

#### MADFKS AUTOMATION OF THE FKS SUBTRACTION METHOD

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HP2.3rd, Florence, September 14 - 17, 2010



#### NEXT-TO-LEADING ORDER





## WHY AUTOMATE?

#### To save time

NLO calculations can take a long time. It would be nice to spend this time doing phenomenology instead.

To reduce the number of bugs in the calculation Having a code that does everything automatically will be without\* bugs once the internal algorithms have been checked properly.

\* To have all processes within one framework To learn how to use a new code for each process is not something all our (experimental) colleagues are willing to do.



#### IR DIVERGENCE

$$\sigma^{\text{NLO}} = \int_{m+1} d^{(d)} \sigma^R + \int_m d^{(d)} \sigma^V + \int_m d^{(4)} \sigma^B$$

- Real emission -> IR divergent
- % (UV-renormalized) virtual corrections
  -> IR divergent
  - After integration, the sum of all contributions is finite (for infrared-safe observables)
  - \* To see this cancellation the integration is done in a noninteger number of dimensions: Not possible with a Monte-Carlo integration



# **SUBTRACTION TERMS** $\sigma^{\text{NLO}} = \int_{m+1} d^{(d)} \sigma^{R} + \int_{m} d^{(d)} \sigma^{V} + \int_{m} d^{(4)} \sigma^{B}$



 $\sigma^{\text{NLO}} = \int_{m+1} d^{(d)} \sigma^R + \int_m d^{(d)} \sigma^V + \int_m d^{(4)} \sigma^B$  $\sigma^{\text{NLO}} = \int_{m+1} \left[ d^{(4)} \sigma^R - d^{(4)} \sigma^A \right] + \int_m \left[ d^{(4)} \sigma^B + \int_{\text{loop}} d^{(d)} \sigma^V + \int_1 d^{(d)} \sigma^A \right]_{\epsilon=0}$ 

Include subtraction terms to make real emission and virtual contributions separately finite

#All can be integrated numerically



In collaboration with Stefano Frixione, Fabio Maltoni e<sup>3</sup> Tim Stelzer, arXiv: 0908.4247



$$\mathbf{FKS SUBTRACTION}$$

$$\sigma^{\text{NLO}} = \int_{m+1} \left[ d^{(4)}\sigma^R - d^{(4)}\sigma^A \right] + \int_m \left[ d^{(4)}\sigma^B + \int_{\text{loop}} d^{(d)}\sigma^V + \int_1 d^{(d)}\sigma^A \right]_{\epsilon=0}$$

- FKS subtraction: Frixione, Kunszt & Signer 1996.
- \* Also known as "residue subtraction"
- Sased on partitioning the phase space such that each partition has at most one soft and/or collinear divergence
- \* Use simple plus-distributions to regulate the divergences
- \*\* Are relevant formulae can be found in our MadFKS paper, arXiv:0908.4247



## FKS -- TECHNICALITIES

- \* Naive scaling of the number of subtraction terms is n<sup>2</sup> (as opposed to n<sup>3</sup> of CS dipoles). Can be greatly reduced by using symmetry of the matrix elements
  - Adding additional gluons does not lead to more phase-space partitions
- In a given phase space partition, Born amplitudes need be computed only once for each real-emission event, and can be used for the Born and collinear, soft and soft-collinear counter events (and their remainders)
- Trivially extended to BSM physics. Massive particles have only soft singularity which is independent of the spin



#### MADFKS

- Automatic FKS subtraction for QCD within the MadGraph/MadEvent framework
- Siven the (n+1) process, it generates the real, all the subtraction terms and the Born processes
- Completely general & all automatic, using the same user-friendly interface as MadGraph
- MadFKS works also for any BSM physics model implemented in MadGraph, e.g. MSSM
- \* Color-linked Borns generated by MadDipole *RF, Gehrmann & Greiner*
- MC-ing over helicities possible; only more efficient for high-multiplicity final states
- Phase-space generation for the (n)-body is the same as in standard MG. It has been heavily adapted to generate (n+1)-body emission events at the same time



## FULL NLO

- Solution Structure Stru
- Interface to link with the virtual corrections following the Binoth-Les Houches Accord
  - Standardized way to link MC codes to one-loop programs
- We are also working on an interface to CutTools In collaboration with Hirschi, Garzelli & Pittau

# 5 JETS AT LEP @ NLO



#### RF, Frixione, Melnikov, Zanderighi arXiv:1008.5313



- Durham jet algorithm
- Scale dependence: +45% -30% at LO; ±20% at NLO
- \* LO and NLO bands overlap (LO uses  $\alpha_s(M_Z)=0.130$ )
- Point-by-point agreement with BlackHat (*Berger et al.*) for the virtuals Rikkert Frederix, September 14, 2010



#### HADRONIZATION CORRECTIONS

- \* Historically, for LEP, hadronization corrections have been estimated by using Ariadne, Herwig & Pythia, which are tuned to data
- <sup>∞</sup> However, they are based on 2→2 and 2→3 matrix elements. Therefore, for these 5 jet observables, they rely strongly on the showers.
- We have estimated the hadronization corrections using Sherpa, with CKKW matching up to 5 hard partons
- ☆ Corrections are mild up to -ln(y<sub>45</sub>) ~ 6



#### $\alpha_s$ extraction



LEP1, hadr. LEP1, no hadr $\sigma_{\text{tot}}^{-1} d\sigma/dy_{45}, R_5 = \sigma_{\text{tot}}^{-1} d\sigma/dy_{45}, R_5$ stat.	
+0.0002 +0.0002	5
stat.	0
Stat.	
-0.0002 $-0.0002$	
syst. $+0.0027 +0.0027$	
-0.0029 -0.0029	
pert. $+0.0062 +0.0068$	
-0.0043 $-0.0047$	
fit range $+0.0014$ $+0.0005$	
-0.0014 -0.0005	
+0.0012 –	
-0.0012	
$\alpha_s(M_Z) = 0.1159 \frac{+0.0070}{0.0055} = 0.1163 \frac{+0.0073}{0.0055}$	}
$\alpha_s(M_Z) = 0.1159 \begin{array}{c} + 0.0016 \\ -0.0055 \end{array} \begin{array}{c} 0.1163 \\ -0.0055 \end{array} \begin{array}{c} -0.0055 \end{array}$	<u> </u>
LEP2, no hadr. LEP2, no hadr. LEP2,	no hadr
	$\mathrm{d}y_{45},R_{4}$
+0.0020 $+0.0022$ $+0.$	0015
stat. $-0.0022$ $-0.0025$ $-0.$	0016
syst. $+0.0008 +0.0012 +0.$	0008
-0.0009 -0.0012 -0.	0008
pert. $+0.0049 +0.0029 +0.$	0029
-0.0034 $-0.0020$ $-0.$	0020
ht range	0028
-0.0038 $-0.0030$ $-0.$	0028
$\alpha_{e}(M_{Z}) = 0.1189 = 0.1120 = 0.1155$	+0.0044
-0.0057 $-0.0047$ $-0.0047$	-0.0039

- Statistical uncertainties negligible at LEP1; larger at LEP2
- Systematic and Perturbative uncertainties
   larger at LEP1 than LEP2, fit range
   uncertainties are opposite
- Uncertainties from hadronization corrections already negligible at LEP1, not even considered for LEP2
- Correlations between bins and LEP energies taken into account conservatively

$$\alpha_s(M_Z) = 0.1156^{+0.0041}_{-0.0034}$$



#### **α**<sub>s</sub> extraction -correlations

- Statistical uncertainties are uncorrelated between different center-of-mass energies. At a given c.o.m. energy, y<sub>45</sub> is uncorrelated, while R<sub>5</sub> is assumed to be fully correlated.
- Systematic uncertainties are assumed to be fully correlated at a given c.o.m. energy and between all LEP2 energies, however completely uncorrelated between LEP1 and LEP2.
- \* Perturbative uncertainties are assumed to be fully correlated.
- # Hadronization uncertainties (considered only at LEP1) are assumed to be fully correlated.



#### COMPARISON WITH OTHER MEASUREMENTS

Observable	$\alpha_s(M_Z)$	Ref.
au decays	$0.1197 \pm 0.0016$	S. Bethke
$\Upsilon$ decays	$0.119 \pm 0.0055$	N. Brambilla <i>et al.</i>
3 jet observables	$0.1224 \pm 0.0039$	G. Dissertori et al.
jets in DIS	$0.1198 \pm 0.0032$	H1 collaboration
DIS	$0.1142 \pm 0.0021$	J. Blümlein
$\operatorname{thrust}$	$0.1135 \pm 0.0011$	R. Abbate <i>et al.</i>
lattice	$0.1183 \pm 0.0008$	HPQCD collaboration
EW fits	$0.1193 \pm 0.0028$	H. Flacher <i>et al.</i>
world average	$0.1184 \pm 0.0007$	S. Bethke
$e^+e^- \rightarrow \text{five jets}$	$0.1156 \pm 0.0038$	RF, S. Frixione, K. Melnikov & G. Zanderighi

- # Uncertainty competitive with other measurements
- Slightly smaller than world average, but consistent within uncertainties

#### VIRTUAL CORRECTIONS

In collaboration with Valentin Hirschi and Maria-Vittoria Garzelli e<sup>3</sup> Roberto Pittau



#### VIRTUAL CORRECTIONS

- Interface using the Binoth-LHA is available
- \* For more flexibility (e.g. massive particles & BSM) we also started working on generating the virtual corrections ourselves
- Solution Using the OPP method as implemented in CutTools
  Ossola, Papadopoulos & Pittau



#### IMPLEMENTATION

- MadGraph generates the loop diagrams by cutting one of the particles in the loop: simple tree-level diagrams remain
   → Passed to CutTools
- \*\* NLOComp filters duplicates and sets-up the interference with the Born diagrams: computation of the color factors.
- Ghosts also needed
- R2 terms are computed using tree-level Feynman Rules Draggiotis, Garzelli, Papadopoulos & Pittau
- Point-by-point agreement found with MCFM (and private codes): Drell-Yan, 2-jet production, top pair, W/Z+1 jet...
- \* Not yet optimized in any way. This will be done only in the MGv5 framework



#### WBB ASSOCIATED PRODUCTION

- # First new results with MadFKS+NLOComp/CutTools
- $\gg pp \to W^+ (\to e^+ \nu_e) b\bar{b}$ , with massive b's
- \*\* Similar calculation by Febres Cordero, Reina & Wackeroth
- \* However, here the W boson decay is included, and stable enough to generate results without cuts on the bottom quarks

## Febres Cordero et al. WBB RESULTS



- Transverse momentum of the hardest and 2nd hardest b-jets regulated by the b quark mass
- W boson decay included
- Unfortunately, slight disagreement between calculations when cuts are also applied to our results

work in progress





#### MATCHING TO A PARTON SHOWER

In collaboration with Stefano Frixione & Paolo Torrielli



## **AUTOMATION OF MC@NLO**

$$d\sigma_{\text{mconlo}}^{(\mathbb{H})} = d\phi_{n+1} \left( \mathcal{M}^{(r)}(\phi_{n+1}) - \mathcal{M}^{(\text{mc})}(\phi_{n+1}) \right)$$

$$d\sigma_{\text{MCONLO}}^{(\mathbb{S})} = \int_{+1}^{} d\phi_{n+1} \Big( \mathcal{M}^{(b+v+rem)}(\phi_n) - \mathcal{M}^{(c.t.)}(\phi_{n+1}) + \mathcal{M}^{(\text{MC})}(\phi_{n+1}) \Big)$$

- In black: pure NLO, fully tested in MadFKS
- In red: already implemented for Herwig 6;
  Pythia and Herwig++ are work in progress
  - FKS is based on a collinear picture, so are the MC counter terms: branching structure is for free
  - % Automatic determination of color partners
- \*\* Automatic computation of leading-color matrix elements
   \*\* Works also when MC-ing over helicities
   Rikkert Frederix, September 14, 2010



#### MADFKS MATCHED TO PARTON SHOWER

In MadFKS many process fully tested and working (e.g. e<sup>+</sup>e<sup>-</sup> to jets, Drell-Yan, top pair production, ...)

\* New result: t-channel single top production



#### **T-CHANNEL SINGLE TOP**

\*\* Already implemented in MC@NLO and POWHEG
Frixione, Laenen, Motylinski & Webber (2006);
Alioli, Nason, Oleari & Re (2009)

\* However, due to the massless initial state b quark in the fixed order calculation, some strange behavior at low pT and for forward B hadrons





## **INITIAL STATE B QUARK**

"Standard" way of looking at this process



leading order



5-flavor scheme

4-flavor

scheme

(contribution to) NLO

But there is an equivalent description with no bottom PDF and an explicit gluon splitting to b quark pairs





(part of) leading order



#### THE TWO SCHEMES





5-flavor scheme: " $2 \rightarrow 2$ "

At all orders both description should agree;
 otherwise, differ by:

- \* evolution of logarithms in PDF: they are resummed
- \* available phase space
- \*\* approximation by large logarithm



## FOUR-FLAVOR SCHEME

Campbell, RF, Maltoni, Tramontano (2009)

- <sup></sup> We the 4-flavor (2 → 3) process as the Born and calculate NLO
  - Much harder calculation due to extra mass and extra parton



- Spectator b for the first time at NLO
- Process implemented in the MCFM-v5.7 parton-level NLO code

<sup>∗</sup> Starting point for future NLO+PS beginning at (2 → 3)



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#### ALSO BSM



% squark-gluino associated production

% real emission corrections included, but virtual correction not (yet)





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## **TO CONCLUDE**

- Sor any QCD NLO computation (SM & BSM) MadFKS takes care of:
  - Senerating the Born, real emission, subtraction terms, phase-space integration and overall management of symmetry factors, subprocess combination etc.
- Substitution Structure Structure
- With the shower subtraction terms, interface to showers to generate automatically unweighted events with NLO precision is working with Herwig and work in progress with Pythia and Herwig++
- First physics results at NLO are being produced within the MadGraph/MadEvent framework using the MadFKS code
  Rikkert Frederix, September 14, 2010