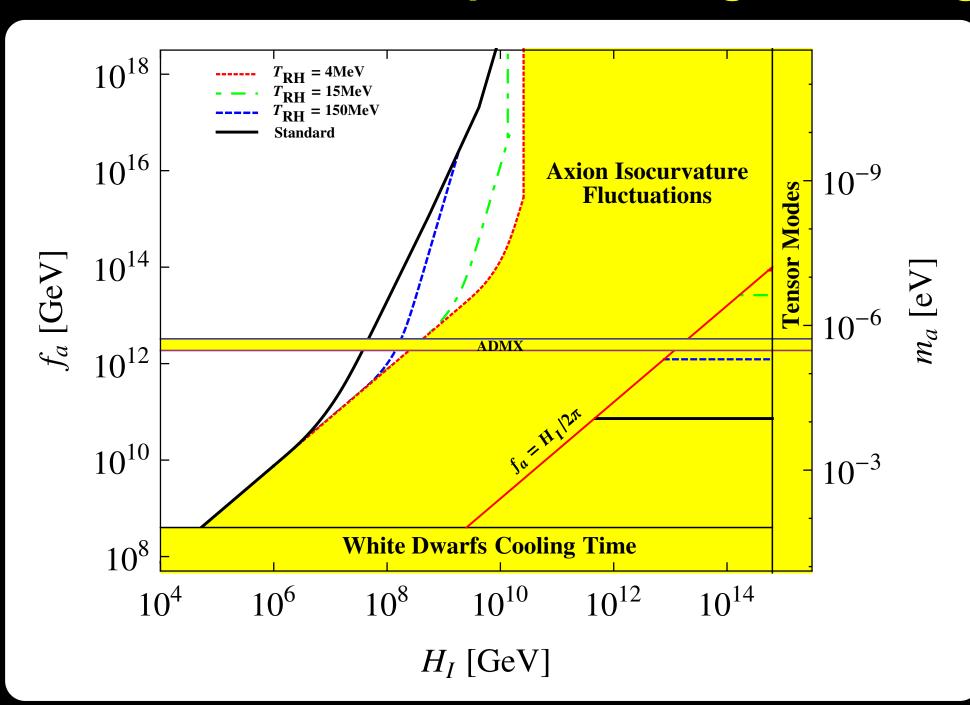
#### Non-standard cosmology

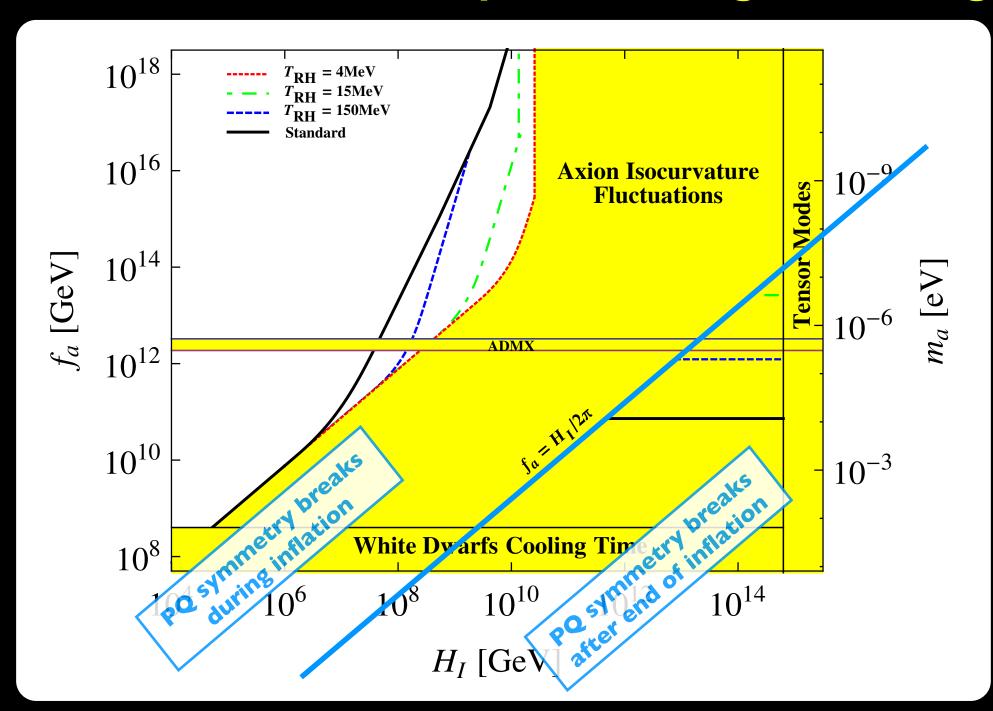


#### Low Temperature Reheating cosmology



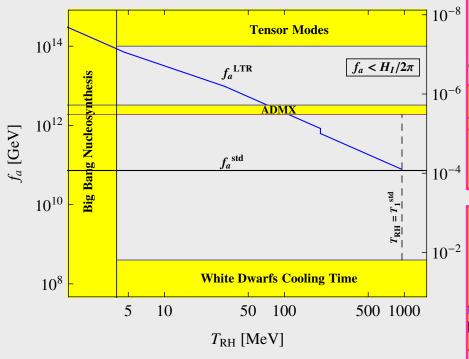
Turner 1983, Scherrer, Turner 1983, Dine, Fischler 1983

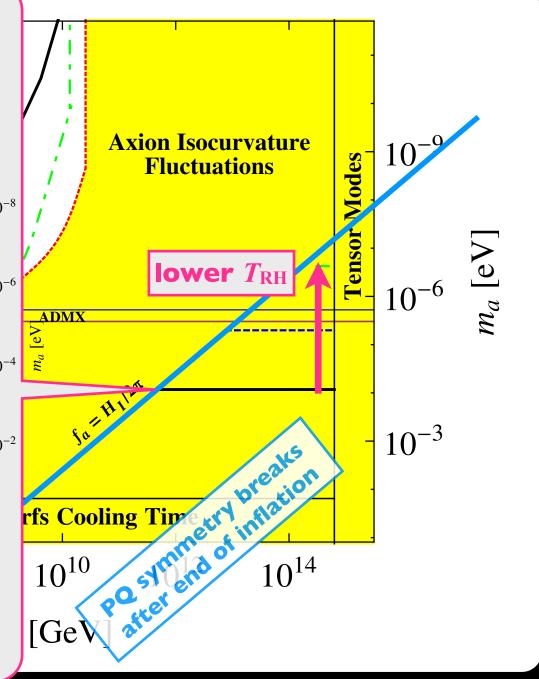


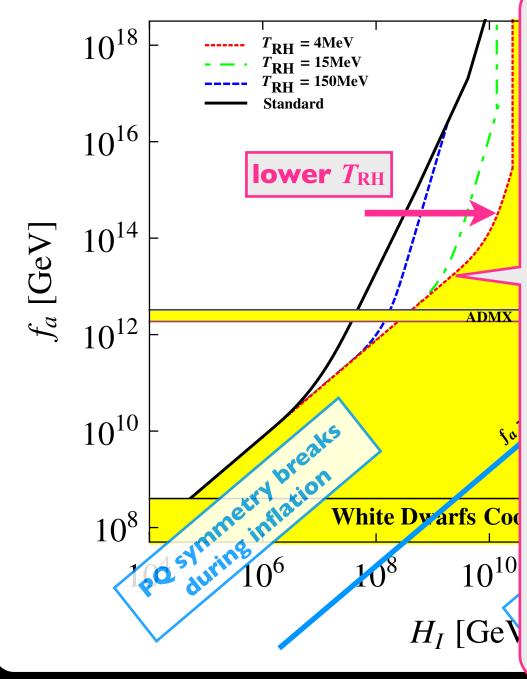


# PQ symmetry breaks after end of inflation

• As  $T_{RH}$  decreases,  $f_a$  must increase and  $m_a$  decrease

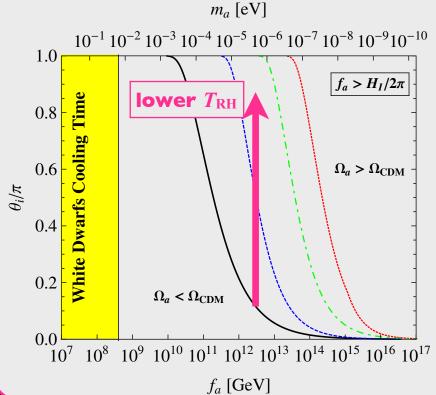




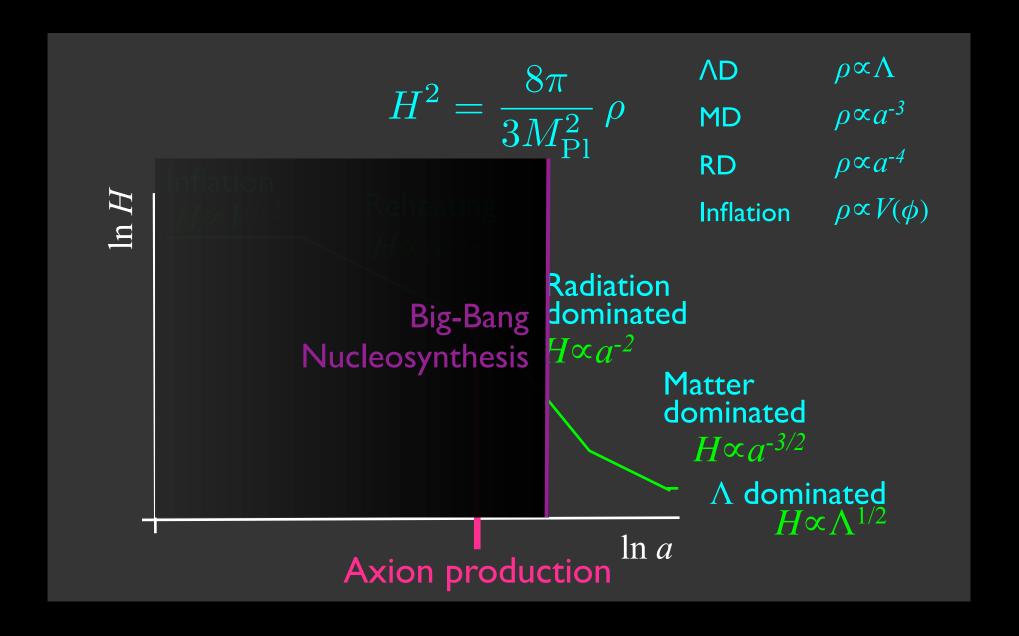


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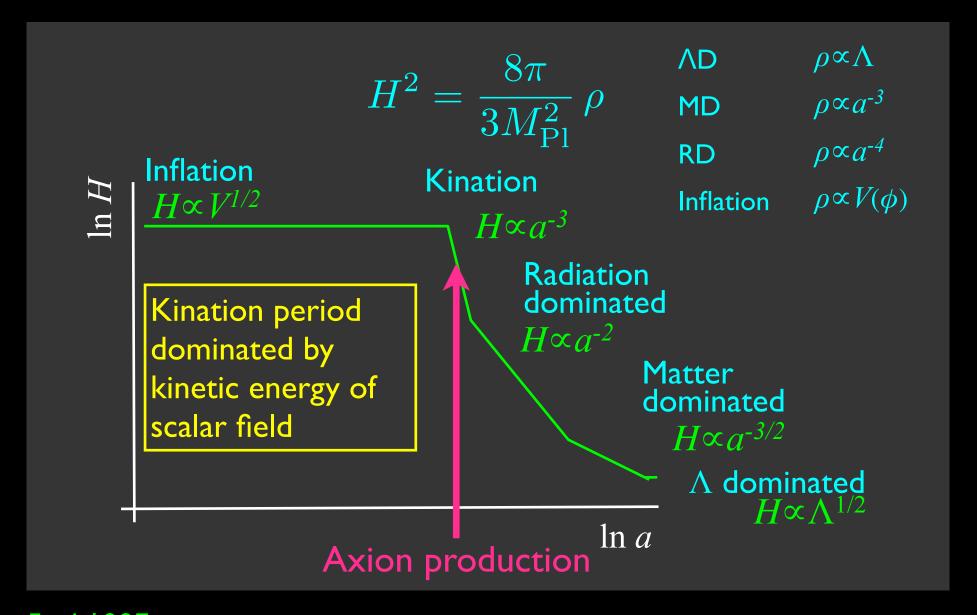
- As  $T_{RH}$  decreases, constraints from non-adiabatic fluctuations become weaker
- And the initial misalignment angle  $\theta_i$  must be larger



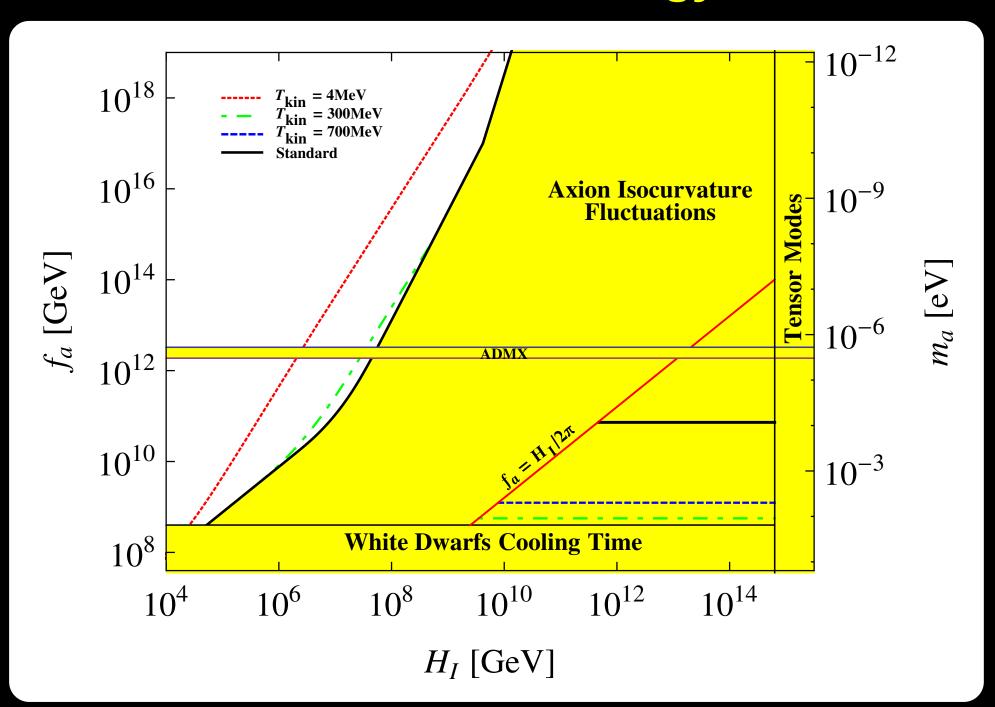
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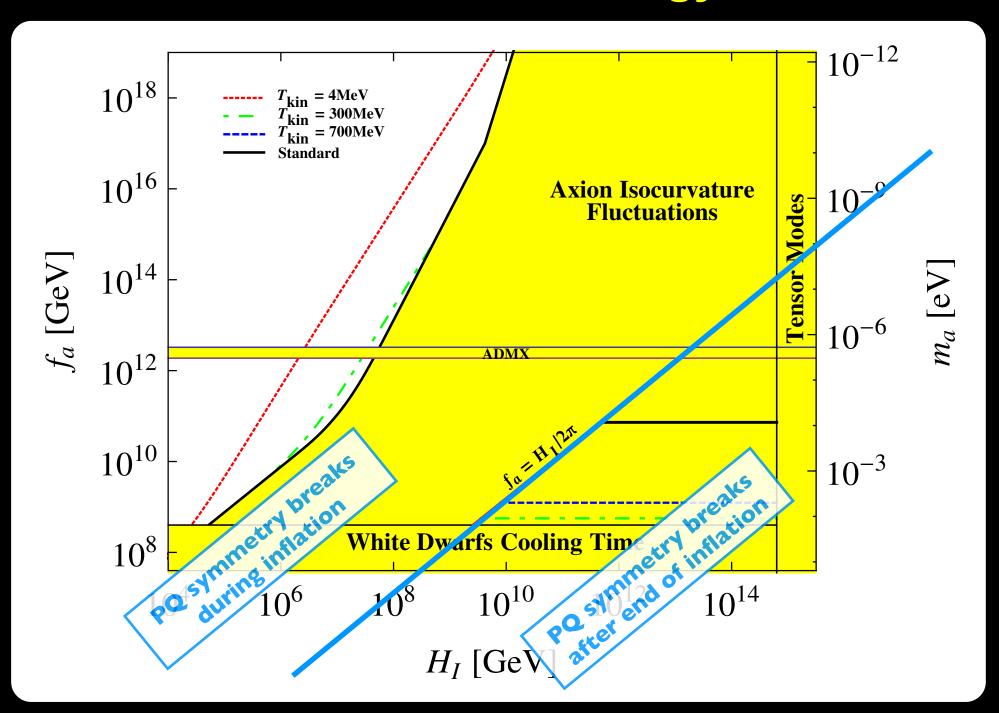


#### **Kination cosmology**



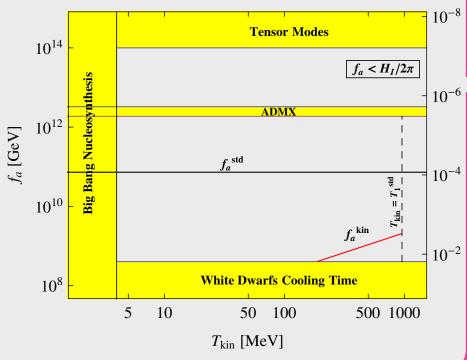
Ford 1987

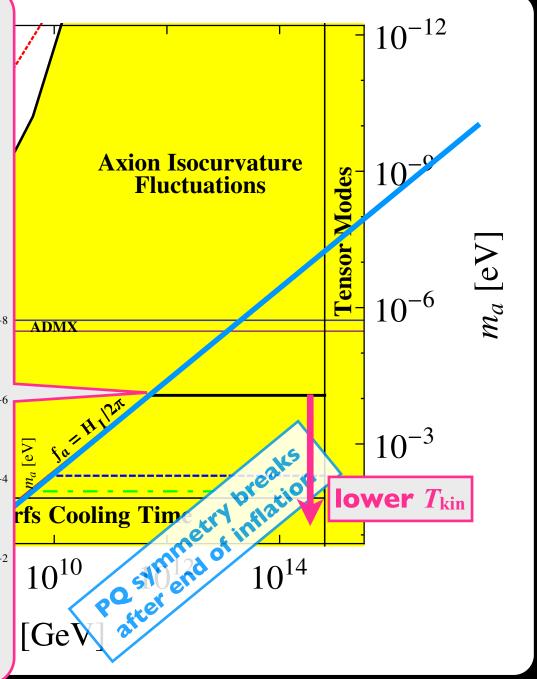


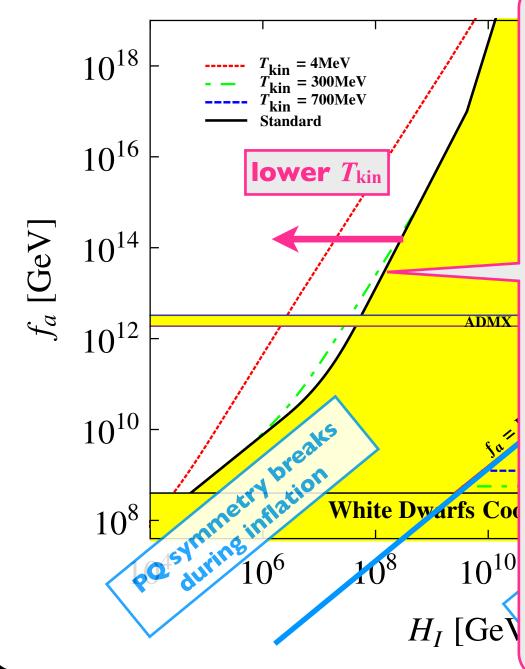


# PQ symmetry breaks after end of inflation

- As  $T_{kin}$  decreases,  $f_a$  must decrease and  $m_a$  increase
- String decay contribution is
   15 × vacuum realignment

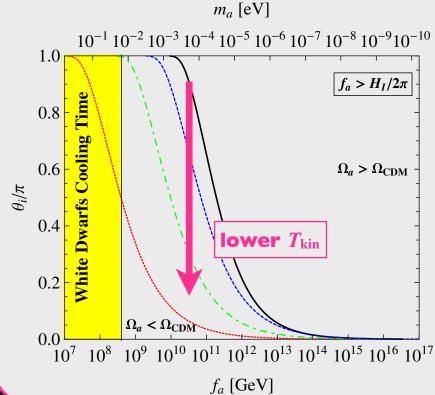






## PQ symmetry breaks during inflation

- As  $T_{\rm kin}$  decreases, constraints from non-adiabatic fluctuations become stronger
- And the initial misalignment angle  $\theta_i$  must be smaller



#### Conclusions

#### For axions to be 100% of cold dark matter.....

- If the Peccei-Quinn symmetry breaks after inflation ends, the axion mass must be  $m_a$ =85±3  $\mu eV$  in standard cosmology
  - much smaller  $m_a$  in LTR cosmology
  - much larger  $m_a$  in kination cosmology
- If the Peccei-Quinn symmetry breaks during inflation, cosmological limits on non-adiabatic fluctuations constrain parameter space and a specific initial misalignment angle  $\theta_i$  must be chosen
  - larger allowed region and larger  $\theta_i$  in LTR cosmology
  - smaller allowed region and smaller  $\theta_i$  in kination cosmology