

Gravitino dark matter and LHC



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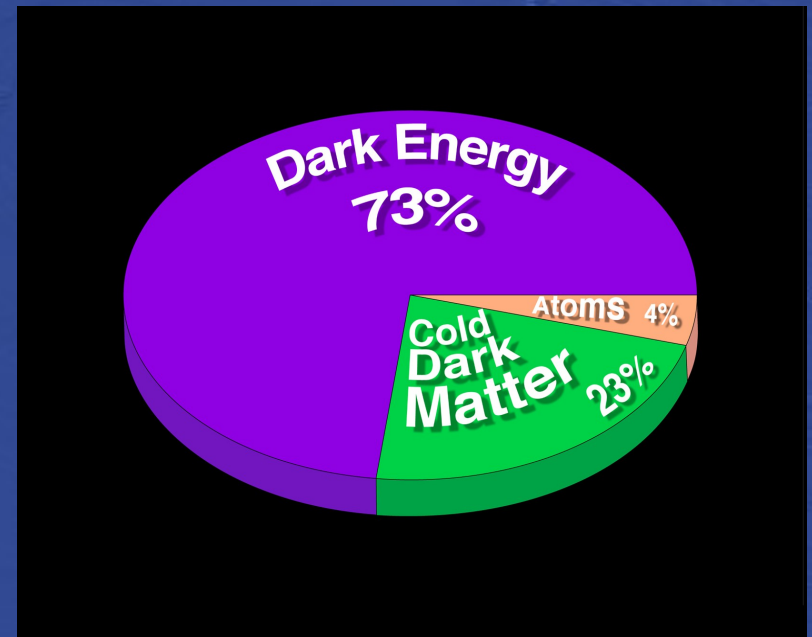
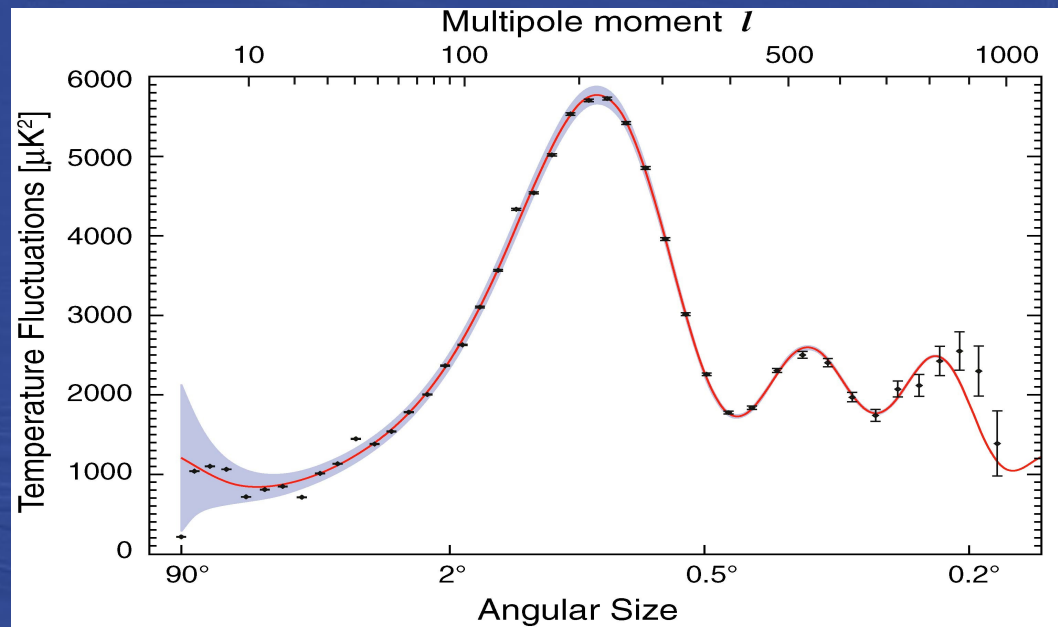
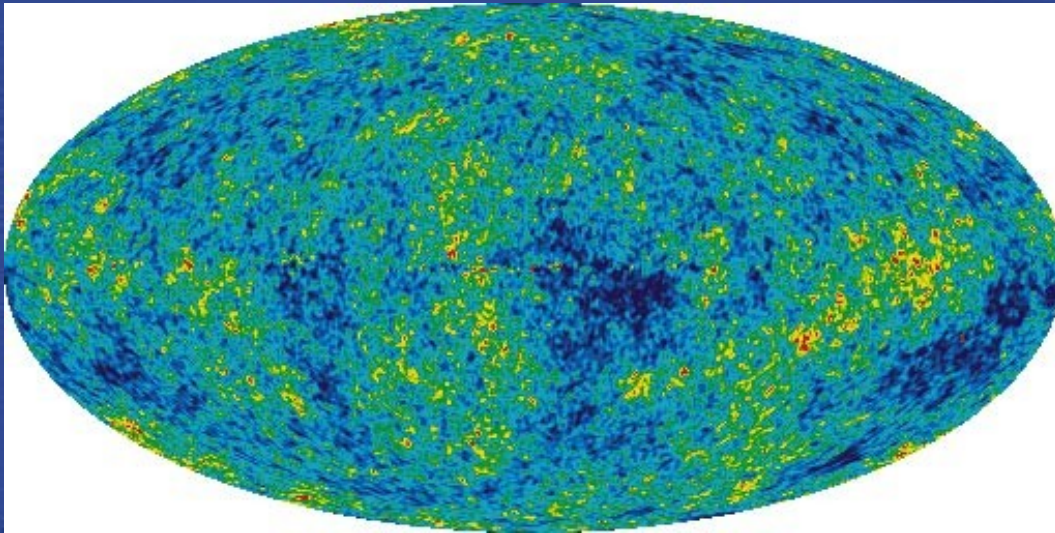
University of Valencia & IFIC

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Outline

- Motivation
- Candidates for DM in SUSY models
- Gravitino mass and interactions
- Cosmology of unstable gravitino
- Cosmology of stable gravitino
- Gravitino at LHC
- Conclusions

CMB Temperature anisotropies



Are there DM candidates in the SM?

THE STANDARD MODEL						
Fermions			Bosons			
Quarks	<i>u</i> up	<i>c</i> charm	<i>t</i> top	γ photon	Force carriers	
	<i>d</i> down	<i>s</i> strange	<i>b</i> bottom	<i>Z</i> Z boson		
Leptons	ν_e electron neutrino	ν_μ muon neutrino	ν_τ tau neutrino	<i>W</i> W boson		
	<i>e</i> electron	μ muon	τ tau	<i>g</i> gluon		
			Higgs[*] boson			

*Yet to be confirmed

Source: AAAS

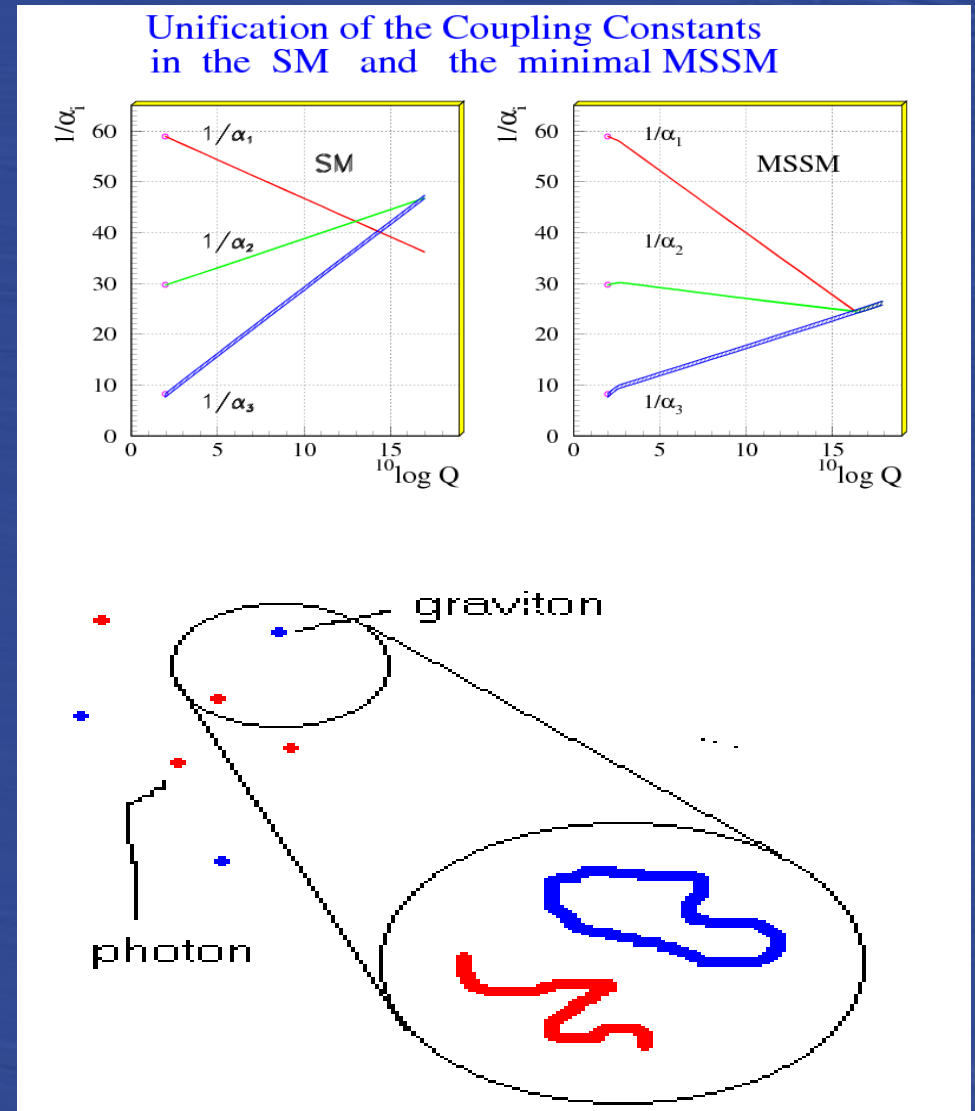
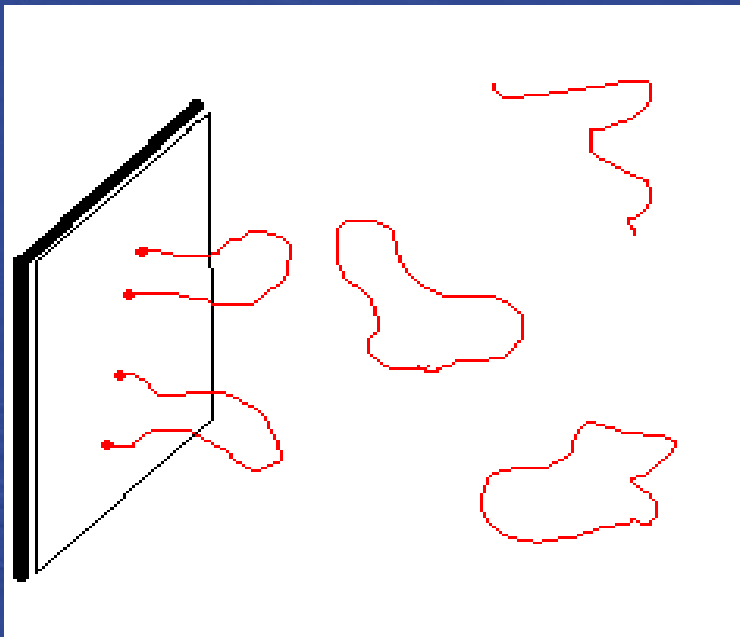
The neutrino cannot
be cold dark matter

From Omega $m=10$ eV

From virial $m > 150$ eV

Supersymmetry: A new symmetry between bosons and fermions

- Basis of superstring theories
- Unification of gauge coupling constants



Particle content in Supersymmetric Models

- The **LSP** is stable in SUSY theories with R-parity conservation. Thus, it will exist as a remnant from the early universe and may account for the observed Dark Matter.

The superpartners

Squarks	$\tilde{u}_{R,L} \quad , \quad \tilde{d}_{R,L}$ $\tilde{c}_{R,L} \quad , \quad \tilde{s}_{R,L}$ $\tilde{t}_{R,L} \quad , \quad \tilde{b}_{R,L}$
Sleptons	$\tilde{e}_{R,L} \quad , \quad \tilde{\nu}_e$ $\tilde{\mu}_{R,L} \quad , \quad \tilde{\nu}_\mu$ $\tilde{\tau}_{R,L} \quad , \quad \tilde{\nu}_\tau$
Neutralinos	$\tilde{B}^0, \quad \tilde{W}^0, \quad \tilde{H}_{1,2}^0$
Charginos	$\tilde{W}^\pm \quad , \quad \tilde{H}_{1,2}^\pm$
Gluino	\tilde{g}
Gravitino	\tilde{G}
Axino	\tilde{a}

Popular candidates for playing the role of cold dark matter

Lightest neutralino: WIMP

Gravitino: Present in Supergravity theories. Can also be the LSP and a good dark matter candidate

Axino: SUSY partner of the axion. Extremely weak interactions

Gauge interactions: Fixed

Yukawa interactions: Superpotential

MSSM:

$$W = \epsilon_{ij} \left(Y_u H_2^j Q^i u + Y_d H_1^i Q^j d + Y_e H_1^i L^j e \right) + \mu \epsilon_{ij} H_1^i H_2^j$$

NMSSM:

$$W = \epsilon_{ij} \left(Y_u H_2^j Q^i u + Y_d H_1^i Q^j d + Y_e H_1^i L^j e \right) - \epsilon_{ij} \lambda S H_1^i H_2^j + \frac{1}{3} \kappa S^3$$

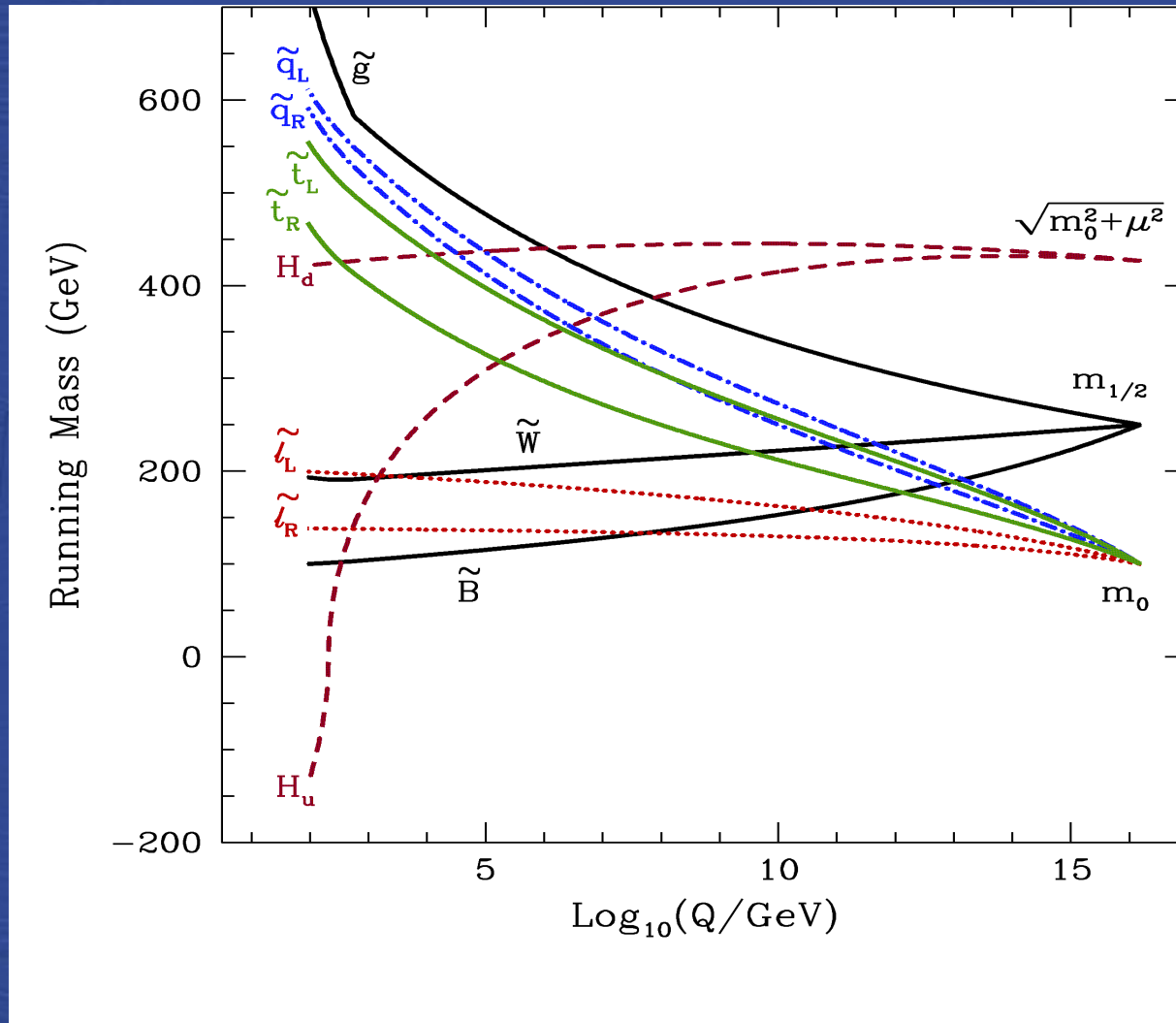
Gut idea: Universal boundary conditions

$$m_{\tilde{f}_i}(M_{GUT}) = m_0$$

$$A_{ij}^u(M_{GUT}) = A_{ij}^d(M_{GUT}) = A_{ij}^l(M_{GUT}) = A_0 \delta_{ij}$$

$$M_1(M_{GUT}) = M_2(M_{GUT}) = M_3(M_{GUT}) = m_{1/2}$$

The running of mass parameters with the mass scale



From 1001.5014
(K. Olive)

Collider constraints on SUSY models

Once the spectrum and couplings are computed, experimental constraints are applied

Masses of superpartners

$$\begin{aligned} m_{\tilde{\chi}_1^\pm} &> 103 \text{ GeV}, \\ m_{\tilde{g}} &> 150 \text{ GeV} \\ m_{\tilde{\tau}} &> 87 \text{ GeV}, \\ &\dots \end{aligned}$$

Mass of the Higgs boson

$$m_h > 114.1 \text{ GeV}$$

Low energy observables that receive SUSY contributions

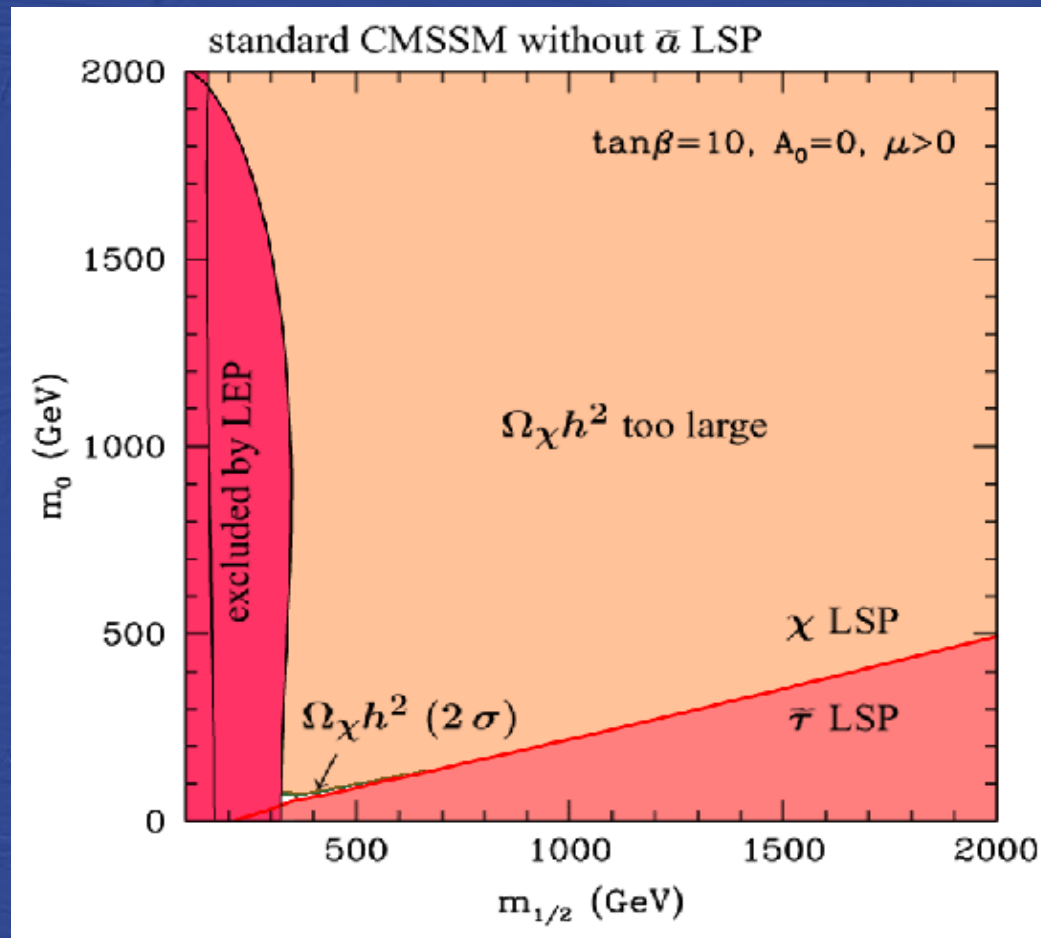
Muon anomalous magnetic moment

$$(g - 2)_\mu$$

Rare decays

$$(b \rightarrow s\gamma)$$

Typical $(m_0 - m_{1/2})$ plane



From hep-ph/0402240
(Covi et al)

Neutralino mass matrix in the MSSM & the NMSSM

Extensions of the MSSM also allow an increase of the Higgs-exchange amplitude. For instance, in the Next-to-MSSM, where a new singlet (and singlino) is included:

$$\mathcal{M}_{\tilde{\chi}^0} = \begin{pmatrix} M_1 & 0 & -M_Z s_\theta c_\beta & M_Z s_\theta s_\beta & 0 \\ 0 & M_2 & M_Z c_\theta c_\beta & -M_Z c_\theta s_\beta & 0 \\ -M_Z s_\theta c_\beta & M_Z c_\theta c_\beta & 0 & -\mu & -\lambda v_2 \\ M_Z s_\theta s_\beta & -M_Z c_\theta s_\beta & -\mu & 0 & -\lambda v_1 \\ 0 & 0 & -\lambda v_2 & -\lambda v_1 & 2\kappa \frac{\mu}{\lambda} \end{pmatrix}$$

The lightest neutralino has now a singlino component

$$\tilde{\chi}_1^0 = \underbrace{N_{11} \tilde{B}^0 + N_{12} \tilde{W}_3^0}_{\text{Gaugino content}} + \underbrace{N_{13} \tilde{H}_d^0 + N_{14} \tilde{H}_u^0}_{\text{Higgsino content}} + \underbrace{N_{15} \tilde{S}}_{\text{Singlino content}}$$

The gravitino can be the LSP in Supergravity

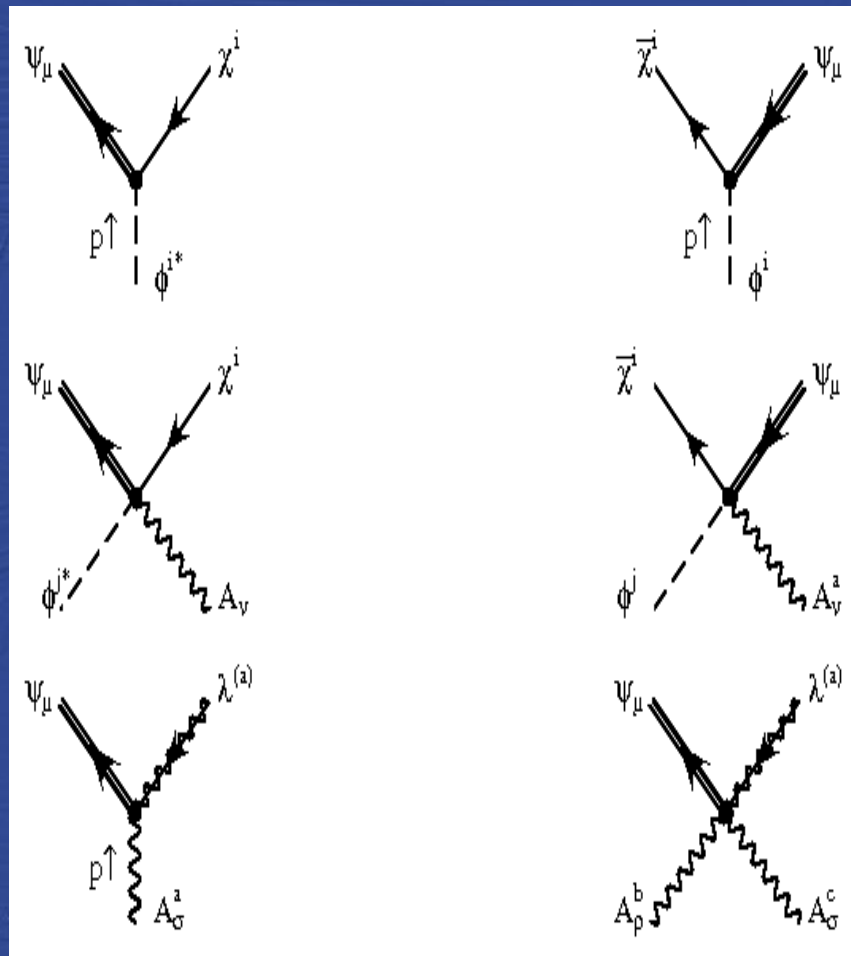
The relation between the gravitino mass and the rest of the soft masses depends on the SUSY-breaking mechanism

The gravitino mass is directly related to the scale of SUSY breaking

Interactions completely determined by the SUGRA Lagrangian

Anomaly-Mediated (AMSB)	$m_{3/2} = O(10^4 - 10^5 \text{ GeV}) \gg m, M$	Gravitino not LSP
Gravity-mediated (GMSB)	$m_{3/2} = O(10^2 - 10^3 \text{ GeV}) \sim m, M$	Gravitino LSP in some regions of the parameter space
Gaugino-Mediated	$m_{3/2} = O(10^{-2} - 10^2 \text{ GeV}) \lesssim m, M$	
Gauge-Mediated	$m_{3/2} = O(10^{-10} - 10^{-8} \text{ GeV}) \ll m, M$	Gravitino LSP

Feynman rules for the gravitino interactions



$$\frac{-1}{\sqrt{2}M_p} \gamma_\nu \gamma_\mu (1 + \gamma_5) p^\nu$$

$$\frac{-1}{\sqrt{2}M_p} \gamma_\mu \gamma_\nu (1 - \gamma_5) p^\nu$$

$$\frac{-1}{2\sqrt{2}M_p} g T_{ji}^a \gamma_\nu \gamma_\mu (1 + \gamma_5)$$

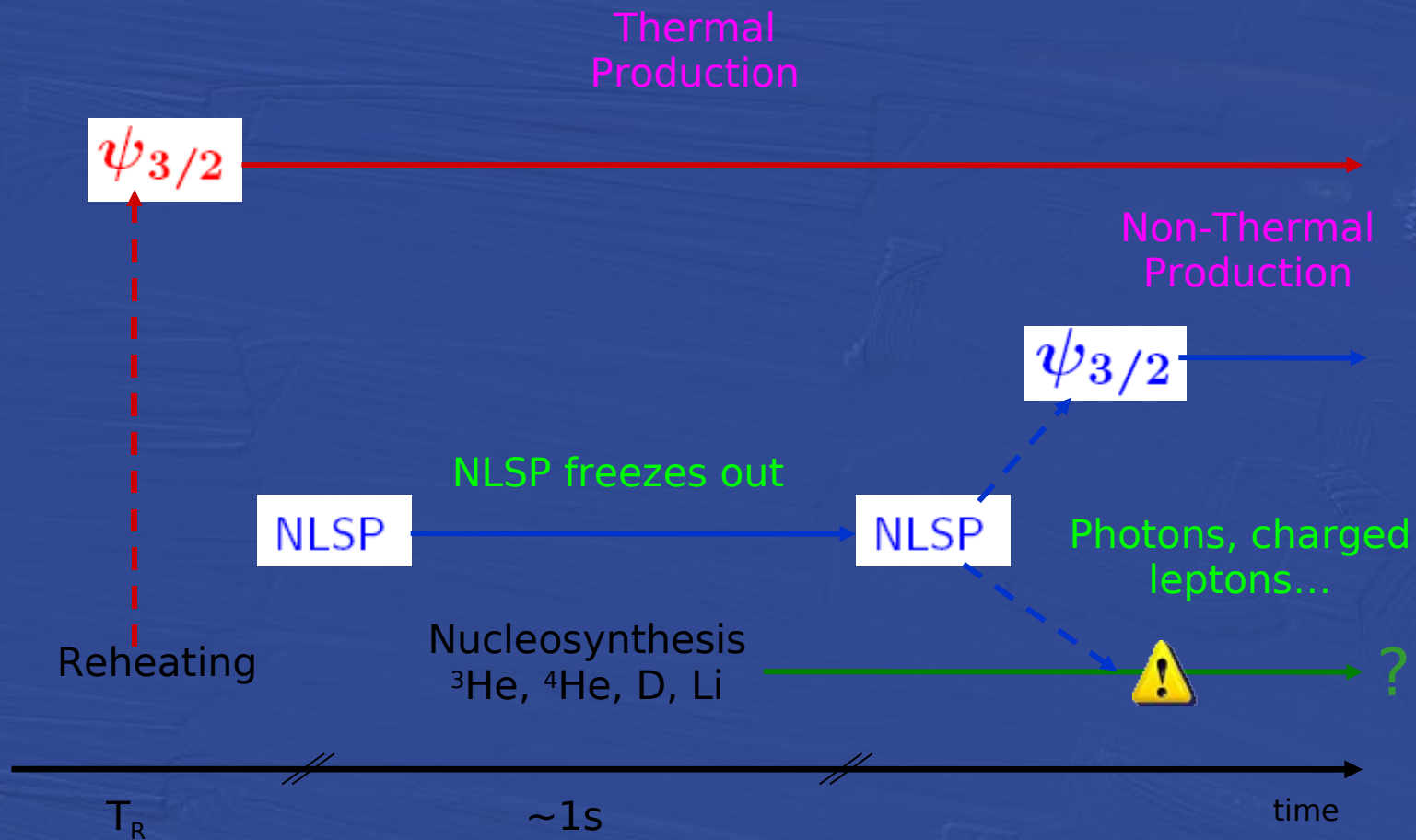
$$\frac{1}{2\sqrt{2}M_p} g T_{ij}^a \gamma_\mu \gamma_\nu (1 - \gamma_5)$$

$$\frac{-i}{4M_p} p_\rho [\gamma^\rho, \gamma_\sigma] \gamma_\mu$$

$$\frac{-1}{4M_p} g f^{abc} [\gamma_\rho, \gamma_\sigma] \gamma_\mu$$

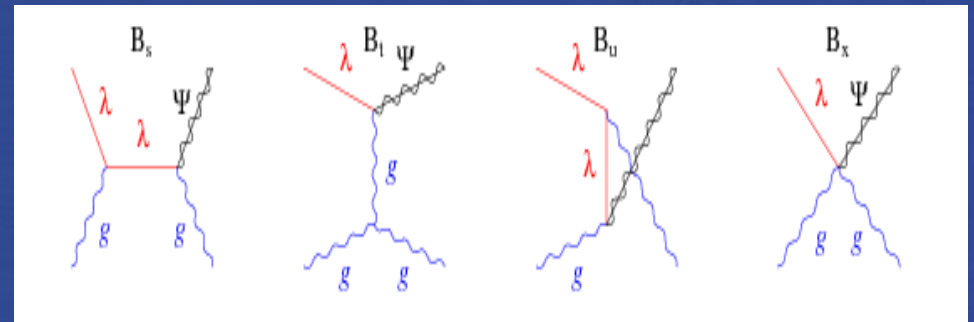
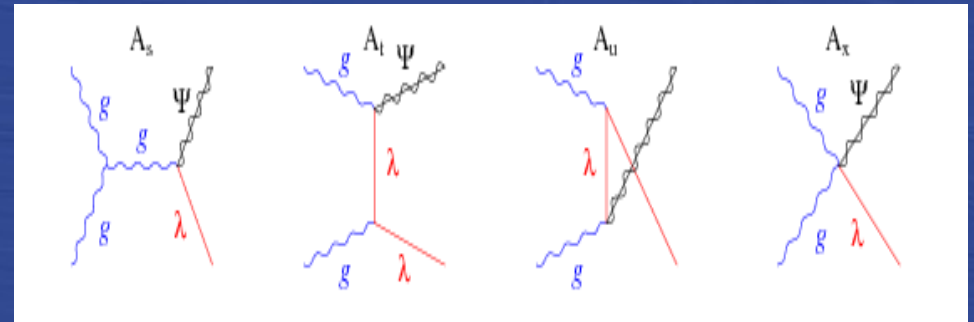
(From Moroi's Thesis hep-ph/9503210)

Gravitino production mechanisms



Thermal production of gravitinos

- SQCD
- Thermal field theory



Thermal production of gravitinos

- Collision term
- Boltzmann equation
- Define the yield
- Hubble parameter & entropy density

$$C(T) \sim \frac{T^6}{M_p^2} \left(1 + \frac{m_{\tilde{g}}^2}{3m_{3/2}^2} \right)$$

$$\dot{n} + 3Hn = C(T)$$

$$Y = n/s, \Omega = mn/\rho_{cr}$$

$$H(T) = 1.66\sqrt{g_*}\frac{T^2}{M_p}$$

$$s(T) = 2\pi^2 h_* T^3/45$$

Assumption: Gravitino LSP & stable <--- R-parity conservation

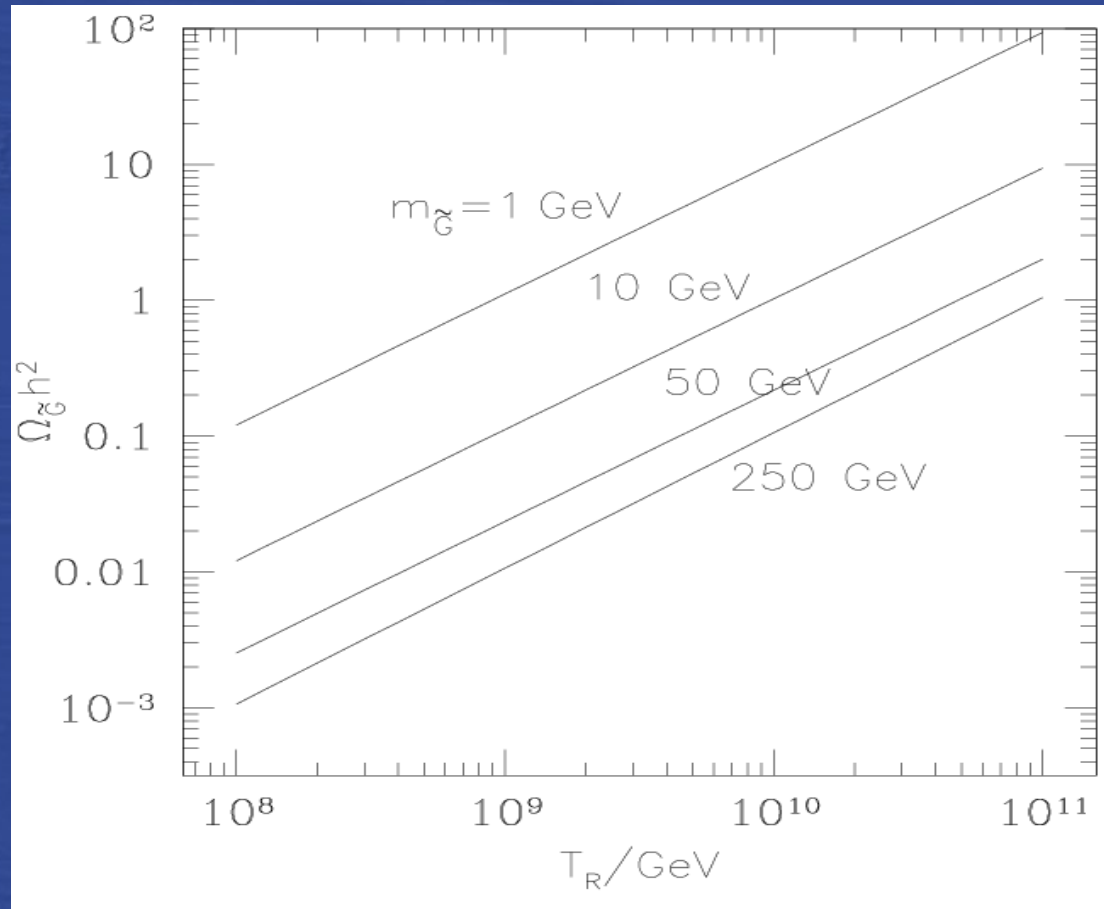
$$\Omega_{3/2} h^2 = \Omega_{3/2}^{TP} h^2 + \Omega_{3/2}^{NTP} h^2$$

$$\Omega_{CDM} h^2 = 0.113$$

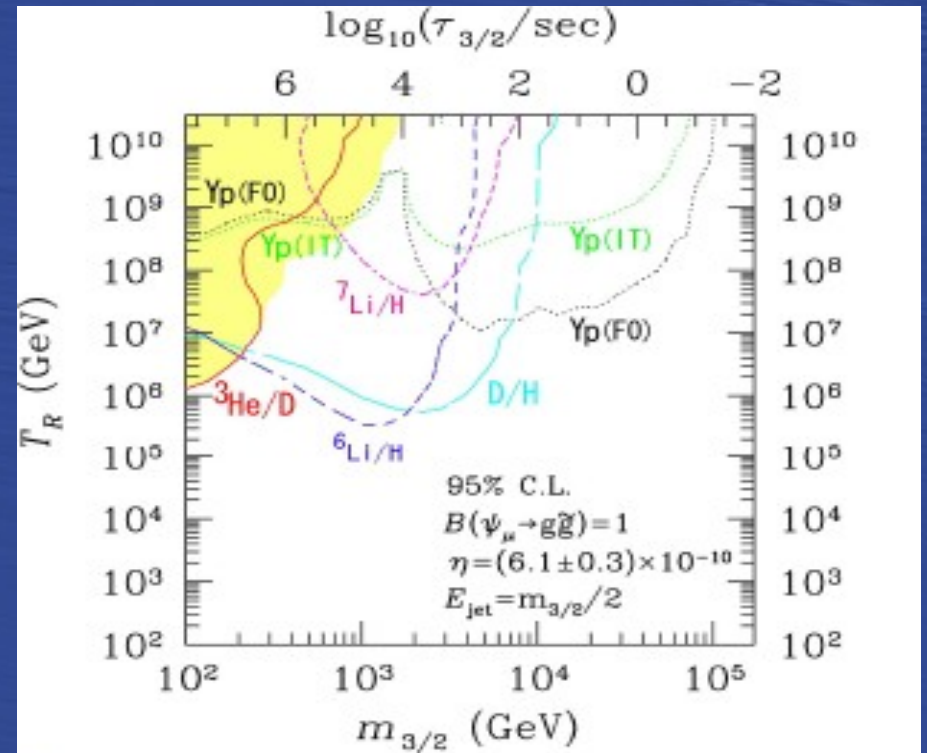
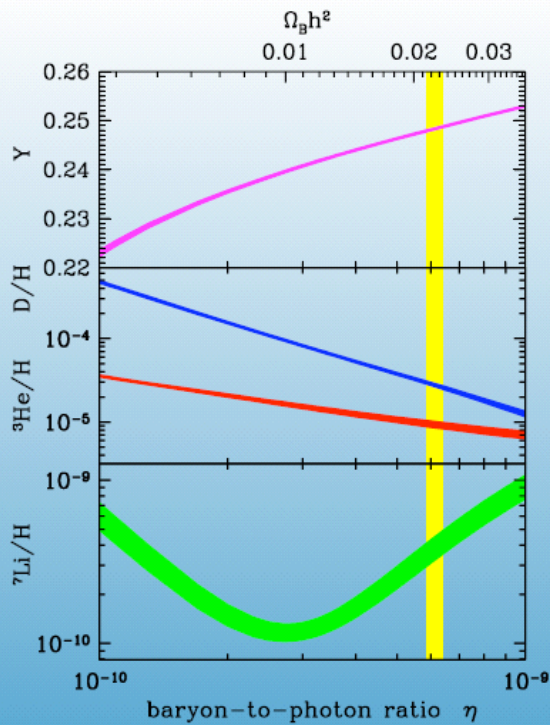
$$\Omega_{3/2}^{NTP} h^2 = \frac{m_{3/2}}{m_{NLSP}} \Omega_{NLSP} h^2$$

$$\Omega_{3/2}^{TP} h^2 = 0.27 \frac{T_R}{10^{10} \text{GeV}} \frac{100 \text{GeV}}{m_{3/2}} \left(\frac{m_{\tilde{g}}}{\text{TeV}} \right)^2$$

Gravitino abundance versus reheating temperature for several gravitino masses



BBN constraints on unstable exotic particles



GRAVITINO PROBLEM (UNSTABLE GRAVITINOS) $T_R < 10^5$ GeV

GDM in R-parity violation

In R-parity conservation \longrightarrow difficult to reconcile:

a) SUSY dark matter

b) BBN

c) Thermal leptogenesis ($T_R \geq 2 \times 10^9 \text{ GeV}$)

$$\Gamma_{NLSP} = \frac{m_{NLSP}^5}{48\pi m_{3/2}^2 M_p^2}$$



very long lifetime

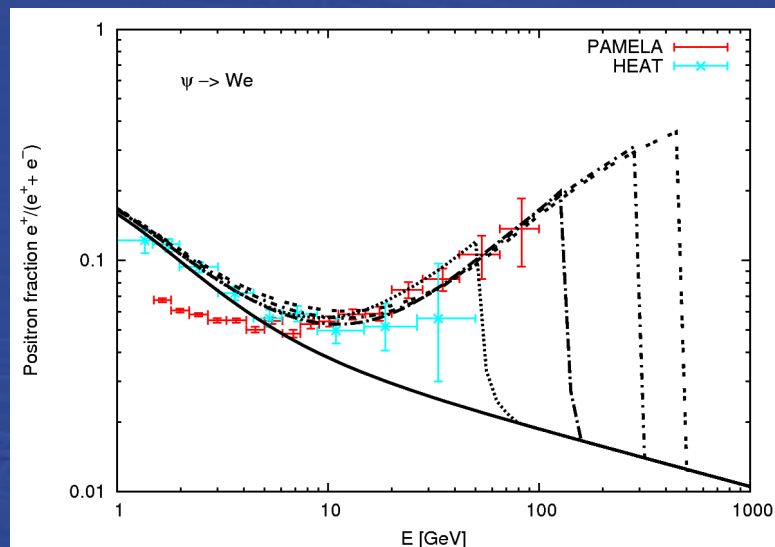
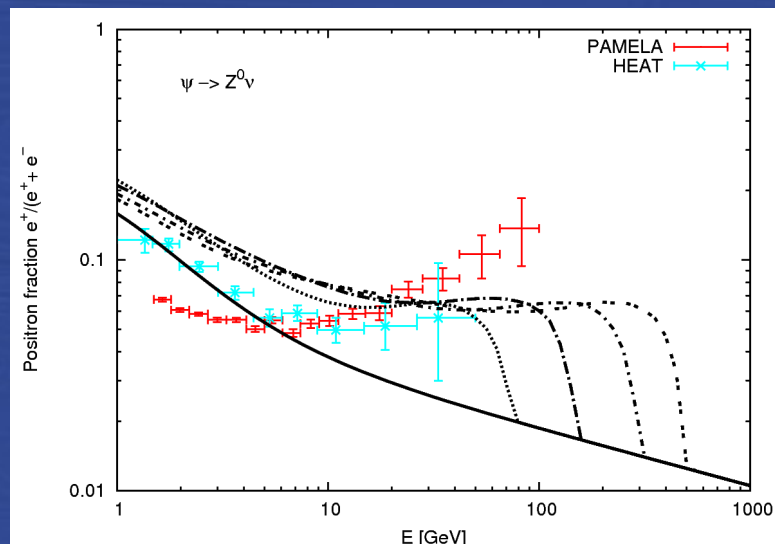
$$\tau_{NLSP} = 2 \text{ days} \left(\frac{m_{3/2}}{5 \text{ GeV}} \right)^2 \left(\frac{150 \text{ GeV}}{m_{NLSP}} \right)^5$$

$$\Omega_{3/2}^{TP} h^2 = 0.27 \frac{T_R}{10^{10} \text{ GeV}} \frac{100 \text{ GeV}}{m_{3/2}} \left(\frac{m_{\tilde{g}}}{\text{TeV}} \right)^2$$

$$m_{3/2} \geq 5 \text{ GeV}$$

GDM in R-parity violation

A. Ibarra et. al.



New formulas for lifetimes

$$\Gamma_{3/2} \sim \lambda^2 \frac{m_{3/2}^3}{M_p^2}$$

$$\tau_{NLSP} = 10^3 \text{ sec} \left(\frac{m_{NLSP}}{100 \text{ GeV}} \right)^{-1} \left(\frac{\lambda}{10^{-14}} \right)^{-2}$$

with tiny couplings

$$10^{-14} < \lambda < 10^{-7}$$

$$m_{3/2} = 150 \text{ GeV} \quad \tau_{3/2} = 10^{26} \text{ sec}$$

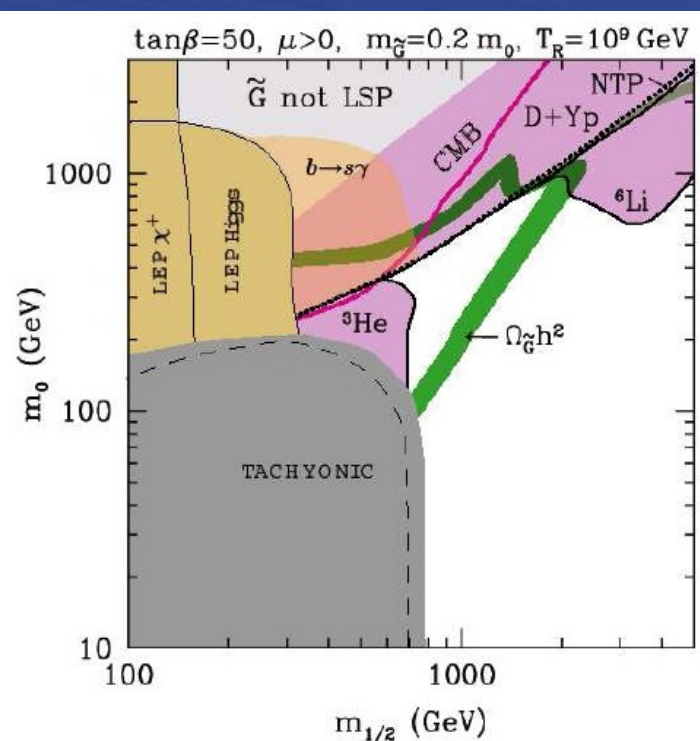
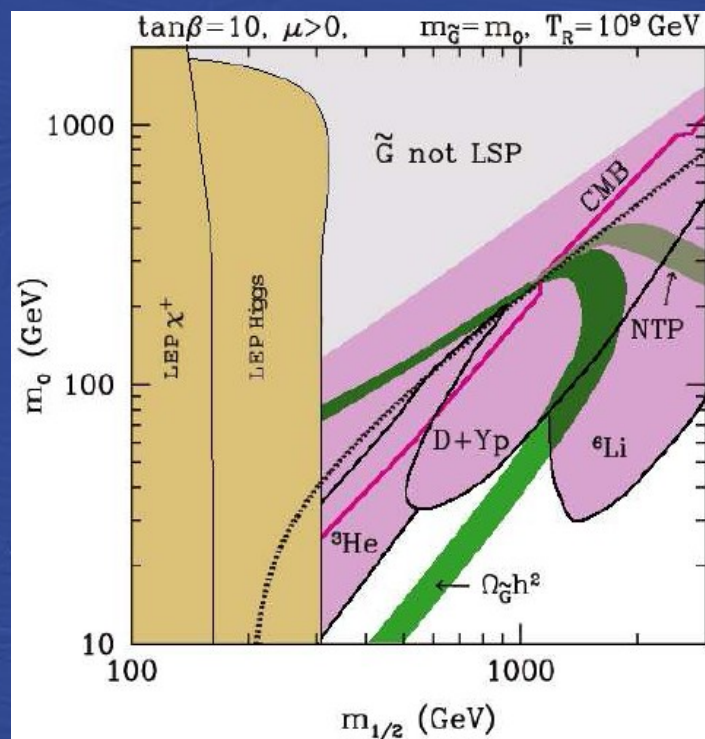
Gravitino dark matter in R-parity conservation

In the CMSSM

Neutralino NLSP areas excluded by BBN constraints. Only part of those with stau NLSP are left.

Non-thermal production alone not sufficient. Large contributions from thermal prod. are necessary.

As long as $T_R \leq 10^9$ GeV sizable regions are found with correct Ω



From hep-ph/0509275

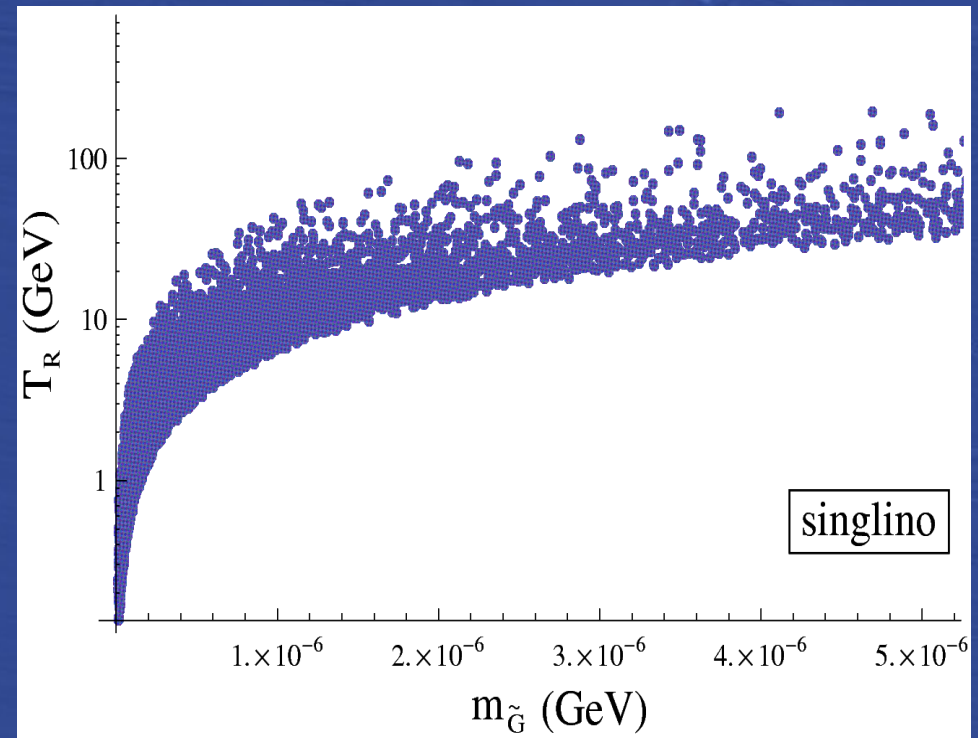
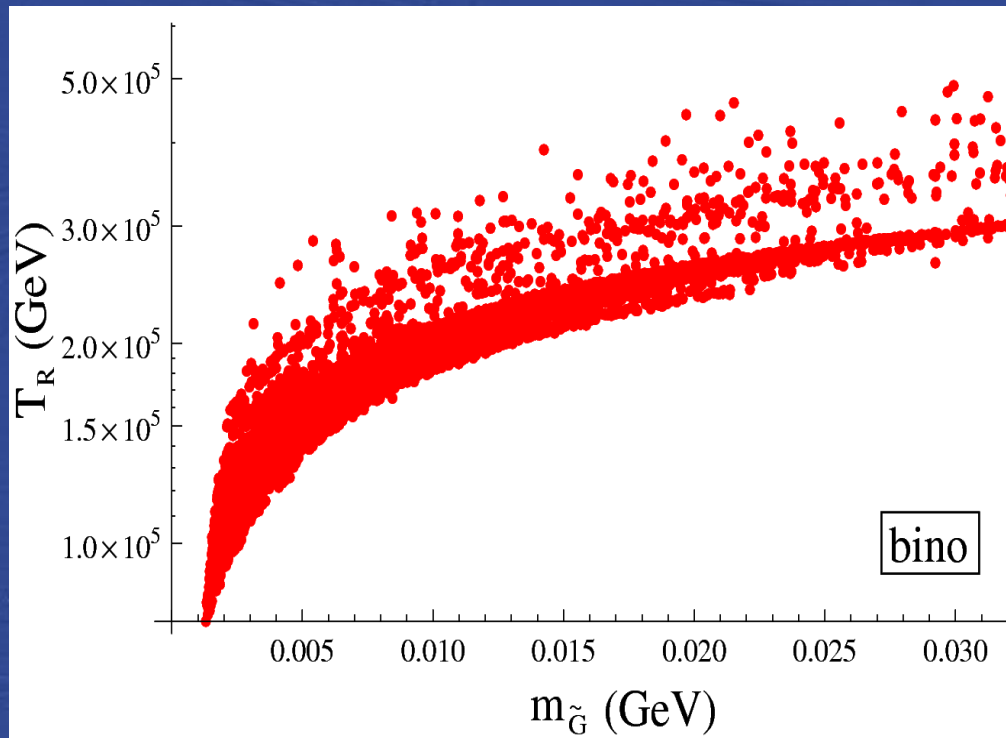
Gravitino dark matter in R-parity conservation

In the cNMSSM

Neutralino NLSP areas NOT excluded by BBN constraints

Non-thermal production alone not sufficient. Large contributions from thermal prod. are necessary.

As long as $T_R \leq 10^7$ GeV & $m_{\tilde{G}} \leq 1$ GeV sizable regions are found with correct Ω .
Singlino case must be excluded



G. Barenboim & GP 1004.4525

Modulus decay

$$X \rightarrow \tilde{G}\tilde{G}$$

$$\Gamma_{tot} \equiv \Gamma(X \rightarrow all) \simeq \Gamma(X \rightarrow gg) + \Gamma(X \rightarrow \tilde{g}\tilde{g}) = \frac{3}{16\pi} \frac{m_X^3}{M_p^2}$$

$$\Gamma_{3/2} = \frac{1}{288\pi} \frac{m_X^3}{M_p^2}$$

$$Br(X \rightarrow \psi_{3/2}\psi_{3/2}) = \frac{\Gamma_{3/2}}{\Gamma_{tot}} = \frac{1}{54} \sim 0.01$$

$$Y_{3/2}^{modulus} = \frac{3}{2} \frac{\Gamma_{3/2}}{\Gamma_{tot}} \frac{T_R}{m_X}$$

$$T_R = 4.9 \times 10^{-3} \left(\frac{10}{g_*(T_R)} \right)^{1/4} \left(\frac{m_X}{10^5 \text{ GeV}} \right)^{3/2} \text{ GeV}$$

Benchmark points in the CMSSM

Model	m_0 (GeV)	$m_{1/2}$ (GeV)	$\tan\beta$	m_χ (GeV)	$\Omega_\chi h^2$
A	200	500	15	205.9	0.64
B	400	800	25	338.6	1.82
C	1000	600	30	252.8	7.81
D	350	450	20	184.9	1.22

Table 1: Four benchmark models considered in the analysis for the neutralino NLSP case.

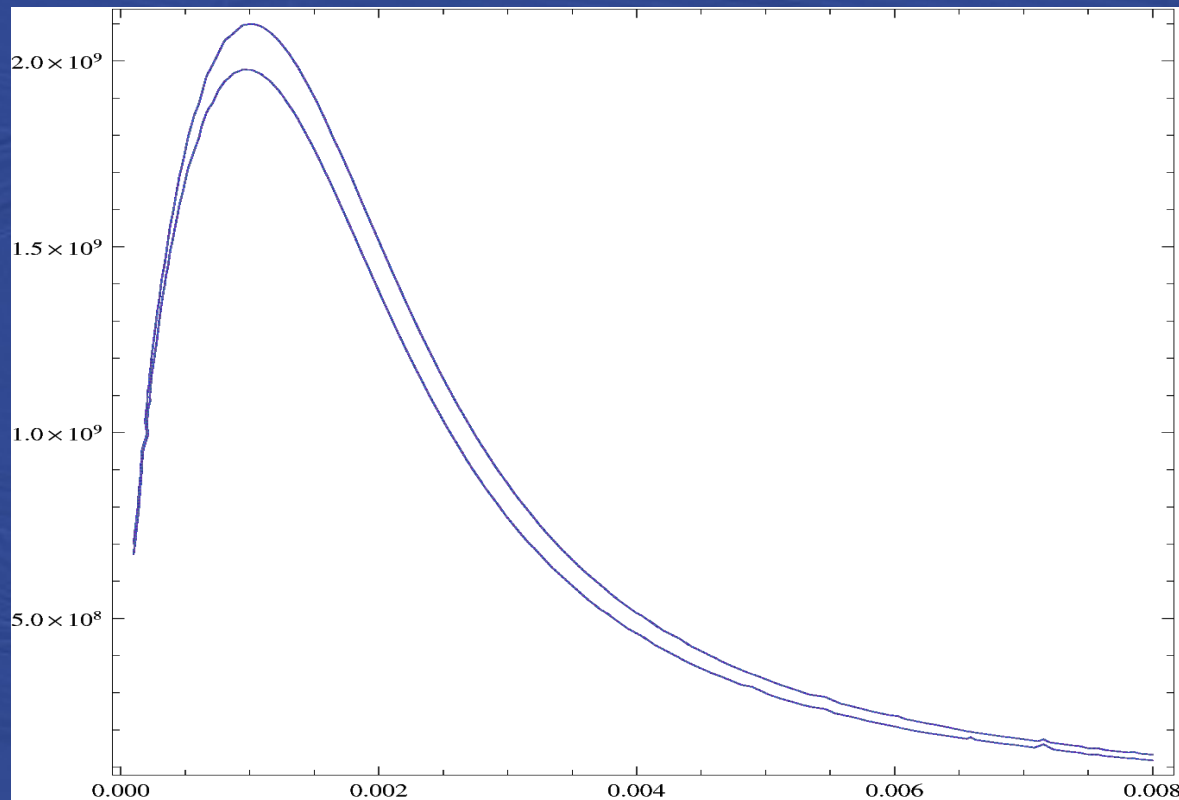
Model	m_0 (GeV)	$m_{1/2}$ (GeV)	$\tan\beta$	$m_{\tilde{\tau}}$ (GeV)	$\Omega_{\tilde{\tau}} h^2$
E	100	1300	5	483.6	0.1
F	50	500	10	186.8	0.01
G	100	800	15	294.3	0.03
H	60	600	20	199.9	0.01

Table 2: Four benchmark models considered in the analysis for the stau NLSP case.

Total gravitino abundance

$$\Omega_{3/2} h^2 = \Omega_{3/2}^{TP} h^2 + \Omega_{3/2}^{NTP} h^2 + \Omega_{3/2}^{modulus} h^2$$

$$\Omega_{CDM} h^2 = 0.113$$

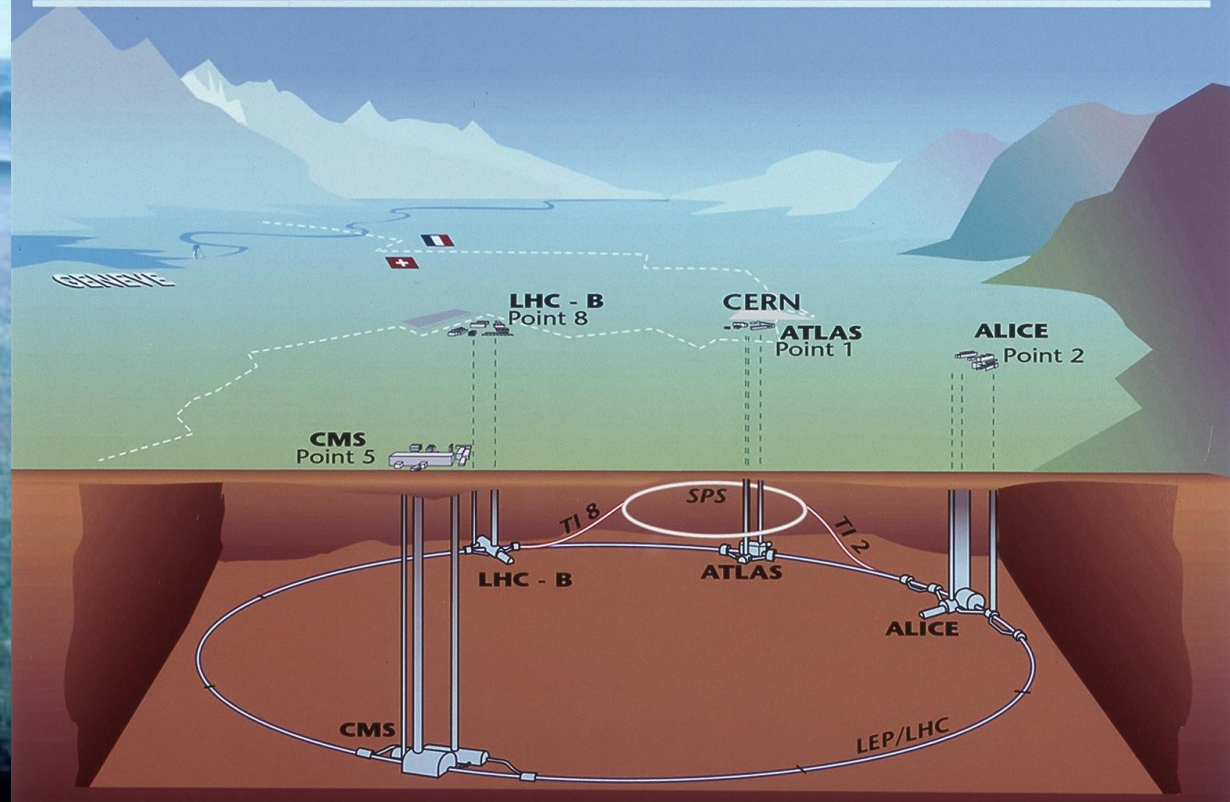


A typical plot for the benchmark points

Gravitino @ LHC



Overall view of the LHC experiments.



Timeline

- 10 September 2008: LHC starts operating
- 20 November 2009: It restarts after the incident
- 23 November 2003: First collisions at 450 GeV
- 30 November 2009: It reaches 1.18 TeV
- 30 March 2010: Two beams collided at 7 TeV at 13:06 CEST

If NLSP=stau

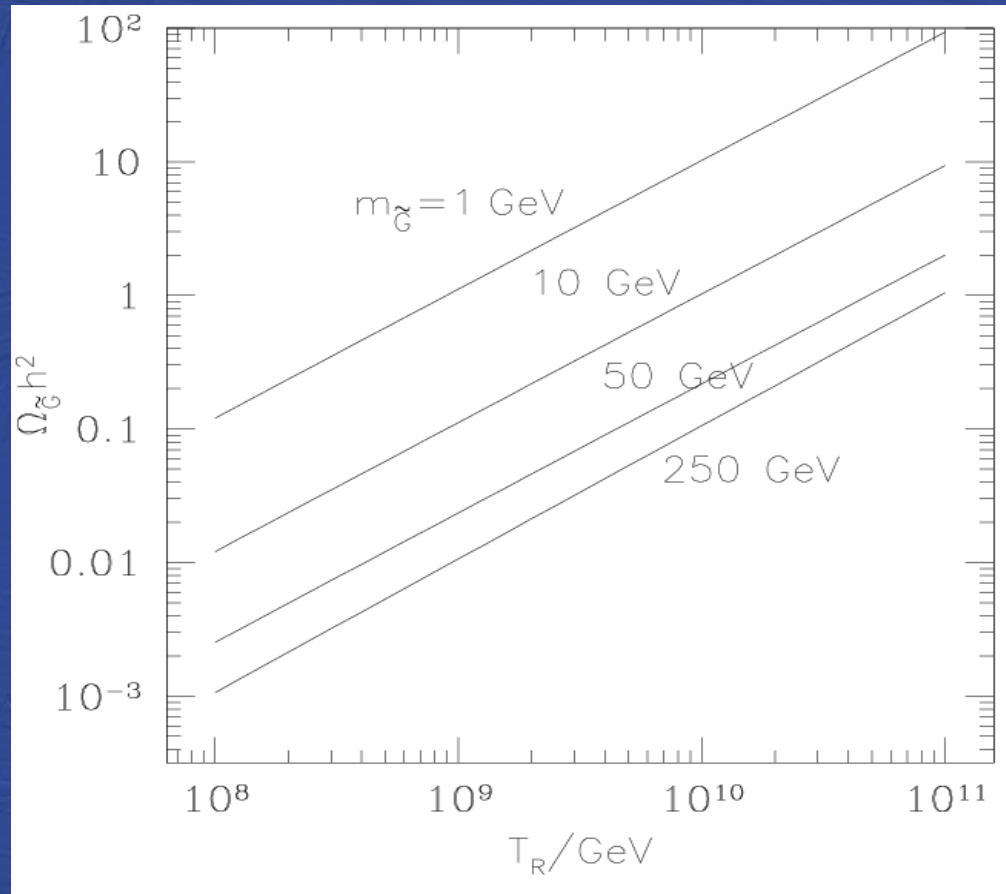
Using gravitino Feynman rules compute the stau lifetime

$$\Gamma(\tilde{f} \rightarrow f \tilde{G}) = \frac{m_{\tilde{f}}^5}{48\pi M_p^2 m_{3/2}^2} \left(1 - \frac{m_{3/2}^2}{m_{\tilde{f}}^2}\right)^4$$

$$\tau_{\tilde{f}} = 6.1 \cdot 10^3 \text{sec} \left(\frac{m_{3/2}}{100 \text{GeV}}\right)^2 \left(\frac{1000 \text{GeV}}{m_{\tilde{f}}}\right)^5 \left(1 - \frac{m_{3/2}^2}{m_{\tilde{f}}^2}\right)^{-4}$$

Measure stau mass and lifetime → determine gravitino mass

Recall the plot Omega versus Reheating temperature



From gravitino mass we can

determine

$$T_R$$

Measure of reheating temperature
at colliders !

Conclusions

- SUSY: Well motivated and most popular journey beyond the SM
- Byproduct: Ideal candidates for CDM (axino, gravitino, neutralino)
- Gravitino: Mass related to SUSY breaking scheme, interactions completely determined by SUGRA Lagrangian
- If light it can play the role of CDM in both R-parity conserving and violating SUSY models
- In R-parity violation: Scenario compatible with thermal leptogenesis, PAMELA anomaly
- In R-parity conservation: Both in CMSSM and cNMSSM there is allowed parameter space where all constraints satisfied
- LHC: Measure T_R in colliders, if stau NLSP from its lifetime